

Charge:

- Identify the most important precision observables that can reveal deviations from the standard model.
- Identify the thresholds of precision that needs to be achieved for each of these observables in order to be definitively sensitive to new physics.
- Study the precision that can be achieved at each proposed facility on these observables, and ask what machine and detector parameters are required to reach the discovery threshold.
- Identify the calculational tools needed to predict standard model rates and distributions in order to perform these measurements at the required precision.

Conveners: Ashutosh Kotwal, Michael Schmitt, DW

Webpage:

<http://www.snowmass2013.org/tiki-index.php?page=Precision+Study+of+Electroweak+Interactions>

- Check of the consistency of the SM by comparing direct with indirect measurements of model parameters, e.g., m_{top} , M_W , $\sin^2 \theta_{eff}$, M_H .
Ayres Freitas (SM), Sven Heinemeyer (MSSM), Alessandro Vicini (M_W)
NC and CC DY: Simone Alioli, Emmanuele Re, Alessandro Vicini
- Search for indirect signals of Beyond-the-SM (BSM) physics in form of small deviations from SM predictions, yielding exclusions of, and constraints on, BSM scenarios using global fits of EW precision observables:
Sven Heinemeyer (MSSM), Jens Erler and Paul Langacker (Z'), GFITTER (global fit and S,T,U)
- Sensitive probe of proton structure (PDFs): Alessandro Vicini, Juan Rojo

- EW gauge boson pair and triple production directly probes the non-abelian gauge structure of the SM.
- Search for non-standard gauge boson self couplings allowed by Lorentz and gauge invariance provide a unique indirect way to look for signals of new physics in a model-independent way.
- Improved constraints on anomalous triple-gauge boson couplings (TGCs) and quartic couplings (QGCs) probe scales of new physics in the multi-TeV range.
- Important backgrounds to Higgs and BSM searches.
- EFT approach (and relation to anomalous coupling parameters) implemented in
Madgraph: Olivier Mattelaer and Celine Degrande
Oscar Eboli (combined fit to LHC EW and Higgs data)
Michael Rauch, Barbara Jaeger (multiboson predictions)
WHIZARD: Jürgen Reuter

- There has been tremendous effort and a lot of progress in calculating higher order EW corrections and in understanding enhanced logarithmic corrections of weak origin.
 NLO EW: $pp, p\bar{p} \rightarrow W; Z \rightarrow l\nu; l^+l^-; VV; Wj; Zj \rightarrow \nu lj; l^+l^-j; t\bar{t}$; single top, and $b\bar{b}, jj$ (weak); and for dominant Higgs production processes, e.g., $gg \rightarrow H; W/ZH; VBF$.
 EW Sudakov logarithms $\alpha_w^l \log^n(Q^2/M^2)$, $n \leq 2l$: results available for W, Z production, $VV, t\bar{t}, bb, cc, jj, VBF$.
 Photon-induced processes (QED PDFs), real W, Z radiation, multiple photon radiation, interplay of QCD/EW corrections, ...
- Their significance strongly depends on details of experimental definition of observable, specifically which kinematic regime is probed.
- Implementation of EW corrections in publicly available MCs in progress, enabling studies of higher-order mixed QED-QCD effects, has only been done for selected processes.
- Important and difficult task: reliable estimates of theoretical uncertainties due to missing higher-order corrections (work in progress for W/Z observables)

EW Sudakov logarithms $\alpha_w^l \log^n(Q^2/M^2)$, $n \leq 2l$

In the high-energy limit, $\frac{Q}{M_{W,Z}} \rightarrow \infty$, EW Sudakov logarithms have been studied in analogy to soft/collinear logarithms in QED,QCD.

- 1-loop: LL and NLL are universal and factorize [Denner, Pozzorini \(2001\)](#)
- Beyond 1-loop: Resummation techniques based on IR evolution equations (IREE) or SCET yield results up to NNLL ($\ln^n(\frac{s}{M_W^2})$, $n = 2, 3, 4$).
 - IREE: EW theory splits into symmetric $SU(2) \times U(1)$ ($M_W = M_Z = M_\gamma = M$ for $\mu > M$) and QED regime and effect of EW symmetry breaking neglected. [Fadin, Lipatov, Martin, Melles \(2000\)](#)
 - SCET: At $\mu = Q$ match full theory to SCET($M = 0$), evolve to $\mu = M$ SCET($M \neq 0$), match to SCET with no gauge bosons.
 - SCET and IREE Sudakov form factors are equivalent. [Chiu, Golf, Kelley, Manohar \(2008\)](#); [Chiu, Fuhrer, Hoang, Kelley, Manohar \(2009\)](#); [Chiu, Fuhrer, Kelley, Manohar \(2010\)](#), [Fuhrer et al \(2011\)](#)

Resummation results at LL and NLL confirmed by explicit diagrammatic one-loop and two-loop calculations.

[Melles \(2000\)](#), [Hori et al \(2000\)](#), [Beenakker, Werthenbach \(2000,2002\)](#), [Pozzorini \(2004\)](#); [Feucht et al \(2003,2004\)](#); [Jantzen et al \(2005,2006\)](#); [Denner et al \(2003,2008\)](#)

- Results available for hadronic cross sections for W, Z production, $VV, t\bar{t}, bb, cc, jj, \text{VBF}$.
- Best studied so far: $f\bar{f} \rightarrow f\bar{f}$
 - up to N³LL for massless fermions ($a = \frac{\alpha}{4\pi s_W^2}, L = \log(s/M_W^2)$):

$$\frac{\delta\sigma(e^+e^- \rightarrow q\bar{q})(s)}{\sigma_{LO}} = a(-2.18L^2 + 20.94L - 35.07) +$$

$$+ a^2(2.79L^4 - 51.98L^3 + 321.20L^2 - 757.35L)$$

$$\approx 2.4\% - 0.4\% \text{ at } 2 \text{ TeV}$$

Note: only LL at 2-loop: +3%

[Jantzen, Kühn, Penin, Smirnov, hep-ph/0509157](#)

- up to NNLL for massive fermions [Denner, Jantzen, Pozzorini \(2008\)](#).
- See also SCET results by [Chiu et al, \(2008\)](#).

$pp \rightarrow Z, \gamma \rightarrow \mu^+ \mu^-$ at $M_{ll} > 2$ TeV: S.Dittmaier and M.Huber, arXiv:0911.2329 [hep-ph].

$q\bar{q}$, weak	$\delta\sigma/\sigma_B[\%]$	QCD	$\delta\sigma/\sigma_B[\%]$
$q\bar{q}$, QED	-11.12	$\gamma\gamma + \gamma q(\bar{q})$	-11.93
$q\bar{q}$, multi- γ	-12.08	h.o. weak	+7.28
h.o. Sudakov	+0.54		-0.32
	+3.38		

$pp \rightarrow W^+ W^-$ at $M_{WW} > 2$ TeV: Bierweiler *et al*, arXiv:1208.3147 [hep-ph].

	$\delta\sigma/\sigma_B[\%]$		$\delta\sigma/\sigma_B[\%]$
$q\bar{q}$, EW	-31.3	QCD	+22.8
$\gamma\gamma$	+21.6	WWV	+4.9