
$|V_{cb}|$ from $B \rightarrow D^{(*)} \ell \nu$ on the lattice

Daping Du

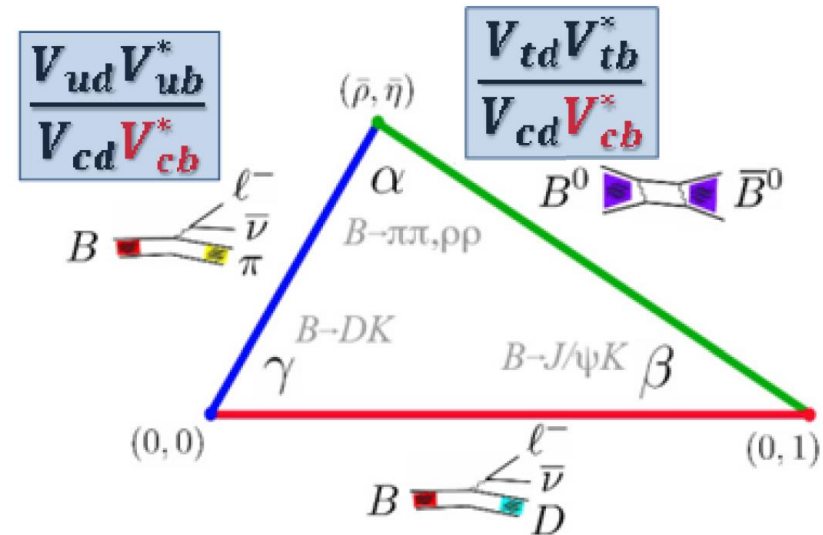
Syracuse University, NY, USA

Sep 8, 2014

CKM 2014, Vienna, Austria

Charmed B semileptonic decays

- CKM matrix element $|V_{cb}|$
 - Supposed to be free from NP effect (dominated by W-exchange)
 - Precision calculation of $|V_{cb}|$ is important to CKM sector of SM.
Normalization of the sides of UT triangle.
 - $|V_{ub}/V_{cb}|$ is of high priority in heavy flavor physics.



- Semileptonic decays:
 - The theory is well studied. HQ expansion Λ/m_b works well.
 - Lots of experimental data are available at B factories (BaBar, Belle,...) . Full-reconstruction of B meson from $\Upsilon(4S)$ decays.
High-precision measurements enable analysis of shapes of q^2 -dependence.
 - Able to probe possible NP contributions lifted by τ mass in processes like $B \rightarrow D^{(*)}\tau\nu$.

The determinations of $|V_{cb}|$

➤ The **inclusive** methods:

Total decay rate + HQE parameters
(1S scheme, kinetic scheme)

Precision : $\sim 2\%$

➤ The **exclusive** methods:

- $B \rightarrow D^* \ell \nu$:

Theory $\mathcal{F}(1)$ + experimental $\mathcal{F}(1)|V_{cb}|$

Precision: $\sim 2\%$

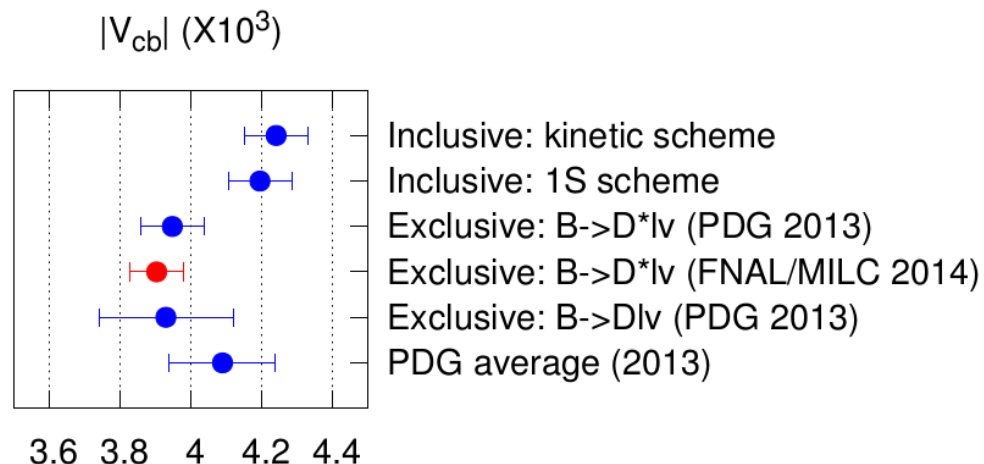
(similar errors from lattice and exp.)

- $B \rightarrow D \ell \nu$:

Theory $\mathcal{G}(w)$ + experimental $\mathcal{G}(w)|V_{cb}|$

Precision: $\sim 5\%$ (dominated by experimental error)

- Discrepancy between the exclusive (lattice) and inclusive $|V_{cb}|$ is more than 2σ . A long standing puzzle.



Exclusive methods: $B \rightarrow D^{(*)} \ell \nu$ decays

$$\frac{d\Gamma}{dw} (B \rightarrow D \ell \nu) =$$

$$w = v \cdot v'$$

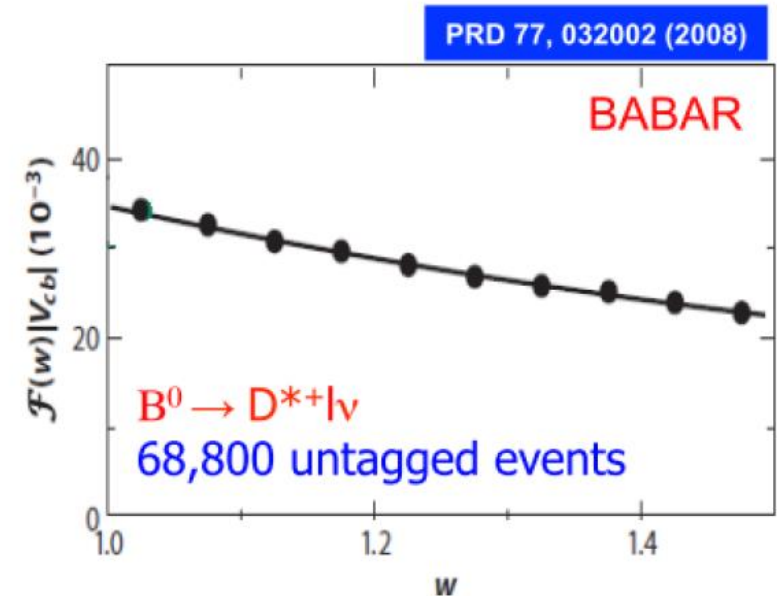
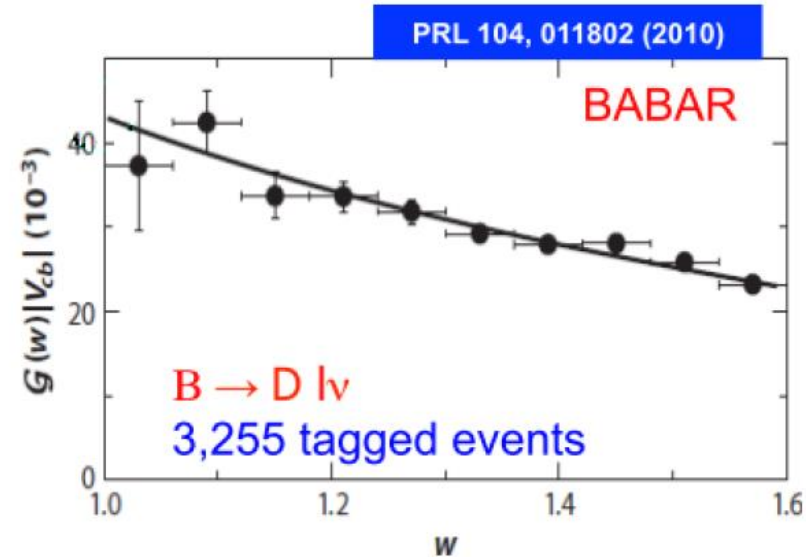
$$\frac{G_F^2 m_B^5}{48\pi^3} |V_{cb}|^2 (1+r)^2 r^3 (w^2 - 1)^{3/2} [\eta_{ew} \mathcal{G}(w)]^2$$

