Tests of Standard Model limits in the light flavor sector

Dinko Počanić

University of Virginia

18 December 2013



5th International Workshop on High Energy Physics in the LHC Era Valparaíso, Chile 16–20 Dec 2013

Light flavors!





Light flavors! Isn't that low energy stuff!?





- Light flavors! Isn't that low energy stuff!?
- What are low energies doing at a meeting on LHC-era HE physics?





- Light flavors! Isn't that low energy stuff!?
- ► What are low energies doing at a meeting on LHC-era HE physics?
- Don't we already know all that's of any use about the low-E, light-flavor sector?





- Light flavors! Isn't that low energy stuff!?
- ► What are low energies doing at a meeting on LHC-era HE physics?
- Don't we already know all that's of any use about the low-E, light-flavor sector?

Not exactly.

The parameter domain of the SM is large — it leaves much room for low-energy tests; they complement HE searches and tests.





- Light flavors! Isn't that low energy stuff!?
- What are low energies doing at a meeting on LHC-era HE physics?
- Don't we already know all that's of any use about the low-E, light-flavor sector?

Not exactly.

The parameter domain of the SM is large — it leaves much room for low-energy tests; they complement HE searches and tests.

 There is vigorous activity on low energy SM tests and searches for BSM physics.





- Light flavors! Isn't that low energy stuff!?
- ► What are low energies doing at a meeting on LHC-era HE physics?
- Don't we already know all that's of any use about the low-E, light-flavor sector?

Not exactly.

The parameter domain of the SM is large — it leaves much room for low-energy tests; they complement HE searches and tests.

- There is vigorous activity on low energy SM tests and searches for BSM physics.
- Rather than giving a comprehensive review, this talk will focus on a small subset: π, μ and n decays to introduce and illustrate the field.

D. Počanić (UVa)

Why low energy tests?



Outline

Overview of π , μ decays and recent experiments

Radiative muon decay, $\mu^+ \to {\rm e}^+ \nu \bar{\nu} \gamma$ (new result)

The π_{e2} decay, $\pi^+ \rightarrow e^+ \nu_e$ (current work)

Correlations in neutron beta decay (experiment being built)





Known and measured pion and muon decays





SM tests with light particles:

Overview of π , μ decays



Recent measurements of π , μ allowed decay

► $\pi^+ \rightarrow \pi^0 e^+ \nu_e$ PIBETA ('99–'01) o SM checks related to CKM unitarity ► $\pi^+ \rightarrow e^+ \nu_e \gamma$ (or $e^+ e^-$) PIBETA ('99–'04), PEN ('06–'10) $\circ \mathbf{F}_{\mathbf{A}}/\mathbf{F}_{\mathbf{V}}, \pi$ polarizability (χ PT calibration) o tensor coupling besides V - A (?) o departures from $\mathbf{V} - \mathbf{A}$ in \mathcal{L}_{weak} ► $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu \gamma$ (or $e^+ e^-$)PIBETA ('04), PEN ('06-'10) o departures from $\mathbf{V} - \mathbf{A}$ in $\mathcal{L}_{\text{weak}}$ $\begin{array}{l} \bullet \ \pi^+ \to {\bf e}^+ \nu_{\bf e} \\ \circ \ {\bf e} \cdot \mu \text{ universality} \end{array} \begin{cases} \begin{array}{l} \mathsf{PEN} \ ('06-'10) \\ \mathsf{PiENu} \ ('06-) \end{array} \end{cases}$ \circ **P**, **S** coupling besides **V** – **A** $\circ \nu$ sector anomalies, Majoron searches, \mathbf{m}_{h+} , PS I-q's, V I-q's, ... o search for signs of SUSY (MSSM)

🛄 D. Počanić (UVa)



The PIBETA/PEN apparatus

- stopped π^+ beam
- active target counter
- 240-detector, spherical pure Csl calorimeter
- central tracking
- beam tracking
- digitized waveforms
- stable temp./humidity





SM tests with light particles:

Apparatus overview



Prior results

- Pion beta decay: $\pi^+
 ightarrow \pi^0 \mathrm{e}^+
 u_\mathrm{e}$
- Radiative pion decay: $\pi^+
 ightarrow {
 m e}^+
 u_{
 m e} \gamma$

(the PIBETA experiment)



SM tests with light particles:



Quark-Lepton (Cabibbo) Universality

The basic weak-interaction V-A form (e.g., μ decay):

$$\mathcal{M} \propto \langle \mathbf{e} | \mathbf{l}^lpha |
u_\mathbf{e}
angle o ar{\mathbf{u}}_\mathbf{e} \gamma^lpha (1-\gamma_5) \mathbf{u}_
u$$

is replicated in hadronic weak decays

 $\mathcal{M} \propto \langle p | h^\alpha | n \rangle \rightarrow \bar{u}_p \gamma^\alpha (G_V - G_A \gamma_5) u_n \quad \text{with} \quad G_{V,A} \simeq 1 \; .$

