Source imaging in RHIC $\sqrt{s_{NN}}$ =200 GeV Au+Au collisions



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INVESTICE DO ROZVOJE VZDĚLÁVÁNÍ

RHIC/PHENIX



RHIC/STAR

RHIC: Broad physics program

- Heavy ions: Au+Au, Cu+Cu, U+U, $\sqrt{s_{NN}}=7.7-200 \text{ GeV}$
- Polarized protons up to $\sqrt{s} = 510 \text{ GeV}$
- Asymmetric systems (d+Au, Cu+Au)

The Solenoidal Tracker at RHIC

- Time Projection Chamber
 - ID via energy loss (dE/dx)
 - Momentum (p)
- Full azimuth coverage

Uniform acceptance for different energies and particles

Both experiments: continuous improvements

to meet evolving physics goals New subsystems for enhanced PID, rates, coverage...



Femtoscopy

Boson emitting source:

Symmetric two-boson wave function

$$N_1(k_1) = \int S(x_1, k_1) |\Psi_1|^2 dx_1$$

$$N_2(k_1, k_2) = \int S(x_1, k_1) S(x_2, k_2) |\Psi_{1,2}|^2 dx_1 dx_2$$



Bose-Einstein Correlation / Hanbury-Brown–Twiss effect

Info about shape and evolution of the particle emitting source

Correlation function:

$$C_2(k_1, k_2) = \frac{N_2(k_1, k_2)}{N_1(k_1)N_1(k_2)} \simeq 1 + \left|\frac{\tilde{S}(q, K)}{\tilde{S}(0, K)}\right|^2 \quad \frac{\tilde{S}(q, K)}{q = k_1 - k_2, K = 0.5(k_1 + k_2)}$$

- Final state interactions
 - Compensating the Coulomb force

$$C_0(q) = C_{\text{raw}}(q) K_{\text{coulomb}}^{-1}$$

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- Strong FSI …
- Solving for the source is difficult \rightarrow assumptions

Gaussian source approximation

If the source is approximated with Gaussian:

$$S(x) \sim \exp\left(-\frac{r_x^2}{2R_x^2} - \frac{r_y^2}{2R_y^2} - \frac{r_z^2}{2R_z^2}\right)$$



C(q)

Then the correlation function is also Gaussian:

$$C(q) - 1 \sim \exp\left(-q_x^2 R_x^2 - q_y^2 R_y^2 - q_z^2 R_z^2\right)$$

- These radii are the so-called HBT radii
- Often specified in the LCMS system (not invariant)
 - Out: direction of the mean transverse momentum of the pair
 - Side: orthogonal to out
 - Long: beam direction

$$C(q) = 1 + \lambda \exp\left(-q_o^2 R_o^2 - q_s^2 R_s^2 - q_l^2 R_l^2\right)$$

Do not necessarily reflect the geometrical size

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Source imaging

Physics in shape: dynamics, resonance decays, rescattering...

Koonin-Pratt equation (1D)

$$C(q) - 1 = 4\pi \int dr r^2 \frac{K(q, r)S(r)}{S(r)}$$

- Imaging: Obtain S(r) directly
 - No assumptions for the shape of source
 - Kernel includes all interactions (QM, FSI)
- Numerical inversion of the equation
 - No analytical solution, hence some limitations and approximations (integral cutoff, finite resolution ...)
 - Assumptions (e.g. weak dependence in single particle sources)
 - Needs statistics, stability is a question

D. A. Brown, P. Danielewicz, Phys.Lett. B398, 252 (1997)



Pion images

PHENIX Year 2002 data

- low $k_T = (p_{T,1} + p_{T,2})/2$
- C from data ~ C restored from image
 → Imaging process can be trusted
- A heavy, non-Gaussian tail is present in the 1D pion source
- Several interpretations suggested
 - Non-zero emission duration
 - Anomalous diffusion due to rescattering in the hadronic phase
 - Conribution of long-lived resonance decays



