

# $\Lambda_b \rightarrow p \ell \nu$ on the lattice

**MATTHEW WINGATE**  
**DAMTP, UNIVERSITY OF CAMBRIDGE**

**STEFAN MEINEL**  
**UNIVERSITY OF ARIZONA**

**CKM 2014**  
**WGII:  $V_{UB}$  &  $V_{CB}$**

# $\Lambda_b \rightarrow p l \nu$ on the lattice

**MATTHEW WINGATE\***  
**DAMTP, UNIVERSITY OF CAMBRIDGE**

**STEFAN MEINEL**  
**UNIVERSITY OF ARIZONA**

**CKM 2014**  
**WGII:  $V_{UB}$  &  $V_{CB}$**

\* **SPEAKER, RESPONSIBLE FOR ANY ERRORS IN SLIDES**

# Outline

- ✦ Motivation
- ✦ Background
- ✦ Published results ( $m_b \rightarrow \infty$  limit)
- ✦ Preliminary results (physical  $m_b$ )

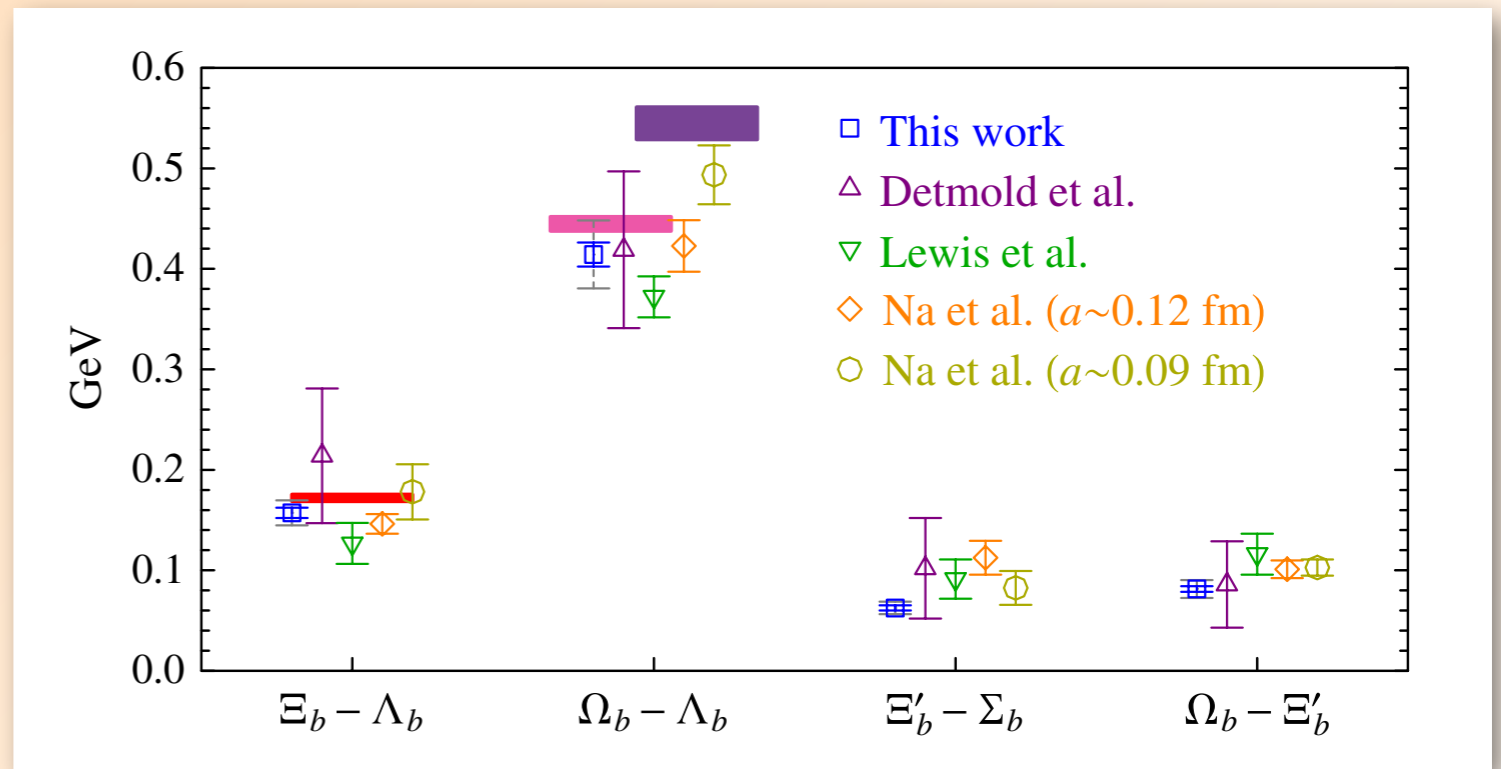
# Motivation

- ✿ [Reminder:  $\Lambda_b$  is the isospin=0, beauty\* =1 baryon]
- ✿  $\Lambda_b \rightarrow p l^- \nu$  is being measured at the LHC
- ✿ Test for new physics in  $b$  to  $u$ :  $B \rightarrow \pi l \nu$  and  $B_s \rightarrow K l \nu$  only expose the vector part of  $(V \pm A)$
- ✿  $\Lambda_b$  form factors are “simple” lattice quantities, in contrast to  $B \rightarrow \rho l \nu, B_s \rightarrow K^* l \nu$
- ✿  $\Lambda_b \rightarrow p$  form factor calculation done in parallel with those for the rare  $b \rightarrow s$  decay  $\Lambda_b \rightarrow \Lambda l^+ l^-$

