

# World average and experimental overview of $\gamma$

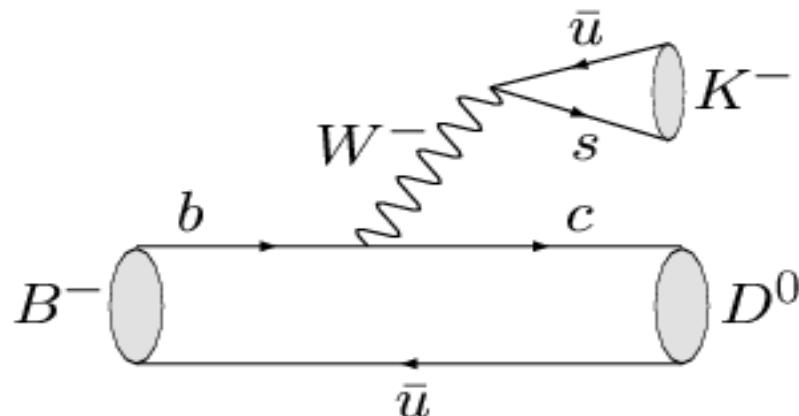
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09/09/2014



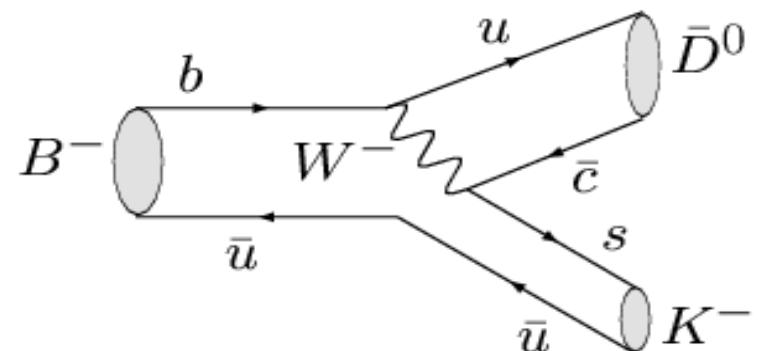
# $\gamma$ measurements from $B^\pm \rightarrow D\bar{K}^\pm$

[see J.Brod's talk]

- Theoretically pristine  $B \rightarrow D\bar{K}$  approach
- Access  $\gamma$  via interference between  $B^- \rightarrow D^0 K^-$  and  $B^- \rightarrow \bar{D}^0 K^-$



color allowed  
 $B^- \rightarrow D^0 K^- \sim V_{cb} V_{us}^*$   
 $\sim A \lambda^3$



color suppressed  
 $B^- \rightarrow \bar{D}^0 K^- \sim V_{ub} V_{cs}^*$   
 $\sim A \lambda^3 (\rho + i \eta)$

relative magnitude of suppressed amplitude is  $r_B$

$$r_B = \frac{|A_{\text{suppressed}}|}{|A_{\text{favoured}}|} \sim \frac{|V_{ub} V_{cs}^*|}{|V_{cb} V_{us}^*|} \times [\text{color supp}] = 0.1 - 0.2$$

relative weak phase is  $\gamma$ , relative strong phase is  $\delta_B$

# $\gamma$ measurements from $B^\pm \rightarrow D K^\pm$

- Reconstruct D in final states accessible to both  $D^0$  and  $\bar{D}^0$ 
  - $D = D_{CP}$ , CP eigenstates as  $K^+ K^-$ ,  $\pi^+ \pi^-$ ,  $K_S \pi^0$

**GLW method (Gronau-London-Wyler)**

- $D = D_{sup}$ , Doubly-Cabbibo suppressed decays as  $K\pi$

**ADS method (Atwood-Dunietz-Soni)**

- Three-body decays as  $D \rightarrow K_S \pi^+ \pi^-$ ,  $K_S K^+ K^-$

**GGSZ (Dalitz) method (Giri-Grossman-Soffer-Zupan)**

- Largest effects due to

- charm mixing



small and can be taken into account

Y.Grossman, A.Soffer, J.Zupan [PRD 72, 031501 (2005)]

- charm CP violation

M.Rama [PRD 89, 014021 (2014)]

- ultimate theoretical error  $< 10^{-7}$  [J.Brod and J.Zupan, arXiv:1308.5663]

- **Different B decays ( $D K$ ,  $D^* K$ ,  $D K^*$ )**

- **different hadronic factors ( $r_B$ ,  $\delta_B$ ) for each**

## The small $r_B$ issue

clearly in the  $r_B \rightarrow 0$  limit the interference disappears and there is no sensitivity to the phase  $\gamma$

when the true value of  $r_B$  is small, then the distribution of  $\hat{r}_B$  best fit values for randomly generated data is biased towards larger values, until the experimental errors are sufficiently small to exclude the  $r_B \sim 0$  region

on the other hand the error on  $\gamma$  is roughly proportional to  $1/r_B$ , hence for small  $r_B$  it is biased towards smaller values

in the language of frequentist statistics it means that the usual  $\Delta \ln \mathcal{L} = 1/2$  rule does not work here, the 68%CL interval extracted from it does not cover the true value of  $\gamma$  at 68% frequency (undercoverage)

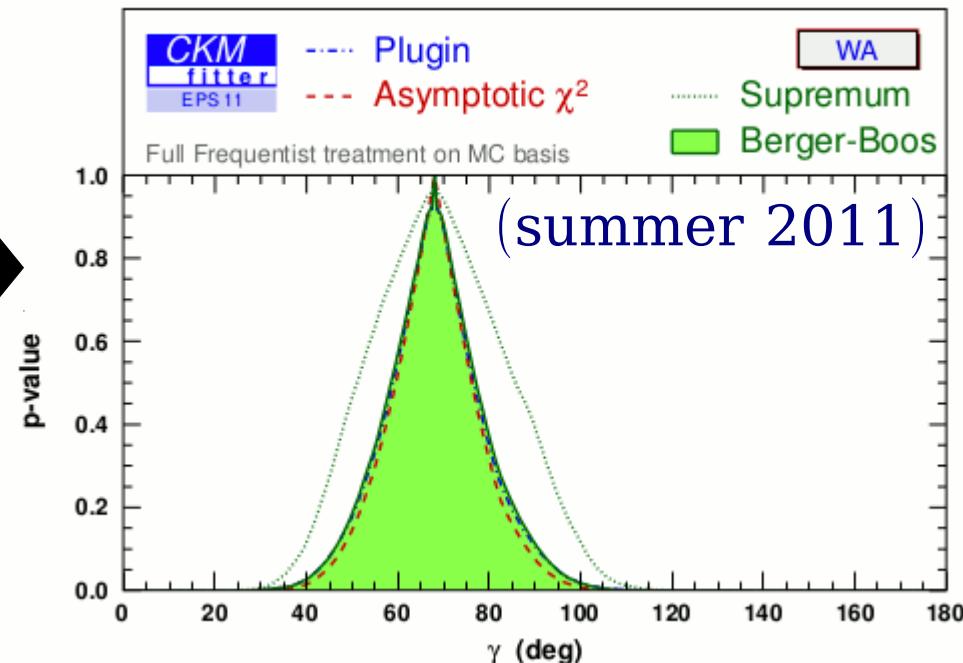
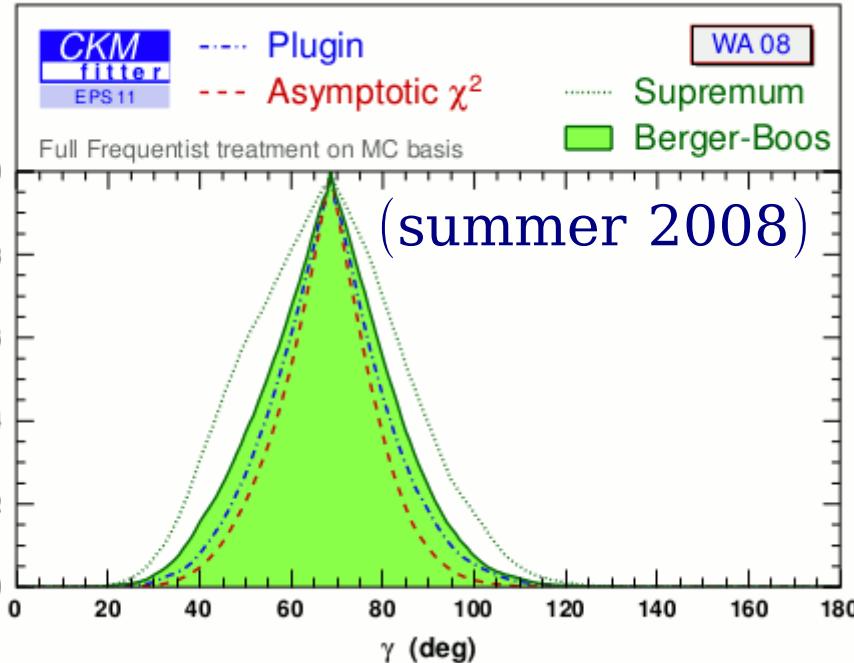
to correct for this effect one has to compute the actual distribution of the profile log-likelihood, and from that distribution deduce a p-value or a CL interval

problem: as soon as the log-likelihood is not distributed as a  $\chi^2$ , its distribution *a priori* depends on the *nuisance parameters*, namely  $r_B$ ,  $\delta_B$  etc.

# Different treatments of the nuisance parameters

to compute the distributions of the log-likelihood depending on the nuisance parameters  $\nu$ :

1. use the best fit estimate,  $\nu = \hat{\nu}$  (plugin method): coverage not guaranteed if the true value of  $\nu$  is different from  $\hat{\nu}$
2. use the worst-case distribution: maximize the p-value over all possible values of  $\nu$  (supremum method): coverage or overcoverage guaranteed by construction
3. maximize the p-value-a well - chosen subspace  $\nu \in N$  (constrained supremum method, e.g. Berger-Boos)



$r_B$ 's away from zero , dependence wrt nuisance parameters is now weak

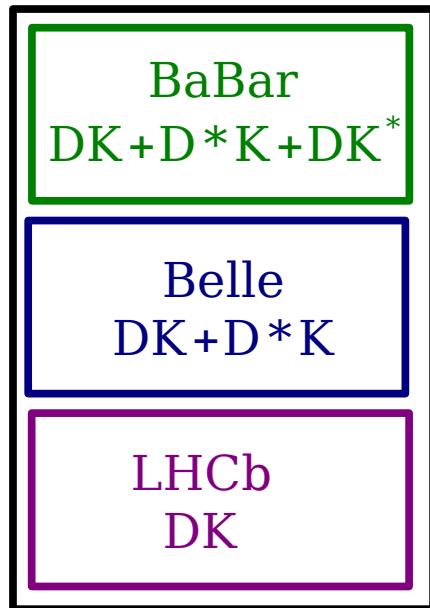
## Outline:

(preliminary results)  
(plugin method)

- start by inputs from  $\gamma$  observables only:
  - GGSZ observables of diff. exp.,  $\gamma$  from GGSZ
  - prediction of GLW observables,  $\gamma$  from GGSZ+GLW
- "charm" analysis: get the  $K\pi$ ,  $K\pi\pi^0$ ,  $K3\pi$  strong phases and other hadronic parameters
  - comparison to CLEO-c results
  - pull of charm observables
- add ADS observables:
  - prediction of ADS observables,  $\gamma$  from GGSZ+GLW+ADS
  - $\delta_D$  discussion (from charm observables, from  $\gamma$  observables...)
- new observables from LHCb:  $DK^{*0}$  (GLW, ADS)...
- Summary

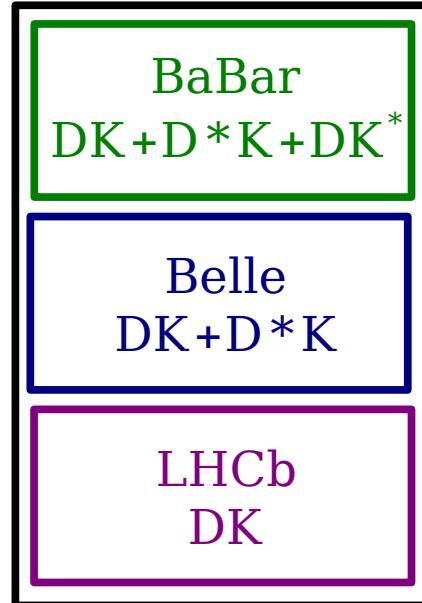
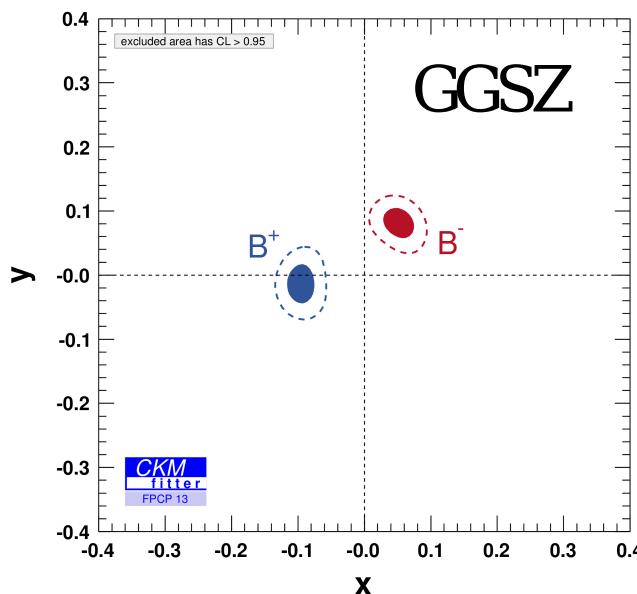
(r discussion:  $DK$ ,  $D\pi$ ,  $D^*\rho$ ,  $D_s K$ )

# GGSZ, GLW, ADS in charged B decays

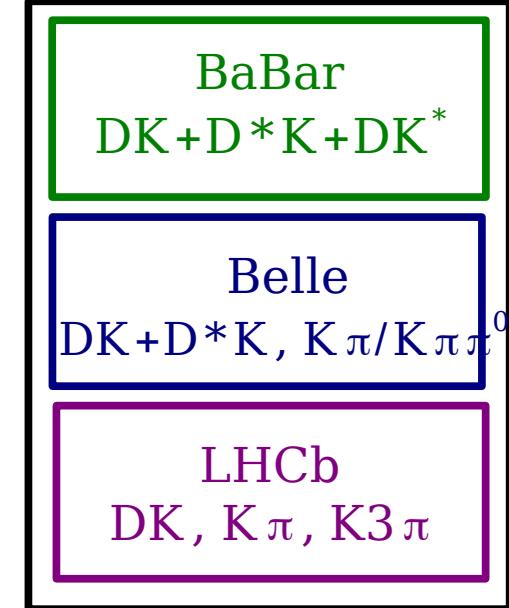
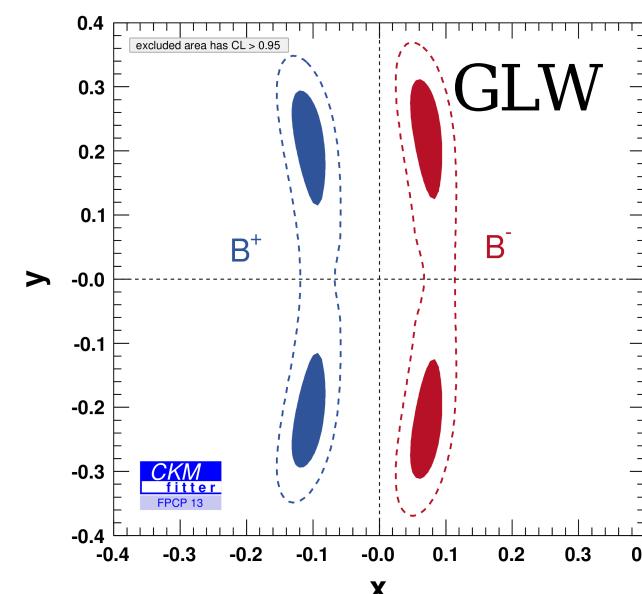


**GGSZ**

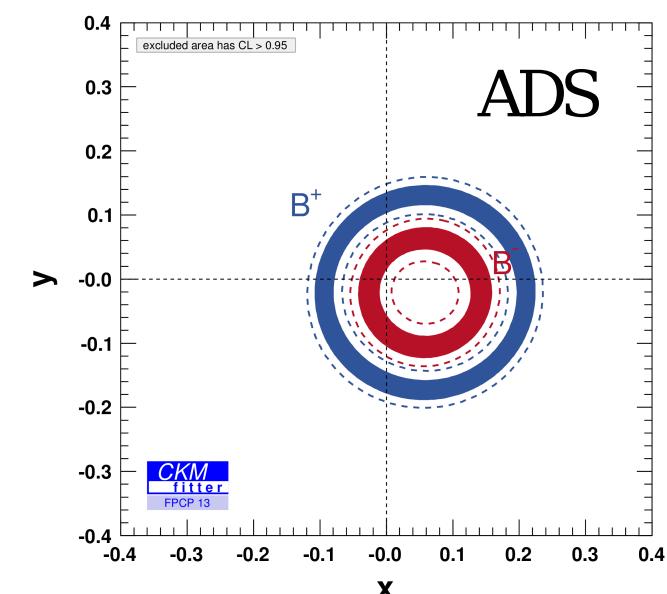
$$x_{\pm} = r_B \cos(\delta_B \pm \gamma) \text{ and } y_{\pm} = r_B \sin(\delta_B \pm \gamma) \quad [\text{just for illustration}]$$



**GLW**



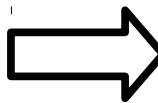
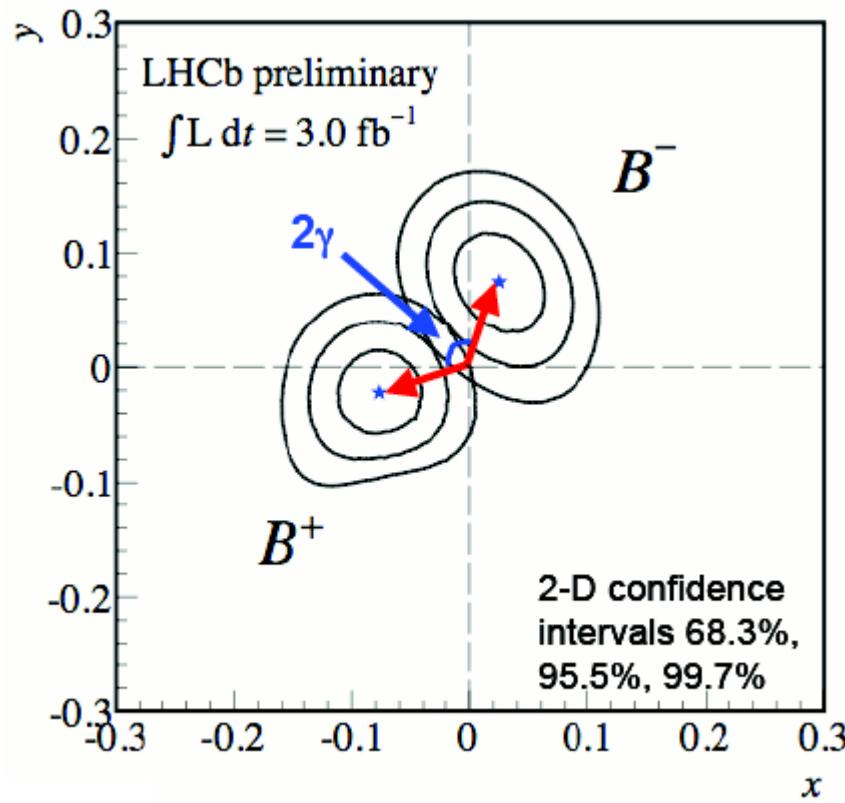
**ADS**



# GGSZ update from LHCb

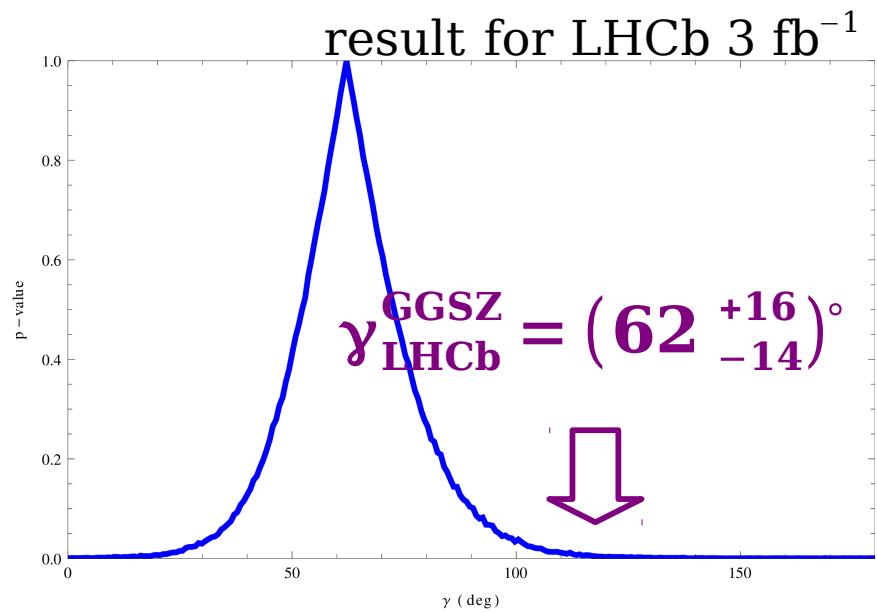
[see M.Karbach's talk]

DK [arXiv:1408.2748]



lower and lower...