Departure from $G_V = 1$ (CVC) comes from weak quark (Cabibbo) mixing: $G_V = G_\mu \cos \theta_C (= G_\mu V_{ud}) \quad \cos \theta_C \simeq 0.97$

3 q generations lead to the Cabibbo-Kobayashi-Maskawa (CKM) matrix (1973): $\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$

CKM unitarity cond.: $\Delta V^2 = 1 - (|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2) \stackrel{?}{=} 0$, stringently tests the SM. Until 2004 appeared violated by $\sim 3\sigma!$



PIBETA result for $\pi^+ \rightarrow \pi^0 e^+ \nu$ (π_β) decay [PRL **93**, 181803 (2004)]

$$\begin{split} \mathcal{B}_{\pi\beta}^{\text{exp-t}} &= [1.040 \pm 0.004 \, (\text{stat}) \pm 0.004 \, (\text{syst})] \times 10^{-8} \,, \\ \mathcal{B}_{\pi\beta}^{\text{exp-e}} &= [1.036 \pm 0.004 \, (\text{stat}) \pm 0.004 \, (\text{syst}) \pm 0.003 \, (\pi_{\text{e2}})] \times 10^{-8} \,, \end{split}$$

McFarlane et al. [PRD 1985]: $B = (1.026 \pm 0.039) \times 10^{-8}$

SM Prediction (PDG):

$$B = 1.038 - 1.041 \times 10^{-8}$$
 (90% C.L.)
 $(1.005 - 1.007 \times 10^{-8}$ excl. rad. corr.)

 \Rightarrow Most sensitive test of CVC/radiative corr. in a meson to date!

PDG 2012: $V_{ud} = 0.97425(22)$ PIBETA: $V_{ud} = 0.9748(25)$ or $V_{ud} = 0.9728(30)$.











SM tests with light particles:

Pion beta decay





Summary of PIBETA results on $\pi
ightarrow e
u \gamma$ [PRL 103, 051802 (2009)]

$F_V = 0.0258 \pm 0.0017$	(14×)
${\sf F}_{\sf A} = 0.0119 \pm 0.0001^{ m exp}_{({\sf F}_{ m V}^{ m CVC})}$	(16×)
$a=0.10\pm0.06~(q^2~\text{dep}~\text{of}~\text{F}_{V})$	(∞)
$-5.2 imes 10^{-4} < F_T < 4.0 imes 10^{-4}$	90 % C.L.

 $\mathsf{B}_{\pi_{\mathrm{e}2\gamma}}(\mathsf{E}_{\gamma}>10\,\mathrm{MeV}, heta_{\mathrm{e}\gamma}>40^{\circ})=73.86(54) imes10^{-8}~(17 imes)$



SM tests with light particles:

Radiative pion decay



Summary of PIBETA results on $\pi \rightarrow e \nu \gamma$ [PRL 103, 051802 (2009)]



 $\mathsf{B}_{\pi_{\mathrm{e}2\gamma}}(\mathsf{E}_{\gamma}>10\,\mathrm{MeV}, heta_{\mathrm{e}\gamma}>40^{\circ})=73.86(54) imes10^{-8}~(17 imes)$

Above results will be improved with the new PEN data analysis.



SM tests with light particles:

Radiative pion decay



Summary of PIBETA results on $\pi \rightarrow e \nu \gamma$ [PRL 103, 051802 (2009)]



$$\mathsf{B}_{\pi_{\mathrm{e}2\gamma}}(\mathsf{E}_{\gamma}>10\,\mathrm{MeV}, heta_{\mathrm{e}\gamma}>40^{\circ})=73.86(54) imes10^{-8}~(17 imes)$$

Above results will be improved with the new PEN data analysis.

At L.O. $(l_9 + l_{10})$, F_V is related to pion polarizability and π^0 lifetime $\alpha_E^{\text{LO}} = -\beta_M^{\text{LO}} = (2.783 \pm 0.023_{\text{exp}}) \times 10^{-4} \text{ fm}^3$ $\tau_{\pi^0} = (8.5 \pm 1.1) \times 10^{-17} \text{ s}$ $\begin{cases} \text{current PDG avg: } 8.52 (12) \\ \text{PrimEx PRL '10: } 8.32 (23) \end{cases}$





Experimental History of Pion F_A and F_V





SM tests with light particles:

Radiative pion decay



Radiative muon decay:

$$\mu^+
ightarrow {
m e}^+
u ar{
u} \gamma$$

PIBETA: 2004 runs (PEN: 2008–2010 runs)



SM tests with light particles:



Radiative muon decay, $\mu^+ \rightarrow e^+ \nu \bar{\nu} \gamma$, (new analysis of 2004 data)





SM tests with light particles:

Radiative muon decay



16/37

RMD preliminary results, cont'd.