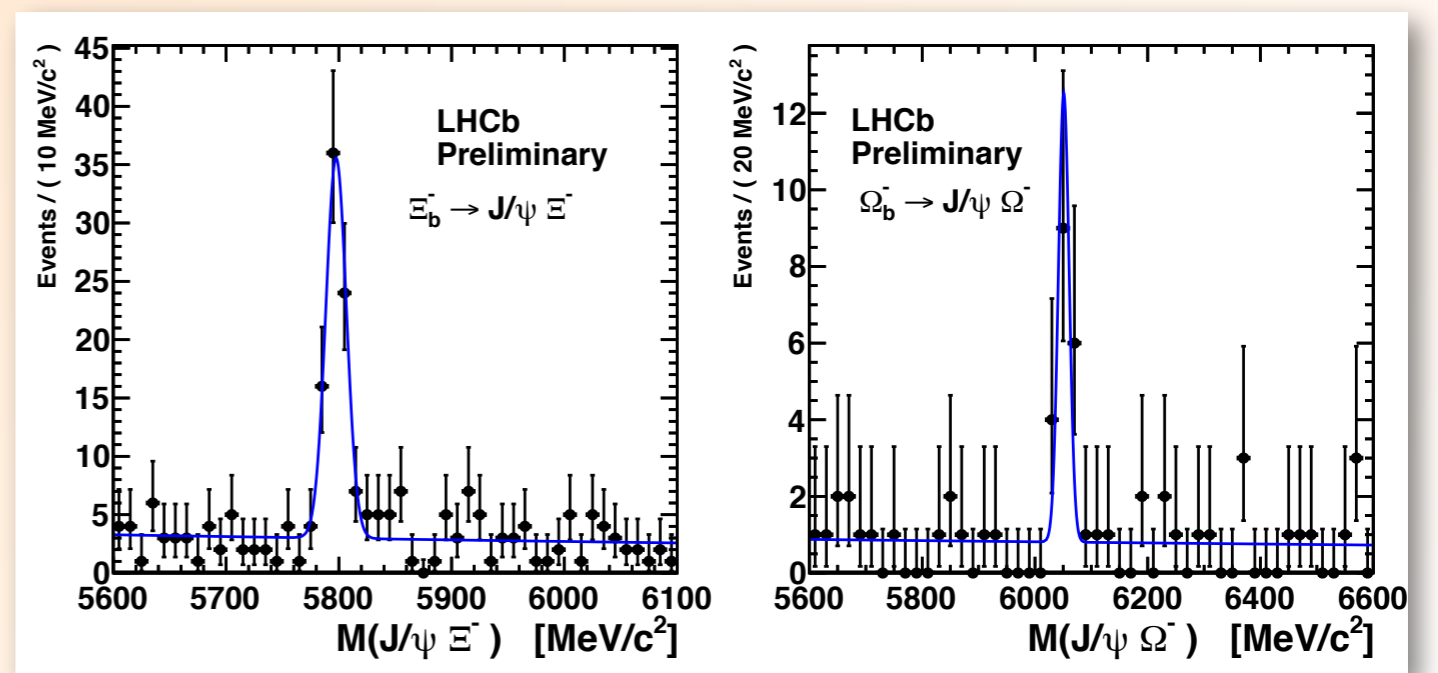
\* beauty = -bottomness

# Beautiful baryons on the lattice

- ❖ Unquenched LQCD  $b$  baryon mass splittings
- ❖ LQCD predicted smaller  $\Omega_b$  mass than originally determined by DØ, later confirmed by CDF & LHCb
- ❖ LQCD calculations of  $\Sigma_b^* \Sigma_b \pi$  and  $\Sigma_b^{(*)} \Lambda_b \pi$  couplings (Detmold, Lin, Meinel, 2011-12)



from Lin, Cohen, Mathur, Orginos, PRD80 (2009)



LHCb-CONF-2011-60

# $\Lambda_b \rightarrow p$ form factors

In general

$$\langle p | \bar{u} \gamma^\mu b | \Lambda_b \rangle = \bar{u}_p \left[ f_1^V \gamma^\mu - f_2^V \frac{i\sigma^{\mu\nu} q_\nu}{m_{\Lambda_b}} + f_3^V \frac{q^\mu}{m_{\Lambda_b}} \right] u_{\Lambda_b}$$

$$\langle p | \bar{u} \gamma^\mu \gamma_5 b | \Lambda_b \rangle = \bar{u}_p \left[ f_1^A \gamma^\mu - f_2^A \frac{i\sigma^{\mu\nu} q_\nu}{m_{\Lambda_b}} + f_3^A \frac{q^\mu}{m_{\Lambda_b}} \right] \gamma_5 u_{\Lambda_b}$$

In the  $m_b \rightarrow \infty$  limit

$$\langle p(k', s') | \bar{u} \Gamma Q | \Lambda_Q(v, s) \rangle = \bar{u}_p(k', s') [F_1(k' \cdot v) + \psi F_2(k' \cdot v)] \Gamma u_{\Lambda_b}(v, s)$$



# $\Lambda_b \rightarrow \Lambda$ form factors

In general

$$\langle \Lambda | \bar{s} \gamma^\mu b | \Lambda_b \rangle = \bar{u}_\Lambda \left[ f_1^V \gamma^\mu - f_2^V \frac{i\sigma^{\mu\nu} q_\nu}{m_{\Lambda_b}} + f_3^V \frac{q^\mu}{m_{\Lambda_b}} \right] u_{\Lambda_b}$$

$$\langle \Lambda | \bar{s} \gamma^\mu \gamma_5 b | \Lambda_b \rangle = \bar{u}_\Lambda \left[ f_1^A \gamma^\mu - f_2^A \frac{i\sigma^{\mu\nu} q_\nu}{m_{\Lambda_b}} + f_3^A \frac{q^\mu}{m_{\Lambda_b}} \right] \gamma_5 u_{\Lambda_b}$$

$$\langle \Lambda | \bar{s} i\sigma^{\mu\nu} q_\nu b | \Lambda_b \rangle = \bar{u}_\Lambda \left[ f_1^{TV} \frac{\gamma^\mu q^2 - q^\mu \not{q}}{m_{\Lambda_b}} - f_2^{TV} \frac{q^\mu}{m_{\Lambda_b}} \right] u_{\Lambda_b}$$

$$\langle \Lambda | \bar{s} i\sigma^{\mu\nu} q_\nu \gamma_5 b | \Lambda_b \rangle = \bar{u}_\Lambda \left[ f_1^{TA} \frac{\gamma^\mu q^2 - q^\mu \not{q}}{m_{\Lambda_b}} - f_2^{TA} \frac{q^\mu}{m_{\Lambda_b}} \right] \gamma_5 u_{\Lambda_b}$$

In the  $m_b \rightarrow \infty$  limit

$$\langle \Lambda(k', s') | \bar{s} \Gamma Q | \Lambda_Q(v, s) \rangle = \bar{u}_\Lambda(k', s') \left[ F_1(k' \cdot v) + \not{v} F_2(k' \cdot v) \right] \Gamma u_{\Lambda_b}(v, s)$$

# Lattice actions & parameters

RBC/UKQCD lattices (2+1 domain wall)