- Noisy experimental data near zero recoil due to suppression at $w=1$. Some model dependence in extrapolation.
- Require non-zero recoil lattice calculation even for zero recoil $G(1)$. Combined fit of lattice and experiments.

$$\frac{d\Gamma}{dw} (B \rightarrow D^* \ell \nu) =$$

$$\frac{G_F^2 m_B^5}{48\pi^3} |V_{cb}|^2 (w^2 - 1)^{1/2} P(w) [\eta_{ew} \mathcal{F}(w)]^2$$

- Lattice calculation at zero recoil benefits from Luke theorem ($\propto 1/m_Q^2$). Experimental extrapolation to $\mathcal{F}(1)|V_{ub}|$ is relatively well under control.
- Non-zero recoil lattice calculation is also desired. An interesting test for consistency of form factor shape.

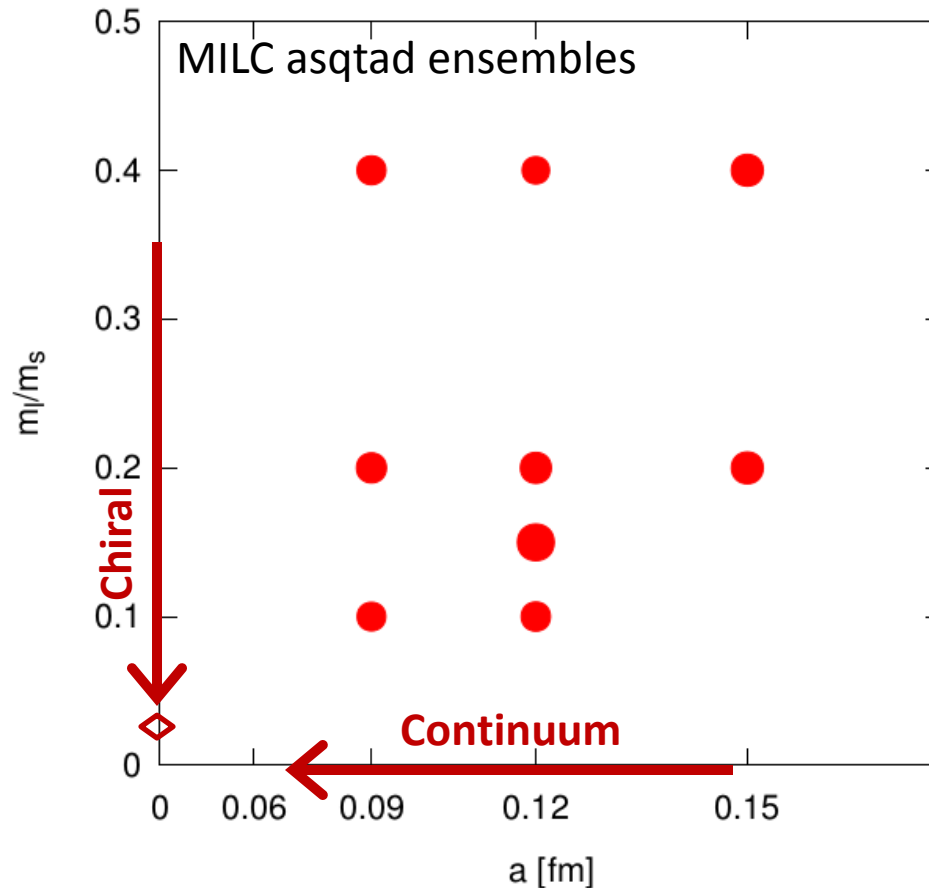


Lattice QCD calculations and errors

- Gauge fields and sea quarks.
 - Number of gauge configurations → **statistical errors**.
 - Implementing dynamical sea quarks (unquenched) is expensive, cost $\propto 1/m_\ell^{1\sim 2}$
Use unphysical heavier masses → **chiral extrapolation error**
 - Small lattice spacings are expensive, cost $\propto 1/a^{4\sim 6}$ → **discretization error**
 - Big volumes are expensive, cost $\propto L^5$ → **finite volume effect**
- Valence quarks and currents
 - Light quarks have the same difficulty as the sea quarks
 - Matching to the continuum currents → **matching error (renormalization)**
 - Input parameters: $g_{DD^*\pi} \dots$ → **systematic errors from inputs**
 - Heavy quarks: $am_c \lesssim 1, am_b \gtrsim 1$ → **heavy quark discretization error**
Fermilab method, NRQCD: error is suppressed $1/m_Q$
Heavier-than-charm staggered?