Preliminary result for RMD branching ratio (thesis E. Munyangabe):

$$\begin{split} B_{\text{exp}} &= 4.365 \, (9)_{\text{stat.}} \, (42)_{\text{syst.}} \times 10^3 \,, \qquad \boxed{\textbf{29}\times} \\ B_{\text{SM}} &= 4.342 \, (5)_{\text{stat-MC}} \times 10^3 \qquad (\text{for } E_{\gamma} > 10 \, \text{MeV}, \ \theta_{\text{e}\gamma} > 30^\circ) \end{split}$$





RMD preliminary results, cont'd.

Preliminary result for RMD branching ratio (thesis E. Munyangabe):

$$\begin{split} & \textit{B}_{\text{exp}} = 4.365\,(9)_{\text{stat.}}\,(42)_{\text{syst.}} \times 10^3 \,, \qquad \boxed{\textbf{29}\times} \\ & \textit{B}_{\text{SM}} = 4.342\,(5)_{\text{stat-MC}} \times 10^3 \qquad (\text{for } E_{\gamma} > 10\,\text{MeV}, \; \theta_{\text{e}\gamma} > 30^\circ) \end{split}$$



NB: preliminary results!

Analysis of PS subset: 13 MeV $< E_{\gamma} < 45$ MeV, and 10 MeV $< E_{e^+} < 43$ MeV, yields

$$\begin{split} \bar{\eta} &= 0.006\,(17)_{\rm stat.}\,(18)_{\rm syst.}, \text{ or} \\ \bar{\eta} &< 0.028 \quad (68\%{\rm CL})\,. \end{split}$$

 $\sim 4 \times$ better than best previous experiment (Eichenberger et al, 84).





What to do with Michel parameters?

For
$$\mu \to e \nu_{\mu} \bar{\nu}_{e} \gamma$$
: $\left(x = \frac{E_{e}}{E_{max}} \text{ and } y = \frac{E_{\gamma}}{E_{max}}\right)$

$$\frac{\mathrm{d}^{3}B(x,y,\theta)}{\mathrm{d}x\,\mathrm{d}y\,2\pi\,\mathrm{d}(\cos\theta)} = f_{1}(x,y,\theta) + \bar{\eta}f_{2}(x,y,\theta) + (1-\frac{4}{3}\rho)f_{3}(x,y,\theta)$$





What to do with Michel parameters?

For
$$\mu \to e \nu_{\mu} \bar{\nu}_{e} \gamma$$
: $\left(x = \frac{E_{e}}{E_{max}} \text{ and } y = \frac{E_{\gamma}}{E_{max}}\right)$

$$\frac{\mathrm{d}^{3}B(x,y,\theta)}{\mathrm{d}x\,\mathrm{d}y\,2\pi\,\mathrm{d}(\cos\theta)} = f_{1}(x,y,\theta) + \bar{\eta}f_{2}(x,y,\theta) + (1-\frac{4}{3}\rho)f_{3}(x,y,\theta)$$

$$\rho = \frac{3}{4} - \frac{3}{4} \left[|g_{LR}^{V}|^{2} + |g_{RL}^{V}|^{2} + 2|g_{LR}^{T}|^{2} + 2|g_{RL}^{T}|^{2} + \Re(g_{RL}^{S}g_{RL}^{T*} + g_{LR}^{S}g_{LR}^{T*}) \right] \stackrel{\text{SM}}{\equiv} \frac{3}{4},$$

$$\vec{\eta} = \left(|g_{RL}^{V}|^{2} + |g_{LR}^{V}|^{2} \right) + \frac{1}{8} \left(|g_{LR}^{S} + 2g_{LR}^{T}|^{2} + |g_{RL}^{S} + 2g_{RL}^{T}|^{2} \right) + 2 \left(|g_{LR}^{T}|^{2} + |g_{RL}^{T}|^{2} \right)$$

$$\stackrel{\text{SM}}{\equiv} \quad \mathbf{0} \,.$$

18/37

SM tests with light particles:

D. Počanić (UVa)

What to do with Michel parameters?