Static ( $m_b = \infty$ ) heavy quarks

Set	$\beta$	$N_s^3 \times N_t \times N_5$	$am_5$	$am_s^{(\text{sea})}$	$am_{u,d}^{(\text{sea})}$	$a$ (fm)	$am_{u,d}^{(\text{val})}$	$m_\pi^{(\text{val})}$ (MeV)	$m_N^{(\text{val})}$ (MeV)	$N_{\text{meas}}$
C14	2.13	$24^3 \times 64 \times 16$	1.8	0.04	0.005	0.1119(17)	0.001	245(4)	1090(21)	2672
C24	2.13	$24^3 \times 64 \times 16$	1.8	0.04	0.005	0.1119(17)	0.002	270(4)	1103(20)	2676
C54	2.13	$24^3 \times 64 \times 16$	1.8	0.04	0.005	0.1119(17)	0.005	336(5)	1160(19)	2782
F23	2.25	$32^3 \times 64 \times 16$	1.8	0.03	0.004	0.0849(12)	0.002	227(3)	1049(25)	1907
F43	2.25	$32^3 \times 64 \times 16$	1.8	0.03	0.004	0.0849(12)	0.004	295(4)	1094(18)	1917
F63	2.25	$32^3 \times 64 \times 16$	1.8	0.03	0.006	0.0848(17)	0.006	352(7)	1165(23)	2782

1-loop operator matching: T Ishikawa *et al.*, JHEP 1105, 040 (2011)



# Form factor shape

- ❖ In static limit,  $z$ -expansion is not applicable
- ❖ Instead, try monopole, dipole, etc. (Latter is a better fit to the data)
- ❖ Incorporate discretization and quark mass effects

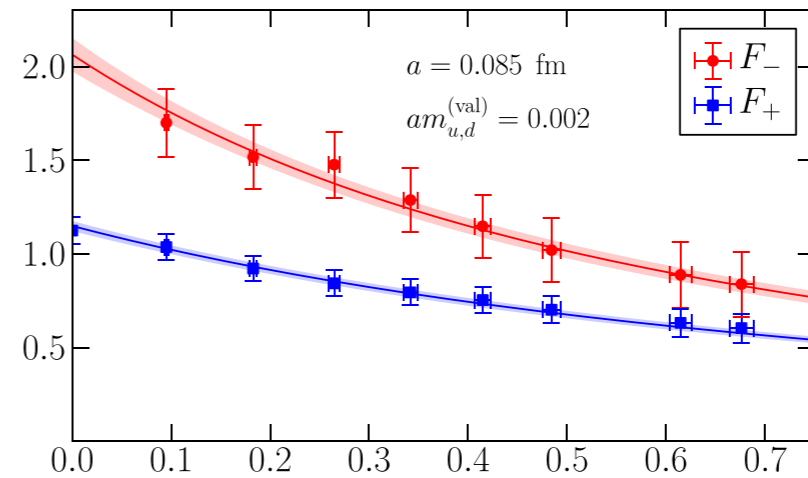
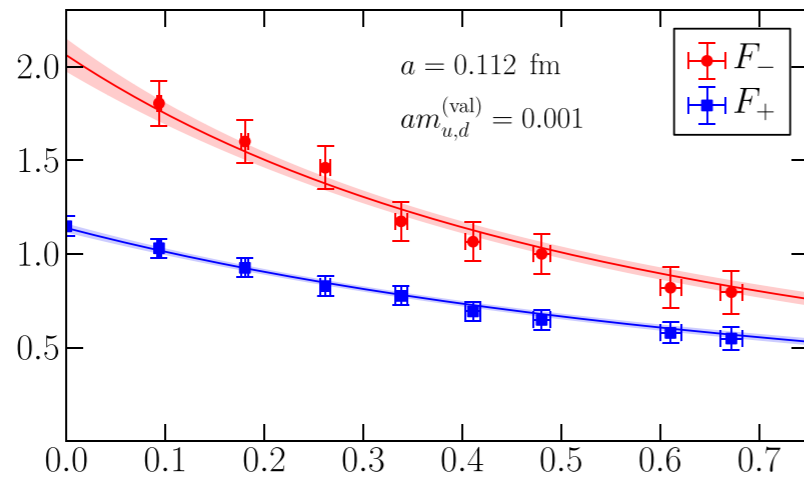
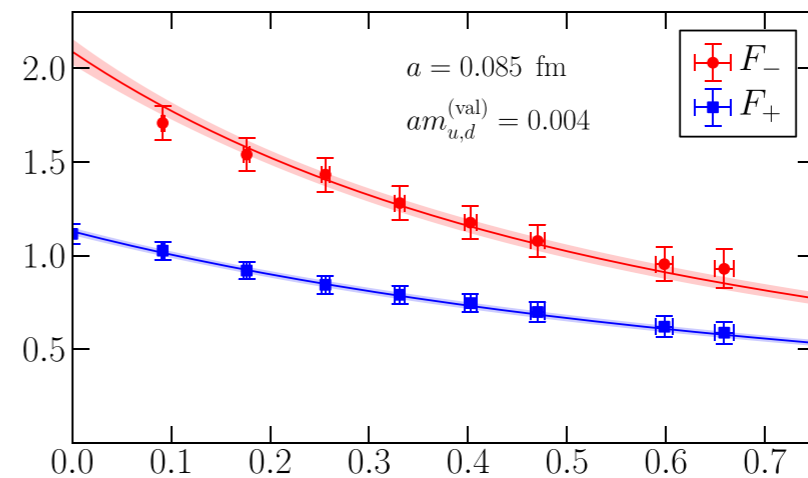
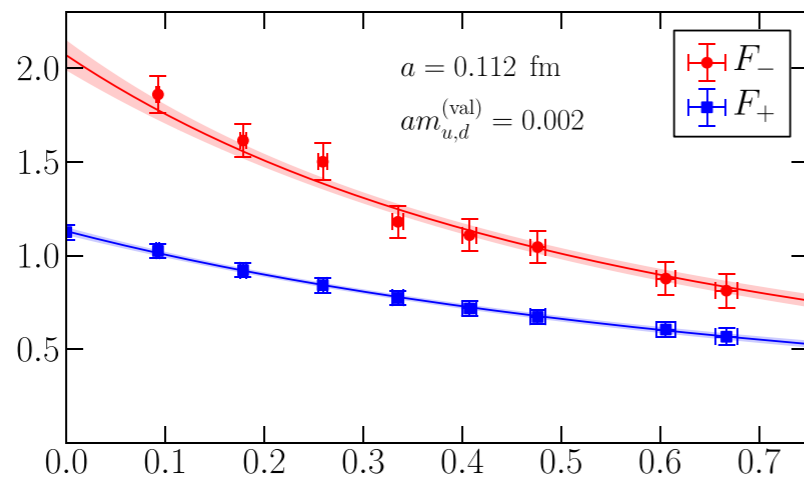
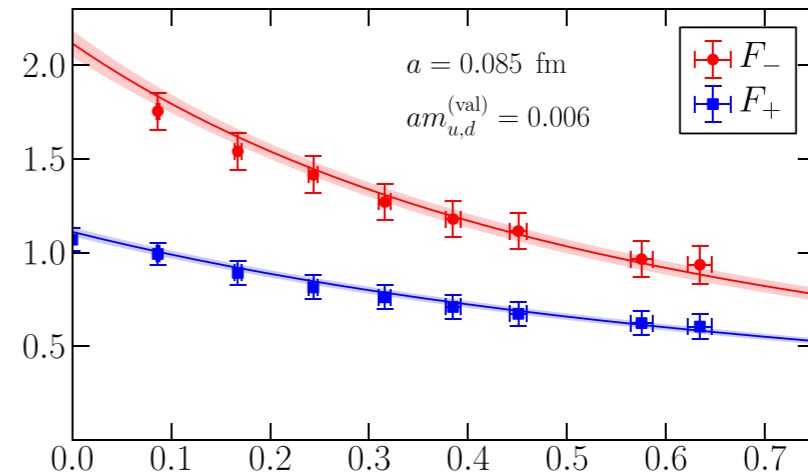
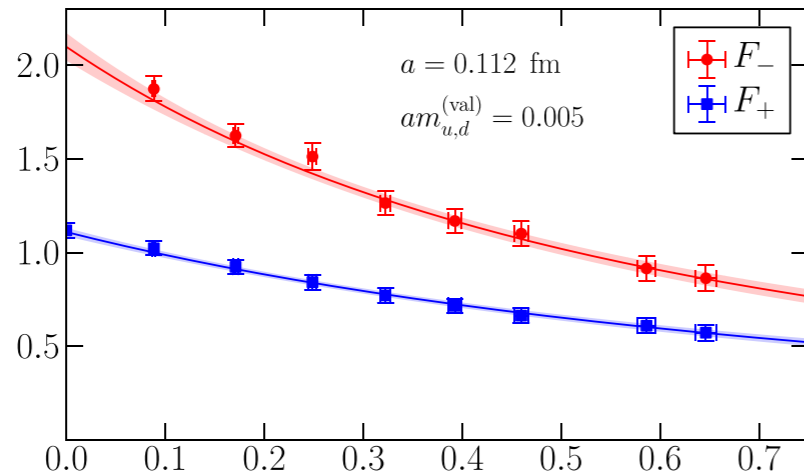
$$F = \frac{Y}{(X + E_p - m_p)^2} [1 + d(aE_p)^2]$$