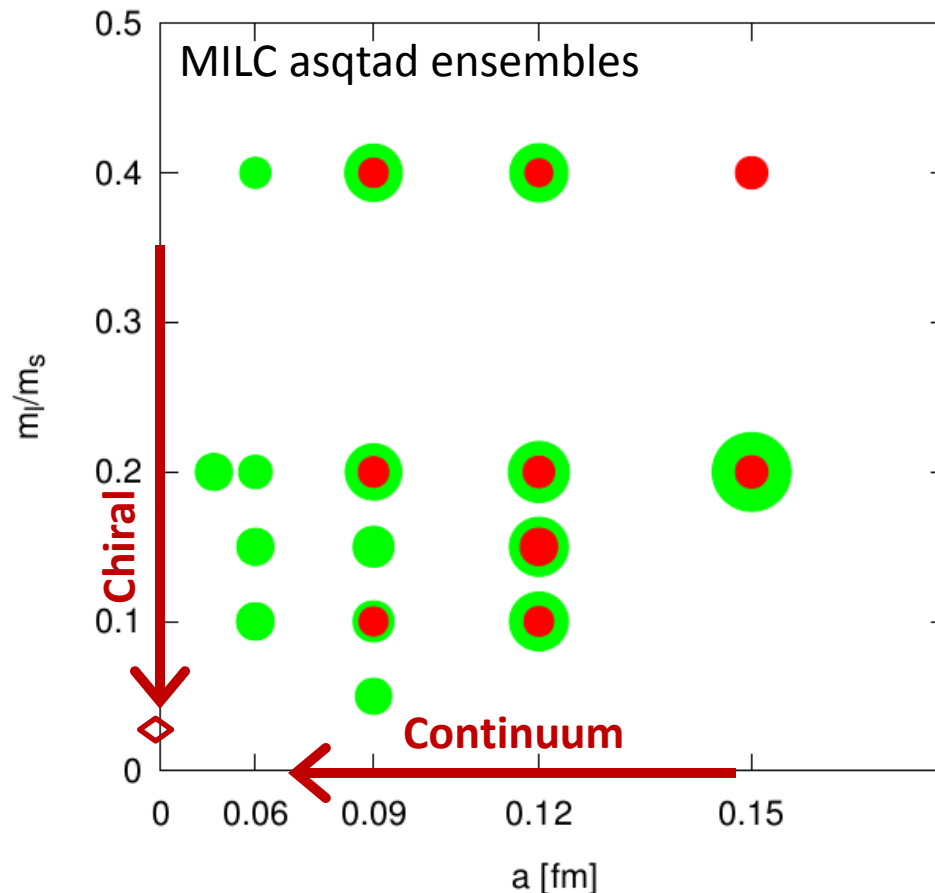
Lattice calculations of $B \rightarrow D^* \ell \nu$

- 2001, quenched calculation, Hashimoto (FNAL) [PRD66\(014503\)](#)
- 2008, quenched calculation at non-zero recoil, de Divitis et al, [NPB807\(373\)](#)
- 2008, **unquenched** (2+1 sea quarks), Laiho (FNAL/MILC), [PRD79\(014506\)](#)



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- 2014, update, Laiho (FNAL/MILC) [PRD89\(114504\)](#)



Lattice calculations of $B \rightarrow D^* \ell \nu$ (FNAL/MILC)

- In the zero recoil limit, only one form factor matters

$$\mathcal{F}(1) = h_{A_1}(1)$$

- Use Fermilab action for heavy b and c quarks.
- Use a double ratio [PRD66\(014503\)](#) that removes all wave-function normalization and the majority of current renormalization

$$\mathcal{R}_{A_1} = \frac{\langle D^* | \bar{c} \gamma^j \gamma^5 b | \bar{B} \rangle \langle \bar{B} | \bar{b} \gamma^j \gamma^5 c | D^* \rangle}{\langle D^* | \bar{c} \gamma^4 c | D^* \rangle \langle \bar{B} | \bar{b} \gamma^4 b | \bar{B} \rangle} = |h_{A_1}(1)|^2$$

- Current renormalization

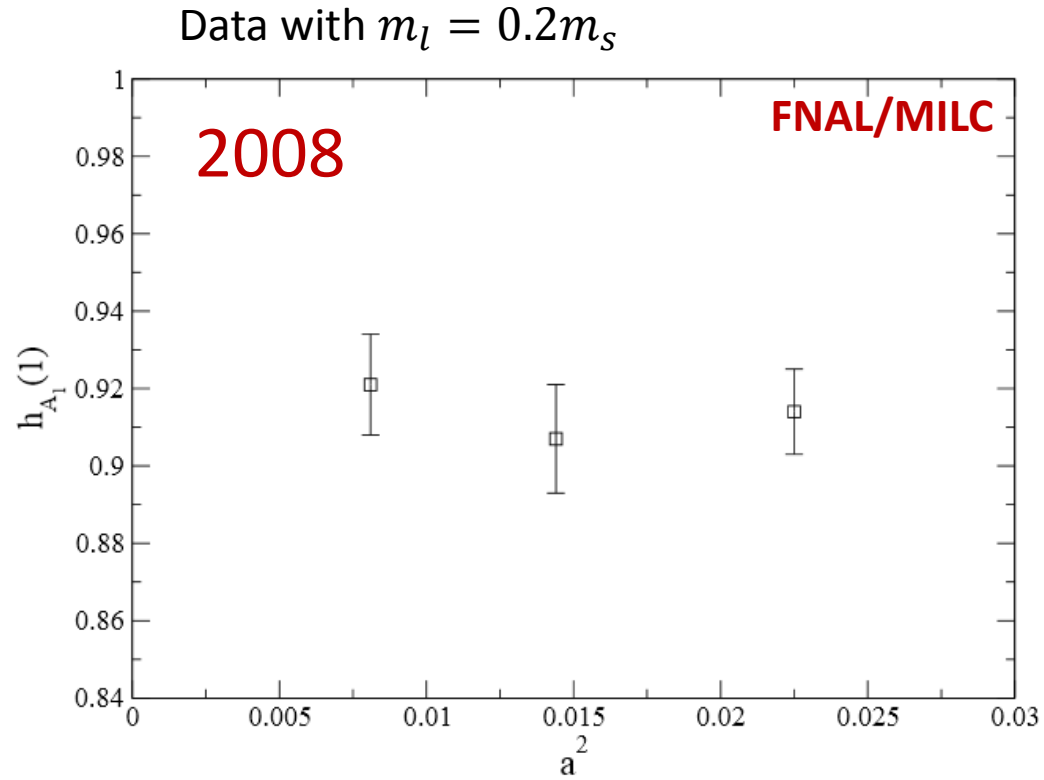
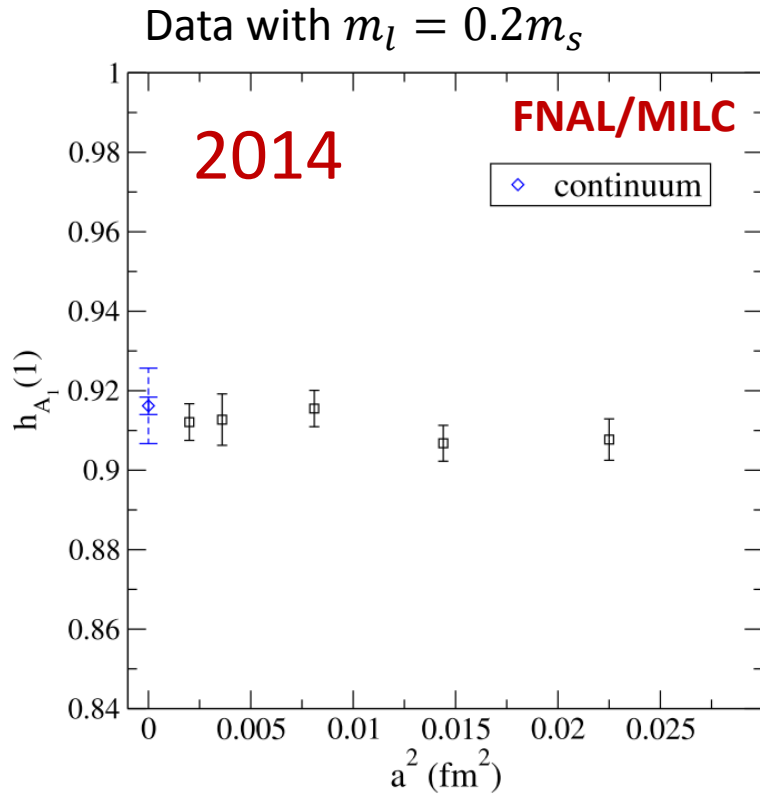
$$\mathcal{J}_{cb}^{\text{con}} = Z_{J_{cb}^\mu} J^{\text{lat}}$$

