

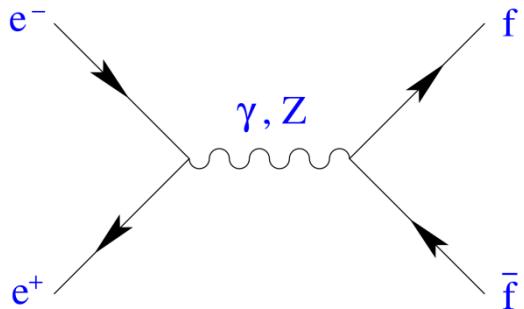
Tau Physics @ TeraZ

A. Pich

IFIC, Valencia

Sixth TLEP Workshop, CERN, 16-18 October 2013

High-Energy $\tau^+\tau^-$ Pairs



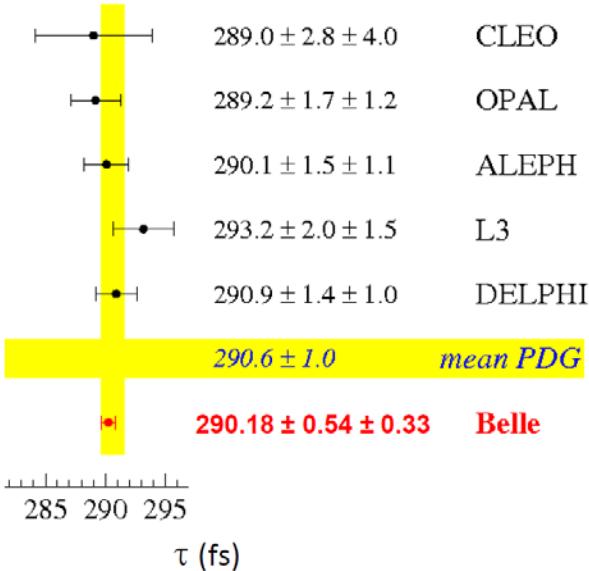
- Boosted $\tau^+\tau^-$ pairs
- Clean signature
- Manageable backgrounds

Many τ results are still dominated by LEP data

Statistics much higher at the B-Factories (backgrounds also)

Let's assume a high-statistics (10^{10}) $\tau^+\tau^-$ data sample at high energies

τ Lifetime



- Not a single lifetime measurement published by Babar/Belle
- Preliminary Babar value @ TAU 2004:
 $\tau_\tau = 289.40 \pm 0.91 \pm 0.90$ (unpublished)
- Preliminary Belle value @ TAU 2012:
 $\tau_\tau = 290.18 \pm 0.54 \pm 0.33$ (unpublished)

Charged-current
universality

$$B_{\tau \rightarrow e} = \frac{B_{\tau \rightarrow \mu}}{0.972559 \pm 0.000005} = \frac{\tau_\tau}{(1632.9 \pm 0.6) \times 10^{-15} \text{ s}}$$

$$\left(B_{\tau \rightarrow \mu} / B_{\tau \rightarrow e} \right)_{\text{exp}} = 0.9761 \pm 0.0028$$

Non-BF: 0.9725 ± 0.0039

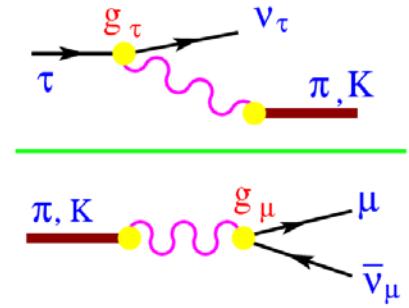
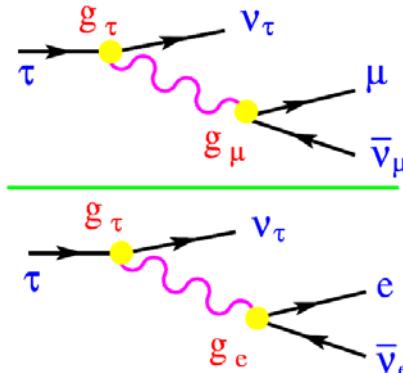
BaBar '10: 0.9796 ± 0.0039