Rescattering

srdist0 Source function from HR model, 0-20%, 0.20-0.36GeV Entries 670343 Hadronic Rescattering Code Mean 13.27 S(r) Core-Core 6.044 * Core-ω Simple but smart cascade model Core-Halo 10⁻¹ ω-ω Only a few resonances ω-Halo △ Halo-Halo (ρ, Δ, Κ* ; ω ; η, η', Φ, Λ) → Sum of all S(r) Causality kept in all scatterings p-dependent cross sections 10^{-2} Gauss Shown to be working Describes spectra, v₂, HBT radii for both SPS and RHIC **PH** Insensitive to initial conditions 10^{-3} Similar predictions to exact hydro 30 40 r[fm] Sensitive to PID (π , K, p) Csanád, Csörgő, Nagy, Braz.J.Phys. 37 (2007) T. J. Humanic, Int. J. Mod. Phys. E 15 (2006)

- HRC able to describe the observed 1D pion source
 Note: model limitations lead to breakdown for higher k_T bin (not shown)
- Underlying mechanism: anomalous diffusion
 - Diffusion with fixed mean free path: Central Limit Theorem \rightarrow Gaussian distrib.
 - Expanding system, changing x-section: Gnedenko–Kolmogorov → Lévy distrib.

Resonances



R.V. (PHENIX), WWND 2007 proc. [arXiv:0706.4409]

- Single FO with resonances: also yields a relatively good description
 - Parameters tuned for PHENIX HBT

Note: model limitations cause problems at $r \rightarrow 0$ (not shown)

- Underlying mechanism: many long lived resonances
 - Different contributions die out gradually
 - Continuously increasing mean lifetimes provide a random variable with timedependent mean and variance → similar effect to anomalous diffusion

Kaons: A cleaner probe

- Less feed-down, less rescattering
 - Interpretation more straightforward
 - More difficult due to ~10 less statistics
- PHENIX 1D Kaon source: an even larger non-Gaussian component
 - Seemingly favors rescattering explanation against resonances
- Interpretation caveat: wide k_T (N_{part}) bin
 - Different k_T → Gaussians with different radii → convolute to non-Gaussian





Kaons: STAR vs. PHENIX



STAR preliminary 1D source in narrow k_T bin consistent with Gaussian

0.20<k_T<0.36 GeV , compared to 0.3<k_T<0.9 GeV

3D source shapes

Expansion of R(q) and S(r) in Cartesian Harmonic basis

Danielewicz and Pratt, Phys.Lett. B618, 60 (2005)

$$R(\mathbf{q}) = \sum_{l} \sum_{\alpha_{1}...\alpha_{l}} R_{\alpha_{1}...\alpha_{l}}^{l}(q) A_{\alpha_{1}...\alpha_{l}}^{l}(\Omega_{q}) \quad (1) \qquad \qquad \alpha_{i} = \mathbf{x}, \mathbf{y} \text{ or } \mathbf{z} \\ \mathbf{x} = \text{out-direction} \\ \mathbf{y} = \text{side-direction} \\ \mathbf{y} = \text{side-direction} \\ \mathbf{z} = \text{long-direction} \end{cases}$$

3D Koonin-Pratt:
$$R(\mathbf{q}) = C(\mathbf{q}) - 1 = 4\pi \int dr^3 K(\mathbf{q}, \mathbf{r}) S(\mathbf{r})$$
 (3)

Plug (1) and (2) into (3)
$$\Rightarrow R^{l}_{\alpha_{1}...\alpha_{l}}(q) = 4\pi \int dr^{3} K_{l}(q,r) S^{l}_{\alpha_{1}...\alpha_{l}}(r)$$
 (4)

Invert (1)
$$\Rightarrow$$
 $R_{\alpha_{1}...\alpha_{l}}^{l}(q) = \frac{(2l+1)!!}{l!} \int \frac{d\Omega_{\mathbf{q}}}{4\pi} A_{\alpha_{1}...\alpha_{l}}^{l}(\Omega_{\mathbf{q}}) R(\mathbf{q})$
Invert (2) \Rightarrow $S_{\alpha_{1}...\alpha_{l}}^{l} = \frac{(2l+1)!!}{l!} \int \frac{d\Omega_{\mathbf{r}}}{4\pi} A_{\alpha_{1}...\alpha_{l}}^{l}(\Omega_{\mathbf{r}}) S(\mathbf{r})$

3D images: figuring out more about freeze-out dynamics!

