

# On regularizing the infrared singularities in QCD NLO splitting functions with the new Principal Value prescription

O. Gituliar, S. Jadach, A. Kusina, M. Skrzypek

**IFJ PAN, Kraków, Poland**

Partly supported by grants:

DEC-2011/03/B/ST2/02632, UMO-2012/04/M/ST2/00240,  
PITN-GA-2010-264564 (LHCPhenoNet), DE-FG02-13ER41996

LHCPhenoNet Annual Meeting  
CERN, 2-5 Dec. 2013

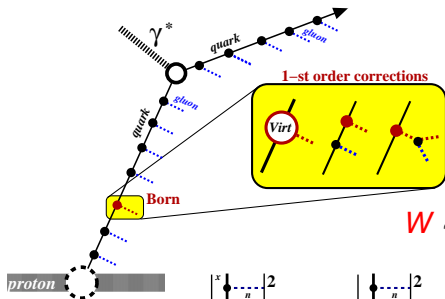


- ▶ Introduction
- ▶ New Principal Value prescription
- ▶ Results for NLO splitting functions
- ▶ Summary

# NLO-corrected middle-of-the-ladder kernel, $C_F^2$



from S. Jadach talk



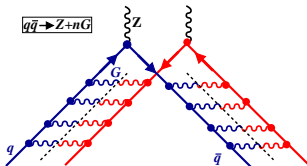
$$W \sim \left| \begin{array}{c} 2 \\ \text{---} \\ 1 \end{array} \right|^2 = \left| \begin{array}{c} 2 \\ \text{---} \\ 2 \\ \text{---} \\ 1 \end{array} \right|^2 + \left| \begin{array}{c} 2 \\ \text{---} \\ 1 \\ \text{---} \\ 2 \end{array} \right|^2 - \left| \begin{array}{c} 2 \\ \text{---} \\ 1 \end{array} \right|^2$$

$$\bar{D}_B^{[1]}(x, Q) = e^{-S_{ISR}} \sum_{n=0}^{\infty} \left\{ \begin{array}{l} \left| \begin{array}{c} x \\ \text{---} \\ a \\ \text{---} \\ n-1 \\ \text{---} \\ \vdots \\ \text{---} \\ 2 \\ \text{---} \\ 1 \end{array} \right|^2 \\ + \sum_{p=1}^n \left| \begin{array}{c} \text{---} \\ a \\ \text{---} \\ n-1 \\ \text{---} \\ \vdots \\ \text{---} \\ p \\ \text{---} \\ 2 \\ \text{---} \\ 1 \end{array} \right|^2 \\ + \sum_{p=1}^n \sum_{j=1}^{p-1} \left| \begin{array}{c} \text{---} \\ n \\ \text{---} \\ p \\ \text{---} \\ \vdots \\ \text{---} \\ j \\ \text{---} \\ 1 \end{array} \right|^2 \end{array} \right\} = e^{-S_{ISR}} \left\{ \delta_{x=1} + \right.$$

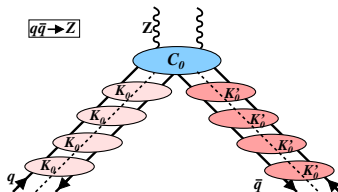
$$\left. + \sum_{n=1}^{\infty} \left( \prod_{i=1}^n \int_{Q > a_i > a_{i-1}} d^3 \eta_i \rho_{1B}^{(1)}(k_i) \right) \left[ 1 + \sum_{p=1}^n \beta_0^{(1)}(z_p) + \sum_{p=1}^n \sum_{j=1}^{p-1} W(\tilde{k}_p, \tilde{k}_j) \right] \delta_{x=\prod_{j=1}^n x_j} \right\}.$$

# Construction of the evolution kernel in collinear factorization

LO cascade and  
construction of the ladder

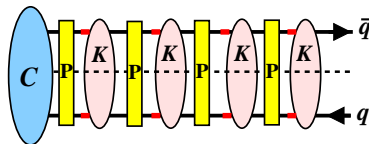


Include NLO and  
group graphs in the ladder into  
"two-particle-irreducible" sets  $K$