$\gamma = (62^{+15}_{-14})^\circ$   
 $r_B = (8.0^{+1.9}_{-2.1}) \times 10^{-2}$   
 $\delta_B = (134^{+14}_{-15})^\circ$

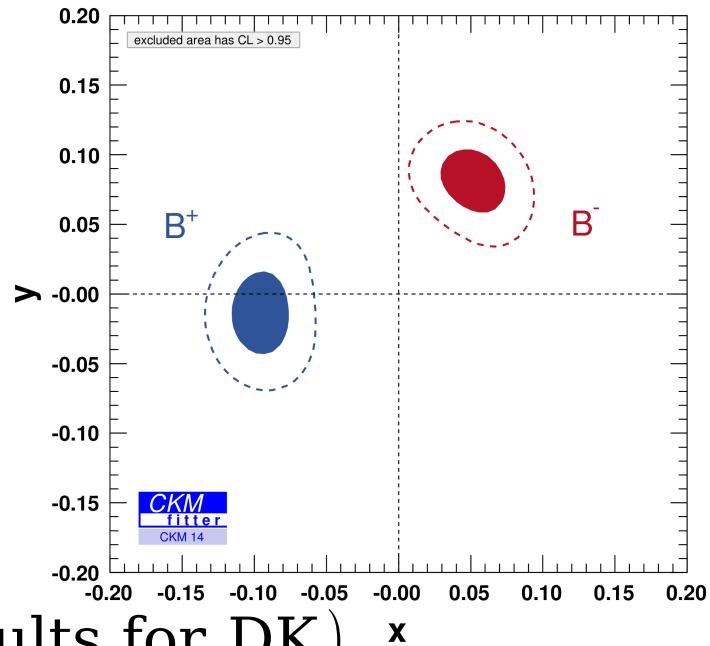


result for LHCb 1+2  $\text{fb}^{-1}$ , LHCb-CONF-2013-004

$\gamma_{\text{LHCb}}^{\text{GGSZ}} = (57 \pm 16)^\circ$

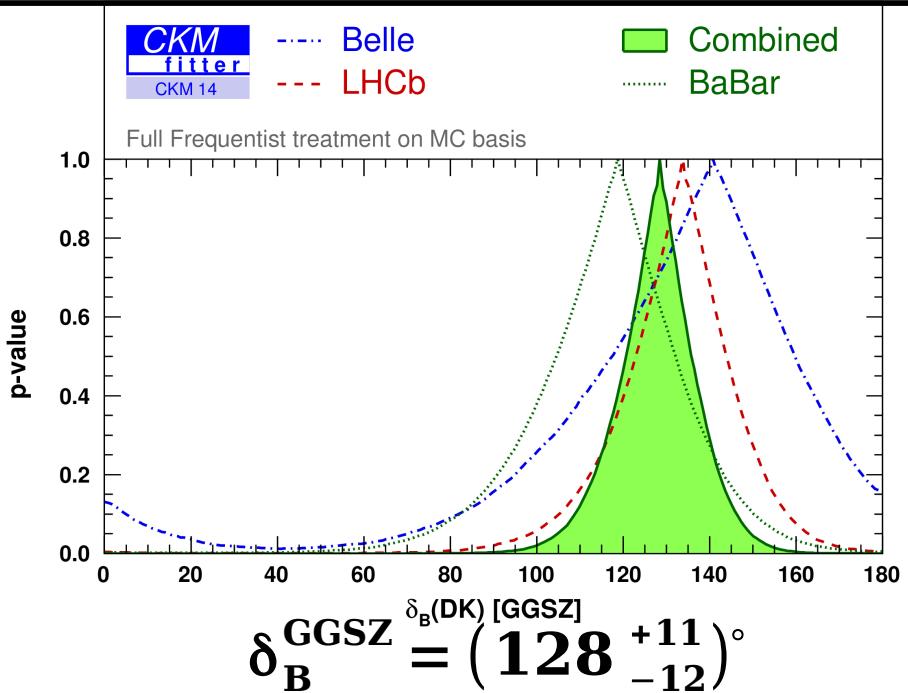
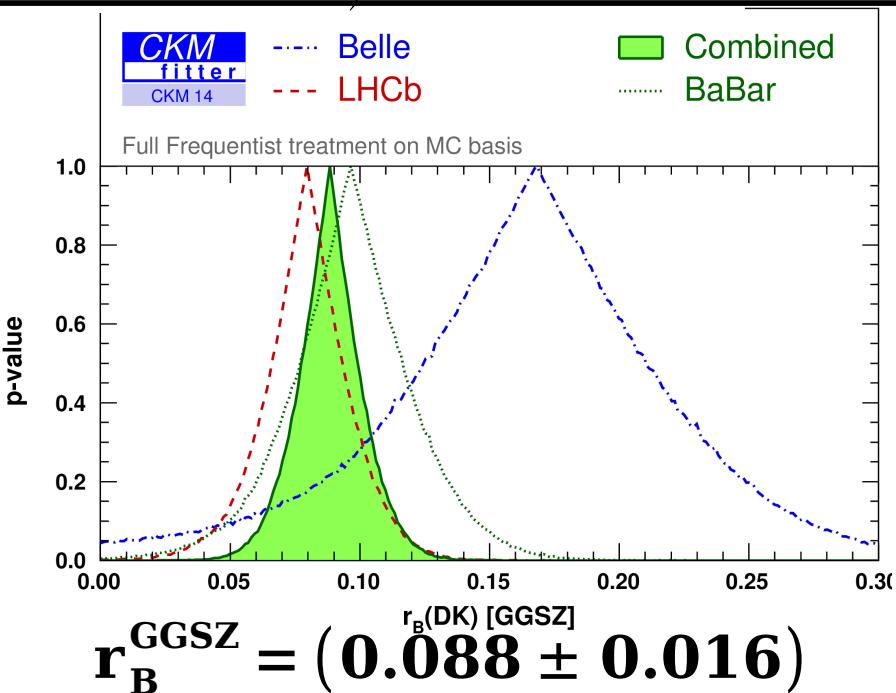
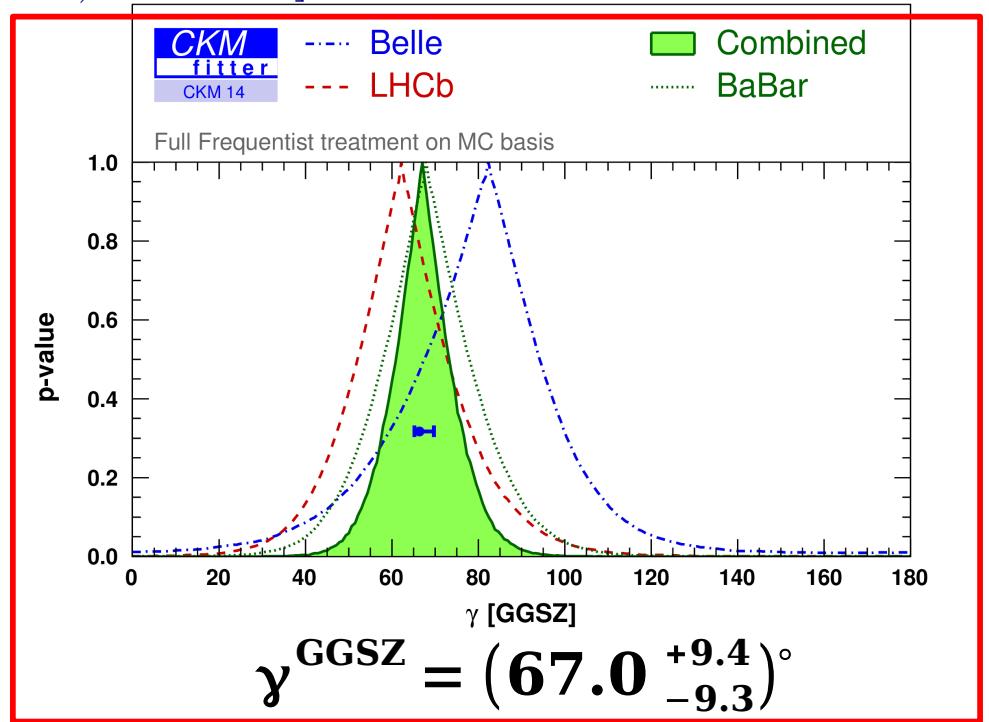
# GGSZ

24 obs.



(results for DK)

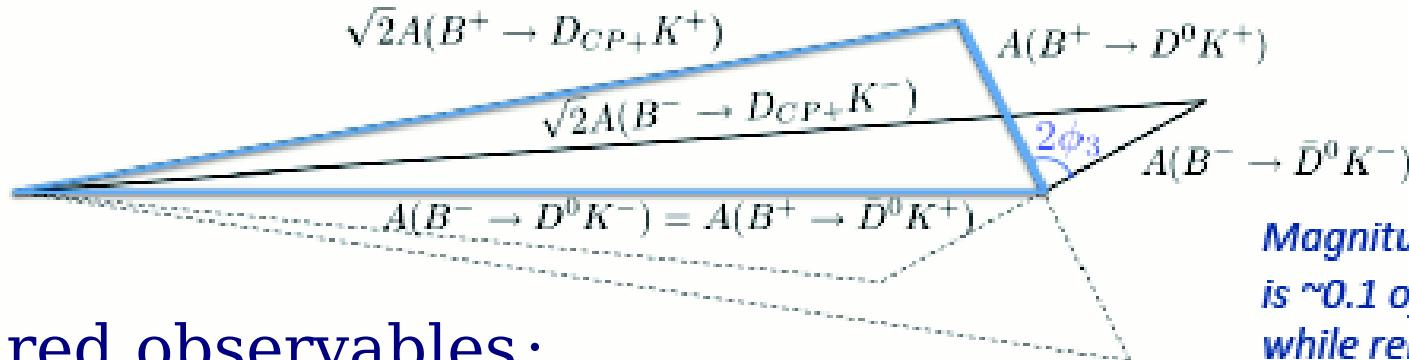
$\left. \begin{array}{l} \text{DK, } D^*K \text{ and } DK^* \text{ [PRL 105 (2010) 121801]} \\ \text{DK and } D^*K \text{ [PRD 81 (2010) 112002]} \\ \text{DK [arXiv:1408.2748]} \end{array} \right\}$  model dep.



# GLW with $D_{CP}^{(*)} K$

D decays to CP eigenstates

➤ Amplitude triangle:



measured observables:

$$R_{CP\pm} \equiv \frac{\text{Br}(B^- \rightarrow D_{CP\pm} K^-) + \text{Br}(B^+ \rightarrow D_{CP\pm} K^+)}{\text{Br}(B^- \rightarrow D^0 K^-) + \text{Br}(B^+ \rightarrow \bar{D}^0 K^+)}$$

$$A_{CP\pm} \equiv \frac{\text{Br}(B^- \rightarrow D_{CP\pm} K^-) - \text{Br}(B^+ \rightarrow D_{CP\pm} K^+)}{\text{Br}(B^- \rightarrow D_{CP\pm} K^-) + \text{Br}(B^+ \rightarrow D_{CP\pm} K^+)}$$

Relation between  $(R_{CP+}, R_{CP-}, A_{CP+}, A_{CP-})$  and  $(\gamma, r_B, \delta_B)$

$$R_{CP+} = 1 + r_B^2 + 2r_B \cos \delta_B \cos \gamma$$

$$R_{CP-} = 1 + r_B^2 - 2r_B \cos \delta_B \cos \gamma$$

$$A_{CP+} = \frac{+2r_B \sin \delta_B \sin \gamma}{1 + r_B^2 + 2r_B \cos \delta_B \cos \gamma}$$

$$A_{CP-} = \frac{-2r_B \sin \delta_B \sin \gamma}{1 + r_B^2 - 2r_B \cos \delta_B \cos \gamma}$$

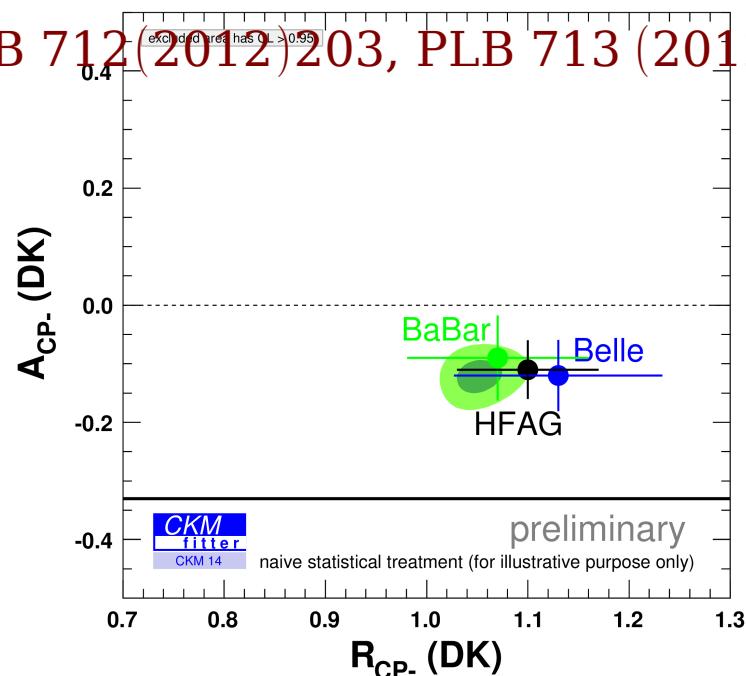
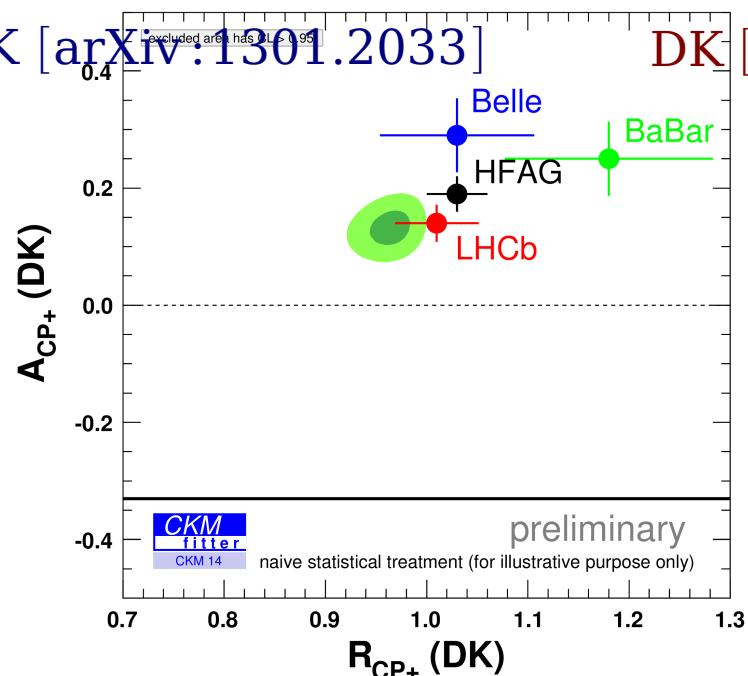
- ⇒ look for  $R_{CP\pm} \neq 1$  and  $A_{CP\pm} \neq 0$
- ⇒  $\neq CP$ ,  $\neq$  sign of asymmetry

Magnitude of one side  
is ~0.1 of the others  
while relative magnitude of  
the others help  $\phi_3$  constraint.

