



# TOP FLAVOR VIOLATING DECAYS

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A. Azatov, G. Panico, G. Perez and YS - arXiv:1408.4525

A. Dery, A. Efrati, Y. Nir, YS and V. Susic - arXiv:1408.1371



# OUTLINE

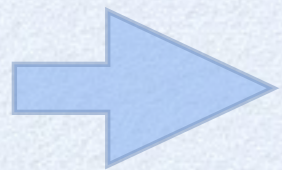
- Introduction and motivation
- Top-flavor violation in composite Higgs models
- Top Higgs flavor violation
- Summary



# INTRODUCTION

Standard Model: 

- diagonal Higgs Yukawa
- universal Z couplings



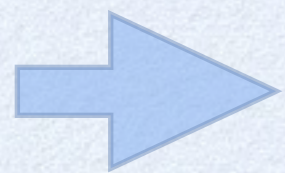
$t \rightarrow qh$  and  $t \rightarrow qZ$  are **highly suppressed**

Experimentally promising channels



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Standard Model: • diagonal Higgs Yukawa  
• universal Z couplings



$t \rightarrow qh$  and  $t \rightarrow qZ$  are **highly suppressed**

Experimentally promising channels

Current bounds:  $\text{BR}(t \rightarrow qh) < 0.56\%$  CMS, ATLAS

(at 95%CL)

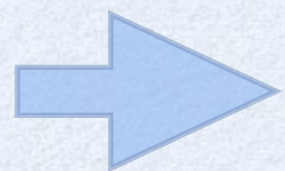
$\text{BR}(t \rightarrow qZ) < 0.05\%$  CMS

well above  
the SM prediction



# INTRODUCTION

Standard Model: • diagonal Higgs Yukawa  
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$t \rightarrow qh$  and  $t \rightarrow qZ$  are **highly suppressed**

Experimentally promising channels

Current bounds:  $\text{BR}(t \rightarrow qh) < 0.56\%$  CMS, ATLAS  
(at 95%CL)  $\text{BR}(t \rightarrow qZ) < 0.05\%$  CMS well above  
the SM prediction

If new physics (NP) at the TeV scale

the NP flavor structure cannot be anarchic

need some alignment with the SM (or universality)





# INTRODUCTION

In many models the SM mixing with the new physics subject to flavor suppression:

- Left-handed fields:  $\lambda_L^t : \lambda_L^c : \lambda_L^u \sim 1 : V_{cb} : V_{ub}$
- Right-handed field:  $\lambda_R^t : \lambda_R^c : \lambda_R^u \sim 1 : \frac{m_c}{m_t V_{cb}} : \frac{m_u}{m_t V_{ub}}$

can use effective field theory to estimate the flavor violating effects



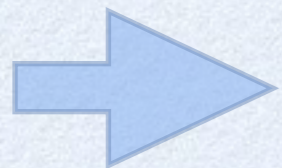
# EFFECTIVE FIELD THEORY

$t \rightarrow cZ$ :

$$\bar{t}_L \gamma^\mu \sigma_3 c_L (H^\dagger \sigma_3 \overleftrightarrow{D}_\mu H)$$

$$\bar{t}_R \gamma^\mu c_R (H^\dagger \overleftrightarrow{D}_\mu H)$$

$$\bar{t}_L \gamma^\mu c_L (H^\dagger \overleftrightarrow{D}_\mu H)$$



$$(g_{tc,L} \bar{t}_L \gamma^\mu c_L + g_{tc,R} \bar{t}_R \gamma^\mu c_R) Z_\mu$$

$$g_{tc,L} \sim \frac{g}{2c_W} \frac{v^2}{M_*^2} V_{cb}$$

$$g_{tc,R} \sim \frac{g}{2c_W} \frac{v^2}{M_*^2} \frac{m_c}{m_t V_{cb}}$$

(assume that the top mixing is an order one)

$M_*$  - the new-physics scale

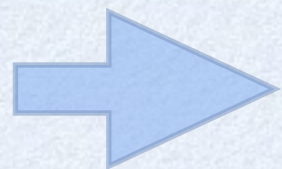


# EFFECTIVE FIELD THEORY

$t \rightarrow ch$ :

$$\bar{t}_L \tilde{H} c_R (H^\dagger H)$$

$$\bar{c}_L \tilde{H} t_R (H^\dagger H)$$



$$(y_{tc,L} \bar{t}_L c_R + y_{tc,R} \bar{c}_L t_R) h$$

$$y_{tc,L} \sim \frac{v^2}{M_*^2} \frac{m_c}{v V_{cb}}$$

$$y_{tc,R} \sim \frac{v^2}{M_*^2} \frac{m_t}{v} V_{cb}$$

$M_*$  - the new-physics scale



# TOP FLAVOR VIOLATION IN COMPOSITE HIGGS MODELS



# COMPOSITE HIGGS MODELS

the framework:

- Global  $SO(5)$  broken to  $SO(4)$ , at scale  $-f$ .
- Solves the gauge/hierarchy problem.
- The Higgs is a composite and pseudo Nambu-Goldstone boson and can be naturally light.
- Non SM contributions to  $Zbb$  are protected by custodial symmetry. Agashe, Contino, Da-Rold, Pomarol  
hep-ph/0605341



# COMPOSITE HIGGS MODELS

The framework (continued):

- Two sectors:
  - **Elementary** - massless SM fields
  - **Composite** - additional massive resonances, no flavor structure (anarchic) - characterized by the scale  $M_*$
- Elementary / composite **mixing**:
  - Left:  $\lambda_L^t : \lambda_L^c : \lambda_L^u \sim 1 : V_{cb} : V_{ub}$
  - Right:  $\lambda_R^t : \lambda_R^c : \lambda_R^u \sim 1 : m_c/m_t V_{cb} : m_u/m_t V_{ub}$
- FCNC are suppressed by the small SM mass and mixing
- Minimizing of the fine-tuning leads to  $\lambda_L^t \sim \lambda_R^t$



$$m_{\text{SM}}^i \sim \frac{v}{f} \frac{\lambda_L^i \lambda_R^i}{M_*}$$

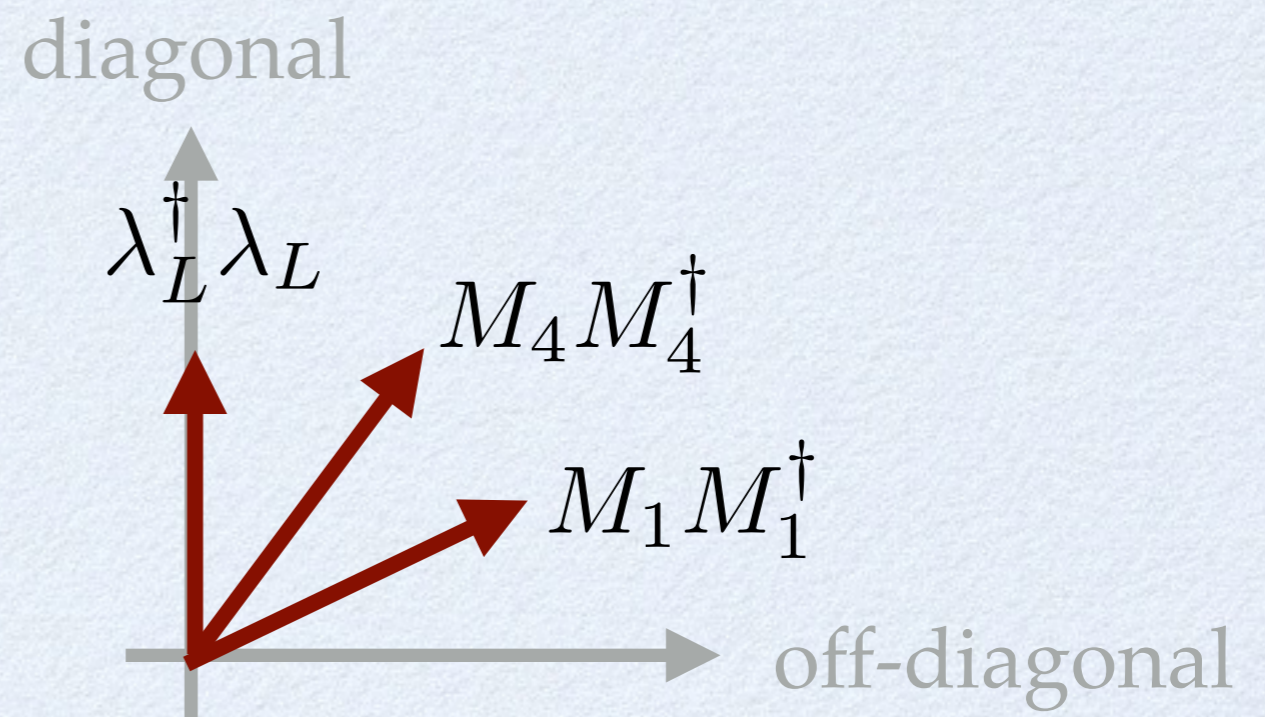


# THE UP FLAVOR STRUCTURE

The flavor parameters:

- Vector-like composite mass matrices -  $M_1, M_4$
- Elementary / composite mixing:
  - Left mixing -  $\lambda_L$
  - Right mixing -  $\lambda_R$

(similar for the right)



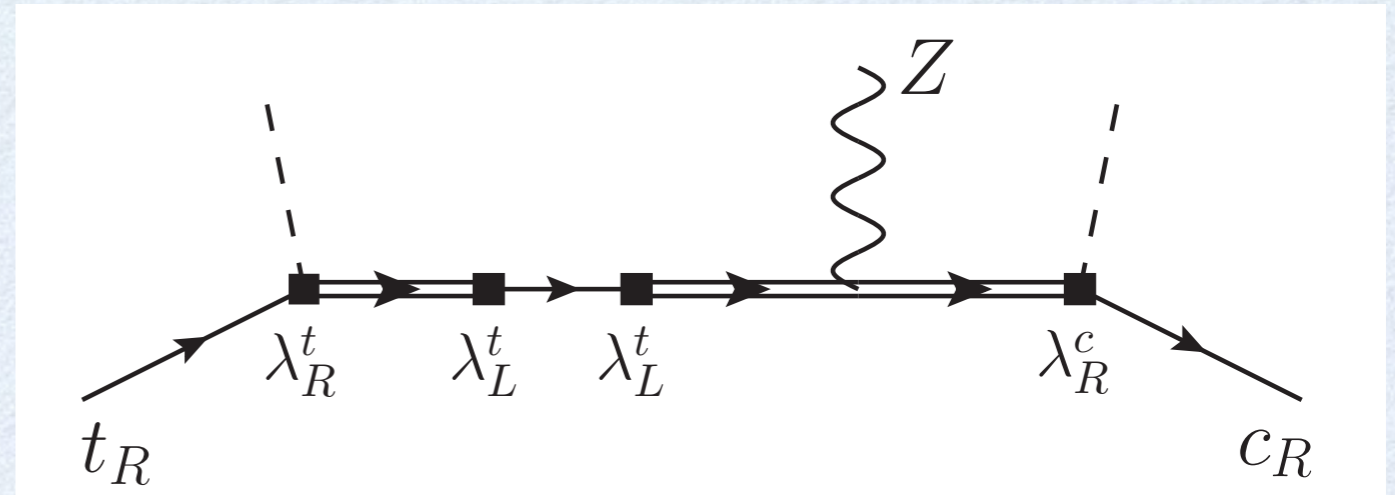
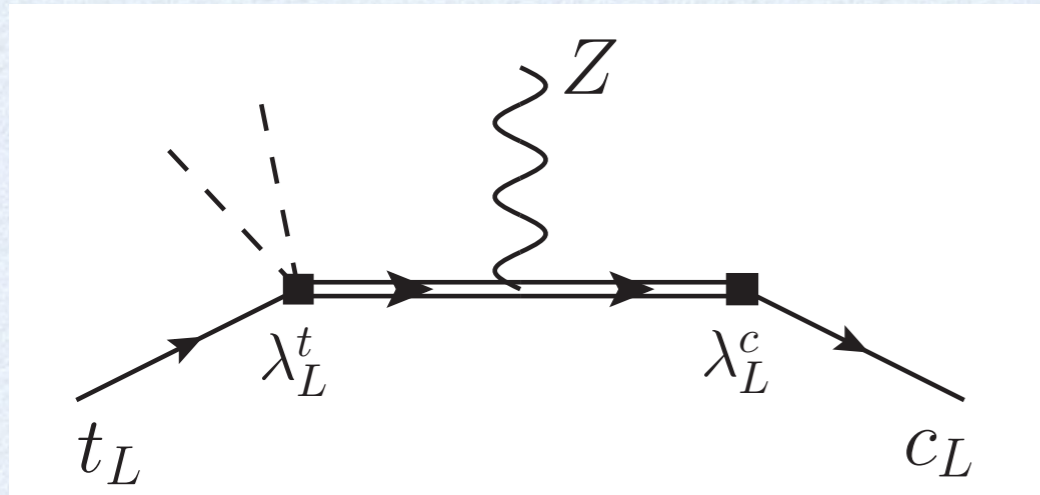
Misalignment in the  
flavor space