0.4% precision

Charged Current Universality

$$|g_\mu / g_e|$$

$B_{\tau \rightarrow \mu} / B_{\tau \rightarrow e}$	1.0018 ± 0.0014
$B_{\pi \rightarrow \mu} / B_{\pi \rightarrow e}$	1.0021 ± 0.0016
$B_{K \rightarrow \mu} / B_{K \rightarrow e}$	0.9978 ± 0.0020
$B_{K \rightarrow \pi\mu} / B_{K \rightarrow \pi e}$	1.0010 ± 0.0025
$B_{W \rightarrow \mu} / B_{W \rightarrow e}$	0.991 ± 0.009



$$|g_\tau / g_e|$$

$B_{\tau \rightarrow \mu} \tau_\mu / \tau_\tau$	1.0024 ± 0.0021
$B_{W \rightarrow \tau} / B_{W \rightarrow e}$	1.023 ± 0.011

$$B_{\tau \rightarrow e} \tau_\mu / \tau_\tau$$

$$\Gamma_{\tau \rightarrow \pi} / \Gamma_{\pi \rightarrow \mu}$$

$$\Gamma_{\tau \rightarrow K} / \Gamma_{K \rightarrow \mu}$$

$$B_{W \rightarrow \tau} / B_{W \rightarrow \mu}$$

$$1.0006 \pm 0.0021$$

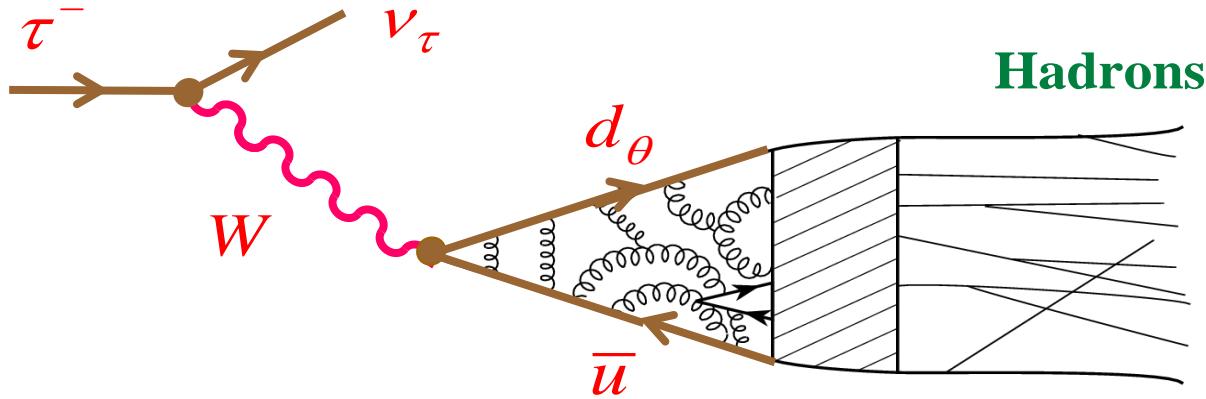
$$0.9956 \pm 0.0031$$

$$0.9852 \pm 0.0072$$

$$1.032 \pm 0.012$$

Precise measurement of $\text{Br}(W \rightarrow \tau) / \text{Br}(W \rightarrow \mu, e)$ needed