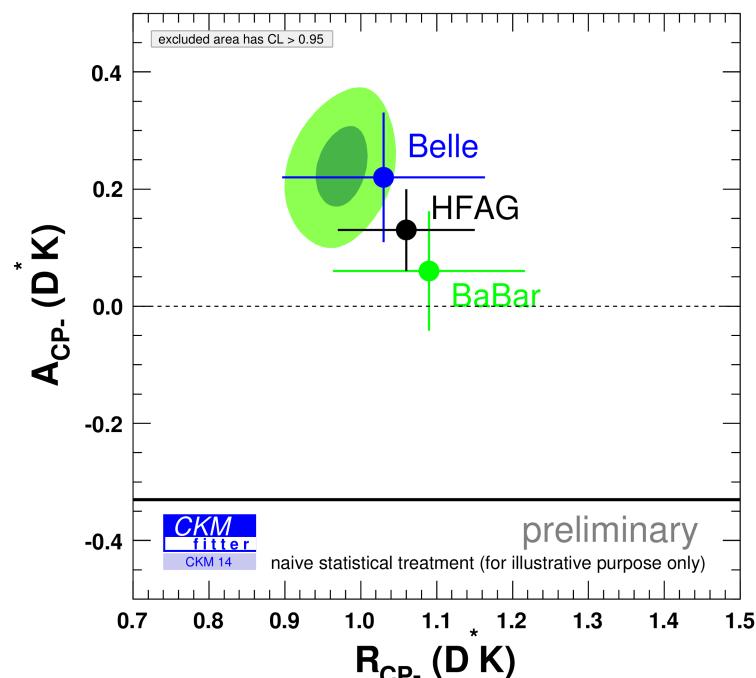
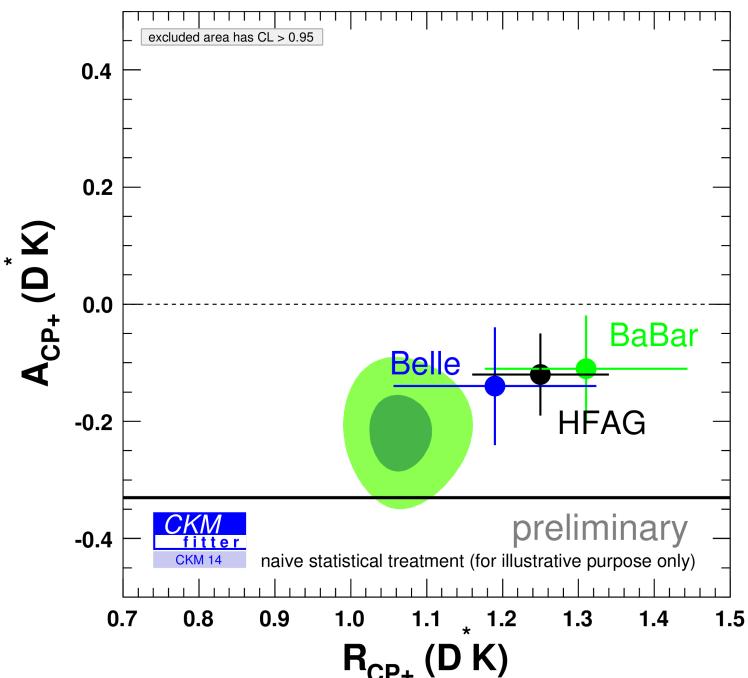
# GLW observables (predictions vs measurements)

DK [PRD82(2010)072004], D<sup>\*</sup>K[PRD78(2008)092002], DK<sup>\*</sup> [PRDD80(2009)092001]

DK and D<sup>\*</sup>K [arXiv:1301.2033] DK [PLB 712(2012)203, PLB 713 (2012) 351]

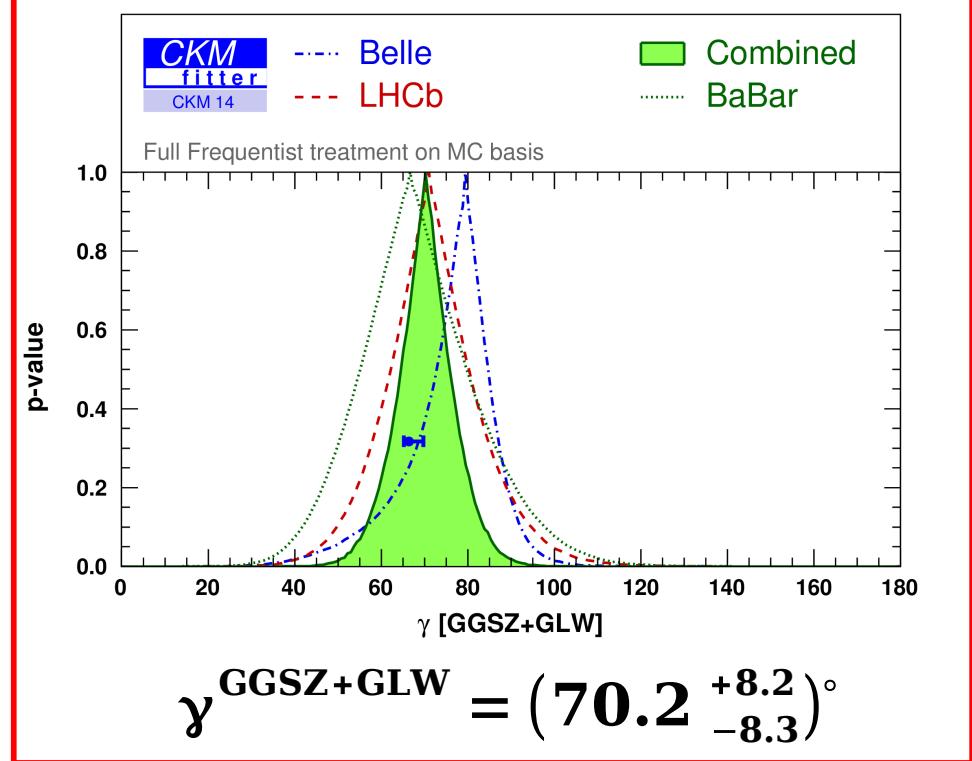


+ 22 obs.

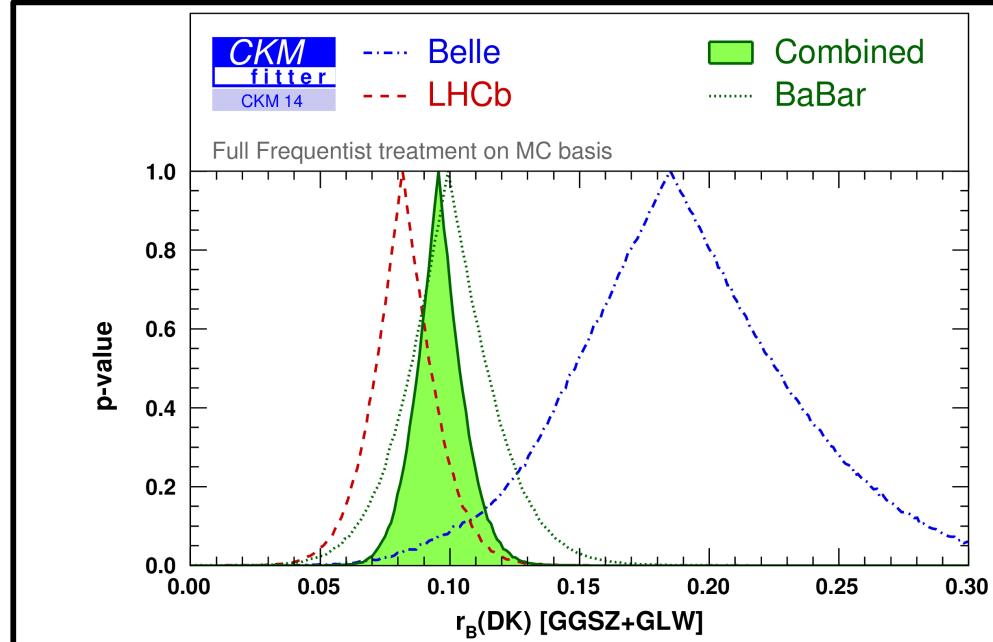


# GGSZ+GLW

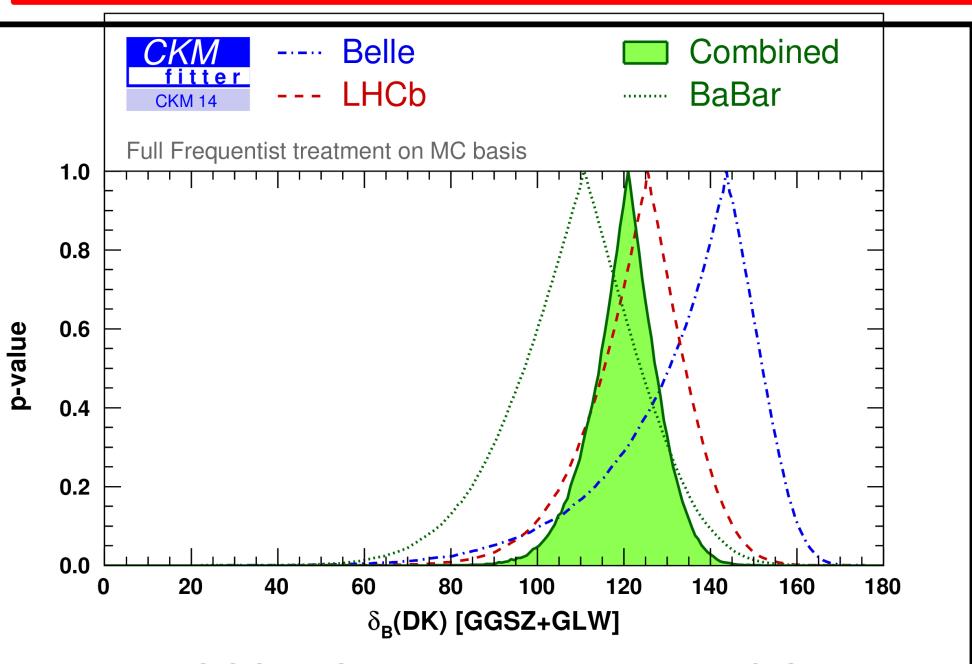
(assuming no DVCP in  $D \rightarrow \pi\pi, KK$ )



(results for DK)

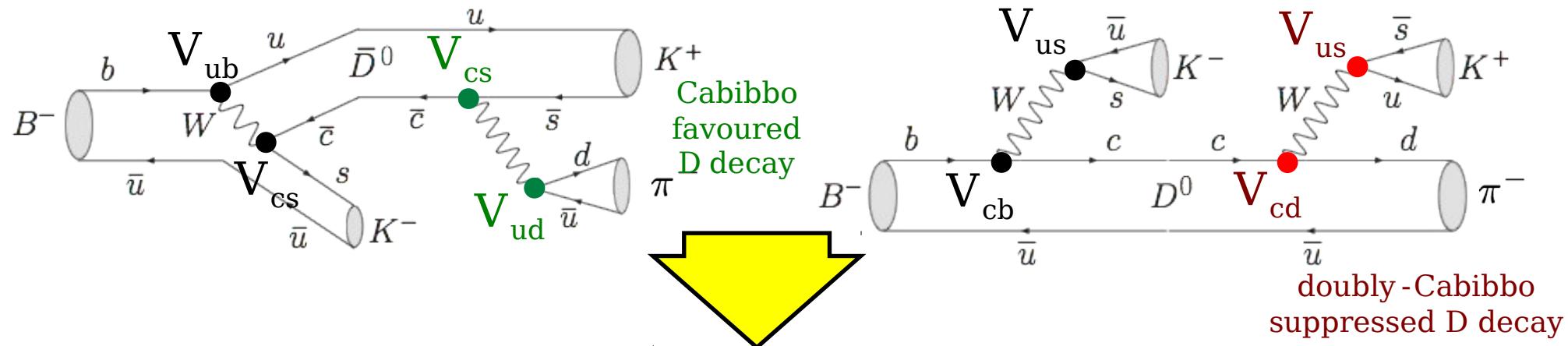


$$r_B^{\text{GGSZ+GLW}} = (0.096 \pm 0.012)$$



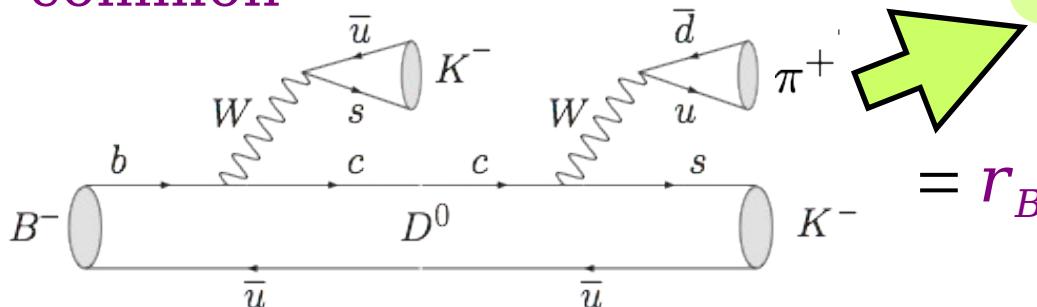
$$\delta_B^{\text{GGSZ+GLW}} = (121.0^{+9.2}_{-10.3})^\circ$$

ADS method:  $\gamma$  via the interference in rare  $B^- \rightarrow [K^+ \pi^-]_D K^-$  decays rate and asymmetry (relative to the common decay):



$$\mathcal{R}_{DK} = \frac{\Gamma([K^+ \pi^-] K^-) + \Gamma([K^- \pi^+] K^+)}{\Gamma([K^- \pi^+] K^-) + \Gamma([K^+ \pi^-] K^+)}$$

common



$$= r_B^2 + r_D^2 + 2 r_B r_D R \cos(\delta_B + \delta_D) \cos \gamma$$

$$\mathcal{A}_{DK} = \frac{\Gamma([K^+ \pi^-] K^-) - \Gamma([K^- \pi^+] K^+)}{\Gamma([K^- \pi^+] K^-) + \Gamma([K^+ \pi^-] K^+)}$$

$$= 2 r_B r_D R \sin(\delta_B + \delta_D) \sin \gamma / R_{DK}$$

where  $r_D = \left| \frac{\mathcal{A}(D^0 \rightarrow K^+ \pi^-)}{\mathcal{A}(\bar{D}^0 \rightarrow K^+ \pi^-)} \right| = 0.0613 \pm 0.0010$

# How to get $\delta_D$ and related (charm) hadronic parameters ?

- dedicated experiments (CLEO-c, BES III) using quantum correlations , running at  $\psi(3770)$ 
  - CLEO-c:  $R_D$ ,  $\cos\delta$ ,  $\sin\delta$  (but also BES III result...)
  - CLEO-c:  $R_{K\pi\pi^0}$ ,  $\delta_{K\pi\pi^0}$ ,  $R_{K3\pi}$ ,  $\delta_{K3\pi}$

$R_f$ : coherence factor , can take any value from 0 to 1

indicates lack coherence between the intermediate states involved in the decay

- mixing/CPV results from BaBar , Belle , CDF , LHCb...

- $D \rightarrow K\bar{K}$  ,  $\pi\pi$ :  $y_{CP}$ ,  $A_\Gamma$  (BaBar , Belle , LHCb)
- $D \rightarrow K_S^0 \pi\pi$ :  $x$  ,  $y$  ,  $|q/p|$  ,  $\varphi$  (BaBar , Belle)
- $D \rightarrow K l \nu$ :  $R_M$  (BaBar , Belle...)
- $D \rightarrow K \pi\pi^0$ :  $x''$  ,  $y''$  (BaBar)
- $D \rightarrow K \pi$ :  $x'$  ,  $y'$  (BaBar , Belle, CDF , LHCb)
- ...

- **CLEO-c/BES III, use external inputs to access the relevant physics parameters**
- **strong phases information in B-factories/LHCb**
- **x , y are also needed for D-mixing corrections in ADS observables**

$$R^\mp = r_B^2 + r_D^2 + 2r_B r_D \cos(\delta_B \mp \gamma + \delta_D)$$

$$\Rightarrow R^\mp = r_B^2 + r_D^2 + 2r_B r_D \cos(\delta_B \mp \gamma + \delta_D) - y r_D \cos \delta_D - y r_B \cos(\delta_B \mp \gamma) + x r_D \sin \delta_D - x r_B \sin(\delta_B \mp \gamma)$$

⇒ **combine charm observables to obtain  $\gamma$  and mixing/CPV charm parameters**

# Charm mixing in $D^0 \rightarrow K^+ \pi^-$ [see M.Charles's talk]

The ratio  $R(t)$  of WS  $D^{*+} \rightarrow D^0 \pi^+_s \rightarrow K^+ \pi^- \pi^+_s$  to RS  $D^{*+} \rightarrow D^0 \pi^+_s \rightarrow K^- \pi^+ \pi^+_s$  decay rates can be approximated (assuming  $|x|, |y| \ll 1$  and no CPV) by:

$$R(t) = R_D + \sqrt{R_D} y t + \frac{x^2 + y^2}{4} t^2$$

DCS to CF ratio    mixing rate

$$\begin{aligned} x &= x \cos \delta_{K\pi} + y \sin \delta_{K\pi} & \delta_{K\pi}: \text{strong phase difference} \\ y &= y \cos \delta_{K\pi} - x \sin \delta_{K\pi} & \text{btw DCS and CF amplitudes} \end{aligned}$$

