

Single top production

(direct measurements of V_{td} , V_{ts} , V_{tb} and new physics related to single top)

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Outline

- Single top production within SM: V_{tb}, V_{ts} & V_{td}?
- <u>Single top production BSM</u>: resonant NP vs EFT, monotops

• Not covered in this talk: th, FCNC top production

(see talks by Brod, Mawatari, Sakurai)

- Disclaimer I: down with a virus infection the past few days
- · Disclaimer II: have not travelled to west-central Africa recently

Single top production within SM



• Theoretical results for single top quark production are available at an ever increasing level of sophistication.

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 G. Bordes and B. van Eijk, Nucl. Phys. B 435, 23 (1995). T. Stelzer, Z. Sullivan and S. Willenbrock, Phys. Rev. D 56, 5919 (1995).
 - NLO QCD & EW predictions
 - resummations

N. Kidonakis, Phys. Rev. D 83, 091503 (2011).
N. Kidonakis, Phys. Rev. D 81, 054028 (2010).
H. X. Zhu et al., JHEP 1102, 099 (2011).
J. Wang et al., arXiv:1010.4509 [hep-ph].

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B.W. Harris et al., Phys. Rev. D 66, 054024 (2002).
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Q.-H. Cao, R. Schwienhorst and C. P. Yuan, Phys. Rev. D 71, 054023 (2005).
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A.S. Papanastasioiu et al., Phys. Lett. B 726, 223 (2013).
R. Schwienhorst et al., Phys. Rev. D 83, 034019 (2011).
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P. Falgari, P. Mellor and A. Signer, Phys. Rev. D 82, 054028 (2010).
P. Falgari et al., Phys. Rev. D 83, 094013 (2011).

matching to parton showers

S. Frixione et al, JHEP 03, 092 (2006).
S. Frixione et al, JHEP 07, 029 (2013).
R. Frederix, E. Re and P. Torrielli, JHEP 1209, 130 (2012).
S. Alioli, P. Nason, C. Oleari and E. Re, JHEP 0909, 111 (2009)

most recently: partial NNLO QCD prediction for t-channel

M. Brucherseifer, F. Caola and K. Melnikov, 1404.7116

8 TeV LHC	$\sigma_{\rm LO},{\rm pb}$	$\sigma_{\rm NLO},{\rm pb}$	$\delta_{ m NLO}$	$\sigma_{\rm NNLO},{\rm pb}$	$\delta_{ m NNLO}$	
top	$53.8^{+3.0}_{-4.3}$	$55.1^{+1.6}_{-0.9}$	+2.4%	$54.2^{+0.5}_{-0.2}$	-1.6%	
anti-top	$29.1^{+1.7}_{-2.4}$	$30.1^{+0.9}_{-0.5}$	+3.4%	$29.7^{+0.3}_{-0.1}$	-1.3%	
smallness of NLO effects appears accidental						

V_{tx} from Single Top Production

• *t*-channel production @ LHC predicted at few % level!

8 TeV LHC $\sigma_{t\,\&\,\bar{t}}^{\rm NNLO} = 83.9^{+0.8}_{-0.3} {\rm pb}$ (does not include ~2% PDF uncertainty)

• proportional to $|V_{tx}|^2$: $\sigma(pp \rightarrow t/X) = A_d |Vtd|^2 + A_s |Vts|^2 + A_b |Vtb|^2$

H. Lacker et al., 1202.4694

 $A_d : A_s : A_b = 5 : 1.8 : 1$ in SM $|V_{td}|^2 : |V_{ts}|^2 : |V_{tb}|^2 = 6 \times 10^{-5} : 1.6 \times 10^{-3} : 1$

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• including top decay $\sigma(pp o (t o bW)/X) \simeq A_b |V_{tb}|^2 R$ J. Alwall et al., hep-ph/0607115

 $R \equiv \frac{|V_{tb}|^2}{|V_{td}|^2 + |V_{ts}|^2 + |V_{tb}|^2}, \quad \text{determined from} \quad \mathcal{B}(t \to bW) / \mathcal{B}(t \to qW)$

- ~8% measurement of $\sigma_{(pp \rightarrow t/X)}$ can be translated into ~4% determination of $|V_{tb}| \Rightarrow$ theoretically $\delta |V_{tb}| < 1\%$ possible
- Access to $V_{ts} \& V_{td}$?