$$X = X_0 + c[m_\pi^2 - (m_\pi^{\text{phys}})^2]$$

- ❖ In practice,  $c$ 's &  $d$ 's small, consistent with zero [except  $c_{l,+} = 0.094(32)$  in the  $\Lambda_b \rightarrow \Lambda$  calculation]

# Form factors on each ensemble

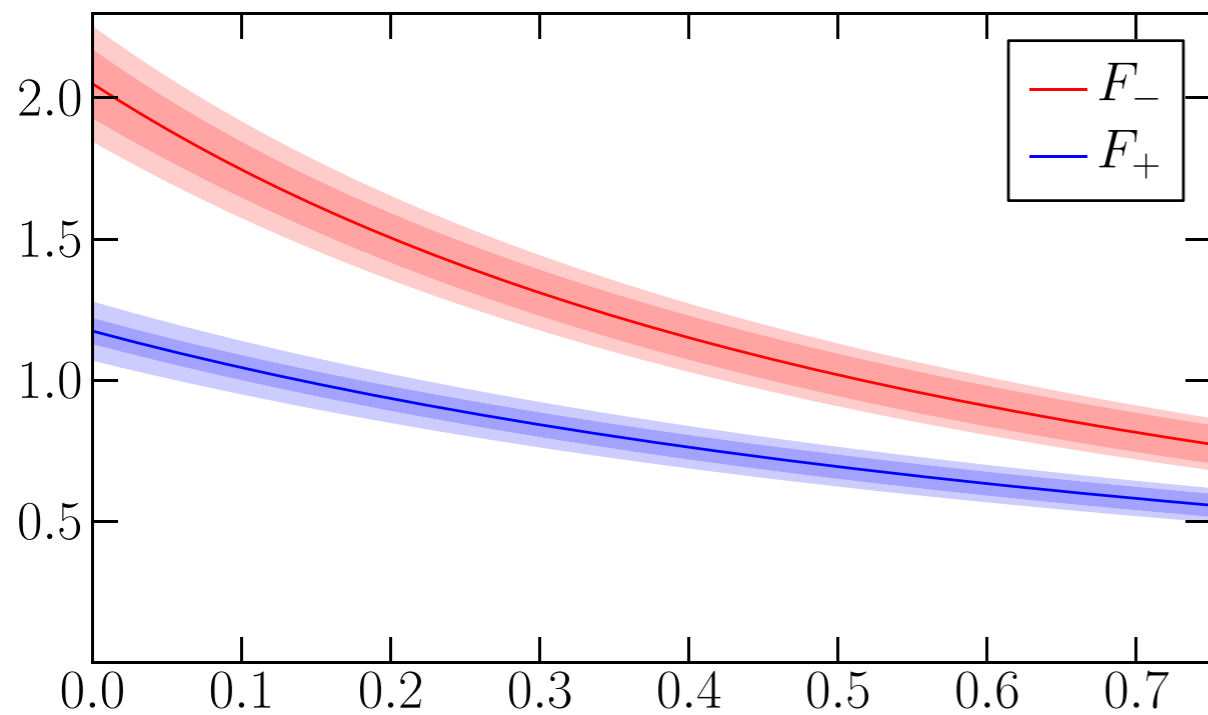