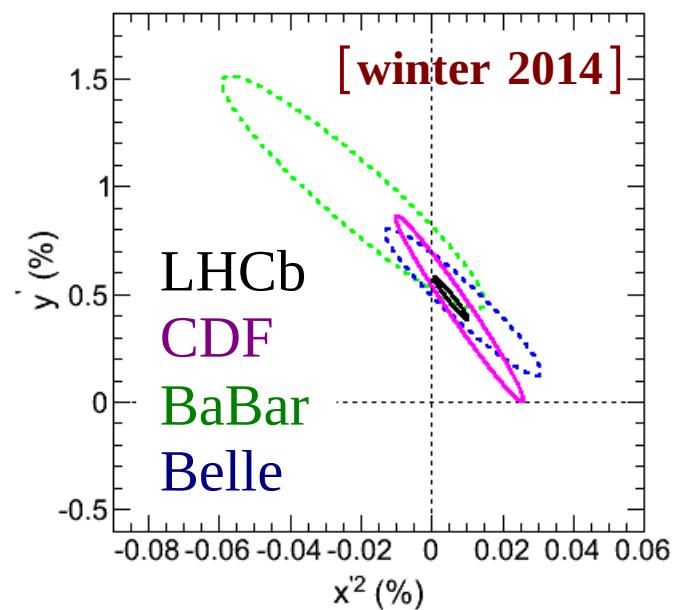
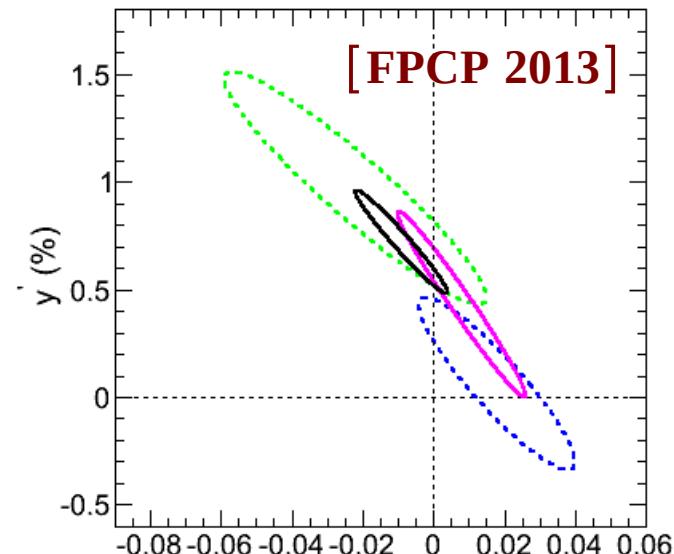
Exp	$R_D$ ( $10^{-3}$ )	$y$ ( $10^{-3}$ )	$x^2$ ( $10^{-3}$ )	$\Sigma$
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<b>Belle</b> PRL112 (2014) 111801	$3.53 \pm 0.13$	$4.6 \pm 3.4$	$+0.09 \pm 0.22$	5.1
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<b>BaBar</b> PRL98 (2007) 211802	$3.03 \pm 0.19$	$9.7 \pm 5.4$	$-0.22 \pm 0.37$	3.9
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<b>LHCb</b> PRL111 (2013) 251801	$3.57 \pm 0.07$	$4.8 \pm 1.0$	$+0.055 \pm 0.049$	?
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<b>CDF</b> preliminary (2013)	$3.51 \pm 0.35$	$4.3 \pm 4.3$	$+0.08 \pm 0.18$	6.1
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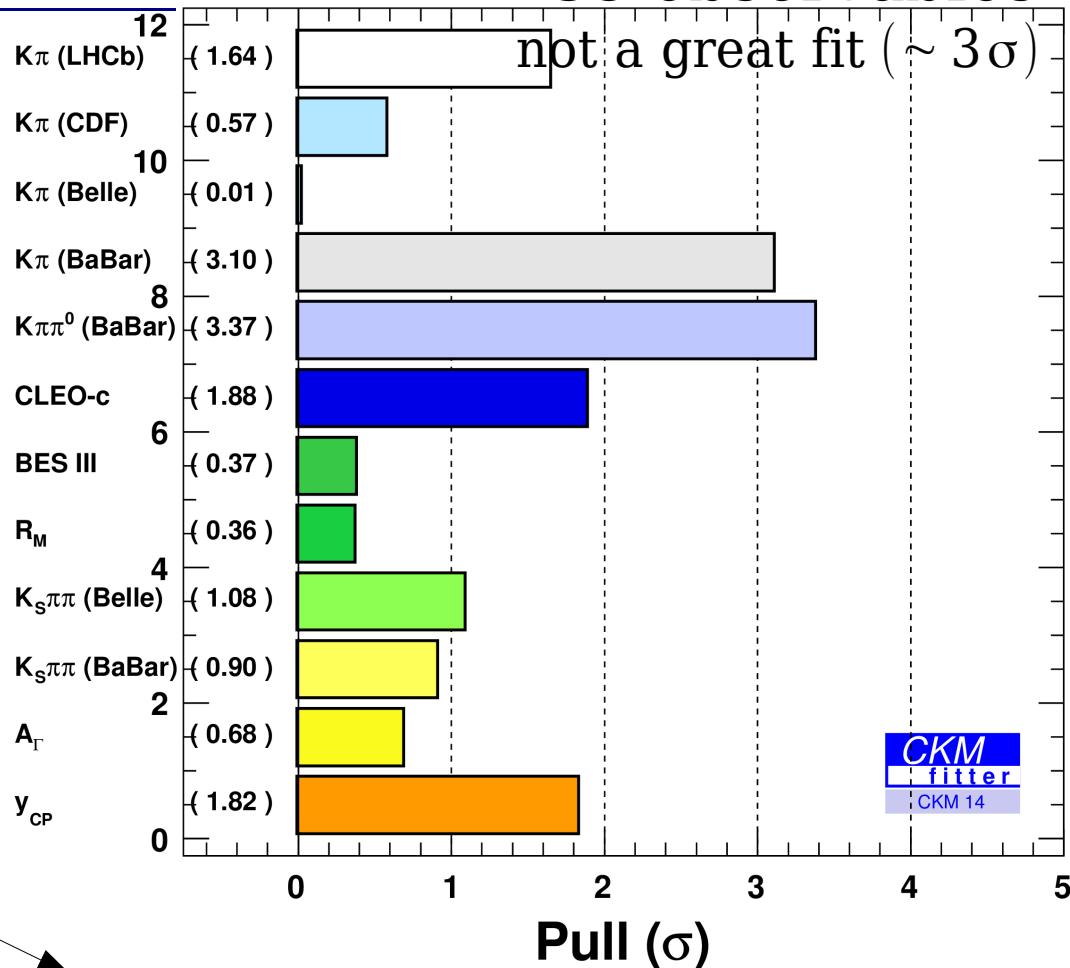
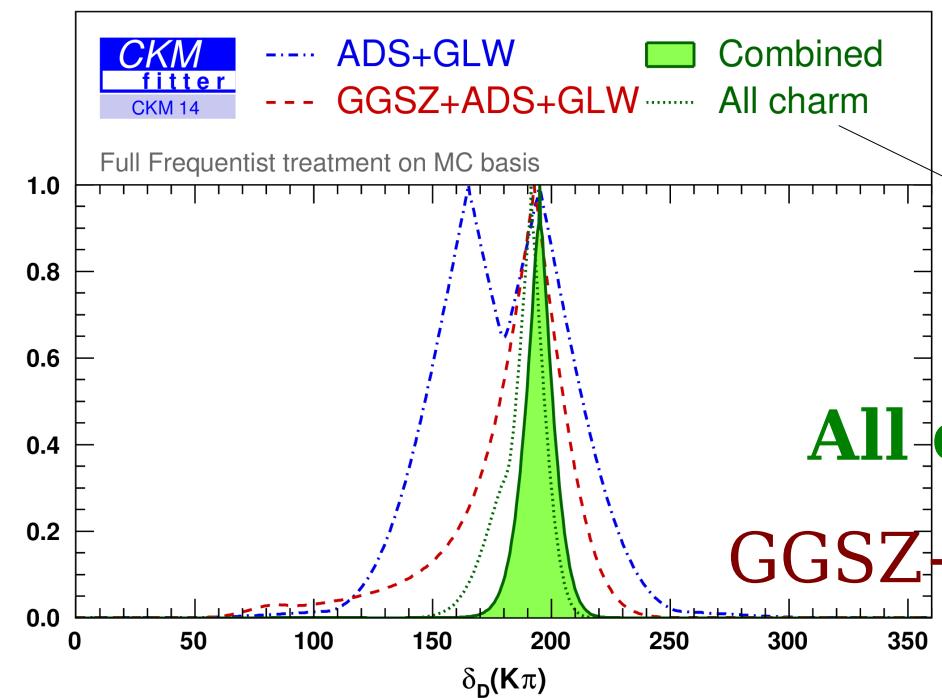


# $\delta_D$ grand combination à la HFAG

~ 35 observables

8 parameters:

$x, y, \delta_D^{K\pi}, r_D, A_D, |q|/|p|, \varphi, \delta_D(K\rho)$



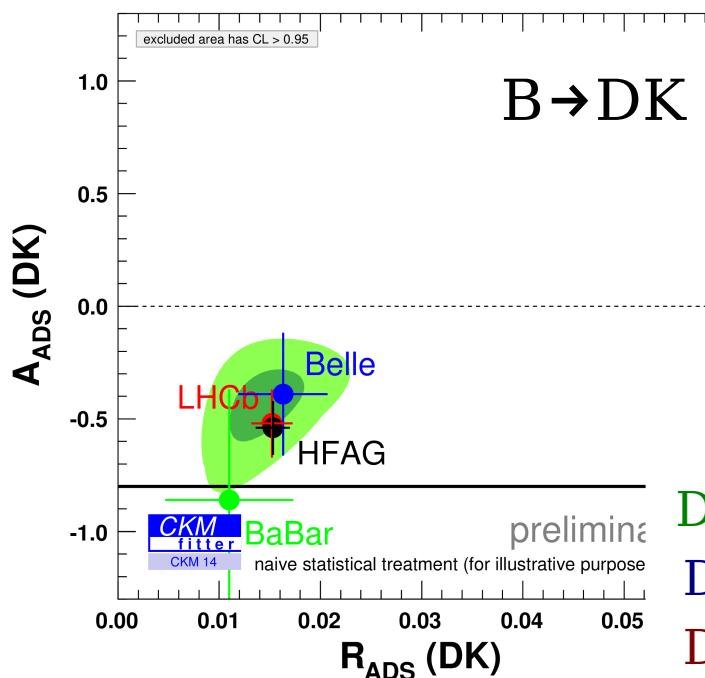
(include  $K3\pi, K\pi\pi^0$  info, see next slides)

All charm:  $\delta_D^{K\pi} = (191.4^{+8.2}_{-11.4})^\circ \quad (^{+16}_{-30})$

GGSZ+GLW+ADS:  $\delta_D^{K\pi} = (193^{+18}_{-23})^\circ \quad (^{+34}_{-77})$

# Comparison of the results obtained for $D^{(*)}K$ with expectations

where "expectations" are derived from the GGSZ observables,  $\delta_D$  and  $r_D$



$$R_{ADS}(DK) = r_B^2 + r_D^2 + 2r_B r_D \cos(\delta_B + \delta_D) \cos \gamma$$

$$A_{ADS}(DK) = 2r_B r_D \sin(\delta_B + \delta_D) \sin \gamma / R_{ADS}(DK)$$

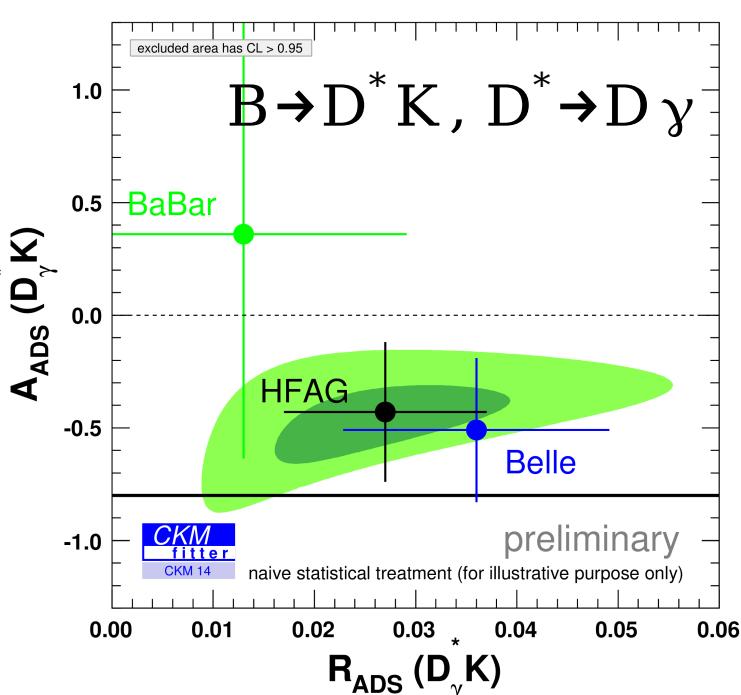
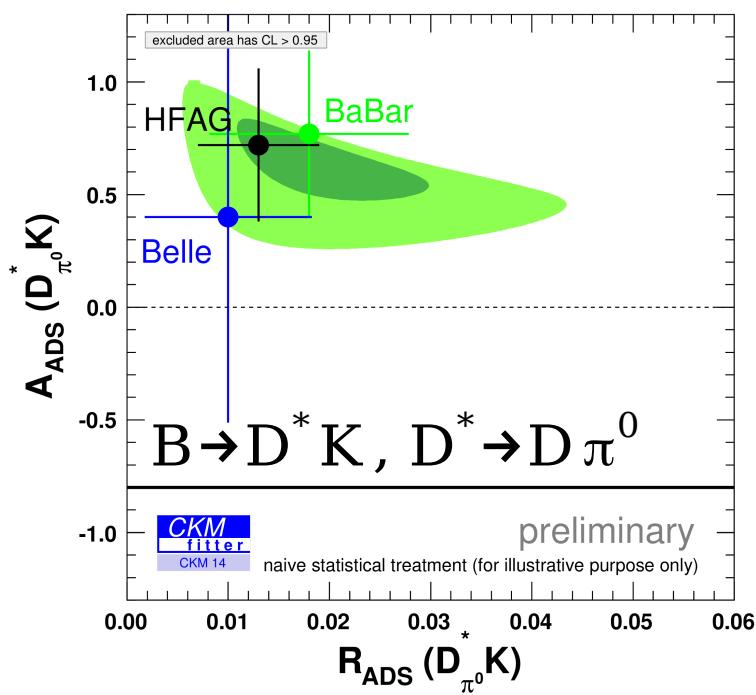
$$R_{ADS}(D_{\pi^0}^* K) = r_B^{*2} + r_D^2 + 2r_B^* r_D \cos(\delta_B^* + \delta_D) \cos \gamma$$

$$A_{ADS}(D_{\pi^0}^* K) = 2r_B^* r_D \sin(\delta_B^* + \delta_D) \sin \gamma / R_{ADS}(D_{\pi^0}^* K)$$

$$R_{ADS}(D_{\gamma}^* K) = r_B^{*2} + r_D^2 - 2r_B^* r_D \cos(\delta_B^* + \delta_D) \cos \gamma$$

$$A_{ADS}(D_{\gamma}^* K) = -2r_B^* r_D \sin(\delta_B^* + \delta_D) \sin \gamma / R_{ADS}(D_{\gamma}^* K)$$

DK,  $D^* K$  [PRD82 (2010) 072006],  $DK^*$  [PRD80 (2009) 092001]  
 DK [PRL 106 (2011) 231803],  $D^* K$  [arXiv:1301.2033]  
 DK [PLB 712(2012)203, PLB 713 (2012) 351]



# ADS observables

- **( $R_+$ ,  $R_-$ ) instead of ( $R_{ADS}$ ,  $A_{ADS}$ ) whenever available**

## Effect of D- $\bar{D}$ mixing on $\gamma$

- M.Rama , arXiv:1307.4384
- $R^\mp = r_B^2 + r_D^2 + 2r_B r_D \cos(\delta_B \mp \gamma + \delta_D)$   
→  $R^\mp = r_B^2 + r_D^2 + 2r_B r_D \cos(\delta_B \mp \gamma + \delta_D) - y r_D \cos \delta_D - y r_B \cos(\delta_B \mp \gamma) +$   
 $x r_D \sin \delta_D - x r_B \sin(\delta_B \mp \gamma)$
- tried on the current LHCb average (DK):  $\sim 1$  degree difference

# $K\pi\pi^0$ , $K3\pi$ from CLEO-c

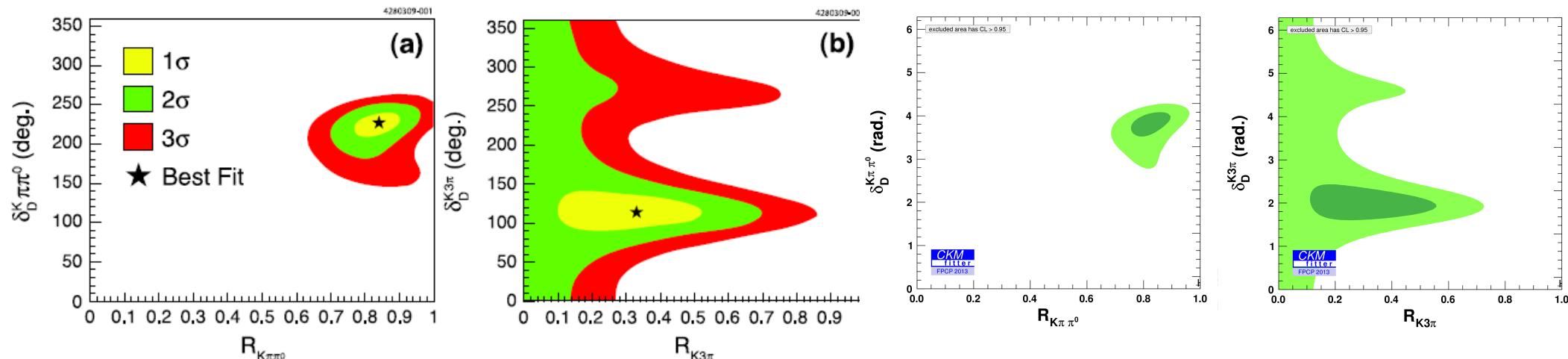
[see R.Briere's talk]

yields of double-tagged events where one meson decays into  $K^- \pi^+ \pi^0$  (or  $K3\pi$ ), and the other meson decays into CP-odd, CP-even and  $K\pi$

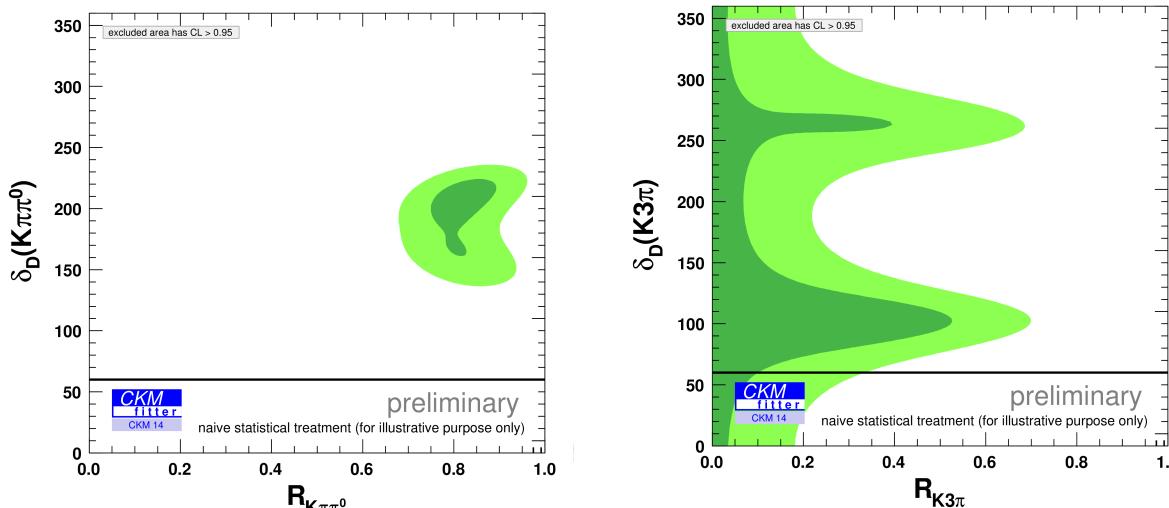
[arXiv:0903.4853, N.Lowrey et al]

(combined with external inputs:  $x$ ,  $y$ ,  $\delta_D$ ...)

that we could reproduce earlier  
extending the charm fitter (+ Br's)



2014 version (currently used in our  $\gamma$  combination):