V_{tx} from Single Top Production

- Access to V_{ts} & V_{td}?
 - In *t*-channel and tW production from initial *s*, *b* sea quarks, the events more central than those resulting from initial *d* quarks.
 J. A. Aguilar-Saavedra and A. Onofre, 1002.4718
 - more pronounced for final state t quarks than for antiquarks.
 - Sensitivity far from expected value of $|V_{td}|$ (by two orders of magnitude)
 - Better prospects for $|V_{ts}|$ determination from *t* decays using characteristic differences in the *b* and *s*-jet profiles

A. Ali, F. Barreiro and Th. Lagouri, (PLB 693(2010) 44-51).

W-prime Single top production beyond SM Single top production beyond SM HANNEL NEW FCHANNEL NEW PHYS • s- t- channer resonant production H W Z'С Z'b Resonance production: 4-th generation quark, vector-like $m_{W'_{R}}$ [GeV] $\mathcal{B}(W' \to tb)$ $\sigma \times \mathcal{B}$ [pb] tion of fermionic t & b partners 500 0.298 54.6 ± 2.1 10.9 ± 0.6 0.319 750 on: 4-th generation q 1000 0.326 2.92 ± 0.18 1250 0.328 0.91 ± 0.07 1500 0.330 0.31 ± 0.03 W 1750 0.331 0.11 ± 0.01 W 2000 0.332 0.04 ± 0.01

• anomalous *tWb* couplings & 4-quark operators



from *D*-<u>*D*</u> mixing

Ahrib et al., hep-ph/0602175

Resonant single t

- The spin analyzing power of top decays resonance chirality
- Tacos (distribution of charged lepton in W decay
 - In top rest frame, lepton always along top spin dire

1 9.0 0.8 0.8 0.8 $\frac{d\Gamma}{d\Gamma} = \frac{1 + \lambda_t \cos \theta_{\rm hel}}{1 + \lambda_t \cos \theta_{\rm hel}}$ In helicity basis $\overline{\Gamma d\cos\theta_{\rm hel}}$ Γ 1); 0.4 $\lambda_t \dot{\lambda}_t \neq 1+1$ (right handed) 0.3 0.2 $\lambda_t \lambda_t = 1 - 1$ (left handed) 0.1 -0.5 0 $\lambda_t = -1$ Kane, Ladinsky & Yuan, Phys.Rev. D45 (1992) 124-141

TOPR

500

750

1000 1250

1750

2000

W

 $\mathcal{B}(W' \to tb)$

0.298

0.319

0.326

0.328 0.330

0.331

0.332

 $\sigma \times \mathcal{B}[pb]$

 0.11 ± 0.01

 0.04 ± 0.01

 H^{+}

top quark rest fra

 $\theta_{\ell} t$

 $d\Gamma$

 $\Gamma d \cos$

to probe

v, <u>q</u>'

(a)

D

 $\cos \theta_{\ell}$





 Fermionic top partners below TeV ubiquitous in models of a light composite Higgs

$$\delta m_h^2 \sim \frac{m_L^2}{2} \Lambda^2 + \frac{m_t^2}{v^2} m_T^2 \log \frac{\Lambda^2}{m_T^2} + \dots$$
 t, T

• top quark partial compositeness implies sizable t-T mixing

 $\mathcal{L} \ni m_L f_L \mathcal{O}_R + m_R f_R \mathcal{O}_L + \lambda \mathcal{O}_L \mathcal{O}_H \mathcal{O}_R, \quad \mathcal{O}_L \sim (3, 2)_{1/6}, \dots$