3D pion correlation and source

PHENIX, PRL100, 232301 (2008) (a) $S^0 + S_{x2}^2 + S_{x4}^4 + S_{x6}^6$ (d) $C^0 + C_{x2}^2 + C_{x4}^4 + C_{x6}^6$ Elongated source in "out" direction 1.1 Data 2(r_x) Imaging Moments up to the 6th order 1.05 (G Hump Fit Au+Au Elliptic and non-Gaussian 10 √s_{NN}=200GeV 1D radii determined by side/long 0<cen<20 % (b) $S^0 + S_{y2}^2 + S_{y4}^4 + S_{y6}^6$ (e) $C^0 + C_{y2}^2 + C_{y4}^4 + C_{y6}^6$ $(x \ 10^{-7} \ fm^{-3})$ 0 $S(r_y)$ 0 $S(r_y)$ $\pi^{+}\pi^{+} \& \pi^{-}\pi^{-}$ Shape described well by a double 0.20<p_<0.36 GeV/c .2 C(q_y -0.35<y_-y_<0.35 Gaussian, or "hump" function (c) $S^0 + S_{z2}^2 + S_{z4}^4 + S_{z6}^6$ $(f) C^0 + C_{z2}^2 + C_{z4}^4 + C_{z6}^6$ $S^{H}(r_{x}, ry, rz) = e^{-F_{s} \left[\left(\frac{r_{x}}{2R_{xs}} \right)^{2} + \left(\frac{r_{y}}{2R_{ys}} \right)^{2} + \left(\frac{r_{z}}{2R_{zs}} \right)^{2} \right]}$ 10 S(r_z) .2 C(q_z) $K_{s} = \frac{1}{1 + (r/r_{0})^{2}}, \quad F_{l} = 1 - F_{s}$ 10 20 20 30 40 0 10 40 r (fm) q (MeV/c) Source profiles **Correlation profiles** $S(r_x) \equiv C(r_x, 0, 0)$ $C(q_x) \equiv C(q_x, 0, 0)$ $\mathbf{S}(\mathbf{r}_{v}) \equiv \mathbf{C}(0,\mathbf{r}_{v},0)$ $C(q_v) \equiv C(0,q_v,0)$ $\mathbf{S}(\mathbf{r}_{7}) \equiv \mathbf{C}(0,0,\mathbf{r}_{7})$ $C(q_z) \equiv C(0,0,q_z)$

3D pion images: STAR vs. PHENIX

Elongated source in "out" direction

- Moments up to the 6th order
- Elliptic and non-Gaussian
- 1D radii determined by side/long
- STAR and PHENIX measurements are consistent
 - Two different detectors with different properties and acceptance
 - Good agreement with same cuts
 - Attests to the reliability of results



Source profiles

3D pion images vs. B/W model

Elongated source in "out" direction

- Moments up to the 6th order
- Elliptic and non-Gaussian
- 1D radii determined by side/long
- Therminator B/W model description
 - Iff resonance contributions ON, and
 - Iff non-zero emission duration
 Δτ~2 fm/c

THERMINATOR Blast-Wave model

- Expansion: $v_r(\rho) = (\rho/\rho_{max})/(\rho/\rho_{max} + v_t)$.
- Thermal emission at proper time τ, ρ=ρ_{max}.
- Freeze-out occurs at $\tau = \tau_0 + a\rho$.
- LAB emission time $t^2 = (\tau_0 + a\rho)^2 + z^2$.
- Finite emission duration $\Delta \tau$ in lab frame



Source profiles

3D kaon correlation and source



 3D Kaon correlation moments and profiles consistent with Gaussian

See talk of **M. Girard** for STAR



- Source Gaussian fit shown
- Uncertainties include shape assumption (error dominated low statistics)