Majority of Z cancels in the ratio;
Small remaining factor is calculated
Using lattice perturbation theory.

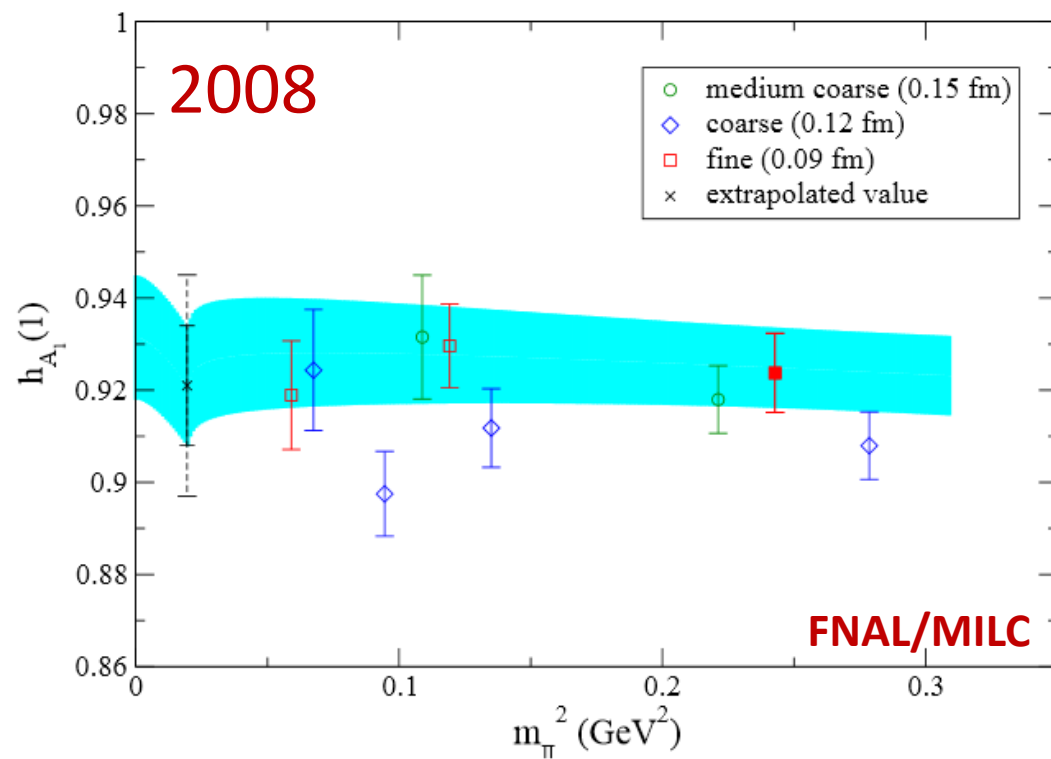
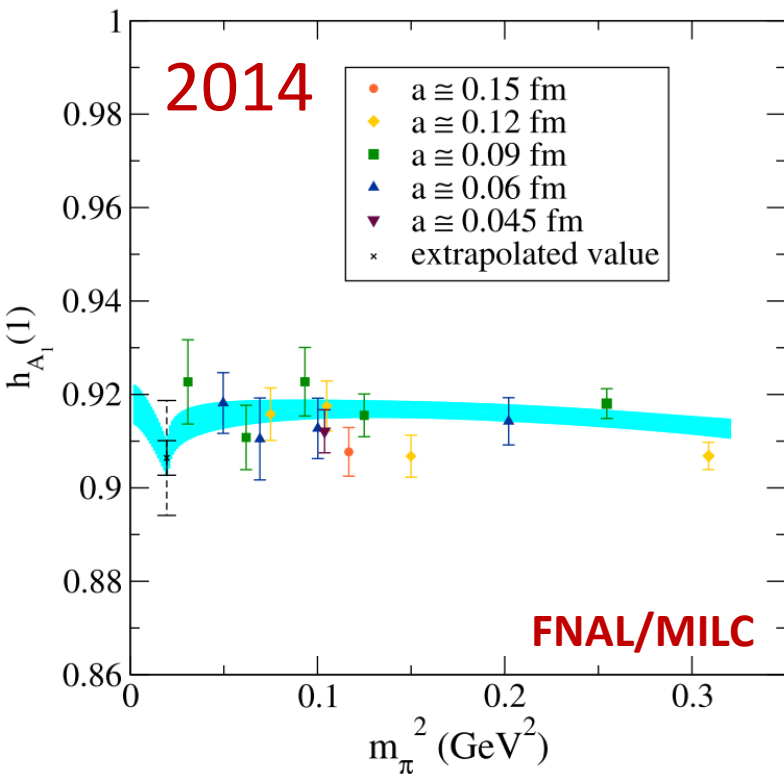
Define

$$\rho_{A^i}^2 = \frac{Z_{A_{cb}^i} Z_{A_{bc}^i}}{Z_{V_{cc}^4} Z_{V_{bb}^4}} \rightarrow \left[1 + \sum_{\ell} \rho_{A^j}^{[\ell]} \alpha_V^\ell(q^*) \right]^2$$