$$F_{\pm} = F_1 \pm F_2$$



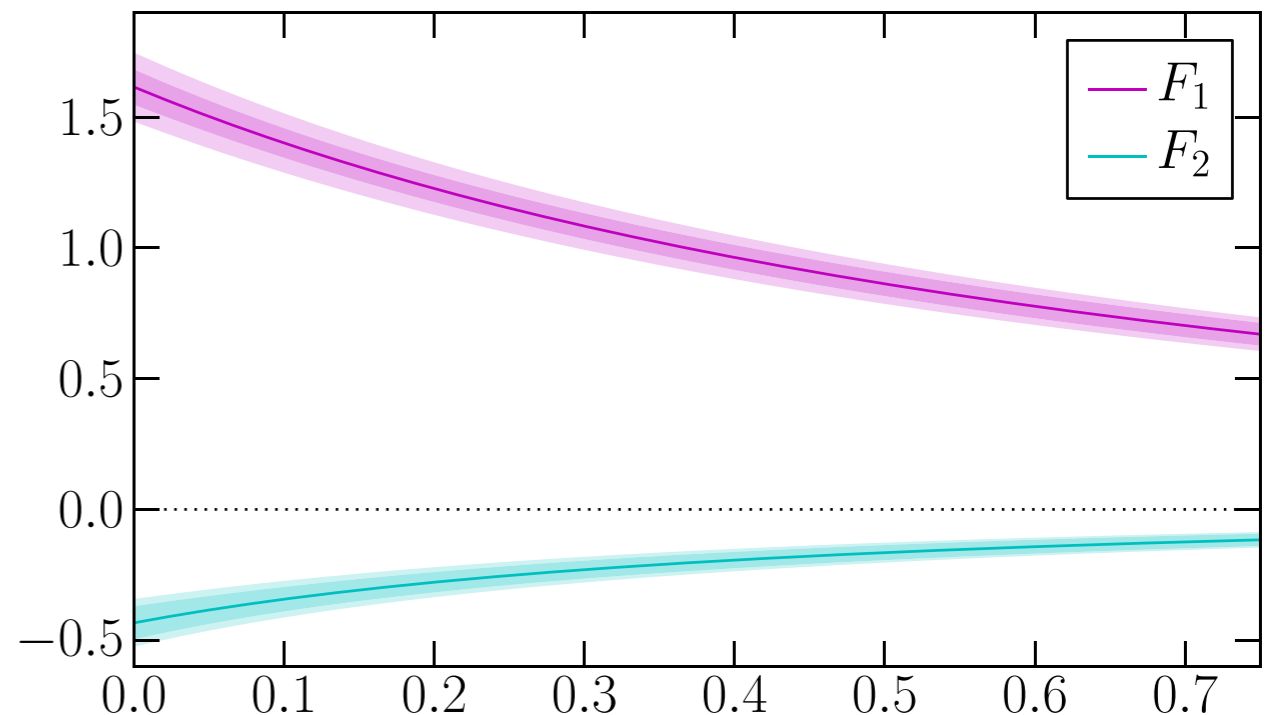
$E_p - m_p$  (GeV)

$E_p - m_p$  (GeV)

# Form factors, physical limit

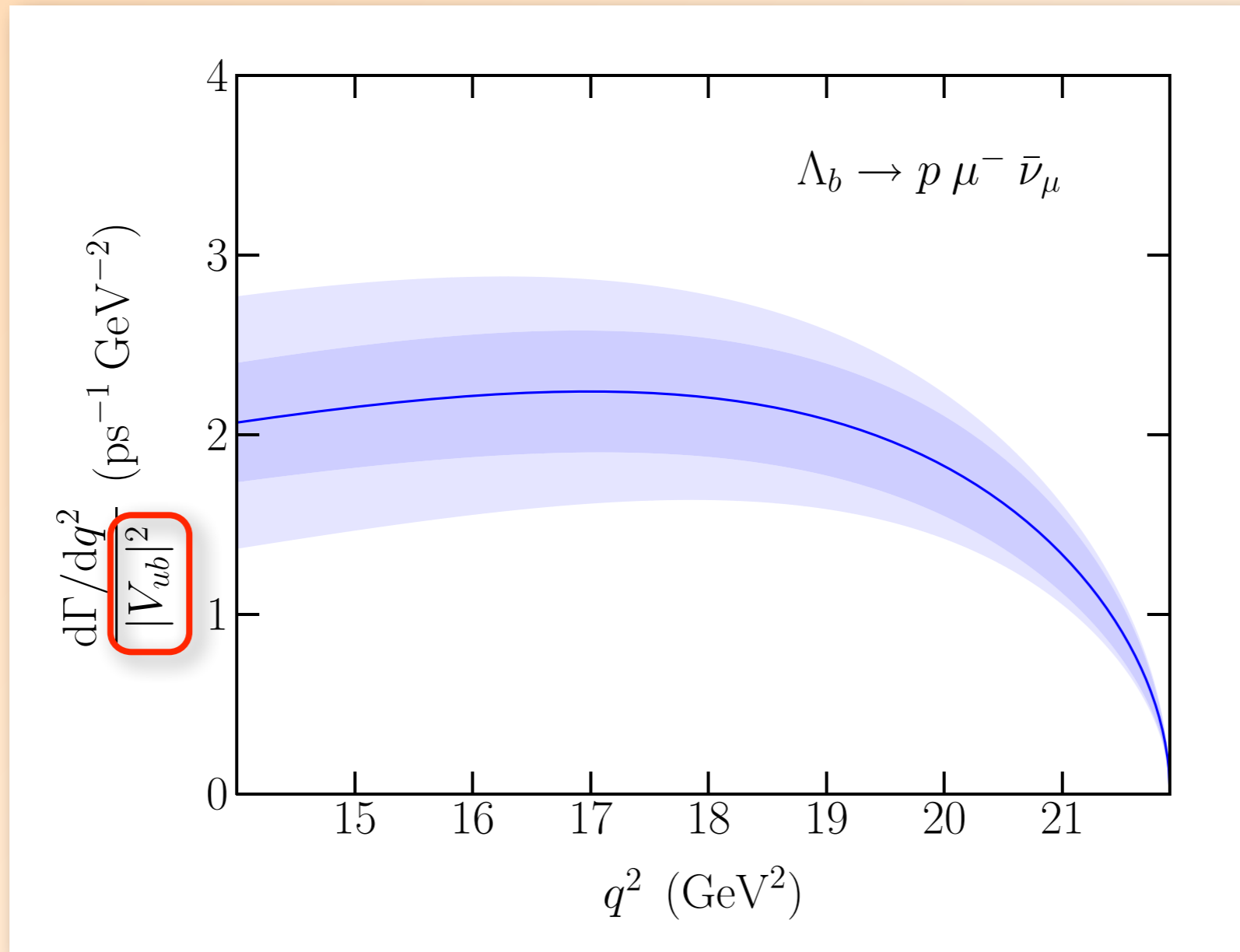


$E_p - m_p$  (GeV)



$E_p - m_p$  (GeV)