Use "projection operators"  $P$  to  
split the ladder and extract kernels

$$\Gamma_{qq} = \text{Tr} \left[ \frac{\hat{n}}{4nq} K \hat{p} \right]$$





# PV prescription

- 😊 Axial gauge = physical interpretation as parton shower
- ☹ Axial gauge = **spurious (unphysical) singularities**

$$\text{gluon propagator: } \frac{1}{l^2} \left( g^{\mu\nu} - \frac{l^\mu n^\nu + n^\mu l^\nu}{nl} \right)$$

**Spurious singularities cancel** in full set of diags, but need regularization  
Curci, Furmanski, Petronzio [80] Ellis, Vogelsang [96], Heinrich, Kunszt [97]:

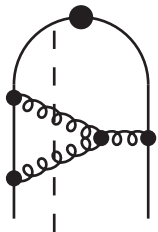
$$\text{Principal Value: } \left[ \frac{1}{nl} \right]_{PV} = \frac{nl}{(nl)^2 + \delta^2(\rho l)^2}$$

PV is more like "phenomenological rule"

Rigorous prescription: Mandelstam [83], Leibbrandt [84].

Difficult in calculations: Bassetto, Heinrich, Kunszt, Vogelsang [97].

## Problem with real emission graph



Standard Heinrich, Kunszt [1998]:

$$N(\epsilon, Q^2) \left[ \frac{P_{qq}(x)}{\epsilon^3} - 2l_0 \frac{P_{qq}(x)}{\epsilon^2} + \frac{p_{qq}(x)}{\epsilon} \left( -2l_1 + 4l_0 + 2l_0 \ln x - 2l_0 \ln(1-x) \right) \right] + \mathcal{O}\left(\frac{1}{\epsilon}\right)$$

Not good for Parton Shower, needs Real-Virt. cancel.

Parton Shower oriented Jadach, et.al. [2011]:

$$\frac{p_{qq}(x)}{\epsilon} \left( +2l_1 + 4l_0 + 2l_0 \ln x - 2l_0 \ln(1-x) \right) + \mathcal{O}\left(\frac{1}{\epsilon}\right)$$

Good for PS,  $\delta$  is a cut-off in 4-dimensions, easy to generate in MC:

$$l_0 = \int_0^1 \frac{dx}{[x]_{PV}} \sim \int_\delta^1 \frac{dx}{x} = -\ln \delta, \quad l_1 = \int_0^1 \frac{dx \ln x}{[x]_{PV}} \sim -\frac{1}{2} \ln^2 \delta,$$

$$P_{qq}(x) = p_{qq} + \epsilon(1-x), \quad p_{qq} = \frac{1+x^2}{1-x}$$



## New use of PV prescription

**Standard:** regularize with PV only the gluon propagator

leave other singularities in (+)-component of integration momenta

$$\frac{d^m l}{l_+^{1-\epsilon}}, \quad l_+ = \frac{nl}{np}$$

**New proposal:** regularize with PV all singularities of the integrand  
in (+)-component of integration momenta, real & virtual

$$\frac{d^m l}{l_+^{1-\epsilon}} \rightarrow d^m l \left[ \frac{1}{l_+} \right]_{PV} \left( 1 + \epsilon \ln l_+ + \epsilon^2 \frac{1}{2} \ln^2 l_+ + \dots \right)$$

All (+)-singularities cancel in the final expression (kernel), so extension of "phenomenological PV rule" of Curci-Furmanski-Petronzio possible