 $\Rightarrow y_t \sim s_L s_R \lambda$

Kaplan, Nucl.Phys. B365 (1991) 259-278

 single T production important search channel (especially for heavier T masses)

recent progress: Azatov et al., 1308.6601 Gutierrez et al., 1403.7490 Gripaios et al., 1406.5957 Matsedonskyia, Panico & Wulzer, 1409.0100 Backovic et al., 1409.0409

c.f. De Simone et al., 1211.5663

 \mathcal{O}_H



Fermionic top partners below TeV ubiquitous in models of a light

W

- composite Higgs Aguilar-Saavedra et al., 1306.0572 10 10 $\delta m_h^2 \sim \frac{r}{r}$ $Q\overline{Q}$ 8 TeV LHC 10 Ybj (B Y) (B Y) (T)Tbj 10 Bbj Tītj top quark σ_{max} (fb) σ_{max} (fb) Xtj $\mathcal{L} \ni m_L f_L \mathcal{O}$ (XT)(X T) 10 10^{-2} 13 TeV 8 TeV 10^{-3} 10 400 600 800 1000 1200 1400 1600 1800 2000 400 600 800 1000 1200 m_Q (GeV) m_Q (GeV)
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W

BSM single top production in decoupling limit

If NP degrees of freedom kinematically inaccessible can describe
 BSM effects in single top production in terms of EFT

$$\mathcal{C}_{\text{int}} = \sum_{a} \frac{C_a}{\Lambda^{n_a}} \mathcal{O}_a$$

• at lowest (
$$n_a=2$$
) order

(1) anomalous tWq, (tZq) couplings

$$\hat{O}_{\varphi q}^{(3)} = (\varphi^{\dagger} i \overleftrightarrow{D}_{\mu}^{I} \varphi) (\bar{q}_{L} \sigma^{I} \gamma^{\mu} q_{L})
\hat{O}_{\varphi t b} = i (\tilde{\varphi}^{\dagger} D_{\mu} \varphi) (\bar{t}_{R} \gamma^{\mu} b_{R})
\hat{O}_{t W} = \bar{q}_{L} \sigma^{\mu \nu} \sigma^{I} t_{R} \tilde{\varphi} W_{\mu \nu}^{I}
\hat{O}_{b W} = \bar{q}_{L} \sigma^{\mu \nu} \sigma^{I} b_{R} \varphi W_{\mu \nu}^{I}$$

 $\mathcal{L}_{Wtb} = -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (\mathcal{F}_L P_L + \mathcal{F}_R P_R) b W_{\mu}^+ \\ -\frac{g}{\sqrt{2}} \bar{b} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \{ \bar{t} \gamma^{\mu} (V_L P_L^{O_{Wtb}} = \frac{g}{\sqrt{P_R}} \} \} \}$

Cao, Wudka & Yuan, 0704.2809

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Cao, Wudka & Yuan, 0704.2809

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- (1) anomalous tWq, (tZq) couplings
- (2) four quark operators

$$\mathcal{O}_{qu}^{(1)} = (\bar{q}_l t_R) (\bar{u}_R q_l),$$

$$\mathcal{O}_{qq}^{(1)} = (\bar{q}_l^i t_R) (\bar{q}_l^j b_R) \epsilon_{ij}, \qquad \Longrightarrow$$

$$\mathcal{O}_{qq}^{(3)} = \frac{1}{2} (\bar{q}_l \gamma_\mu \tau^I q_l) (\bar{q}_h \gamma^\mu \tau^I q_h),$$



BSM single top production in decoupling limit

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Can be constrained with combination of single top production and decay observables:

- modified rates

- effects on $V_L = 1, g_R = 0, g_L = 0$ 95% CL allowed regions $F_{0L}+\sigma_t$ (CMS 2012) 0.15

C.f. Aquilar Saavedra, 0803.3810 Bach & Ohl, 1209.4564 Bernardo et al., 1408.7063



Monotop production

DM Pair Production at Hadron Colliders

• General discussion in terms of EFT

$$\mathcal{L}_{\rm int} = \sum_{a} \frac{C_a}{\Lambda^{n_a}} \mathcal{O}_a$$

• With B preservation, O_a need to be bilinear in quark fields

$$\mathcal{O}_{1a}^{ij} = \left(\bar{Q}_L^i \gamma_\mu Q_L^j\right) \mathcal{J}_a^\mu, \qquad \mathcal{O}_{2a}^{ij} = \left(\bar{u}_R^i \gamma_\mu u_R^j\right) \mathcal{J}_a^\mu, \qquad \mathcal{O}_{3a}^{ij} = \left(\bar{d}_R^i \gamma_\mu d_R^j\right) \mathcal{J}_a^\mu, \\ \mathcal{O}_{4a}^{ij} = \left(\bar{Q}_L^i H u_R^j\right) \mathcal{J}_a, \qquad \mathcal{O}_{5a}^{ij} = \left(\bar{Q}_L^i \tilde{H} d_R^j\right) \mathcal{J}_a,$$