# $K\pi\pi^0$ , $K3\pi$ from CLEO-c

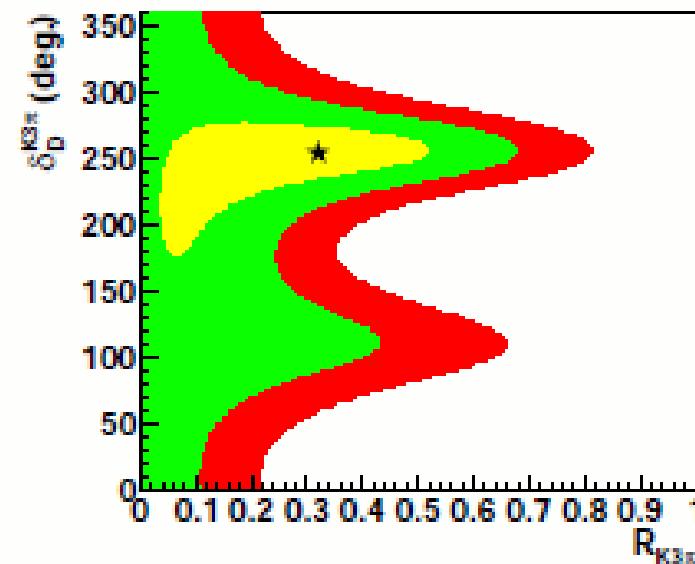
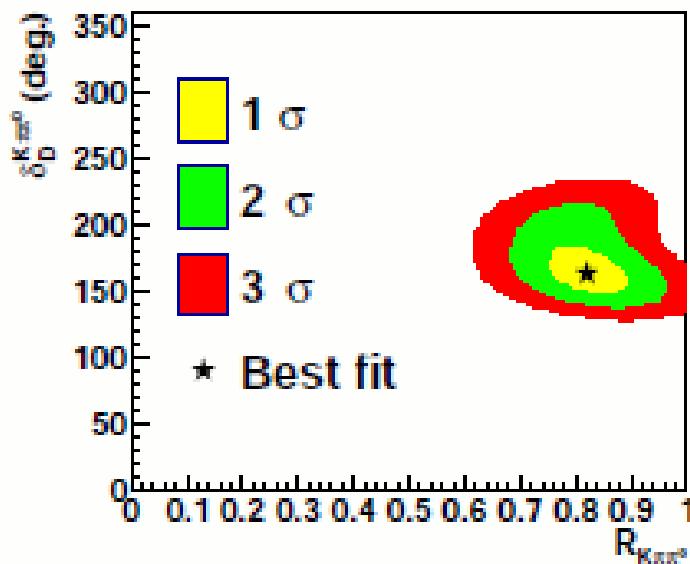
[J.Libby et al ,arXiv:1401.1904]

yields of double-tagged events where one meson decays into  $K^-\pi^+\pi^0$  (or  $K3\pi$ ), and the other meson decays into  $K_S^0\pi^+\pi^-$

$$Y_i = H_{K\pi\pi^0} \left( K_i + (r_D^{K\pi\pi^0})^2 K_{-i} - 2r_D^{K\pi\pi^0} \sqrt{K_i K_{-i}} R_{K\pi\pi^0} [c_i \cos \delta_D^{K\pi\pi^0} + s_i \sin \delta_D^{K\pi\pi^0}] \right),$$

measure by CLEO-c

$K_i$  : fractional yield of  $D^0$  decays that fall into bin i



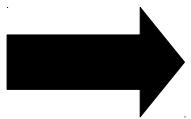
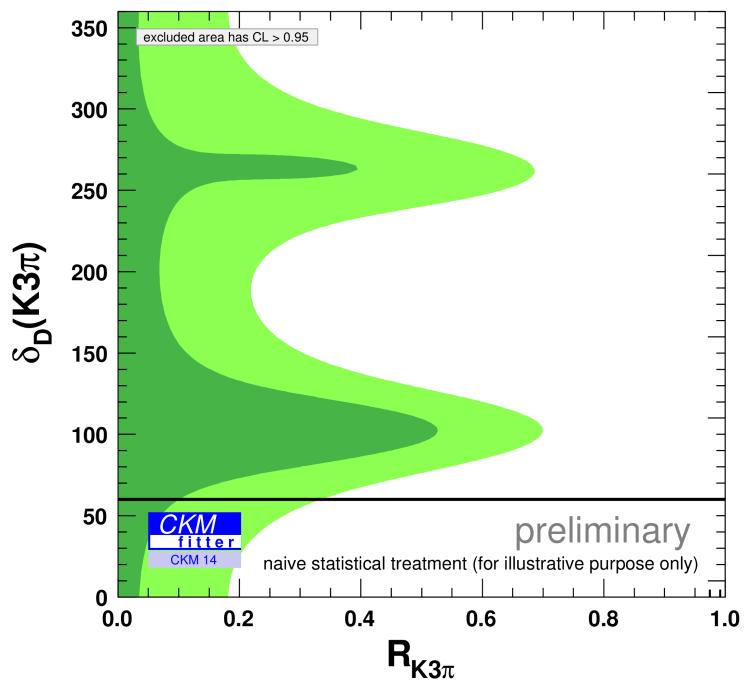
⇒ will soon include this information

$K3\pi$  charm information is limited:

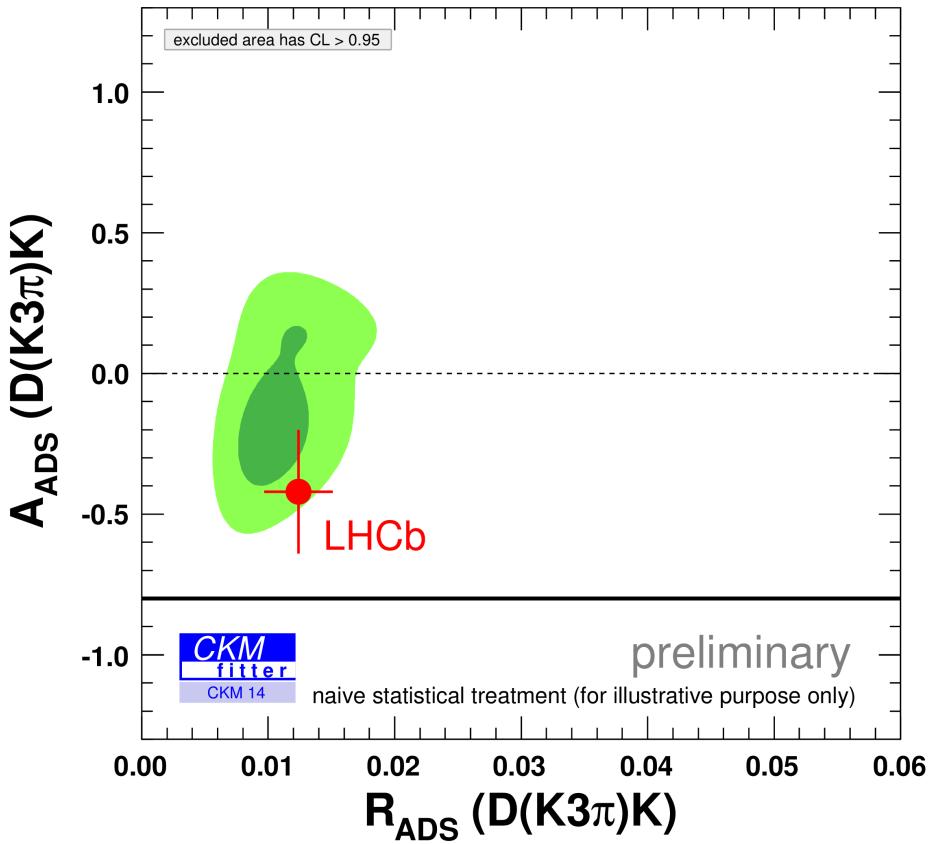
- possible additional inputs from BES III
- B factories/LHCb [S.Harnew and J.Rademacker , arXiv:1309.0134]

# ADS $B \rightarrow D(K3\pi)K$

where "expectations" derived from the GGSZ observables,  $\delta_D$ ,  $r_D$  and  $R$  (for  $K3\pi$ )



D( $K3\pi$ )K [PLB 723 (2013) 44]

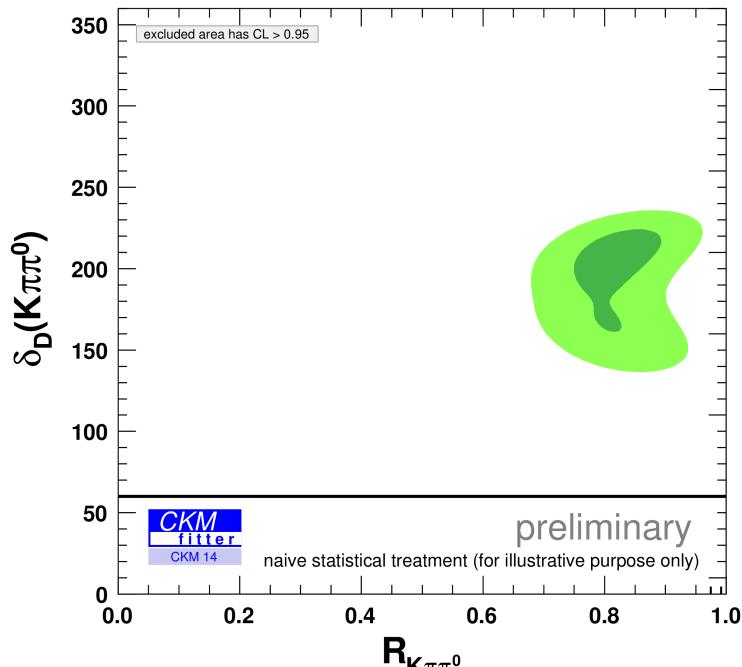


⇒  $D(K3\pi)K$  LHCb result included in the  $\gamma$  combination

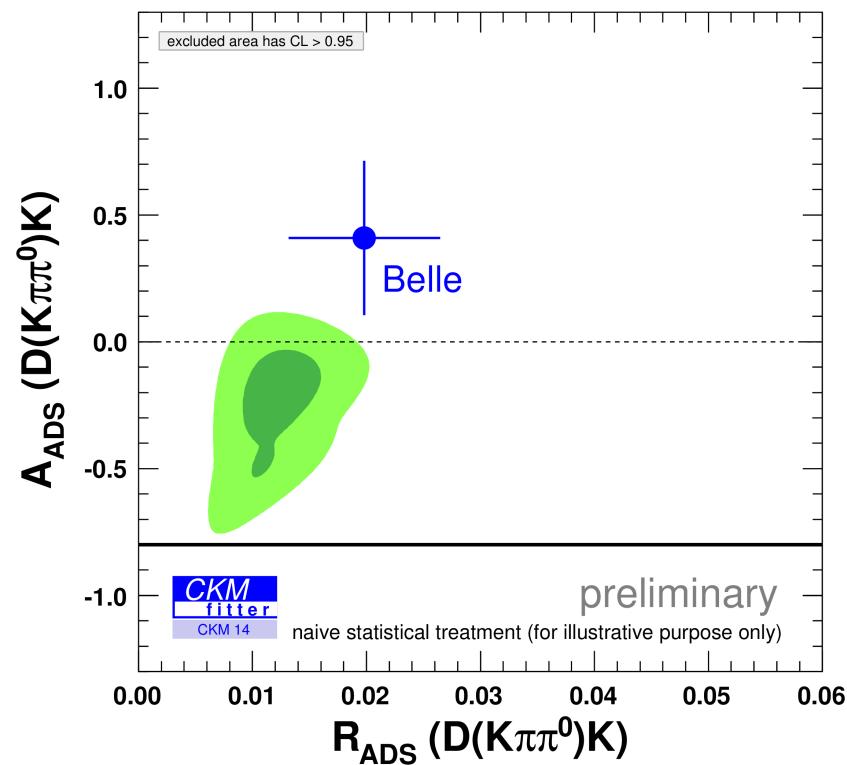
# ADS $B \rightarrow D(K\pi\pi^0)K$

[see J.Libby's talk]

where "expectations" derived from the GGSZ observables,  $\delta_D$ ,  $r_D$  and  $R$  (for  $K\pi\pi^0$ )



Evidence for the suppressed decay  $B^- \rightarrow DK^-, D \rightarrow K^+\pi^-\pi^0$



DK [PRD 88, 091104(R) (2013)]

M. Nayak,<sup>16</sup> J. Libby,<sup>16</sup> K. Trabelsi,<sup>12</sup> I. Adachi,<sup>12</sup> H. Aihara,<sup>55</sup> D. M. Asner,<sup>42</sup> T. Aushev,<sup>20</sup> A. M. Bakich,<sup>49</sup> A. Bala,<sup>43</sup> P. Behera,<sup>16</sup> K. Belous,<sup>18</sup> V. Bhardwaj,<sup>34</sup> G. Bonvicini,<sup>60</sup> A. Bozek,<sup>38</sup> M. Bracko,<sup>27,21</sup> T. E. Browder,<sup>11</sup> D. Červenkov,<sup>5</sup> M.-C. Chang,<sup>8</sup> P. Chang,<sup>37</sup> V. Chekelian,<sup>28</sup> A. Chen,<sup>35</sup> B. G. Cheon,<sup>10</sup> R. Chistov,<sup>20</sup> I.-S. Cho,<sup>62</sup> K. Cho,<sup>24</sup> V. Chobanova,<sup>28</sup> Y. Choi,<sup>48</sup> D. Cinabro,<sup>60</sup> J. Dalseno,<sup>28,51</sup> M. Danilov,<sup>20,30</sup> Z. Doležal,<sup>5</sup> Z. Drásal,<sup>5</sup> D. Duggan,<sup>15</sup> S. Eidelman,<sup>4</sup> A. Esipenko,<sup>60</sup> J. Gao,<sup>42</sup> M. Gershon,<sup>7</sup> Y. Goh,<sup>30</sup> N. Golutvin,<sup>14</sup> G. Gritsan,<sup>15</sup> G.

We report a study of the suppressed decay  $B^- \rightarrow DK^-, D \rightarrow K^+\pi^-\pi^0$ , where  $D$  denotes either a  $D^0$  or a  $D^0$  meson. The decay is sensitive to the  $CP$ -violating parameter  $\phi_3$ . Using a data sample of  $772 \times 10^6$   $B\bar{B}$  pairs collected at the  $\Upsilon(4S)$  resonance with the Belle detector, we measure the ratio of branching fractions of the above suppressed decay to the favored decay  $B^- \rightarrow DK^-, D \rightarrow K^-\pi^+\pi^0$ . Our result is  $R_{DK} = [1.98 \pm 0.62(\text{stat.}) \pm 0.24(\text{syst.})] \times 10^{-2}$ , which indicates the first evidence of the signal for this suppressed decay with a significance of 3.2 standard deviations. We measure the direct  $CP$  asymmetry between the suppressed  $B^-$  and  $B^+$  decays to be  $A_{DK} = 0.41 \pm 0.30(\text{stat.}) \pm 0.05(\text{syst.})$ . We also report measurements for the analogous quantities  $R_{D\pi}$  and  $A_{D\pi}$  for the decay  $B^- \rightarrow D\pi^-, D \rightarrow K^+\pi^-\pi^0$ .

G. Varner,<sup>11</sup> K. E. Varvell,<sup>49</sup> M. N. Wagner,<sup>9</sup> C. H. Wang,<sup>36</sup> M.-Z. Wang,<sup>37</sup> Y. Watanabe,<sup>22</sup> K. M. Williams,<sup>59</sup>

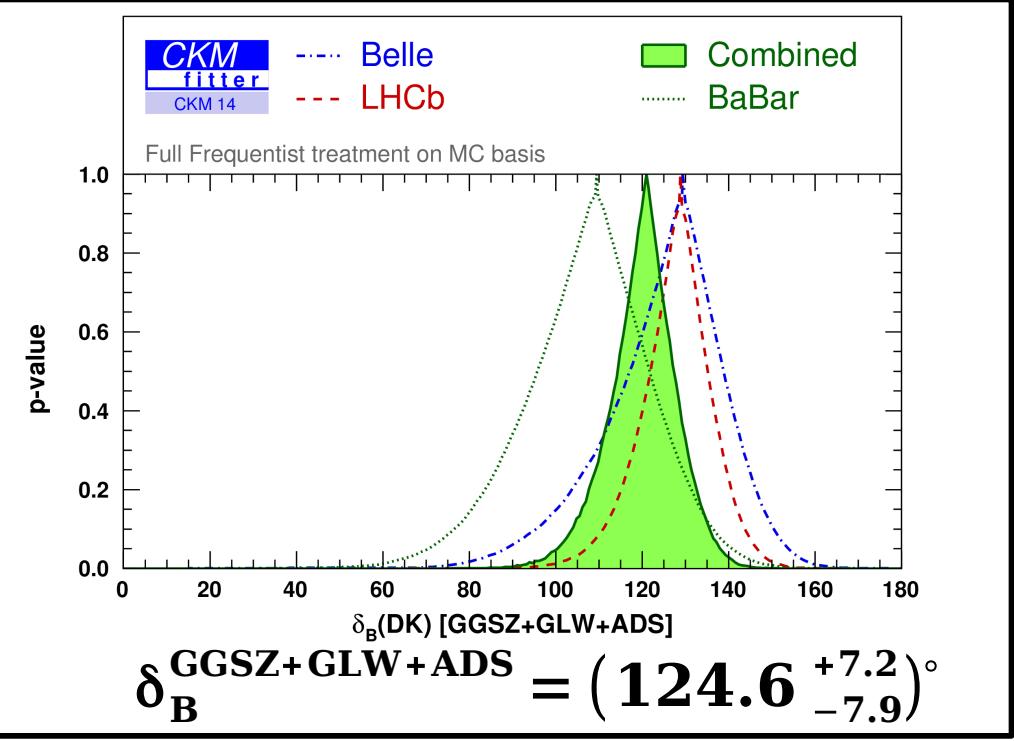
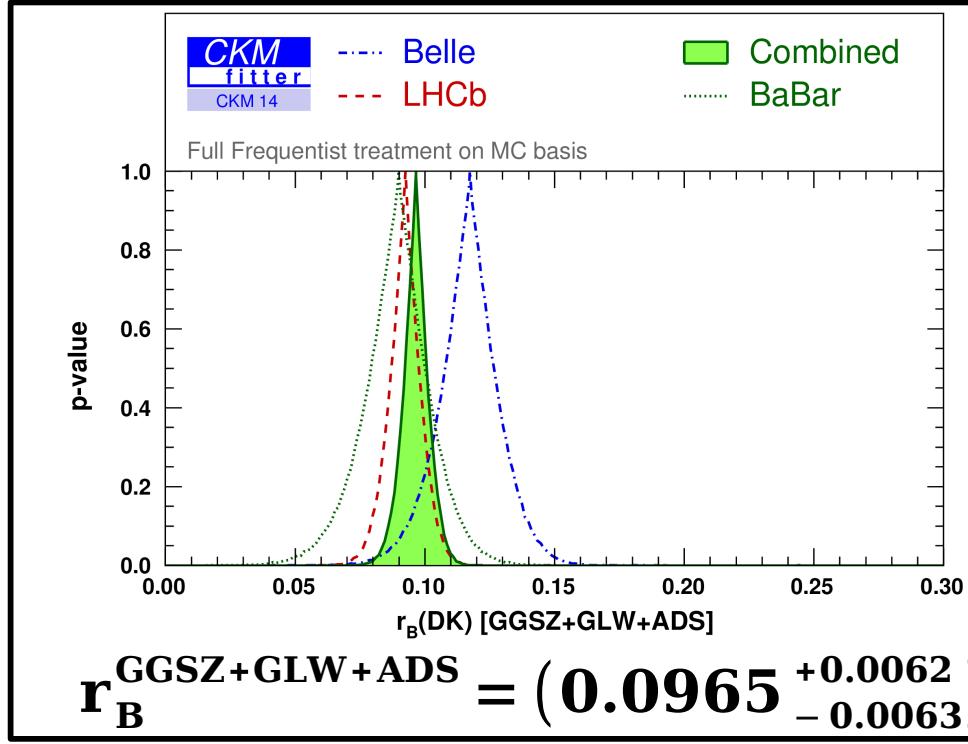
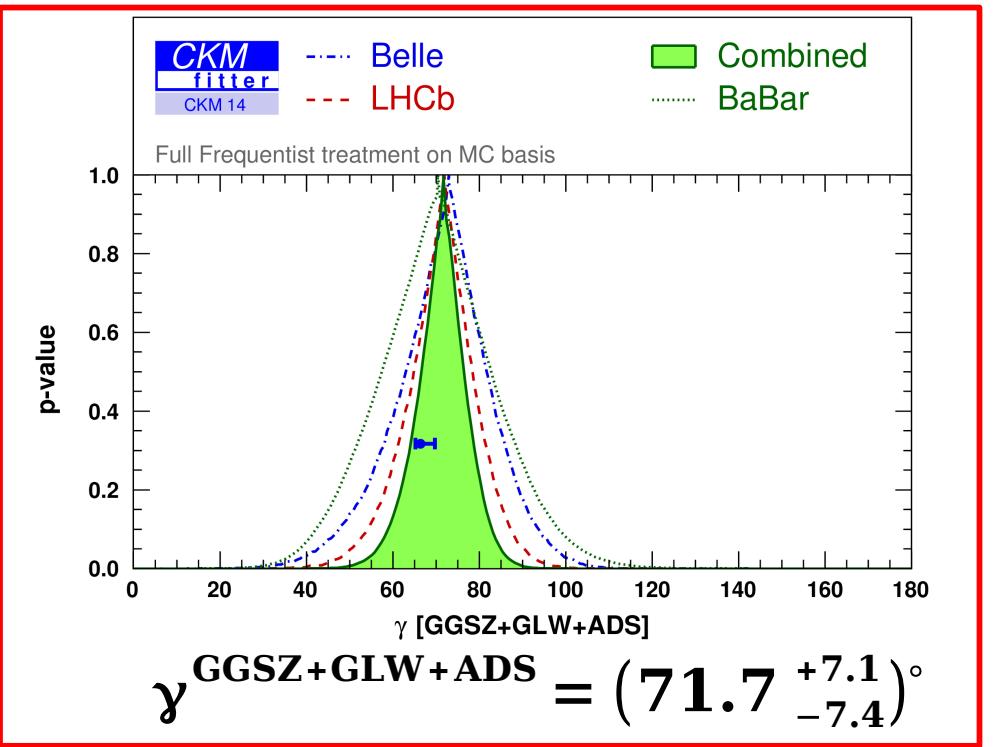
E. Won,<sup>25</sup> Y. Yamashita,<sup>39</sup> S. Yashchenko,<sup>7</sup> Y. Yusa,<sup>40</sup> V. Zhilich,<sup>4</sup> V. Zhulanov,<sup>4</sup> and A. Zupanc<sup>23</sup>

$\Rightarrow$  Belle (and BaBar)  $D(K\pi\pi^0)K$  results included in the  $\gamma$  combination

# GGSZ+GLW+ADS

+20 obs.