## Example: virtual three point integral

Must perform (+)-integral as the last one, Ellis, Vogelsang [1996]:

$$\begin{aligned}
& \int \frac{d^m l}{(2\pi)^m} \frac{f(l_+)}{l^2(l-q)^2(l-p)^2} = \\
& = \frac{-i}{16\pi^2 q^2} \left(\frac{4\pi}{-q^2}\right)^{-\epsilon} \frac{\Gamma(1-\epsilon)}{-\epsilon} \left[ \int_0^x dy f(l_+) \frac{z^\epsilon(1-z)^\epsilon}{1-y} \left(1 + 2\epsilon \ln \frac{1-y}{1-z}\right) \right. \\
& \quad \left. + 2 \frac{\Gamma^2(1+\epsilon)}{\Gamma(1+2\epsilon)} (1-x)^{-\epsilon} \int_x^1 dy f(l_+) (1-y)^{-1+2\epsilon} \right],
\end{aligned}$$

$$x = q_+/p_+, \quad y = l_+/p_+, \quad z = y/x, \quad p^2 = (p-q)^2 = 0, \quad m = 4 + 2\epsilon,$$

Singularities at  $y = 0$  and  $y = x$ : only from gluon propagator.

Singularity at  $y = 1$ : not from gluon propagator!

**Proposal: treat all (+)-singularities on equal footing**

Note: (+)-singularities lead to  $1/\epsilon^3$  poles in kernel





## Example: scalar non-axial integral

kinematics:  $p^2 = (p - q)^2 = 0$

$$J_3^F = \int \frac{d^m l}{(2\pi)^m} \frac{1}{l^2 (q - l)^2 (p - l)^2}$$

The PV prescription:

$$J_3^F = C \left( -\frac{1}{\epsilon^2} + \frac{\pi^2}{6} \right), \quad C = i \frac{\Gamma(1 - \epsilon)}{(4\pi)^2 |q^2|} \left( \frac{4\pi}{|q^2|} \right)^{-\epsilon}$$

New PV prescription:

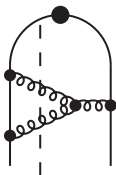
$$J_3^F = C \left( -\frac{2l_0 + \ln(1 - x)}{\epsilon} - 4l_1 + 2l_0 \ln(1 - x) + \frac{\ln^2(1 - x)}{2} \right),$$

$\frac{1}{\epsilon^2}$  replaced by  $l_1$  and  $\frac{1}{\epsilon} l_0$

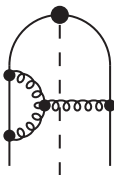
# NLO kernels $P_{qq}$ and $P_{gg}$ in New PV scheme



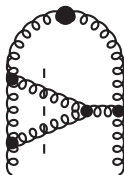
There are only four graphs with  $1/\epsilon^3$  singularity:



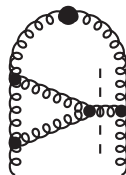
$$\tilde{\Gamma}_{qq}^{(d_R)}(x, \epsilon)$$



$$\tilde{\Gamma}_{qq}^{(d_V)}(x, \epsilon)$$



$$\tilde{\Gamma}_{gg}^{(d_R)}(x, \epsilon)$$

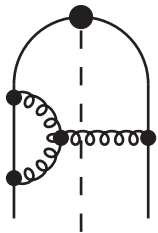


$$\tilde{\Gamma}_{gg}^{(d_V)}(x, \epsilon)$$

related to  $P_{qq}$  and  $P_{gg}$  splitting functions in a standard way:

$$\tilde{\Gamma}_{qq(gg)}(x, \epsilon) = \delta_{1-x} + \frac{1}{\epsilon} \left( \frac{\alpha_S}{2\pi} P_{qq(gg)}^{LO}(x) + \frac{1}{2} \left( \frac{\alpha_S}{2\pi} \right)^2 P_{qq(gg)}^{NLO}(x) + \dots \right) + \mathcal{O}\left(\frac{1}{\epsilon^2}\right)$$