# Decay rate $\Lambda \rightarrow p \mu \nu$

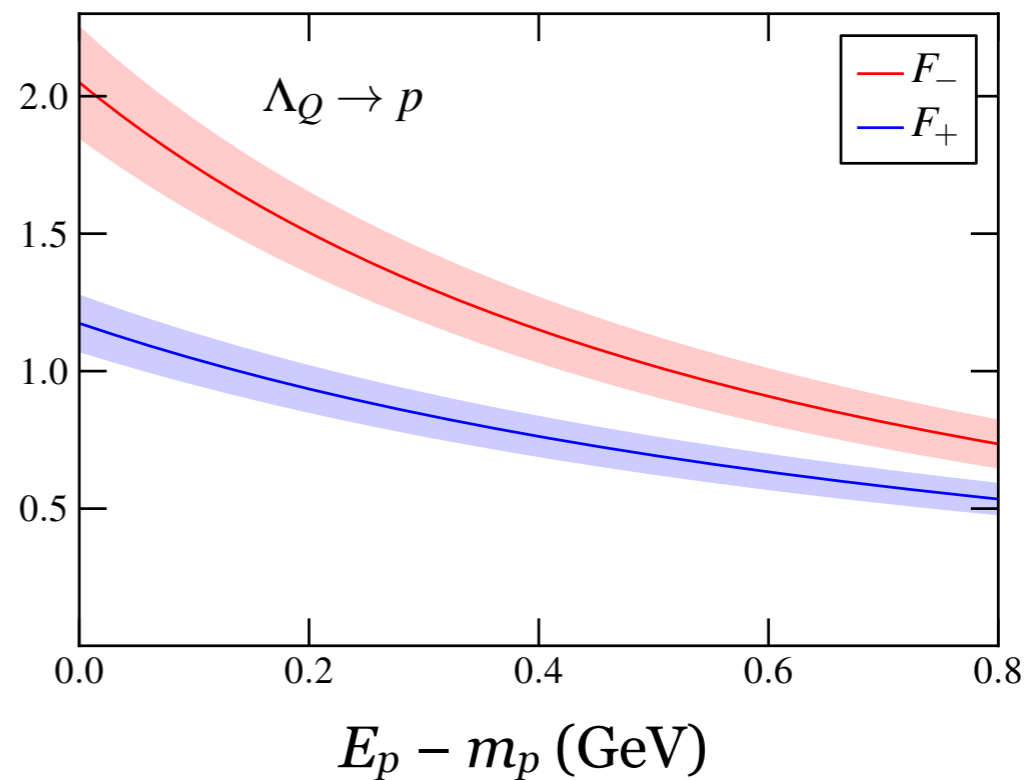
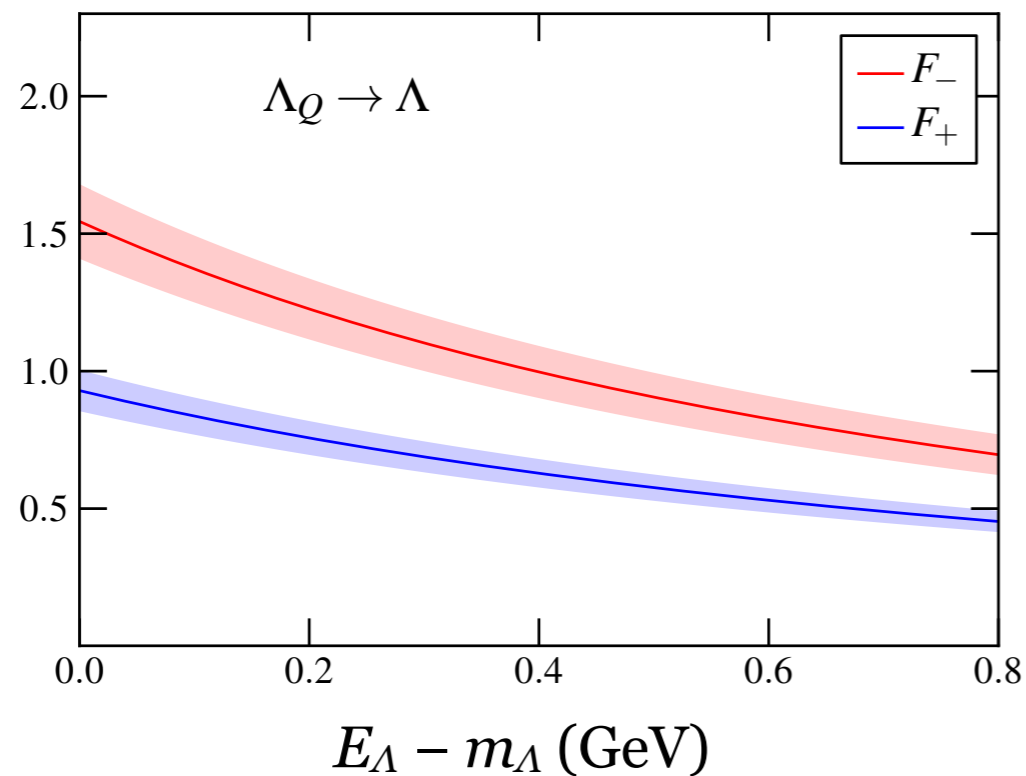


Outer error band includes estimate of terms beyond LO static approximation:  $\sim \sqrt{\frac{\Lambda_{\text{QCD}}^2}{m_b^2} + \frac{|\mathbf{k}'|^2}{m_b^2}}$

Detmold, Lin, Meinel, Wingate, PRD88 (2013)