The discretization effect



Chiral/continuum extrapolation



Error budget

FNAL/MILC

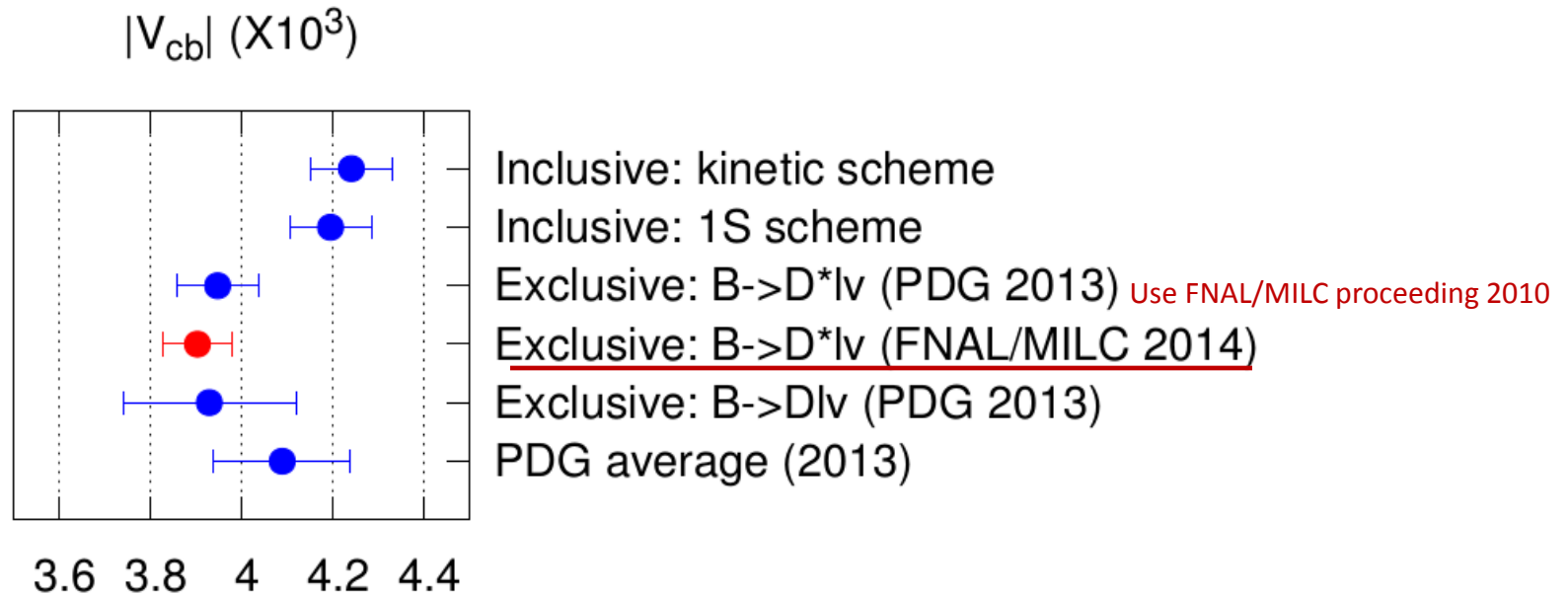
Error of $h_{A_1}(1)$ (%)	2008	2014
Statistics	1.4	0.4
χ PT fits	0.9	0.5
$g_{DD^*\pi}$	0.9	0.3
HQ Discretization errors	1.5	1.0
HQ mass tuning	0.7	-
Perturbation theory	0.3	0.4
Isospin	-	0.1
u_0 tuning	0.4	-
Scale r_1	-	0.1
Total	2.6	1.4

- Significant improvement in the statistical uncertainty.
- HQ discretization error is now the dominant source of error.
- The coupling $g_{DD^*\pi}$ used is 0.53(8) based on [PLB721\(94\)](#), [PLB719\(103\)](#)
- HQ mass tuning is included in “Statistics”.
- No “ u_0 tuning” error in the newer data.

Update of $|V_{cb}|$ (FNAL/MILC)

➤ Summary:

- Full data set analysis is published [PRD89\(114504\)](#)
- Result: $|V_{cb}| = (39.04 \pm 0.49_{\text{expt}} \pm 0.53_{\text{QCD}} \pm 0.19_{\text{QED}}) \times 10^{-3}$
- The discrepancy between inclusive and exclusive methods remains.