For
$$\mu \to e \nu_{\mu} \bar{\nu}_{e} \gamma$$
: $\left(x = \frac{E_{e}}{E_{max}} \text{ and } y = \frac{E_{\gamma}}{E_{max}}\right)$

$$\frac{\mathrm{d}^{3}B(x,y,\theta)}{\mathrm{d}x\,\mathrm{d}y\,2\pi\,\mathrm{d}(\cos\theta)} = f_{1}(x,y,\theta) + \bar{\eta}f_{2}(x,y,\theta) + (1-\frac{4}{3}\rho)f_{3}(x,y,\theta)$$

$$\boldsymbol{\rho} = \frac{3}{4} - \frac{3}{4} \Big[|g_{LR}^V|^2 + |g_{RL}^V|^2 + 2|g_{LR}^T|^2 + 2|g_{RL}^T|^2 + \Re(g_{RL}^S g_{RL}^{T*} + g_{LR}^S g_{LR}^{T*}) \Big] \quad \stackrel{\text{SM}}{\equiv} \quad \frac{3}{4} \,,$$

$$\bar{\boldsymbol{\eta}} = \left(|g_{RL}^{V}|^{2} + |g_{LR}^{V}|^{2} \right) + \frac{1}{8} \left(|g_{LR}^{S} + 2g_{LR}^{T}|^{2} + |g_{RL}^{S} + 2g_{RL}^{T}|^{2} \right) + 2 \left(|g_{LR}^{T}|^{2} + |g_{RL}^{T}|^{2} \right)$$

$$\stackrel{\text{SM}}{\equiv} \quad \boldsymbol{0} \,.$$

Combined with η , δ , ρ , parameters of OMD (TWIST), a global fit will yield model-independent limits on non–(V – A) couplings.



SM tests with light particles:



The $\pi_{
m e2}$ decay: $\pi^+
ightarrow {
m e}^+
u$

Primary motivation for PEN (data runs 2008–10)



SM tests with light particles:

The π_{e2} decay



- Early evidence for V A nature of weak interaction.
- Modern SM calculations:

$${\cal R}^{\pi}_{{
m e}/\mu} = rac{ \Gamma(\pi o e ar
u(\gamma)) }{ \Gamma(\pi o \mu ar
u(\gamma))}_{{
m CALC}} =$$

 $\begin{cases} 1.2352 (5) \times 10^{-4} & \text{Marciano and Sirlin, [PRL$ **71** $(1993) 3629]} \\ 1.2354 (2) \times 10^{-4} & \text{Finkemeier, [PL B$ **387** $(1996) 391]} \\ 1.2352 (1) \times 10^{-4} & \text{Cirigliano and Rosell, [PRL$ **99** $(2007) 231801]} \end{cases}$





- Early evidence for V A nature of weak interaction.
- Modern SM calculations:

$${\cal R}^{\pi}_{{
m e}/\mu} = rac{\Gamma(\pi o e ar{
u}(\gamma))}{\Gamma(\pi o \mu ar{
u}(\gamma))}_{{
m CALC}} =$$

 $\begin{cases} 1.2352 (5) \times 10^{-4} & \text{Marciano and Sirlin, [PRL$ **71** $(1993) 3629]} \\ 1.2354 (2) \times 10^{-4} & \text{Finkemeier, [PL B$ **387** $(1996) 391]} \\ 1.2352 (1) \times 10^{-4} & \text{Cirigliano and Rosell, [PRL$ **99** $(2007) 231801]} \end{cases}$

Experimental world average [PDG] is about $|40\times|$ less accurate:

$$rac{\Gamma(\pi
ightarrow ear{
u}(\gamma))}{\Gamma(\pi
ightarrow \muar{
u}(\gamma))} = 1.230(4) imes 10^{-4};$$





- Early evidence for V A nature of weak interaction.
- Modern SM calculations:

$${\cal R}^{\pi}_{{
m e}/\mu} = rac{\Gamma(\pi o e ar{
u}(\gamma))}{\Gamma(\pi o \mu ar{
u}(\gamma))}_{{
m CALC}} =$$

 $\begin{cases} 1.2352 (5) \times 10^{-4} & \text{Marciano and Sirlin, [PRL$ **71** $(1993) 3629]} \\ 1.2354 (2) \times 10^{-4} & \text{Finkemeier, [PL B$ **387** $(1996) 391]} \\ 1.2352 (1) \times 10^{-4} & \text{Cirigliano and Rosell, [PRL$ **99** $(2007) 231801]} \end{cases}$

Experimental world average [PDG] is about $|40\times|$ less accurate:

$$\overline{\Gamma(\pi o e ar{
u}(\gamma))}_{(\pi o \mu ar{
u}(\gamma))} = 1.230(4) imes 10^{-4};$$