# Virtual contribution to NLO $P_{qq}$ kernel in NPV

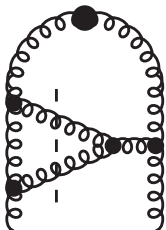


The virtual graph contributes:

$$\begin{aligned}\tilde{\Gamma}_{qq}^{(d_V)}(x, \epsilon) &= -\frac{1}{\epsilon^2} P_{qq} (1 + \epsilon \ln(1-x)) \tilde{Z}_{d_V} \\ &\quad - \frac{1}{\epsilon} P_{qq} \left[ I_0(2 \ln x + 2 \ln(1-x)) - 6I_1 - \text{Li}_2(1-x) \right. \\ &\quad \left. + \ln^2 x - 3 + \frac{8}{12} \pi^2 \right] + \frac{1}{\epsilon} \frac{1}{2} \frac{1+x}{1-x}, \\ \tilde{Z}_{d_V} &= 4I_0 + 2 \ln(1-x) + \ln x - \frac{3}{2},\end{aligned}$$

**Inclusive sum of Real and Virtual graphs  $\tilde{\Gamma}_{qq}^{(d)}$  identical as in standard PV scheme**

# Real contribution to NLO $P_{gg}$ kernel in NPV



Only  $\epsilon^{-1}$  poles

→ the calculation can be done in 4-dimensions,  
→ much simpler than in standard PV

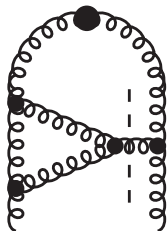
The real graph in New PV prescription:

$$\begin{aligned}\tilde{\Gamma}_{gg}^{(d_R)}(x, \epsilon) = & C_S^{(d_R)} \frac{1}{\epsilon} \left[ P_{gg} \left( -4I_1 + 4I_0 (\ln(1-x) - \ln(x) - 2) \right. \right. \\ & + 2\ln^2(1-x) + 2\ln^2(x) - 4\ln(x)\ln(1-x) - 8\ln(1-x) \\ & + \left. \frac{11}{3}\ln(x) + 2\frac{\pi^2}{6} + 4 \right) + \ln(x) \left( \frac{11}{3}x^2 + \frac{23}{6}x + \frac{23}{6} + \frac{11}{3x} \right) \\ & \left. - \frac{22}{3}x^2 + \frac{24}{3}x - \frac{25}{3} + \frac{22}{3x} \right].\end{aligned}$$

# Virtual contribution to NLO $P_{gg}$ kernel in NPV



The virtual graph in New PV prescription:



$$\begin{aligned}\Gamma_{gg, NPV}^{(d_V)}(x, \epsilon) = & C_S^{(d_V)} P_{gg} \left[ \frac{1}{\epsilon^2} (1 + \epsilon \ln(1-x)) \tilde{Z}_{GS}^V \right. \\ & + \frac{1}{\epsilon} \left( 4l_0 \ln(1-x) + 8l_0 \ln(x) - 16l_1 + 4 \ln^2(x) \right. \\ & \left. \left. + 12 \frac{\pi^2}{6} - \frac{134}{9} \right) \right] - C_S^{(d_V)} \frac{1}{3\epsilon} x\end{aligned}$$

$$\tilde{Z}_{GS}^V = 12l_0 + 4 \ln(1-x) + 4 \ln(x) - \frac{22}{3},$$

Inclusive sum of Real and Virtual graphs  $\tilde{\Gamma}_{gg}^{(d)}$   
identical as in standard PV scheme

**Both schemes, PV and New PV,  
give the same  $P_{qq}$  and  $P_{gg}$  kernels**

# Summary



- ▶ We modified the way of using the PV prescription in the light-cone gauge by applying it to all the singularities in the plus component of the integration momentum.
- ▶ In the New PV prescription the NLO kernels  $P_{qq}$  and  $P_{gg}$  are reproduced correctly.
- ▶ Partial results of four graphs contributing to the kernels, differ in PV and New PV: the  $1/\epsilon^3$  poles, present in PV, are replaced by  $(1/\epsilon) \ln^2 \delta$  etc.
- ▶ Real graphs, free of  $1/\epsilon^3$  and  $1/\epsilon^2$  poles, can be calculated in four dimensions and are usable for the stochastic simulations.
- ▶ The higher order poles cancel separately for real and for virtual components.