# Comparison: $\Lambda$ vs. $p$ final states

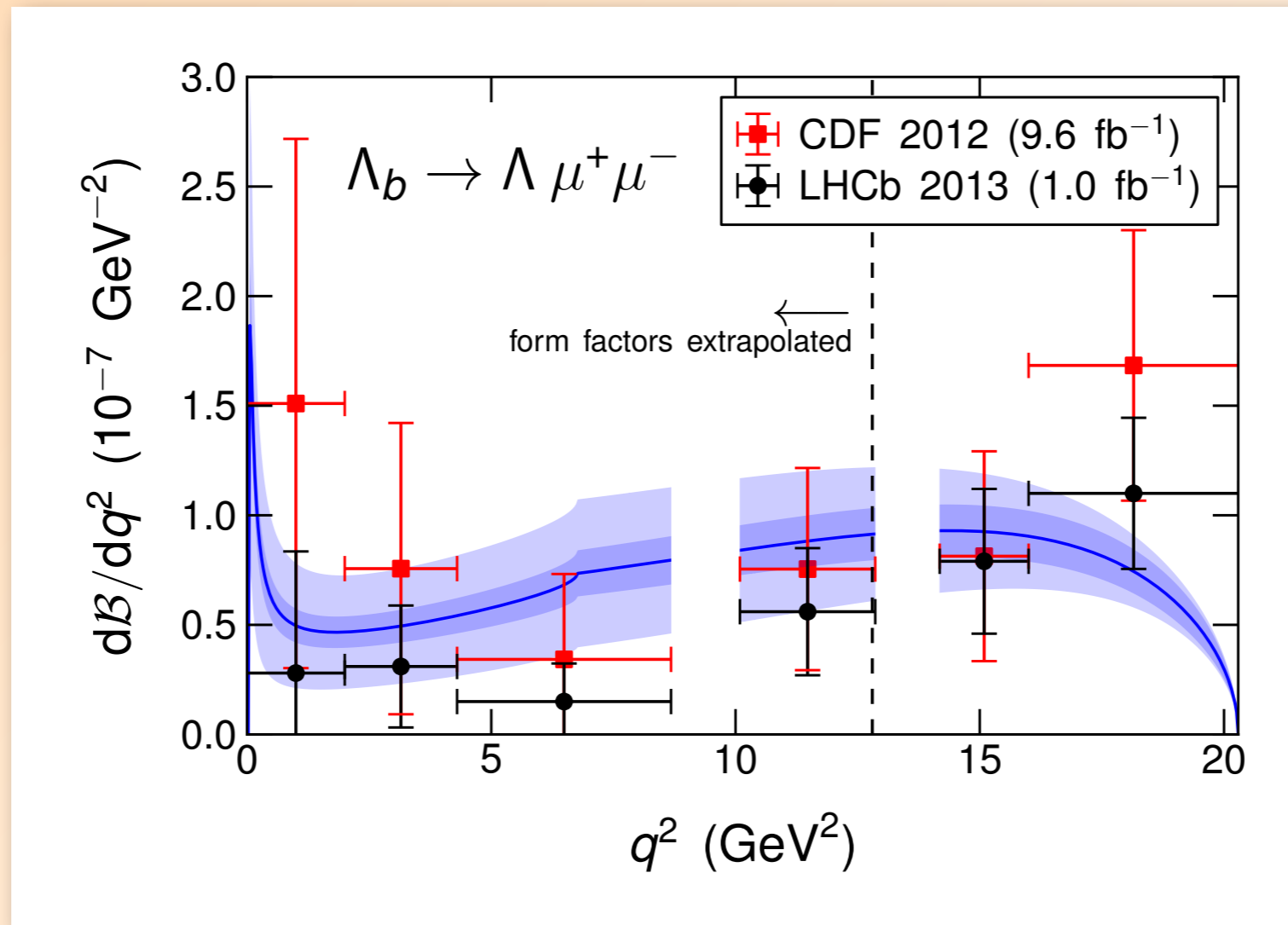
$$F_{\pm} = F_1 \pm F_2$$



Error bands: statistical + 8% lattice systematic uncertainties

figure from S Meinel, Lattice 2013

# Differential b.f. $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$



Note:  $C_9^{\text{NP}} < 0$  would imply a decrease in the b.f.

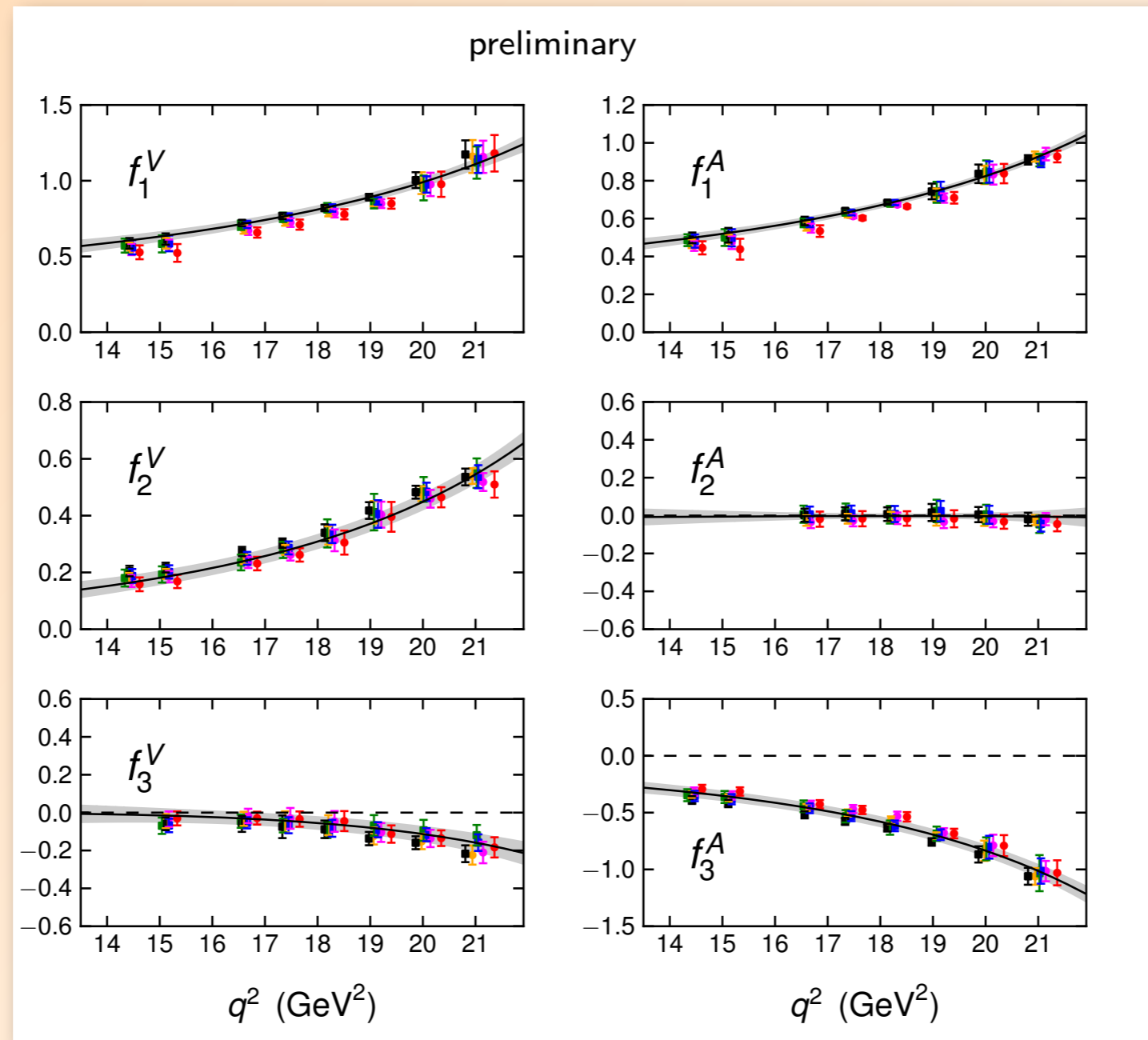


# Beyond the static approximation

- ❖ S Meinel is extending the calculation beyond static action
- ❖ RBC/UKQCD lattice ensembles
- ❖ Relativistic heavy quark action: Fermilab method with nonperturbatively tuned coefficients. (RBC/UKQCD collab)
- ❖ “Mostly nonperturbative” operator matching (Fermilab)
- ❖ Replace pole model with  $z$ -expansion
- ❖ Result: dramatic reduction of systematic uncertainties