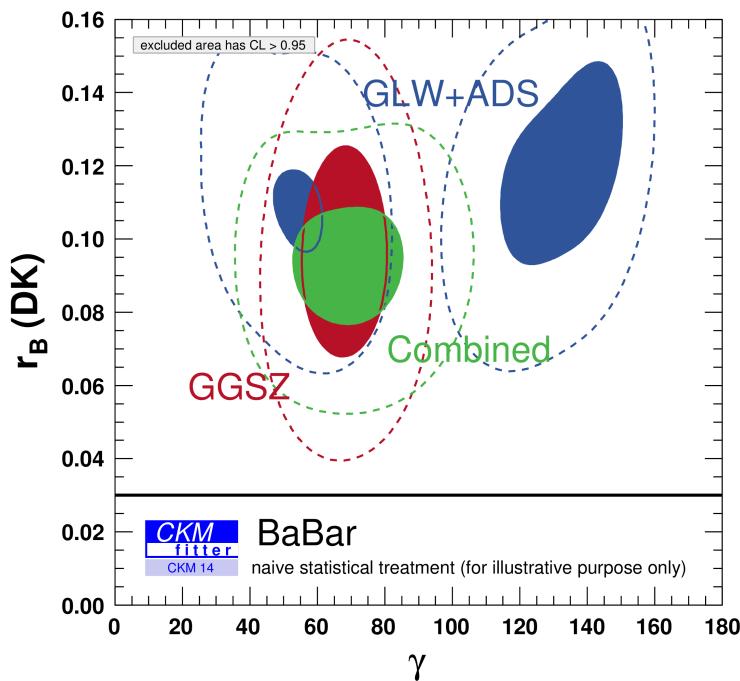
(results for DK)



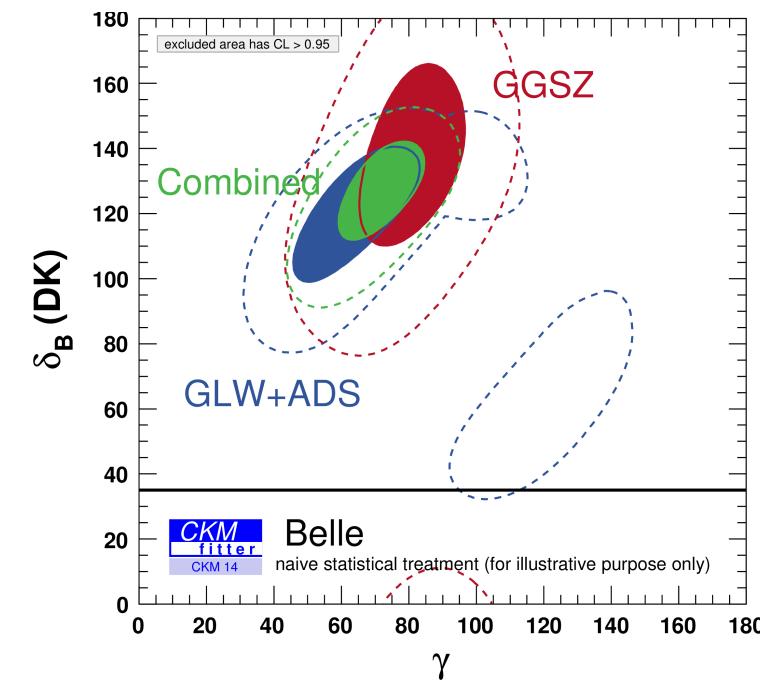
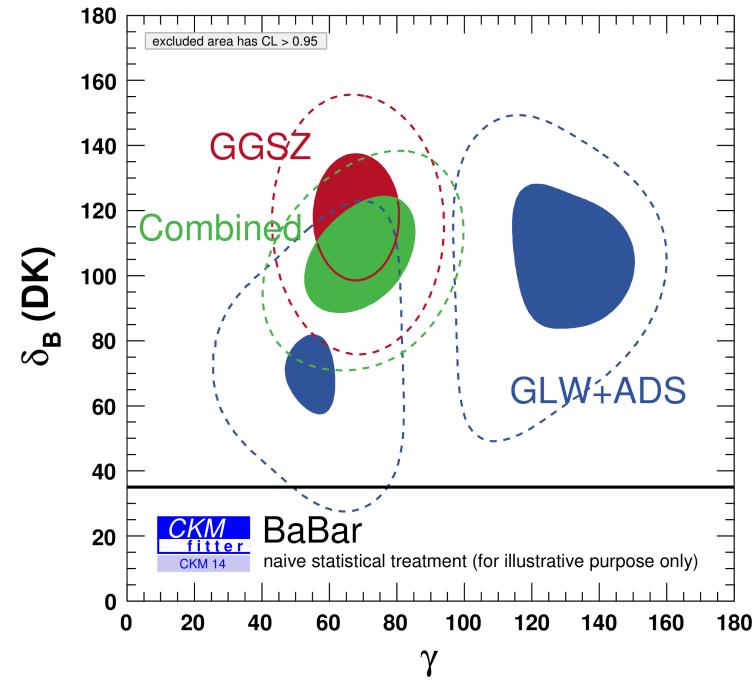
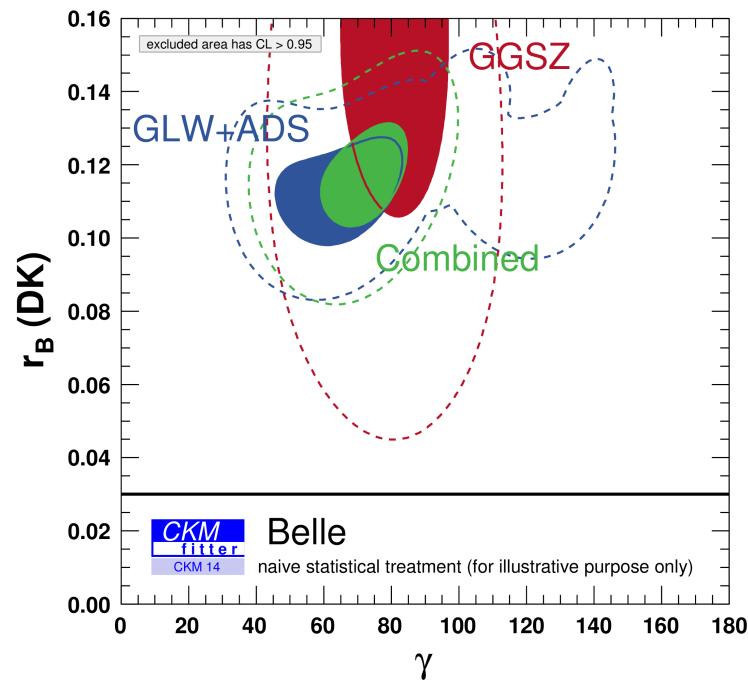
# GGSZ versus GLW+ADS

$(r_B(\text{DK}) \text{ vs } \gamma, \delta_B(\text{DK}) \text{ vs } \gamma)$

BaBar



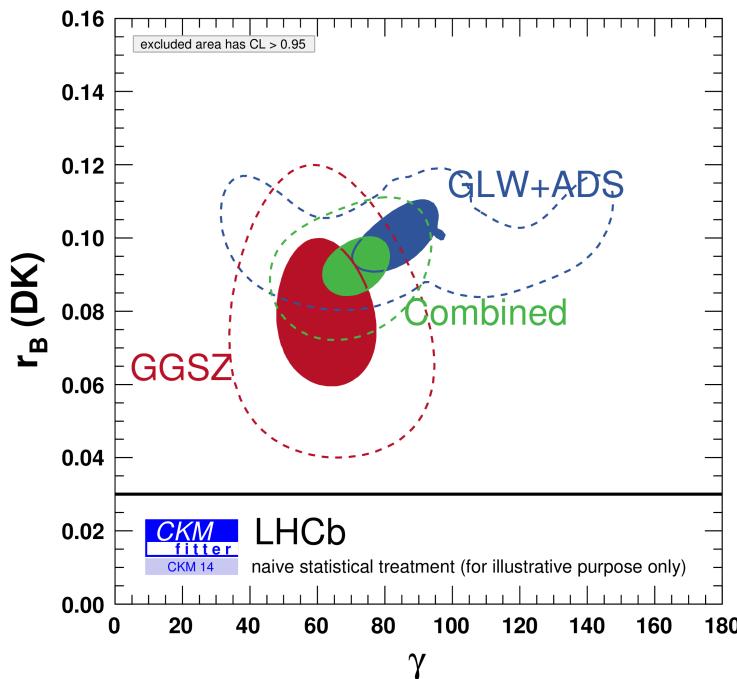
Belle



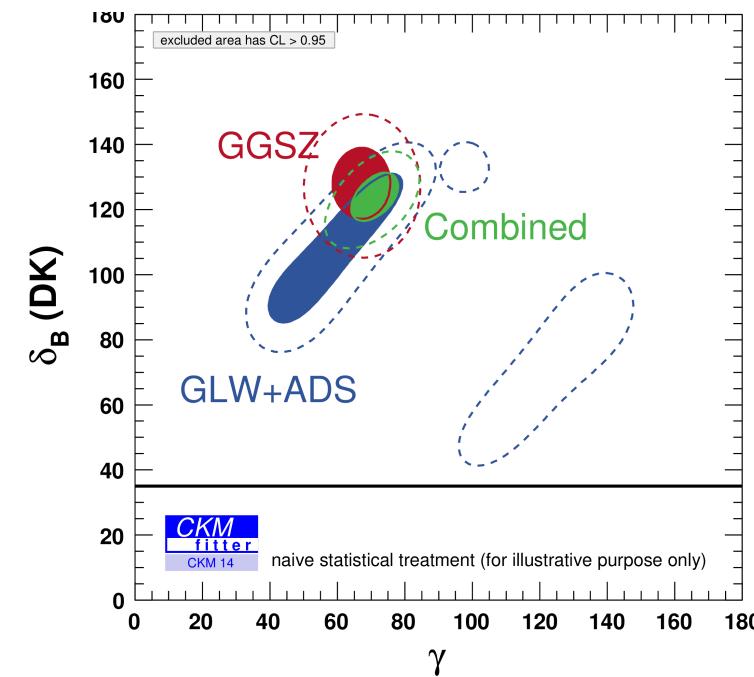
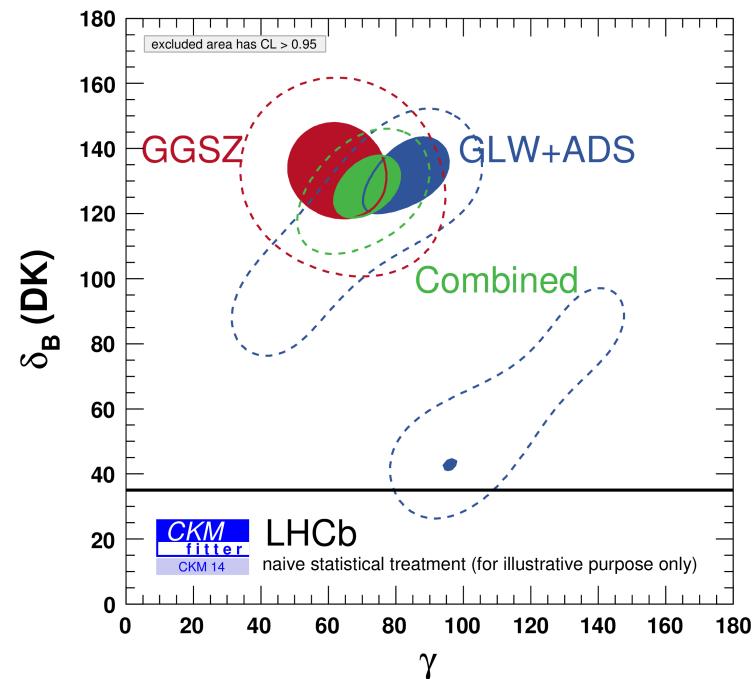
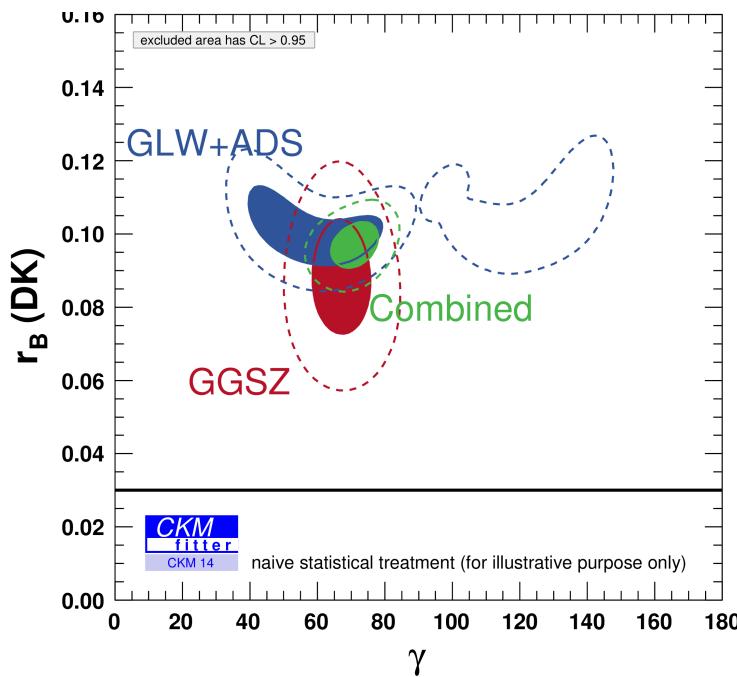
# GGSZ versus GLW+ADS

$(r_B(\text{DK}) \text{ vs } \gamma, \delta_B(\text{DK}) \text{ vs } \gamma)$

LHCb



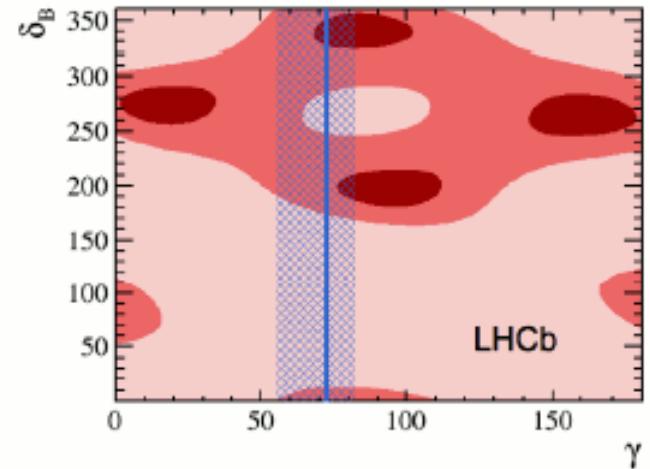
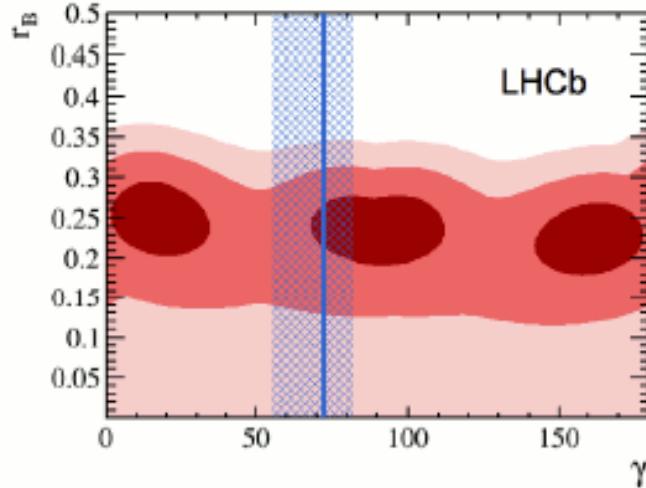
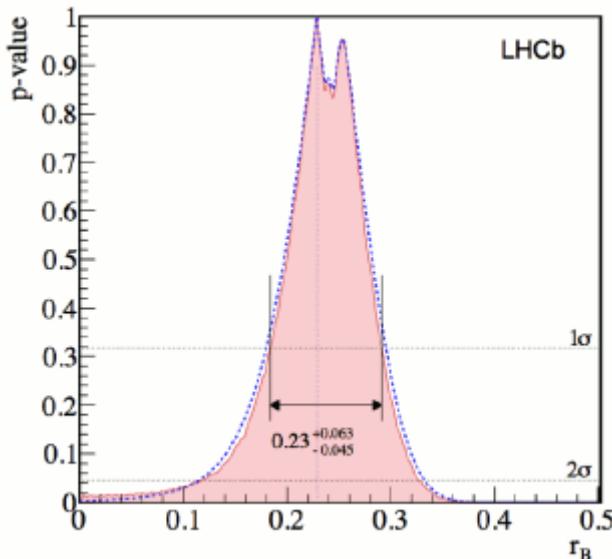
All