PEN goal:
$$\frac{\Delta R}{R} \simeq 5 \times 10^{-4}$$





- Early evidence for V A nature of weak interaction.
- Modern SM calculations:

$${\cal R}^{\pi}_{{
m e}/\mu} = rac{\Gamma(\pi o e ar{
u}(\gamma))}{\Gamma(\pi o \mu ar{
u}(\gamma))}_{{
m CALC}} =$$

 $\begin{cases} 1.2352 (5) \times 10^{-4} & \text{Marciano and Sirlin, [PRL$ **71** $(1993) 3629]} \\ 1.2354 (2) \times 10^{-4} & \text{Finkemeier, [PL B$ **387** $(1996) 391]} \\ 1.2352 (1) \times 10^{-4} & \text{Cirigliano and Rosell, [PRL$ **99** $(2007) 231801]} \end{cases}$

Experimental world average [PDG] is about 40× less accurate:

$$\frac{\Gamma(\pi \to e\bar{\nu}(\gamma))}{\Gamma(\pi \to \mu\bar{\nu}(\gamma))} = 1.230(4) \times 10^{-4}; \quad \text{PEN goal: } \frac{\Delta R}{R} \simeq 5 \times 10^{-4}$$

- Strong SM helicity suppression amplifies sensitivity to PS terms ("door" for New Physics) by factor $2m_{\pi}/m_e(m_{\mu}+m_d) \approx 8000$.
- ▶ $\mathbf{R}_{\mathrm{e}/\mu}^{\pi}$ tests lepton universality: in SM e, μ , τ differ by Higgs couplings only; there could also be new S or PS bosons with non-universal couplings (New Physics).



$$\begin{split} \mathcal{L}_{\mathsf{NP}} &= \left[\pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{V}}^{2}} \bar{u} \gamma_{\alpha} d \pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{A}}^{2}} \bar{u} \gamma_{\alpha} \gamma_{5} d \right] \bar{e} \gamma^{\alpha} (1 - \gamma_{5}) \nu \\ &+ \left[\pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{S}}^{2}} \bar{u} d \pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{P}}^{2}} \bar{u} \gamma_{5} d \right] \bar{e} (1 - \gamma_{5}) \nu \,, \quad (\mathsf{\Lambda}_{i} \dots \mathsf{scale of NP}) \end{split}$$



SM tests with light particles:

The π_{e2} decay



$$\begin{split} \mathcal{L}_{\mathsf{NP}} &= \left[\pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{V}}^{2}} \bar{u} \gamma_{\alpha} d \pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{A}}^{2}} \bar{u} \gamma_{\alpha} \gamma_{5} d \right] \bar{e} \gamma^{\alpha} (1 - \gamma_{5}) \nu \\ &+ \left[\pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{S}}^{2}} \bar{u} d \pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{P}}^{2}} \bar{u} \gamma_{5} d \right] \bar{e} (1 - \gamma_{5}) \nu \,, \quad (\mathsf{\Lambda}_{i} \dots \mathsf{scale of NP}) \end{split}$$

CKM unitarity and superallowed Fermi nuclear decays currently limit:

 $\Lambda_V \ge 20 \text{ TeV}, \quad \text{ and } \quad \Lambda_S \ge 10 \text{ TeV}.$



SM tests with light particles:



$$\begin{split} \mathcal{L}_{\mathsf{NP}} &= \left[\pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{V}}^{2}} \bar{u} \gamma_{\alpha} d \pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{A}}^{2}} \bar{u} \gamma_{\alpha} \gamma_{5} d \right] \bar{e} \gamma^{\alpha} (1 - \gamma_{5}) \nu \\ &+ \left[\pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{S}}^{2}} \bar{u} d \pm \frac{\pi}{2\mathsf{\Lambda}_{\boldsymbol{P}}^{2}} \bar{u} \gamma_{5} d \right] \bar{e} (1 - \gamma_{5}) \nu \,, \quad (\mathsf{\Lambda}_{i} \dots \mathsf{scale of NP}) \end{split}$$

CKM unitarity and superallowed Fermi nuclear decays currently limit:

$$\Lambda_V \ge 20 \text{ TeV}, \quad \text{and} \quad \Lambda_S \ge 10 \text{ TeV}.$$