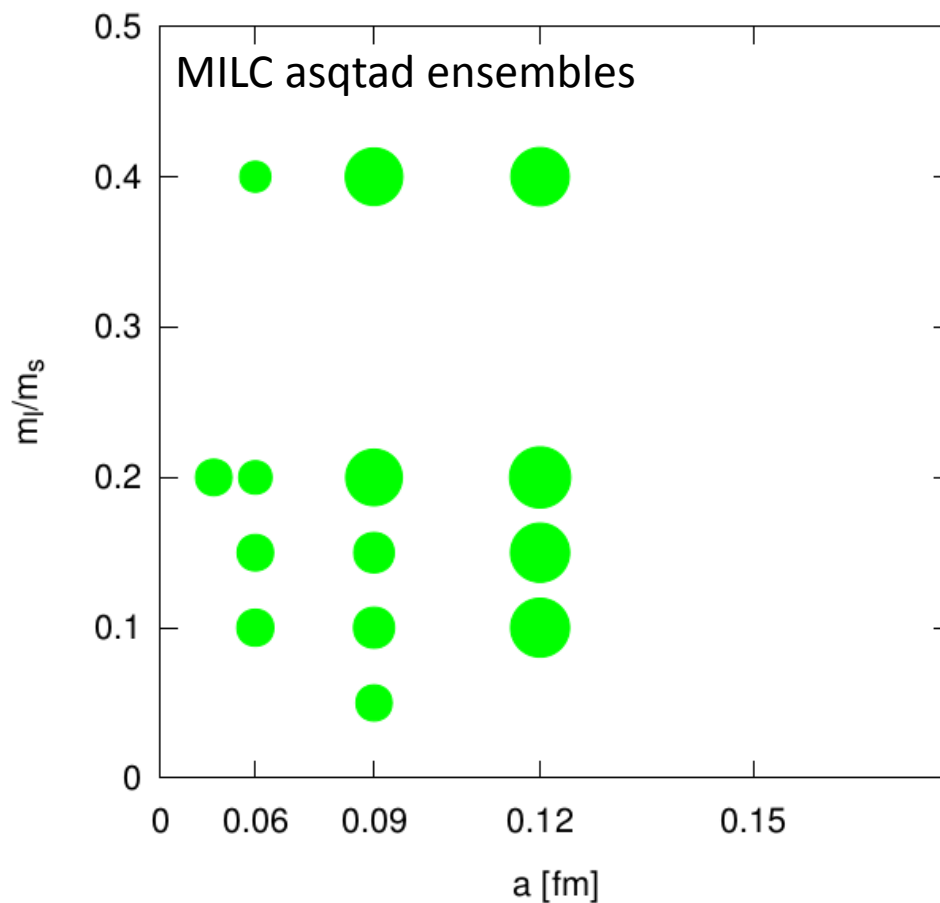


Lattice calculations of $B \rightarrow D\ell\nu$

- 1999, quenched, zero recoil, Hashimoto (FNAL), [PRD61\(014502\)](#)
- 2005, unquenched, zero recoil, Okamoto (FNAL/MILC) [hep-lat/0409116](#)
2+1 asqtad (old) + Fermilab b : $\mathcal{G}(1) = 1.074(18)(16)$
- 2012, unquenched, non-zero recoil, Qiu (FNAL/MILC) [1312.0155](#)
2+1 asqtad(new) + Fermilab b :
Fit using full w -range with experiments to determine $|V_{cb}|$
- 2012, unquenched, non-zero recoil, Na (HPQCD) [Lattice 2014](#)
2+1 asqtad + hisq light valence and charm + NRQCD b :
Ratio $f_0(B_s \rightarrow D_s\ell\nu)/f_0(B \rightarrow D\ell\nu)$ and $|V_{cb}|$
- 2013, unquenched, near zero recoil, Atoui et al, [1311.5071](#)
2+0 twisted mass Wilson, $B_{(s)} \rightarrow D_{(s)}\ell\nu$

Lattice calculations of $B \rightarrow D\ell\nu$ (FNAL/MILC)

- Use a similar set of lattice configuration as the FNAL/MILC $B \rightarrow D^*\ell\nu$
- Use Fermilab bottom and charm quarks; vector current
- Four non-zero momenta for the daughter D meson.



Moving to non-zero recoil

- The form factor

$$\mathcal{G}(w) = h_+(w) + \frac{1-r}{1+r} h_-(w)$$

$h_-(1) \rightarrow 0$, in HQ limit, but not in general
 $h_-(1)$ has to be extrapolated from $h_-(w)$

- Double ratio for $h_+(1)$

$$\mathcal{R}_+ = \rho_{V^4}^2(1) \frac{\langle D|\bar{c}\gamma^4|\bar{B}\rangle\langle\bar{B}|\bar{b}\gamma^4 c|D\rangle}{\langle D|\bar{c}\gamma^4 c|D\rangle\langle\bar{B}|\bar{b}\gamma^4 b|\bar{B}\rangle} = |h_+(1)|^2$$

- Single ratios for $h_+(w)$ and $h_-(w)$

$$h_+(w(\mathbf{p})) = \sqrt{\mathcal{R}_+} Q_+(\mathbf{p}) \left[1 - \mathcal{R}_-(\mathbf{p}) \cdot \mathbf{x}_f(\mathbf{p}) \right],$$

$$h_-(w(\mathbf{p})) = \sqrt{\mathcal{R}_+} Q_+(\mathbf{p}) \left[1 - \frac{\mathcal{R}_-(\mathbf{p}) \cdot \mathbf{x}_f(\mathbf{p})}{x_f^2(\mathbf{p})} \right],$$

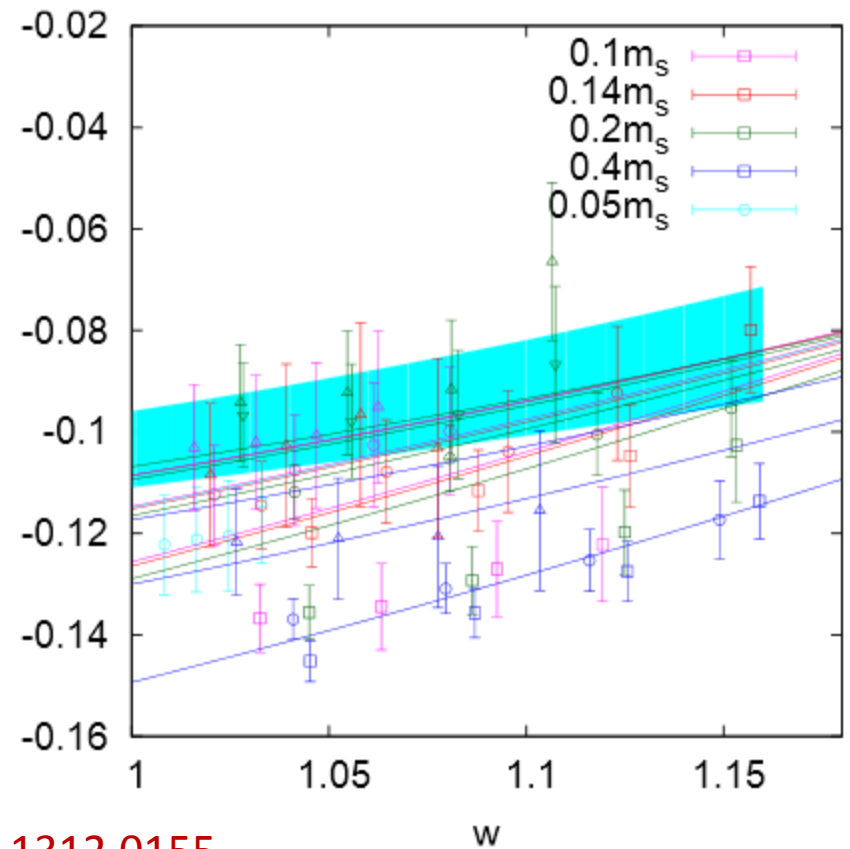
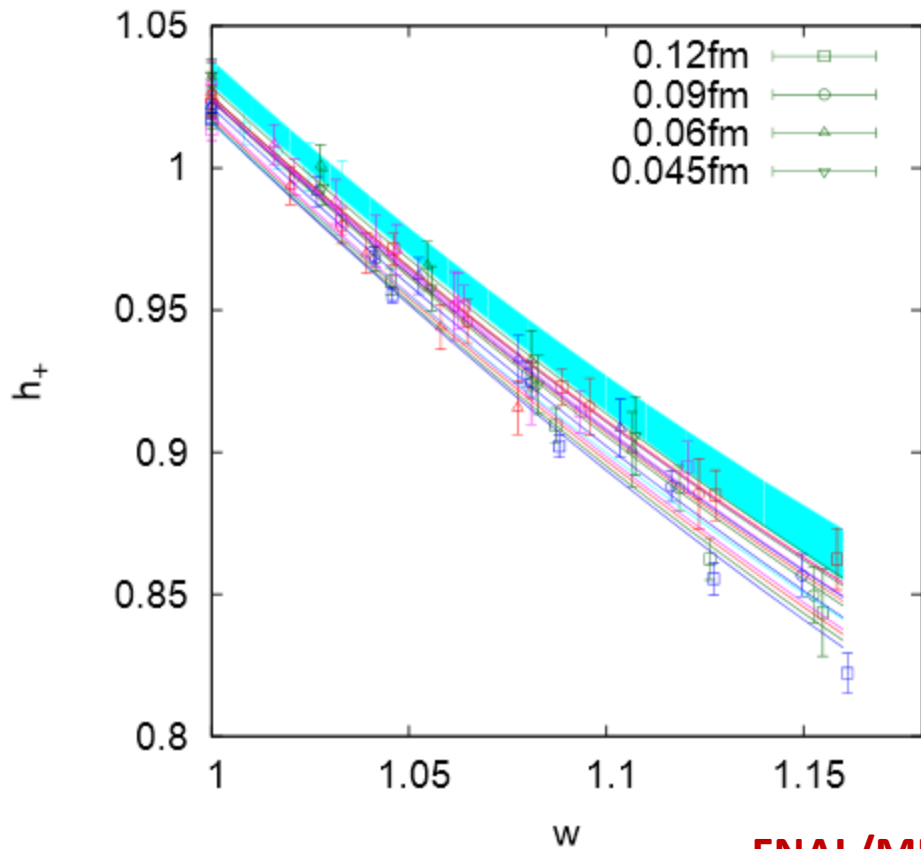
← Renormalized
 Q_+, R_-, x_f