Flavor violation



# TOP/Z FLAVOR VIOLATION



$$g_{tc,L} \sim \frac{g}{2c_W} \frac{v^2}{2f^2} \frac{\lambda_L^t}{M_*} \frac{\lambda_L^c}{M_*}$$

$$g_{tc,R} \sim \frac{g}{2c_W} \frac{v^2}{f^2} \frac{\lambda_R^t}{M_*} \frac{\lambda_R^c}{M_*} \left( \frac{\lambda_L^t}{M_*} \right)^2$$

$$\text{BR}(t \rightarrow cZ) \sim (0.2, 0.8) \times 10^{-5} \left( \frac{700 \text{ GeV}}{M_*} \right)^4$$

Right current is larger than the left - can be tested

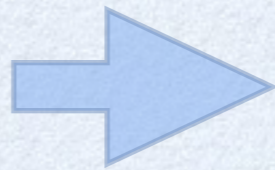


# TOP HIGGS FLAVOR VIOLATION



# FROGGATT-NIELSEN

Approximate U(1)'s  
symmetries



Yukawa are suppressed  
by different powers of the  
breaking parameters

$$\epsilon_i \simeq \lambda = 0.2$$

Consider two possibilities

- SM + nonrenormalizable terms
- Supersymmetric model + nonrenormalizable terms



# FROGGATT-NIELSEN

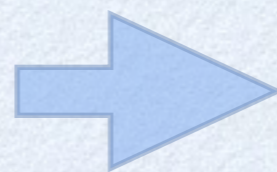
SM + nonrenormalizable terms:

$$\mathcal{L}_Y = -\lambda_{ij}^u Q_i \bar{U}_j \phi - \frac{\lambda'_{ij}{}^u}{\Lambda^2} Q_i \bar{U}_j \phi (\phi^\dagger \phi) + \text{h.c.}$$

→

$$Y_{ij}^u = \frac{\sqrt{2} m_i^u}{v} \delta_{ij} + \frac{v^2}{\Lambda^2} (V_L^u \lambda'^u V_R^{u\dagger})_{ij}$$

$(\phi^\dagger \phi)$  does not carry a FN charge



$\lambda$  and  $\lambda'$  have the same parametric structure

$$\epsilon_i \simeq \lambda = 0.2 \quad Y_{tc} \sim \frac{v^2}{\Lambda^2} \lambda \lesssim 10^{-3} \ll 0.14$$

FCNC bounds (charm mixing)

$$\frac{v^2}{\Lambda^2} \lesssim 10^{-2}$$

← current bound



# FROGGATT-NIELSEN

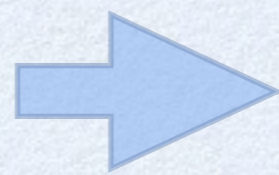
Supersymmetric model + nonrenormalizable terms:

$$\mathcal{L}_Y = \lambda_{ij}^u Q_i \bar{U}_j \phi_u + \frac{\lambda'^u_{ij}}{\Lambda^2} Q_i \bar{U}_j \phi_u (\phi_u \phi_d)$$

➔

$$(Y_h^u)_{ij} = \frac{c_\alpha}{s_\beta} \frac{m_i^u}{v} \delta_{ij} + \frac{v^2 c_{\alpha+\beta} s_\beta}{2\sqrt{2}\Lambda^2} (V_L^u \lambda'^u V_R^{u\dagger})_{ij}$$

If  $(\phi_u \phi_d)$  carry  
a FN charge



$\lambda$  and  $\lambda'$  have different  
parametric structure

Use “holomorphic zeros” to align the new sources of flavor violation and avoid FCNC bounds

Leurer, Nir, Seiberg  
hep-ph/9310320



# FROGGATT-NIELSEN

Supersymmetric model + nonrenormalizable terms:

$$\mathcal{L}_Y = \lambda_{ij}^u Q_i \bar{U}_j \phi_u + \frac{\lambda'^u_{ij}}{\Lambda^2} Q_i \bar{U}_j \phi_u (\phi_u \phi_d)$$