• coupling to suitable dark sector currents, i.e.

$$\mathcal{J}_{V,A}^{\mu} = \bar{\chi}\gamma^{\mu}\{1,\gamma_{5}\}\chi \quad \mathcal{J}_{S,P} = \bar{\chi}\{1,\gamma_{5}\}\chi \qquad \mathcal{J} = \chi^{\dagger}\chi, \ \mathcal{J}^{\mu} = \chi^{\dagger}\partial^{\mu}\chi$$

Fermionic Scalar

DM Pair Production at Hadron Colliders

- Flavor universal contributions (C_{ij} ~ δ_{ij})
- mono[jet, γ] constraints using initial state radiation for tagging



DM Pair Production at Hadron Colliders

- Can flavor violating interactions be competitive?
 - Constraints from ΔF =2 observables



• large monotop ($t+E_{miss}$) signals possible due to chirality flipping operators (also $b+E_{miss}$, but can be due to flavor conserving ops.)

J.F.K. & Zupan, 1107.0623 Boucheneb et al., 1407.7529

• reconstruction using $j_{(b)}jj + E_{miss}$, or $j_{(b)}I + E_{miss}$ (M_T)

(~ 1% signal eff.)

Andrea, Fuks & Maltoni, 1106.6199 Alvarez, Coluccio Leskow, Drobnak & J.F.K., 1310.7600 Agram et al., 1311.6478

Expectations in Models of Flavor

Minimal Flavor Violation

$$C_{2a} = b_1^{(2a)} + b_2^{(2a)} Y_u^{\dagger} Y_u + b_3^{(2a)} Y_u^{\dagger} Y_d Y_d^{\dagger} Y_u + \cdots,$$

$$C_{4a} = \left(b_1^{(4a)} + b_2^{(4a)} Y_d Y_d^{\dagger} + \cdots \right) Y_u.$$

- For $b_1^a \sim b_2^a \sim b_3^a$ C_{2a} almost flavor diagonal and universal
- C_{4a} is highly hierarchical, can have large flavor violation if $y_b \sim 1$

200

Larger effects expected with horizontal symmetries

Single X + t Production

- Corresponds to production of neutral mediators in DM models
- Example: Scalar DM (S) via (heavy h₂) Higgs portal in THDMIII

$$\mathcal{L}_{h_2}^{\tilde{y}} = \sum_{ij} \left(\tilde{y}_u^{ij} \bar{u}^i P_R u^j h_2 + \tilde{y}_d^{ij} \bar{d}^i P_R d^j h_2 \right) + \text{h.c.} + \lambda v_{\text{EW}} h_2 SS,$$

• *D-<u>D</u> mixing*

$$\begin{split} &|\tilde{y}_{u}^{ut}\tilde{y}_{u}^{ct}|, |\tilde{y}_{u}^{tu}\tilde{y}_{u}^{tc}| < 0.030 \times \left(\frac{m_{h_{2}}}{250 \text{GeV}}\right)^{2}, \\ &|\tilde{y}_{u}^{tu}\tilde{y}_{u}^{ct}|, |\tilde{y}_{u}^{ut}\tilde{y}_{u}^{tc}| < 0.0088 \times \left(\frac{m_{h_{2}}}{250 \text{GeV}}\right)^{2}, \\ &\sqrt{|\tilde{y}_{u}^{ut}\tilde{y}_{u}^{tu}\tilde{y}_{u}^{ct}\tilde{y}_{u}^{tc}|} < 0.0036 \times \left(\frac{m_{h_{2}}}{250 \text{GeV}}\right)^{2}, \end{split}$$



Conclusions

- Single top quark production a key part of the top physics program at the LHC
 - fruitful laboratory of EW & QCD effects (e.g. PDF fits)
 - provides direct access to third row of CKM
- Rich BSM phenomenology, complementary to other searches
 - Example 1: single top partner production expected to become dominant discovery channel within composite scenarios
 - Example 2: monotop provides powerful probe of dark sectors coupling to quark flavor