At $\Delta R_{e/\mu}^{\pi}/R_{e/\mu}^{\pi} = 10^{-3}$, π_{e2} decay is directly sensitive to:

$$\label{eq:relation} \begin{split} & \boxed{\Lambda_{\mathcal{P}} \leq 1000 \, \text{TeV}} \quad \text{and} \quad \boxed{\Lambda_{\mathcal{A}} \leq 20 \, \text{TeV}}, \\ \text{and indirectly, through loop effects to} \quad \boxed{\Lambda_{\mathcal{S}} \leq 60 \, \text{TeV}}. \end{split}$$





$$\begin{split} \mathcal{L}_{\mathsf{NP}} &= \left[\pm \frac{\pi}{2 \Lambda_{\boldsymbol{V}}^2} \bar{u} \gamma_\alpha d \pm \frac{\pi}{2 \Lambda_{\boldsymbol{A}}^2} \bar{u} \gamma_\alpha \gamma_5 d \right] \bar{e} \gamma^\alpha (1 - \gamma_5) \nu \\ &+ \left[\pm \frac{\pi}{2 \Lambda_{\boldsymbol{S}}^2} \bar{u} d \pm \frac{\pi}{2 \Lambda_{\boldsymbol{P}}^2} \bar{u} \gamma_5 d \right] \bar{e} (1 - \gamma_5) \nu \,, \quad (\Lambda_i \dots \text{scale of NP}) \end{split}$$

CKM unitarity and superallowed Fermi nuclear decays currently limit:

$$\Lambda_V \ge 20 \text{ TeV}, \qquad \text{and} \qquad \Lambda_S \ge 10 \text{ TeV}.$$

At $\Delta R_{e/\mu}^{\pi}/R_{e/\mu}^{\pi} = 10^{-3}$, π_{e2} decay is directly sensitive to:

 $\begin{array}{l} \hline \Lambda_P \leq 1000 \, {\rm TeV} & {\rm and} & \hline \Lambda_A \leq 20 \, {\rm TeV} \end{array},\\ \mbox{and indirectly, through loop effects to } \hline \Lambda_S \leq 60 \, {\rm TeV} \end{array}.\\ \mbox{In general multi-Higgs models with charged-Higgs couplings}\\ \hline \lambda_{e\nu} \approx \lambda_{\mu\nu} \approx \lambda_{\tau\nu}, \mbox{ at } 0.1 \% \mbox{ precision, } R^{\pi}_{e\mu} \mbox{ probes } \hline m_{\rm H^{\pm}} \leq 400 \, {\rm GeV} \end{array}. \end{array}$





Other processes and limits; status of PEN

From $\pi \rightarrow \mathbf{e}\nu$, additional constraints on:

- pseudoscalar and vector leptoquarks,
- neutrino sector anomalies through lepton universality,
- heavy neutrinos.





Other processes and limits; status of PEN

From $\pi \rightarrow \mathbf{e}\nu$, additional constraints on:

- pseudoscalar and vector leptoquarks,
- neutrino sector anomalies through lepton universality,
- heavy neutrinos.

Current status of PEN:

- ► Data acquisition runs in 2008, 2009, 2010 completed.
- ► Collected: > 22 M π → e events, > 200 M π → μ → e events.
- Comprehensive, blinded maximum likelihood analysis in progress.





Target waveform fitting:

(A. Palladino)

(1) Shape (filter) wf signals, (2) Use predicted $\pi_{stop}(DEG)$ and $e^+(PH)$ wf's, (3) Fit with 2 and 3-peak wf's; compare χ^2 values.



Key PEN systematic: low E "tail" response



Correlations in **n** beta decay: Nab and ABba/PANDA experiments

(apparatus under construction)



SM tests with light particles:





Neutron beta decay observables (SM)

$$\begin{split} \frac{\mathrm{d}w}{\mathrm{d}E_{e}\mathrm{d}\Omega_{e}\mathrm{d}\Omega_{\nu}} &\simeq p_{e}E_{e}(E_{0}-E_{e})^{2} \\ &\times \left[1+a\frac{\vec{p}_{e}\cdot\vec{p}_{\nu}}{E_{e}E_{\nu}}+b\frac{m}{E_{e}}+\langle\vec{\sigma}_{n}\rangle\cdot\left(A\frac{\vec{p}_{e}}{E_{e}}+B\frac{\vec{p}_{\nu}}{E_{\nu}}\right)+\dots\right] \end{split}$$





Neutron beta decay observables (SM)

$$\frac{\mathrm{d}w}{\mathrm{d}E_{e}\mathrm{d}\Omega_{e}\mathrm{d}\Omega_{\nu}} \simeq p_{e}E_{e}(E_{0}-E_{e})^{2} \\ \times \left[1+a\frac{\vec{p}_{e}\cdot\vec{p}_{\nu}}{E_{e}E_{\nu}}+b\frac{m}{E_{e}}+\langle\vec{\sigma}_{n}\rangle\cdot\left(A\frac{\vec{p}_{e}}{E_{e}}+B\frac{\vec{p}_{\nu}}{E_{\nu}}\right)+\dots\right]$$

where in SM:

$$a = rac{1-|\lambda|^2}{1+3|\lambda|^2}$$
 $A = -2rac{|\lambda|^2+Re(\lambda)}{1+3|\lambda|^2}$ $b \equiv 0$

$$B = 2 \frac{|\lambda|^2 - Re(\lambda)}{1 + 3|\lambda|^2} \qquad \lambda = \frac{G_A}{G_V} \text{ (with } \tau_n \Rightarrow \mathsf{CKM} \ V_{ud}\text{)}$$

also proton asymmetry: $C = \kappa (A + B)$ where $\kappa \simeq 0.275$.