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$$(Y_h^u)_{ij} = \frac{c_\alpha}{s_\beta} \frac{m_i^u}{v} \delta_{ij} + \frac{v^2 c_{\alpha+\beta} s_\beta}{2\sqrt{2}\Lambda^2} (V_L^u \lambda'^u V_R^{u\dagger})_{ij}$$

Approximate  $U(1) \times U(1)$  FN symmetry:

$$V_L^u \lambda'^u V_R^{u\dagger} \sim \begin{pmatrix} \lambda^{14} & \lambda^{11} & \lambda^{12} \\ \lambda^7 & \lambda^4 & \lambda^5 \\ \lambda^3 & 1 & \lambda^1 \end{pmatrix} \Rightarrow Y_{tc} \sim \frac{v^2}{\Lambda^2} \sim 0.1$$

no problem with charm mixing



# SUMMARY

- $\text{BR}(t \rightarrow cZ) \sim 10^{-5}$  with dominant right-handed current is generic in composite Higgs models, even in the presence of custodial protection - can be tested at the next run of the LHC.
- In Froggatt-Nielsen models the bounds from meson mixing push the new physics scale such that the  $t \rightarrow ch$  rate is too small to be observed.
- By using further alignment of the new flavor violation sources, as demonstrated in Supersymmetric models, the  $t \rightarrow ch$  rates can be much larger and even saturate the current bound.



# BACKUP SLIDES



# FINE-TUNING

the Higgs potential:  $V(h) = \alpha \sin^2 (h/f) + \beta \sin^4 (h/f)$

$$\xi \equiv \sin^2 (\langle h \rangle / f) = -\alpha / 2\beta \ll 1$$

typical values:  $\beta \sim \frac{N_c}{16\pi^2} (\lambda_L^t \lambda_R^t)^2$      $\alpha_{L,R}^t \sim \frac{N_c}{16\pi^2} (\lambda_{L,R}^t)^2 M_*^2$

the fine-tuning:  $\text{FT}^{-1} \sim \frac{-2\beta \xi}{\alpha_{L,R}^t} \sim \frac{\min [(\lambda_R^t)^2, (\lambda_L^t)^2] v^2}{M_*^2 f^2}$

$$\sim \frac{m_t^2}{\max [\lambda_{L,R}^t]^2} \sim y_t \frac{v^2}{M_* f} \sim \frac{y_t}{g_*} \frac{v^2}{f^2}$$

$\lambda_L^t \sim \lambda_R^t$  minimize the tuning



# COMPOSITE TOP/Z FLAVOR VIOLATION

$$\mathcal{L} \supset \bar{u}_L^i (m_u^{\text{SM}})^{ij} u_R^j + \bar{u}_L^i (g_{Z,L}^{\text{SM}})^{ij} Z_\mu \gamma^\mu u_L^j + \bar{u}_R^i (g_{Z,R}^{\text{SM}})^{ij} Z_\mu \gamma^\mu u_R^j$$