# $\Lambda_b \rightarrow p$ form factors

Status of S Meinel's ongoing calculation

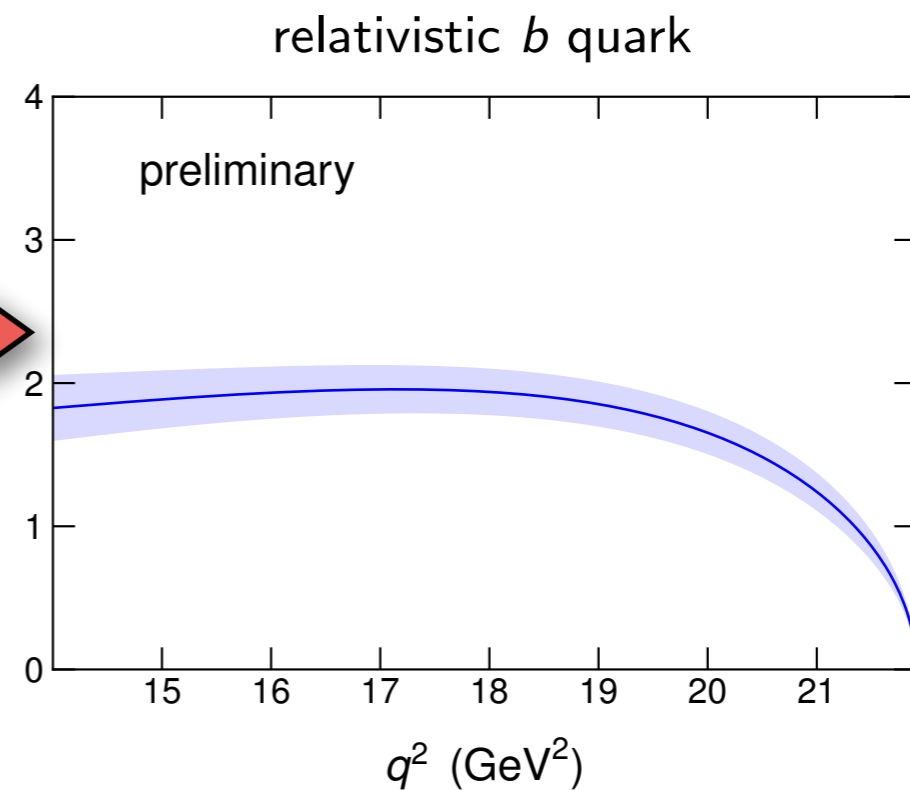
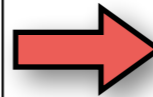
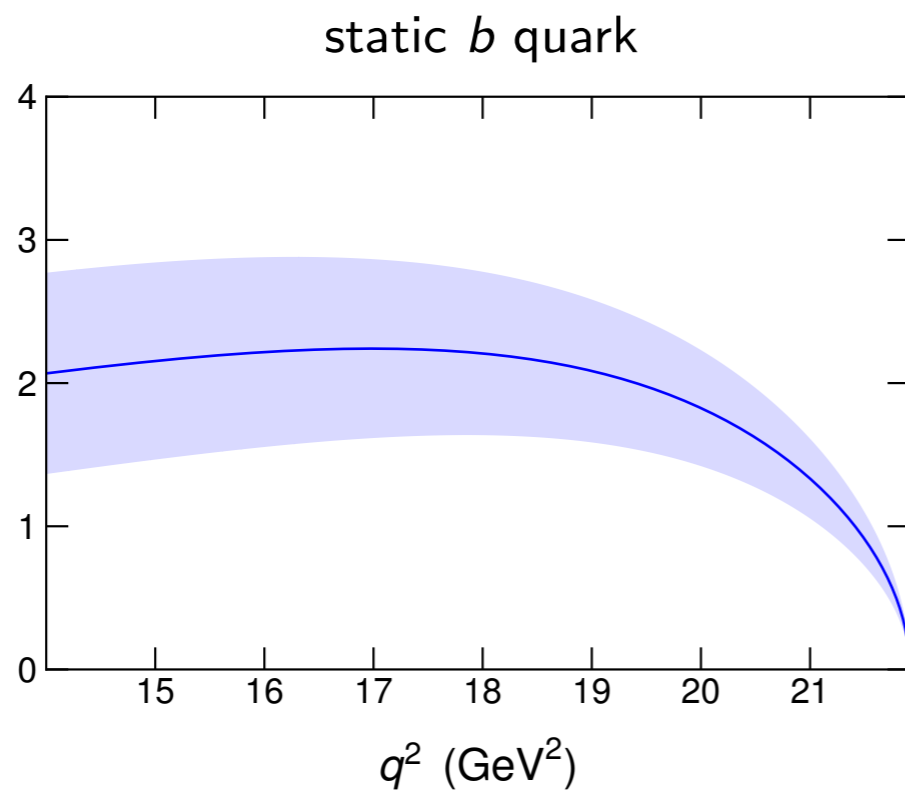


S Meinel at FPCP 2014

# Predicted precision

$$\Lambda_b \rightarrow p \mu^- \bar{\nu}_\mu$$

$$\frac{d\Gamma/dq^2}{|V_{ub}|^2} \text{ (ps}^{-1} \text{ GeV}^{-2}\text{)}$$



Total uncertainties shown.

# Conclusions

- ❖ First lattice calculations of  $\Lambda_b$  form factors
- ❖  $\Lambda_b \rightarrow p l^- \nu$  : Novel method to determine  $|V_{ub}|$
- ❖  $\Lambda_b \rightarrow \Lambda l^+ l^-$  : complements  $B \rightarrow K^{(*)} l^+ l^-$  and  $B_s \rightarrow \varphi l^+ l^-$
- ❖ Look forward to experimental measurements

# References

$$\Lambda_b \rightarrow p$$

W Detmold, C-J D Lin, S Meinel, M Wingate  
Phys. Rev. D 88, 014512 (2013) [arXiv:1306.0446]

$$\Lambda_b \rightarrow p, \text{ etc.}$$

S Meinel  
Lattice 2013 [arXiv:1401.2685], FPCP 2014 slides

$$\Lambda_b \rightarrow \Lambda$$

W Detmold, C-J D Lin, S Meinel, M Wingate  
Phys. Rev. D 87, 074502 (2013) [arXiv:1212.4827]