3D kaon source: Model comparison

Therminator B/W model

- Kaons: Instant freeze-out Δτ = 0 (contrary to pions!)
- Parameters tuned for STAR kaons!
- Resonances are needed

Hydrokinetic model

- Consistent in "side"
- Slightly more tail (r>15fm) in "out" and "long"

Hybrid Hydrokinetic Model (hHKM)

PRC81, 054903 (2010)

- Glauber initial conditions
- Pure hydro expansion
- Hadronic cascade with UrQMD Gets many RHIC observables right



Therminator: Kisiel, Taluc, Broniowski, Florkowski, Comput. Phys. Commun. 174 (2006) 669.

HKM data: Shapoval, Sinyukov, Karpenko , arXiv:1308.6272 [hep-ph]

Summary

Long-range component in pion source of RHIC $\sqrt{s_{NN}}$ =200 GeV Au+Au collisions observed by PHENIX, confirmed by STAR

- Source is elongated in the "out" direction
- "long" and "side" are similar

Contrary to PHENIX, STAR does not observe the heavy tail in source functions for (central) Kaons

The k_T range for PHENIX, however, is 9x wider than for STAR

No one over-simplified model explains the observables simultaneously – multiple contributions

- Resonance contributions are required to describe the source shapes
- Kaons and pions may also be subject of different freeze-out dynamics
- Rescattering taken into account by successful models



Backup slides follow...

Imaging: Inversion procedure

$$C(q) = 4\pi \int dr r^2 K(q, r) S(r)$$

$$S(r) = \sum_{j} S_{j} \cdot B_{j}(r)$$
 Expansion in B-spline basis

$$C_{i}^{Th}(q) = \sum_{j} K_{ij} \cdot S_{j}$$

$$K_{ij} = \int dr \cdot K(q, r) B_{j}(r) \qquad \chi^{2} = \frac{\left(C_{i}^{Expt}(q) - \sum_{j} K_{ij} \cdot S_{j}\right)^{2}}{\Delta^{2} C_{i}(q)^{Expt}}$$

Freeze-out occurs after last scattering Hence only Coulomb & BE effect included in kernel

PHENIX pion moments



PHENIX, PRL100, 232301 (2008)

Peripheral pions in STAR



NA49 pions in Pb+Pb - correlation



NA49 pions in Pb+Pb - sources



Kaon source shape



Fit to correlation moments by kT



Transverse mass dependence

- Radii: rising trend at low m_T
 - Strongest in "long"
- Buda-Lund model
 - Perfect hydrodynamics, inherent m_T-scaling
 - Works perfectly for pions
 - Deviates from kaons in the "long" direction in the lowest m_T bin
- HKM (Hydro-kinetic model)
 - Describes all trends
 - Some deviation in the "out" direction
 - Note the different centrality definition



Buda-Lund: M. Csanád, arXiv:0801.4434v2 HKM: PRC81, 054903 (2010)

Radii vs. mT, pion, kaon

- STAR kaons
- PHENIX pions +,-
- Buda-Lund
- HKM



HKM: PRC81, 054903 (2010)

Kaon femtoscopy analyses

Au+Au @ $\sqrt{s_{NN}}=200 \text{ GeV}$ Mid-rapidity |y|<0.5

- 1. Source shape: 20% most central Run 4: 4.6 Mevts, Run 7: 16 Mevts
- 2. m_T-dependence: 30% most central Run 4: 6.6 Mevts







PID cut applied

- 1. Source shape analysis
 - dE/dx: nσ(Kaon)<2.0 and nσ(Pion)>3.0 and nσ(electron)>2.0 nσ(X) :deviation of the candidate dE/dx from the normalized distribution of partice type X at a given momentum
 - 0.2 < p_T < 0.4 GeV/c
- 2. m_T -dependent analysis
 - -1.5< nσ(Kaon)<2.0

-0.5< nσ(Kaon)<2.0