Neutron beta decay observables (SM)

$$\frac{\mathrm{d}w}{\mathrm{d}E_{e}\mathrm{d}\Omega_{e}\mathrm{d}\Omega_{\nu}} \simeq p_{e}E_{e}(E_{0}-E_{e})^{2}$$
$$\times \left[1+a\frac{\vec{p}_{e}\cdot\vec{p}_{\nu}}{E_{e}E_{\nu}}+b\frac{m}{E_{e}}+\langle\vec{\sigma}_{n}\rangle\cdot\left(A\frac{\vec{p}_{e}}{E_{e}}+B\frac{\vec{p}_{\nu}}{E_{\nu}}\right)+\dots\right]$$

where in SM:

$$a = rac{1-|\lambda|^2}{1+3|\lambda|^2}$$
 $A = -2rac{|\lambda|^2+Re(\lambda)}{1+3|\lambda|^2}$ $b \equiv 0$

$$B = 2 \frac{|\lambda|^2 - Re(\lambda)}{1 + 3|\lambda|^2} \qquad \lambda = \frac{G_A}{G_V} \text{ (with } \tau_n \Rightarrow \mathsf{CKM} \ V_{ud}\text{)}$$

also proton asymmetry: $C = \kappa (A + B)$ where $\kappa \simeq 0.275$.

SM overconstrains a, A, B observables in n β decay! Fierz interf. term b adds more sensitivity to non-SM processes!



SM tests with light particles:



Goals of the Nab experiment (at SNS, ORNL)

Measure the electron-neutrino parameter a in neutron decay

with accuracy of

$$\boxed{\frac{\Delta a}{a} \simeq 10^{-3}}$$

or \sim 50 \times better than:

current results:

- -0.1054 ± 0.0055 Byrne et al '02 -0.1017 ± 0.0051 -0.091 ± 0.039
 - Stratowa et al '78 Grigorev et al '68





Goals of the Nab experiment (at SNS, ORNL)

Measure the electron-neutrino parameter a in neutron decay

with accuracy of

$$\frac{\Delta a}{a}\simeq 10^{-3}$$

or $\sim 50\times$ better than:

	-0.1054 ± 0.0055	Byrne et al '02
current results:	-0.1017 ± 0.0051	Stratowa et al '78
	-0.091 ± 0.039	Grigorev et al '68

Measure the Fierz interference term b in neutron decay

 $\Delta b \simeq 3 \times 10^{-3}$

with accuracy of

current results:

none (not yet measured in n decay)





Goals of the Nab experiment (at SNS, ORNL)

Measure the electron-neutrino parameter a in neutron decay

with accuracy of

$$\frac{\Delta a}{a}\simeq 10^{-3}$$

or \sim 50 \times better than:

	-0.1054 ± 0.0055	Byrne et al '02
current results:	-0.1017 ± 0.0051	Stratowa et al '78
	-0.091 ± 0.039	Grigorev et al '68

Measure the Fierz interference term b in neutron decay

with accuracy of

 $\Delta b \simeq 3 \times 10^{-3}$

current results: none (not yet measured in n decay)

 Nab will be followed by the ABba/PANDA polarized program to measure A, electron, and B/C, neutrino/proton, asymmetries with ² 10⁻³ relative precision, an independent measurement of λ.



Current status of V_{ud} and λ , from n decay

... remains an unresolved mess:



Limits on T, S couplings from beta decay



Measurement of *b* with $\delta b < 10^{-3} \Rightarrow > 4$ -fold improvement on the current limit for ϵ_T from $\pi^+ \to e^+ \nu \gamma$ decay.

From T. Bhattacharya, V. Cirigliano, S.D. Cohen, A. Filipuzzi, M. González-Alonso, M.L. Graesser, R. Gupta, H-W. Lin, Phys. Rev. D 85 (2012) 054512.



SM tests with light particles:



Nab measurement principles: proton phase space



NB: For a given E_e , $\cos \theta_{e\nu}$ is a function of p_p^2 only.



SM tests with light particles:



Nab measurement principles: proton phase space



Numerous consistency checks are built-in!