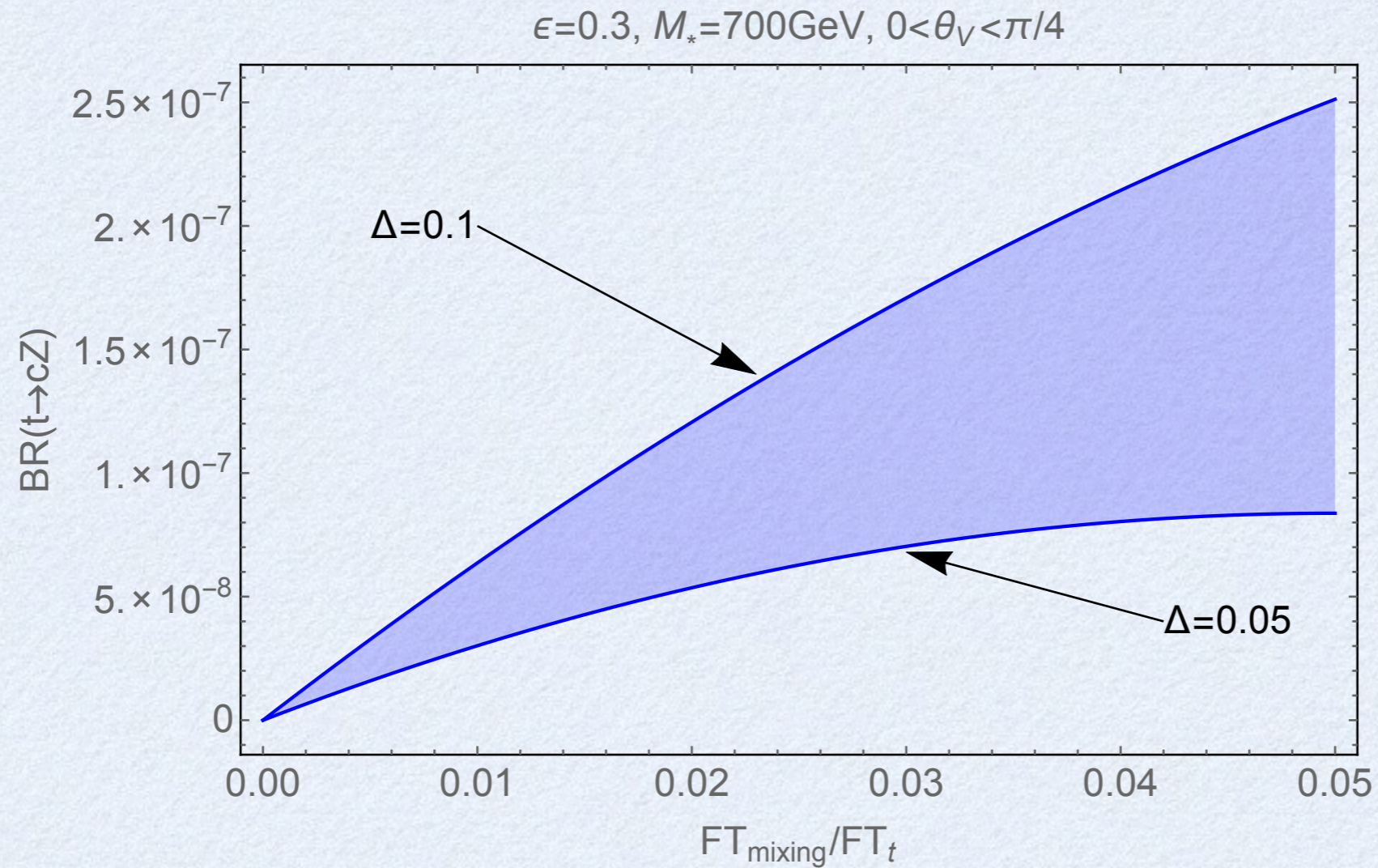
$$m_u^{\text{SM}} = \frac{\epsilon}{\sqrt{2}} \lambda_L [M_4^{-1} - M_1^{-1}] \lambda_R \quad \epsilon \equiv \frac{v}{f}$$

$$g_{Z,L}^{\text{SM}} = -\frac{g}{2c_W} \frac{\epsilon^2}{2} \lambda_L \left[ \left( M_1^\dagger M_1 \right)^{-1} + \left( M_4^\dagger M_4 \right)^{-1} \right] \lambda_L^\dagger$$

$$g_{Z,R}^{\text{SM}} = -\frac{g}{2c_W} \frac{\epsilon}{\sqrt{2}} \left[ m_u^{\text{SM}\dagger} \lambda_L \left( M_4 M_4^\dagger M_4 \right)^{-1} \lambda_R + h.c. \right]$$



# FINE-TUNING VS. FLAVOR VIOLATION





# PARAMETER COUNTING

2 generation limit no CPV:

- $M_1, M_4, \lambda_L, \lambda_R \rightarrow 0: SO(2)^6 = SO(2)_Q \times SO(2)_U \times SO(2)_{Q4,L} \times SO(2)_{Q4,R} \times SO(2)_{U1,L} \times SO(2)_{U1,R}$
- 6 real  $2 \times 2$  matrices  $(M_1, M_4, \lambda_{L1}, \lambda_{L4}, \lambda_{R1}, \lambda_{R4})$
- $\lambda_{L1} = \lambda_{L4}, \lambda_{R1} = \lambda_{R4}$  respect  $SO(5)$  - finite potential. Flavor does commute with the global  $SO(5)$ .
- single elementary (composite) mixing for  $\lambda_L, \lambda_R$  and 1 and 4 have same eigenvalues :  $24 - 6 = 18$  parameters.
- $SO(2)^6$  allows to remove 6 unphysical parameters.
- 12 parameters: 8 eigenvalues  $(M_1, M_4, \lambda_L, \lambda_R) + 4$  mixing angles (misalignment between the matrices)..