# New LHCb results, $\text{DK}^{*0}$ , $D \rightarrow hh$ (ADS, GLW)

$$r_B(\text{DK}^{*0}) = (0.230 \begin{array}{l} +0.063 \\ -0.045 \end{array})$$

[arXiv:1407.8136]  
 [see M.Karbach's talk]

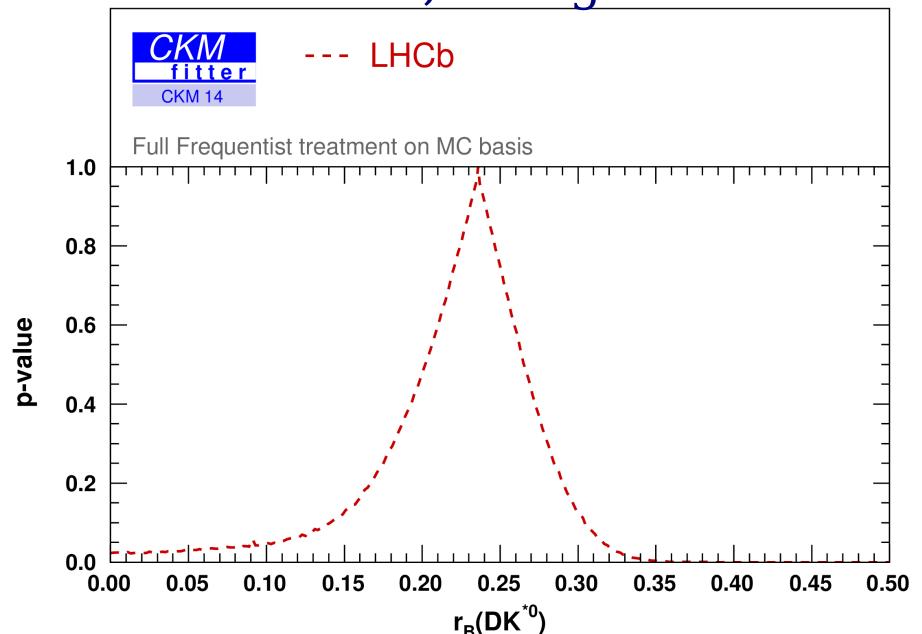


Combining all LHCb  $\gamma$  measurements (GGSZ+GLW+ADS, charged and neutral)

$$r_B(\text{DK}^{*0}) = 0.236 \begin{array}{l} +0.043 \\ -0.052 \end{array}$$

$$\begin{array}{c} +0.081 \\ -0.137 \end{array} \quad (2\sigma)$$

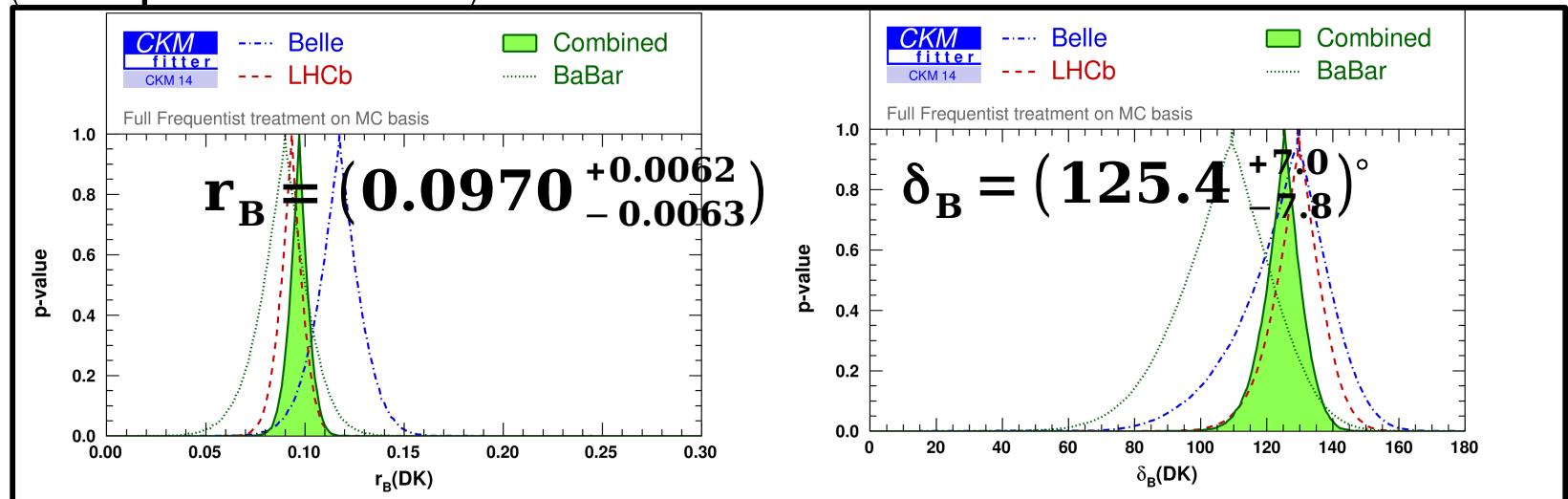
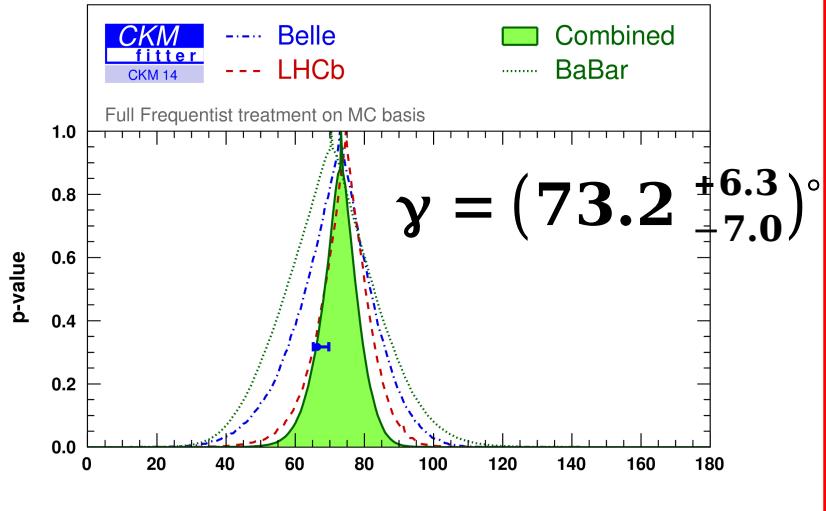
$$\begin{array}{c} +0.12 \\ -\infty \end{array} \quad (3\sigma)$$



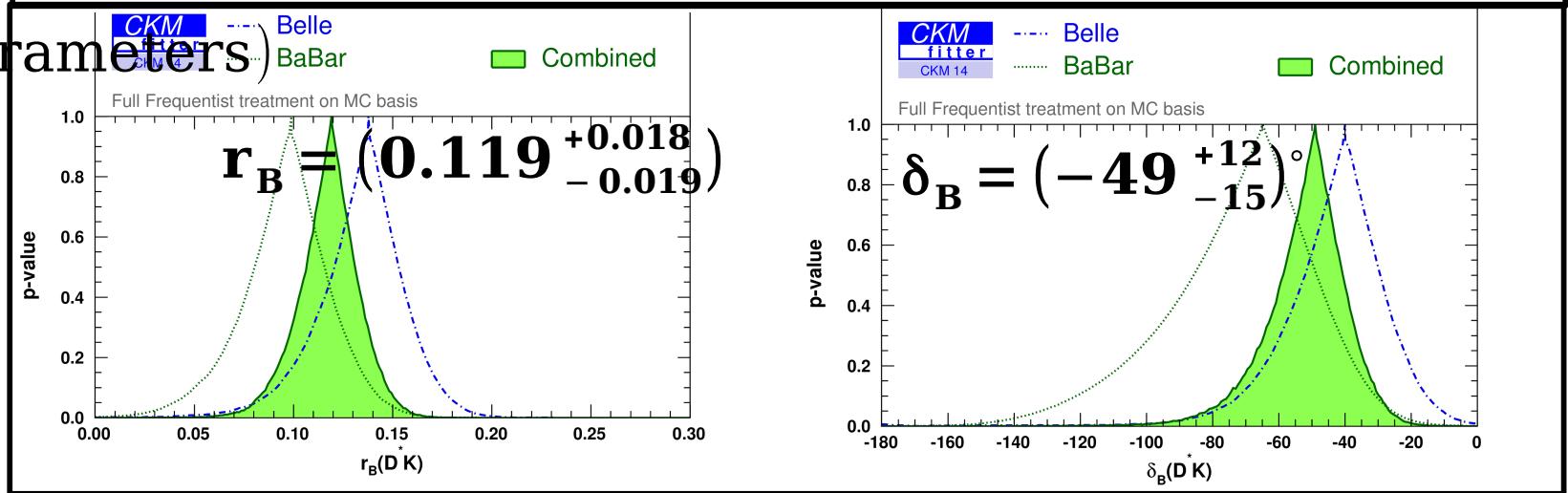
# GGSZ+GLW+ADS

(charged and neutral B decays)

(DK parameters)



(D<sup>\*</sup> K parameters)



# Conclusion

- Combining GGSZ, GLW and ADS observables from charged and neutral B modes obtained by BaBar, Belle and LHCb experiments

$$\gamma = (73.2 \begin{array}{l} +6.3 \\ -7.0 \end{array})^\circ \quad [\text{CKM2014, preliminary}]$$

$$\gamma \text{ [no direct } \gamma \text{ measurements]} = (66.4 \begin{array}{l} +1.3 \\ -2.5 \end{array})^\circ \quad [\text{CKMfitter, Moriond 2014}]$$

- Next: include  $D(K_S K \pi)K$  info (GLS method) from LHCb ?  
⇒ include CLEO-c information on charm observables [arXiv:1203.3804]
- looking forward for additional/updated  $\gamma$  observables:

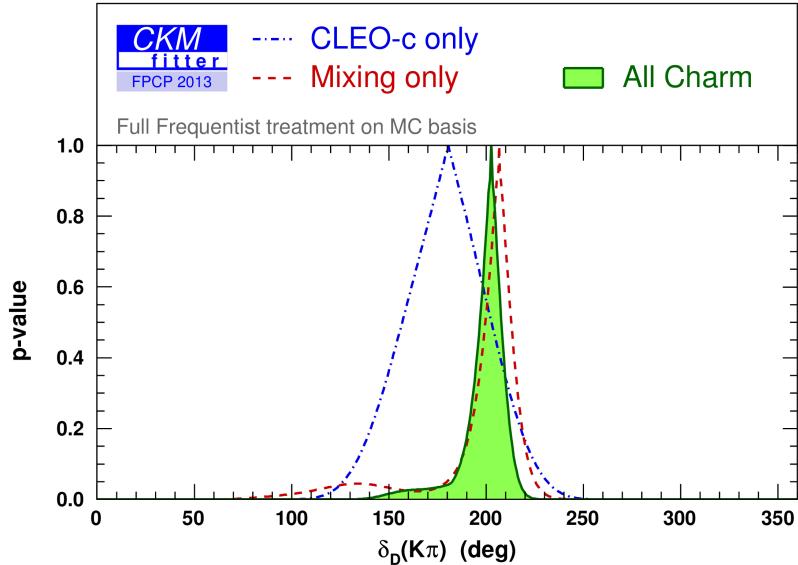
- ADS/GLW observables with  $3 \text{ fb}^{-1}$  from LHCb
- GGSZ observables from  $D\bar{K}^{*0}$
- $(x^\pm, y^\pm)$  from  $D(K_S \pi\pi)\pi$
- ...

and for charm observables:

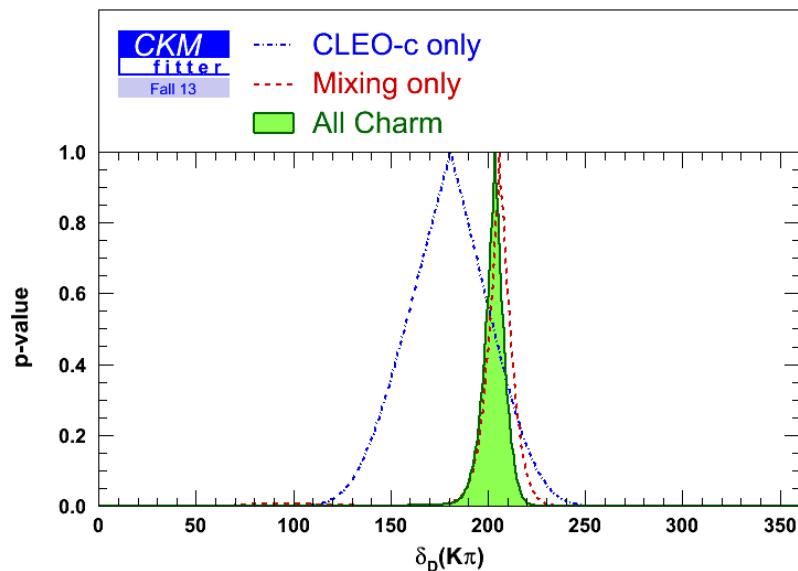
- from BES III
- additional information on  $K3\pi$  from LHCb and B-factories
- ...



# $\delta_D$ grand combination

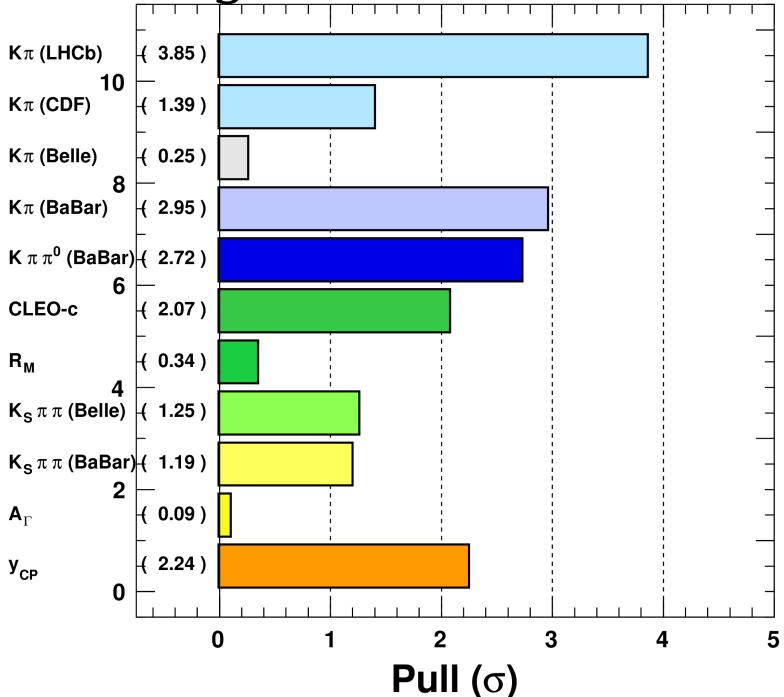


$$\delta_D = (202.6^{+7.6}_{-8.9})^\circ \quad (+15) \quad (+22) \quad (-64)$$

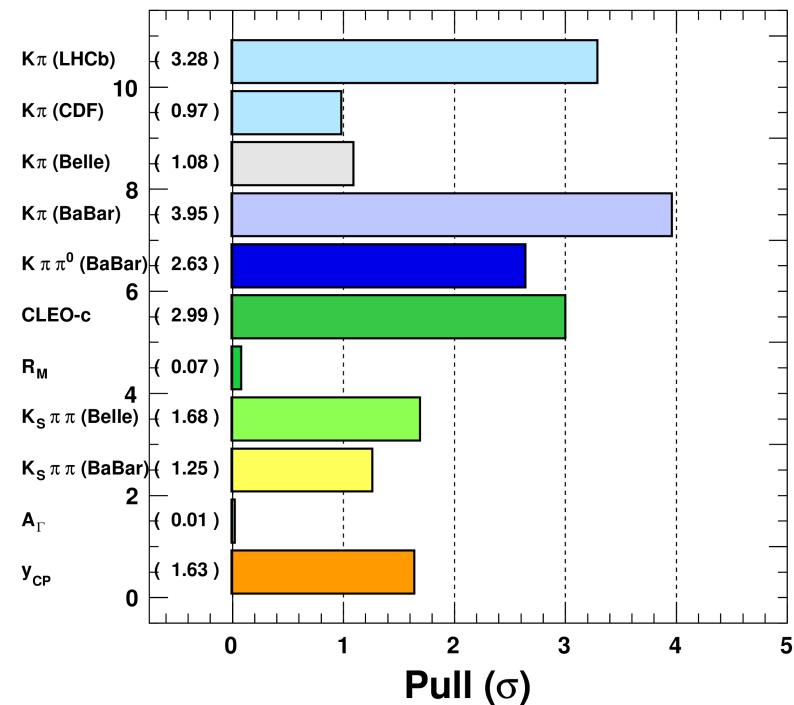


$$\delta_D = (203.4^{+6.4}_{-6.9})^\circ \quad (+13) \quad (+20) \quad (-35)$$

Not a great fit...