$$Q_+(\mathbf{p}) \equiv \frac{\langle D(\mathbf{p})|V^4|B(\mathbf{0})\rangle}{\langle D(\mathbf{0})|V^4|B(\mathbf{0})\rangle}$$

$$x_f^i(\mathbf{p}) \equiv \frac{\langle D(\mathbf{p})|V^i|D(\mathbf{0})\rangle}{\langle D(\mathbf{p})|V^4|D(\mathbf{0})\rangle}$$

$$R_-^i(\mathbf{p}) \equiv \frac{\langle D(\mathbf{p})|V^i|B(\mathbf{0})\rangle}{\langle D(\mathbf{p})|V^4|B(\mathbf{0})\rangle}$$

Chiral/continuum extrapolation



FNAL/MILC 1312.0155

Lattice error budget (FNAL/MILC, preliminary)

- The dominant source of errors:
 - statistical, HQ discretization and (perturbative) renormalization .
- The perturbative renormalization ρ for the spatial component of the vector current is difficult to calculate (in progress).
- Result (2013 lattice conference, **preliminary**):
$$|V_{cb}| = (38.5 \pm 1.9_{\text{expt+lat}} \pm 0.2_{\text{QED}}) \times 10^{-3}$$
- Uncertainty is dominated by experimental uncertainty.

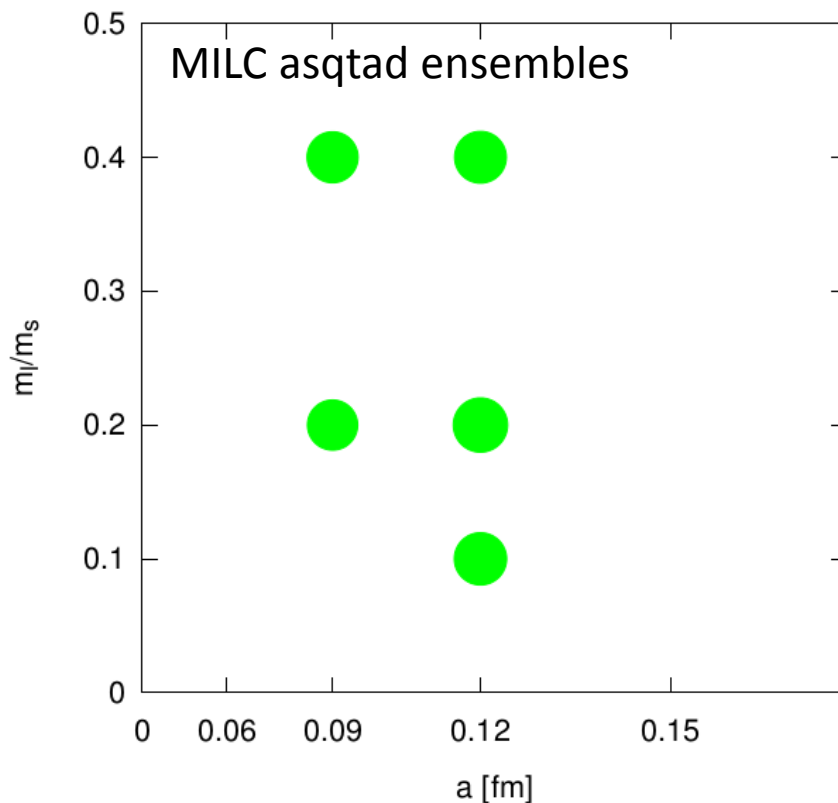
Take only ~5% in $G(w)$, $w \sim 1$

FNAL/MILC 1312.0155

source	$h_+(\%)$	$h_-(\%)$
Statistics+ χ PT	≤ 1	≤ 11
κ -tuning adjustment	≤ 0.1	1.4
Lattice scale r_1	0.2	≤ 0.1
heavy-quark discretization	2.0	10.
ρ factor	0.4	20.
Net systematic error	2.1	22.

Lattice calculations of $B \rightarrow D\ell\nu$ (HPQCD)