🛄 D. Počanić (UVa)

SM tests with light particles:





SM tests with light particles:

Neutron β decay

33/37



- ► Use edges to determine and verify shape of detection function Φ(p_p, 1/t_p);
- Use central part of $P_t(1/t_p^2)$ (~ 70%) to extract **a**.



Nab systematic uncertainties:

(expt not stat limited)

Exper	imental parameter	$(\Delta a/a)_{ m SYST}$
Magnetic field:	curvature at pinch	$5 imes 10^{-4}$
	ratio $r_{ m B}=B_{ m TOF}/B_0$	$2.5 imes10^{-4}$
	ratio $r_{\rm B,DV} = B_{\rm DV}/B_0$	$3 imes 10^{-4}$
L_{TOF} , length of TOF	region	(*)
U inhomogeneity:	in decay / filter region	$5 imes 10^{-4}$
	in TOF region	$1 imes 10^{-4}$
Neutron Beam:	position	$4 imes 10^{-4}$
	width	$2.5 imes10^{-4}$
	Doppler effect	small
	unwanted beam polarization	small
Adiabaticity of proto	on motion	$1 imes 10^{-4}$
Detector effects:	E_{e} calibration	(*)
	$E_{\rm e}$ resolution	$5 imes 10^{-4}$
	Proton trigger efficiency	$2.5 imes10^{-4}$
Accidental coinciden	ces	small
Residual gas		small
Background		small
Sum		$1 imes 10^{-3}$
(*) Free fit paramete	er	
ć (UVa) SM test	ts with light particles: Neutron β decay	18 Dec '13 👹

35/37

Nab systematic uncertainties:

(expt not stat limited)

Exper	imental parameter	$(\Delta a/a)_{ m SYST}$
Magnetic field:	curvature at pinch	$5 imes 10^{-4}$
	ratio $r_{ m B}=B_{ m TOF}/B_0$	$2.5 imes10^{-4}$
	ratio $r_{\rm B,DV}=B_{\rm DV}/B_0$	$3 imes 10^{-4}$
L_{TOF} , length of TOF	region	(*)
U inhomogeneity:	in decay / filter region	$5 imes 10^{-4}$
	in TOF region	$1 imes 10^{-4}$
Neutron Beam:	position	$4 imes 10^{-4}$
	width	$2.5 imes10^{-4}$
	Doppler effect	small
	unwanted beam polarization	small
Adiabaticity of proto	on motion	$1 imes 10^{-4}$
Detector effects:	$E_{\rm e}$ calibration	(*)
	$E_{\rm e}$ resolution	$5 imes 10^{-4}$
	Proton trigger efficiency	$2.5 imes10^{-4}$
Accidental coinciden	ces	small
Residual gas		small
Background		small
Sum		$1 imes 10^{-3}$
(*) Free fit paramet	er	
anić (UVa) SM tes	ts with light particles: Neutron β decay	18 Dec '13 🕮









➤ A significant experimental effort is under way on the low-E precision frontier to make use of the unparalleled theoretical precision in the weak interactions and properties of the lightest particles.





- ➤ A significant experimental effort is under way on the low-E precision frontier to make use of the unparalleled theoretical precision in the weak interactions and properties of the lightest particles.
- Information thus obtained is complementary to the collider results, and essential for building a consistent overall picture.





- ➤ A significant experimental effort is under way on the low-E precision frontier to make use of the unparalleled theoretical precision in the weak interactions and properties of the lightest particles.
- Information thus obtained is complementary to the collider results, and essential for building a consistent overall picture.
- ► LHC will gradually eliminate much (but not all) of the parameter space within reach of low-energy experiments, barring increases in the latters' precision.





- ➤ A significant experimental effort is under way on the low-E precision frontier to make use of the unparalleled theoretical precision in the weak interactions and properties of the lightest particles.
- Information thus obtained is complementary to the collider results, and essential for building a consistent overall picture.
- ► LHC will gradually eliminate much (but not all) of the parameter space within reach of low-energy experiments, barring increases in the latters' precision.
- Great projects for graduate students and postdocs; excellent preparation for subsequent work on any project in subatomic physics.





- ➤ A significant experimental effort is under way on the low-E precision frontier to make use of the unparalleled theoretical precision in the weak interactions and properties of the lightest particles.
- Information thus obtained is complementary to the collider results, and essential for building a consistent overall picture.
- ► LHC will gradually eliminate much (but not all) of the parameter space within reach of low-energy experiments, barring increases in the latters' precision.
- Great projects for graduate students and postdocs; excellent preparation for subsequent work on any project in subatomic physics.

Home pages: http://pen.phys.virginia.edu http://nab.phys.virginia.edu