overall:  $3.6\sigma \rightarrow 3.1\sigma$



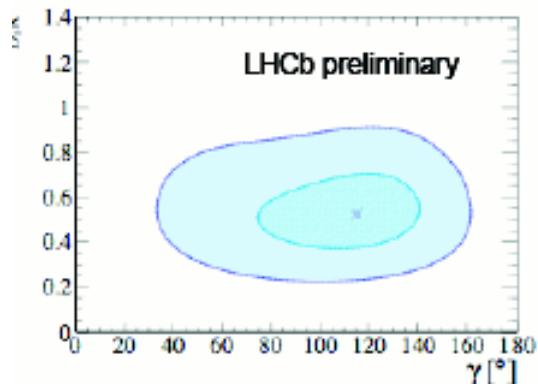
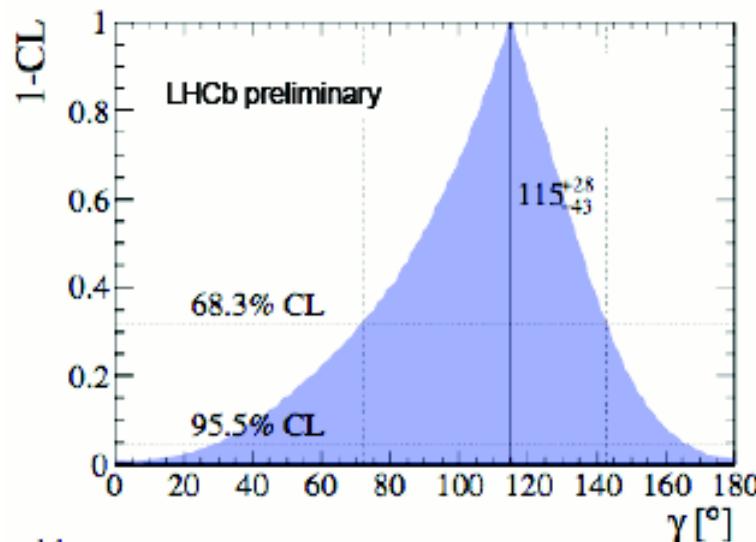
# New LHCb results, $D_s K$ (only 1 $\text{fb}^{-1}$ )

note: would expect a  $r_B \approx 0.3$

$$\gamma = (115^{+28}_{-43})^\circ$$

$$r_{D_s K} = (0.53^{+0.17}_{-0.16})$$

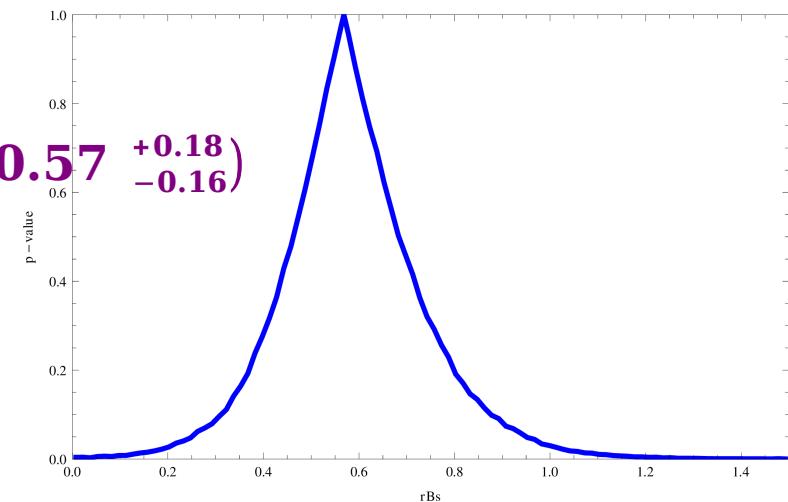
$$\delta_{D_s K} = (3^{+19}_{-20})^\circ$$



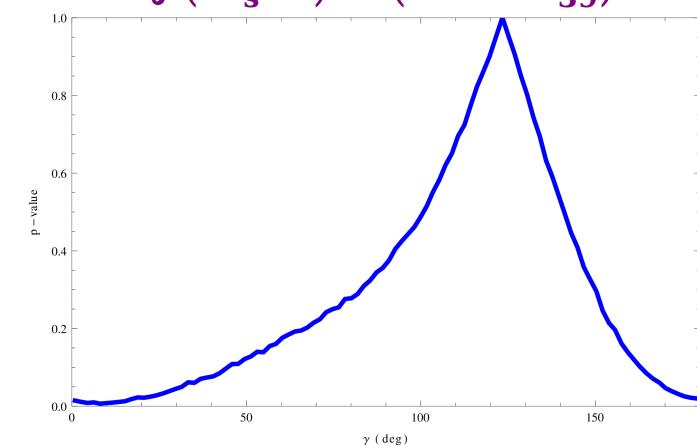
Projection  
of contours  
onto 1-D  
gives

- 68.3%
- 95.5%

$$r_B(D_s K) = (0.57^{+0.18}_{-0.16})$$



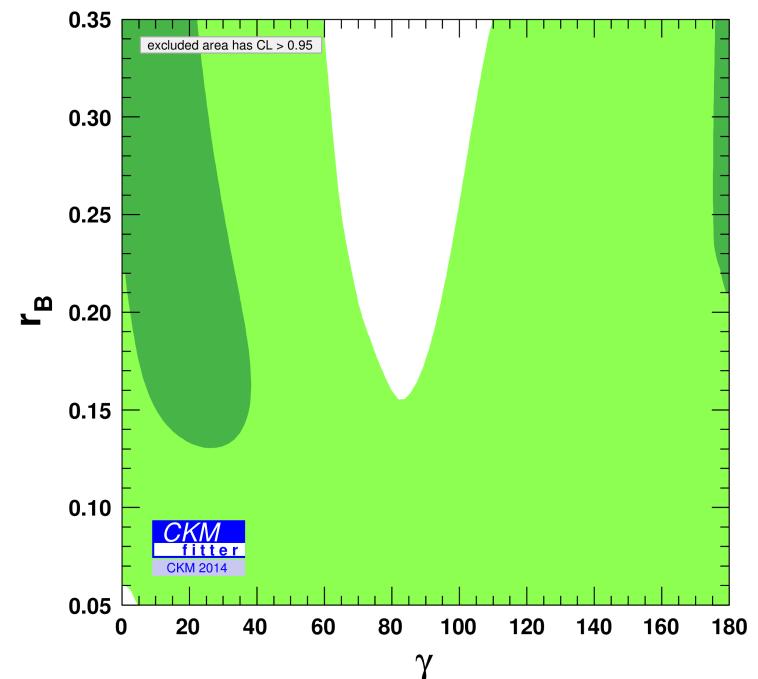
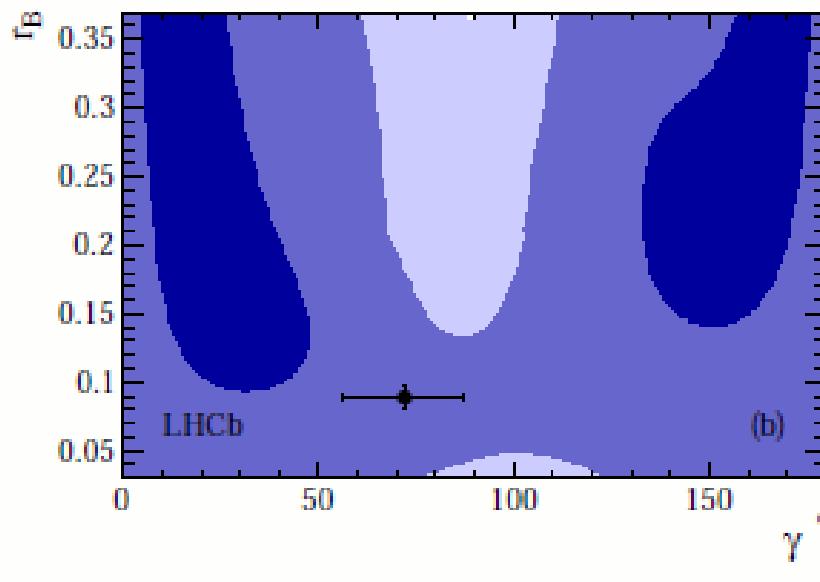
$$\gamma(D_s K) = (123^{+26}_{-39})^\circ$$



# New LHCb results, DK, $D \rightarrow K_s K\pi$

arXiv:1402.2982, 3  $\text{fb}^{-1}$

Observable	Whole Dalitz plot	$K^*(892)^{\pm}$ region
$\mathcal{R}_{\text{SS/OS}}$	$1.528 \pm 0.058 \pm 0.025$	$2.57 \pm 0.13 \pm 0.06$
$\mathcal{R}_{DK/D\pi, \text{SS}}$	$0.092 \pm 0.009 \pm 0.004$	$0.084 \pm 0.011 \pm 0.003$
$\mathcal{R}_{DK/D\pi, \text{OS}}$	$0.066 \pm 0.009 \pm 0.002$	$0.056 \pm 0.013 \pm 0.002$
$\mathcal{A}_{\text{SS, } DK}$	$0.040 \pm 0.091 \pm 0.018$	$0.026 \pm 0.109 \pm 0.029$
$\mathcal{A}_{\text{OS, } DK}$	$0.233 \pm 0.129 \pm 0.024$	$0.336 \pm 0.208 \pm 0.026$
$\mathcal{A}_{\text{SS, } D\pi}$	$-0.025 \pm 0.024 \pm 0.010$	$-0.012 \pm 0.028 \pm 0.010$
$\mathcal{A}_{\text{OS, } D\pi}$	$-0.052 \pm 0.029 \pm 0.017$	$-0.054 \pm 0.043 \pm 0.017$



(not using proper inputs for strong phases/coherence factor)

# The r issue...

- was relevant for  $B^\pm \rightarrow D^{(*)} K^\pm$ , still is for  $B^\pm \rightarrow D^{(*)} \pi^\pm$  (charged B)

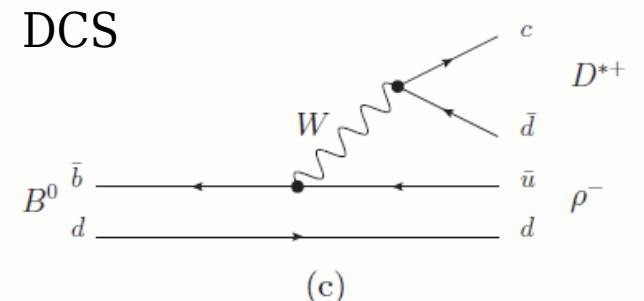
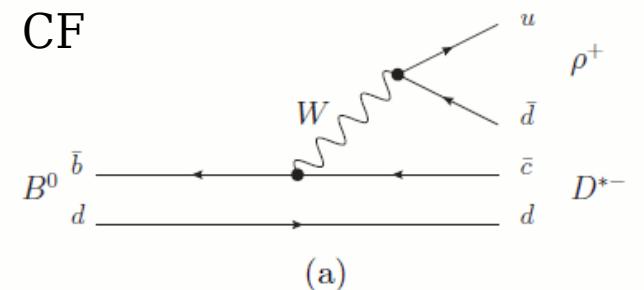
- remember  $B \rightarrow D^{*\pm} \pi^\mp$ ,  $D^{*\pm} \rho^\mp$ ?  
access  $\sin 2\beta + \gamma$  with time-dep analysis

done for  $D^{*\pm} \pi^\mp$  by B-factories  
with external inputs ( $\text{BR}(B \rightarrow D_s^* \pi)$ ) and SU(3)

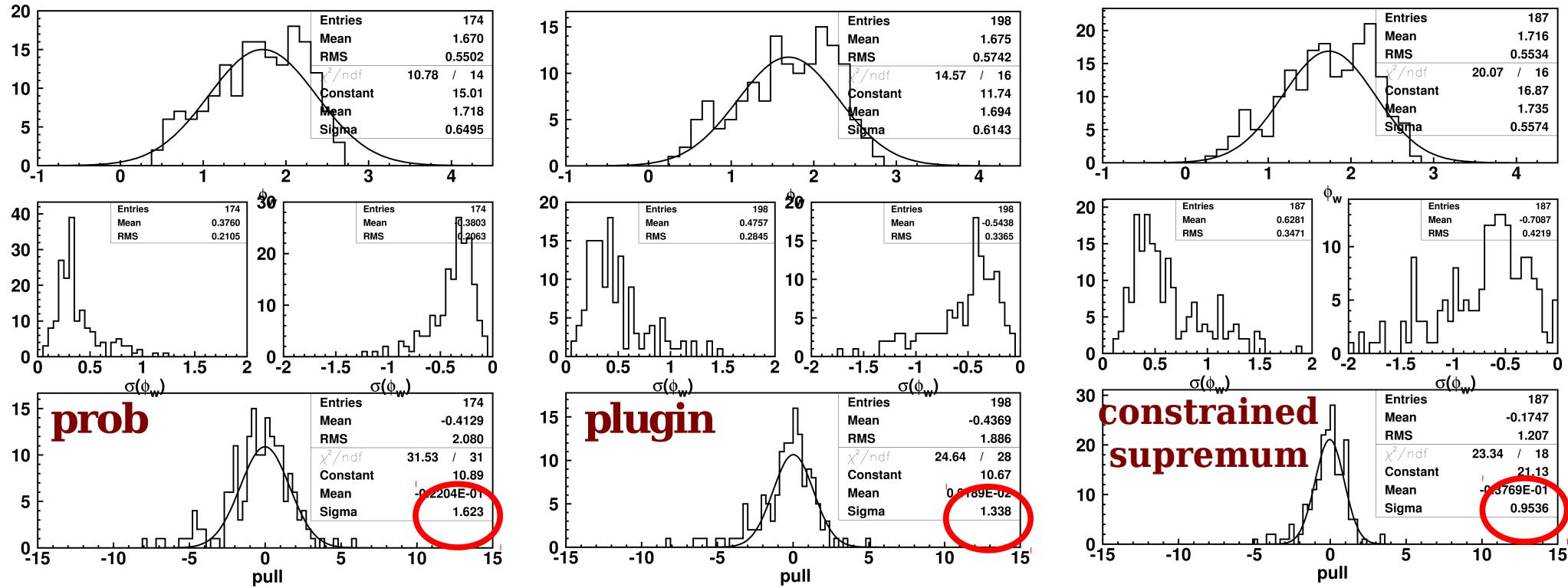
- could be done in  $D^{*\pm} \rho^\mp$  and without external inputs ( $B \rightarrow VV$ )  
recently revisited the analysis, switched to cartesian coordinates:  
 $\{\varphi_w, r_\lambda, \delta_\lambda\} \rightarrow \{x_\lambda, y_\lambda, \bar{x}_\lambda, \bar{y}_\lambda\}$   
 on-going analysis at Belle ( $\rightarrow$  realistic toys)  
 how to extract  $\varphi_w$ ?  
 r is small...  $\sim 0.01$  (DCS/CF)

$$R_{D^{(*)} h} = \frac{|V_{cd}|}{|V_{cs}|} \frac{f_{D^{(*)}}}{f_{D_s^{(*)}}} \sqrt{\frac{\mathcal{B}(B^0 \rightarrow D_s^{(*)+} h^-)}{\mathcal{B}(B^0 \rightarrow D^{(*)-} h^+)}} ,$$

$$R_{D^* \pi} = (1.65 \pm 0.18 \pm 0.04)\%,$$



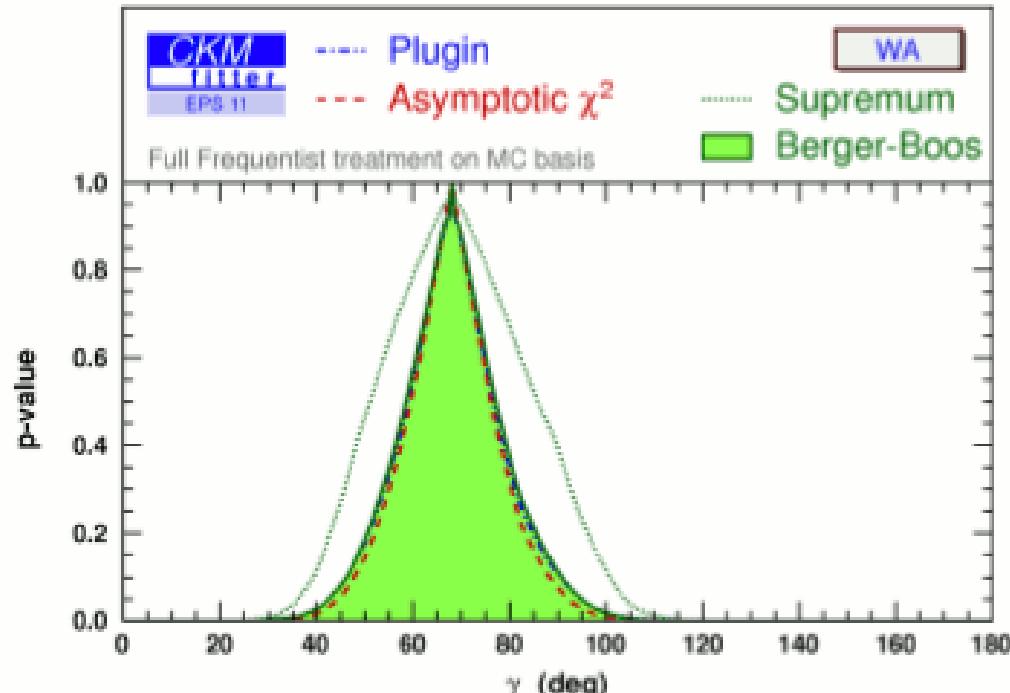
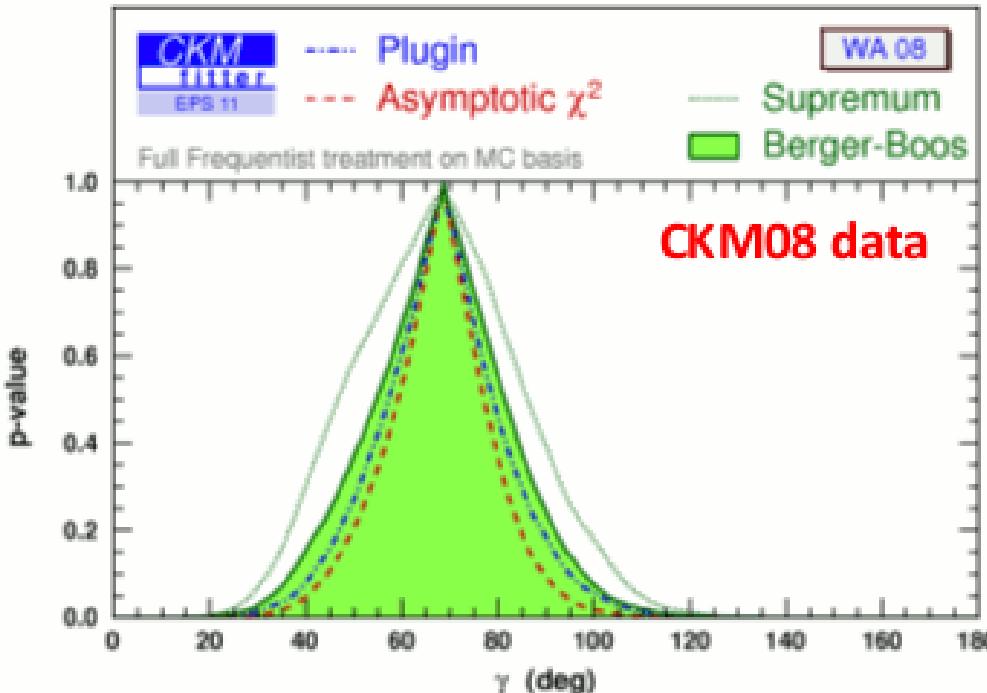
# $B \rightarrow D^{*\pm} \rho^\mp$



seems to work... will pursue the study (r's plots)

# Gamma and the Berger-Boos $p$ value

- The Berger-Boos,  $p_\beta$ ,  $p$ -Value [JASA 89, 427 (1994)] makes a more powerful use of the data than the supremum  $p$  value,  $p_{\text{sup}}$ , by providing control over the nuisance parameters,  $\theta$ . It is a valid / conservative  $p$  value defined as:  $p_\beta = \sup_{\theta \in C_\beta} p(\theta) + \beta$ , where  $C_\beta$  is a level  $1-\beta$  confidence set for the nuisance  $\theta$ .  
⇒ we use the Likelihood under the null hypothesis to infer the confidence region  $C_\beta$ .
- The very increased accuracy on  $\gamma$  not only comes from the new statistical treatment, but also from more accurate measurements, which help constraining the nuisance,  $r_\beta$ . This is illustrated below by re-playing various stat. treatment with CKM08 data.



# Experiment by experiment