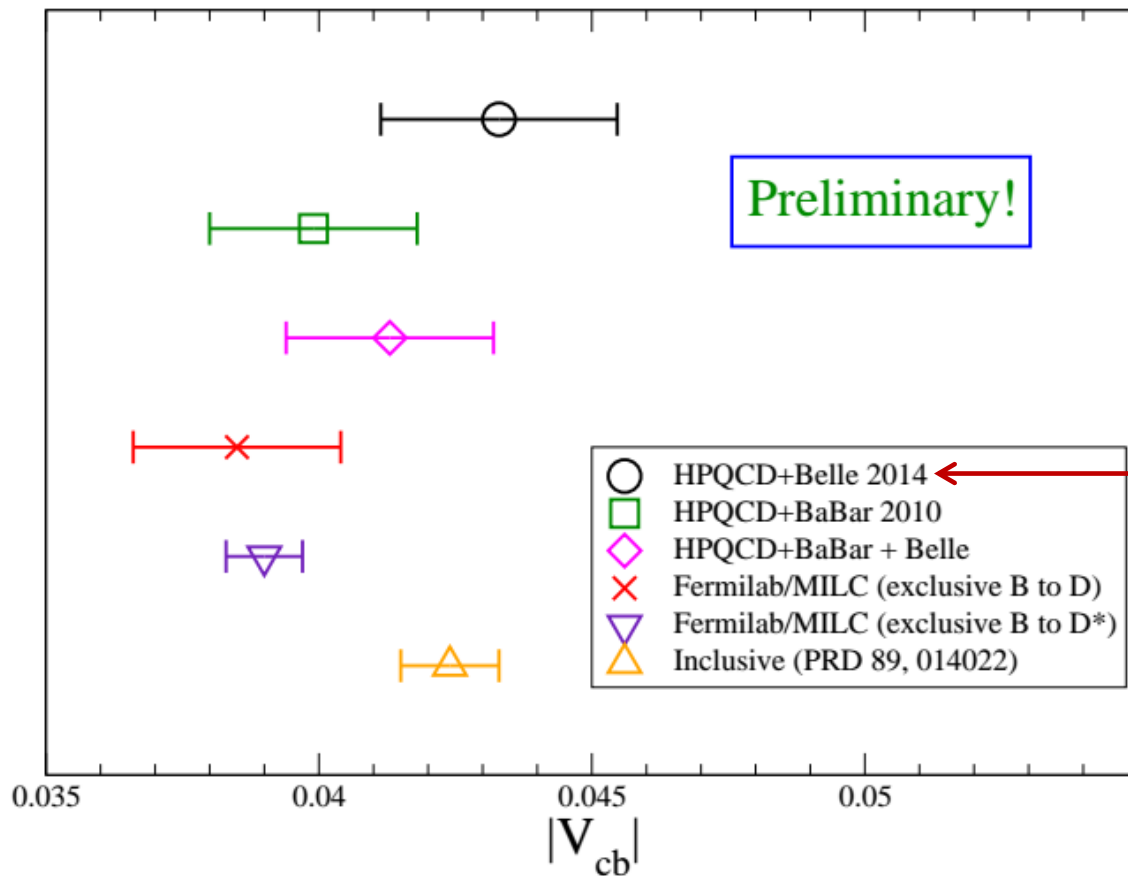
- Use a subset of the MILC asqtad ensembles
- NRQCD bottom quark, hisq light (valence) and **charm** quarks (relativistic).
- Lattice vector current correction: $O(\alpha_S, \Lambda/m_b, \alpha_S/am_b)$ at one-loop perturbation.



Preliminary result and errors

➤ Dominant sources of error:

Current matching error ($\sim 4\%$), statistics, discretization, z-expansion

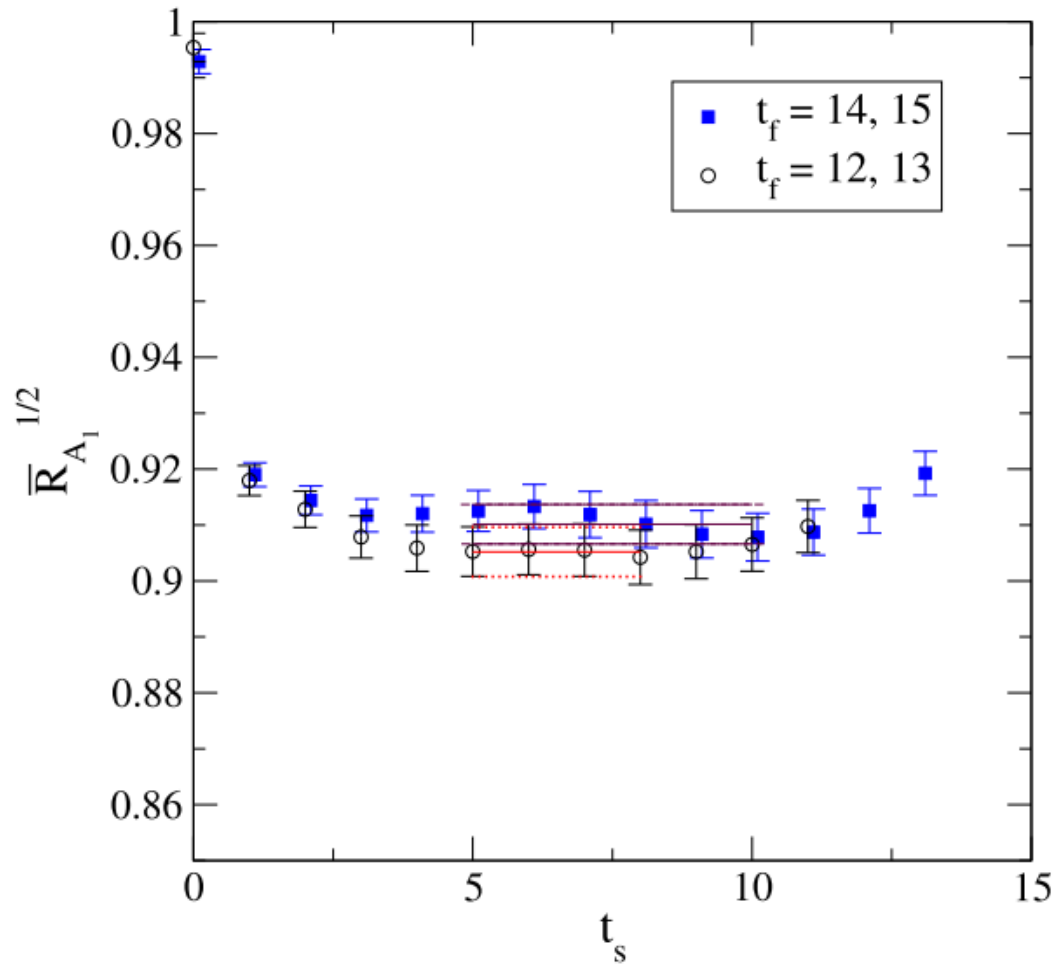


Robin Glattauer, ICHEP 2014

Summary and outlook

- The exclusive result of $|V_{cb}|$ determination is dominated by the $B \rightarrow D^* \ell \nu$ decay channel. Lattice error is commensurate with the experimental error. Using the FNAL/MILC lattice result, the $|V_{cb}|$ from $B \rightarrow D^* \ell \nu$ decay is $0.03905(49)(53)(19)$, with a total uncertainty of 1.9%.
- The determination of $|V_{cb}|$ using $B \rightarrow D \ell \nu$ decay at non-zero recoil is in progress. The preliminary results from FNAL/MILC and HPQCD are consistent with the determination from $B \rightarrow D^* \ell \nu$ decay.
- New analysis results from Belle is important for the crosscheck.
- A lattice calculation of $B \rightarrow D^* \ell \nu$ decay at non-zero recoil is desired.

Backup: excited state effect is small



Backup: hyperfine splittings