HADRONIC TAU DECAY

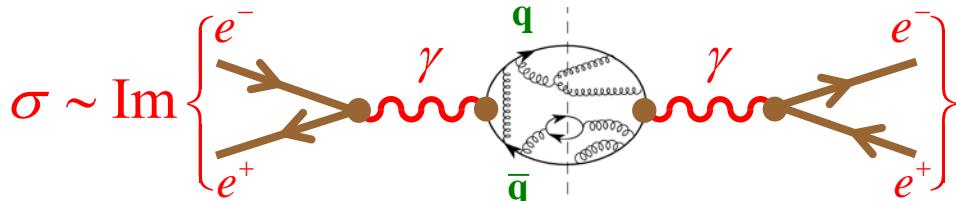


$$d_\theta = V_{ud} \ d + V_{us} \ s$$

Only lepton massive enough to decay into hadrons

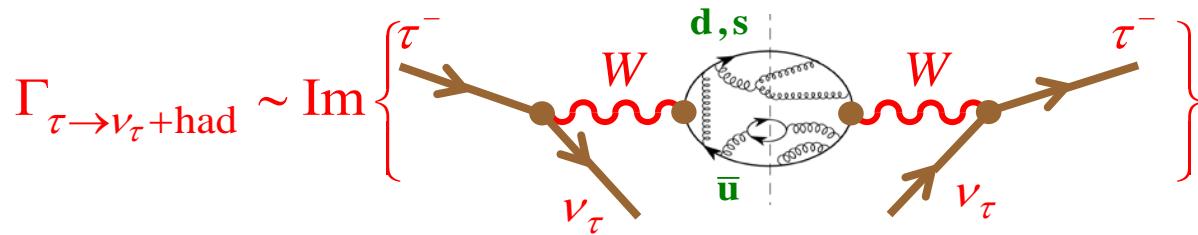
$$R_\tau \equiv \frac{\Gamma(\tau^- \rightarrow \nu_\tau + \text{Hadrons})}{\Gamma(\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e)} \approx N_c \quad ; \quad R_\tau = \frac{1 - B_e - B_\mu}{B_e} = 3.636 \pm 0.011$$

$$R_\tau = \frac{1}{B_e^{\text{univ}}} - 1.97256 = 3.6331 \pm 0.0088 \quad ; \quad R_\tau = \frac{\text{Br}(\tau^- \rightarrow \nu_\tau + \text{Hadrons})}{B_e^{\text{univ}}} = 3.6280 \pm 0.0094$$



$$\frac{\sigma(e^+ e^- \rightarrow \text{had})}{\sigma(e^+ e^- \rightarrow \mu^+ \mu^-)} = 12\pi \operatorname{Im} \Pi_{\text{em}}(s)$$

$$\Pi_{\text{em}}^{\mu\nu}(q) \equiv i \int d^4x e^{iqx} \langle 0 | T[J_{\text{em}}^\mu(x) J_{\text{em}}^\nu(0)] | 0 \rangle = (-g^{\mu\nu} q^2 + q^\mu q^\nu) \Pi_{\text{em}}(q^2)$$



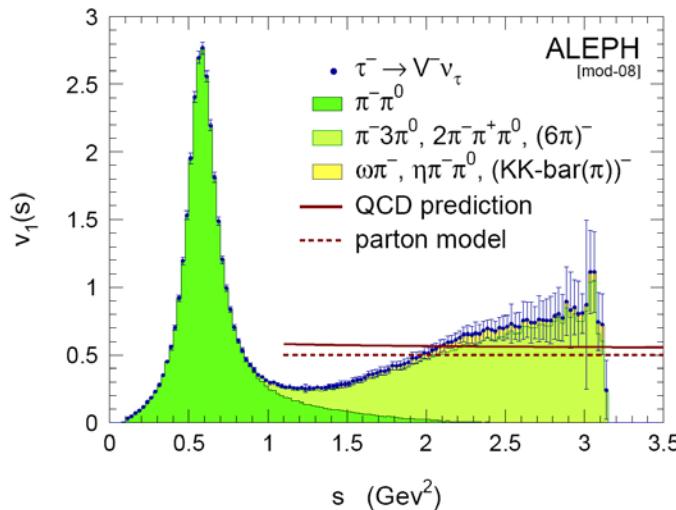
$$R_\tau \equiv \frac{\Gamma(\tau^- \rightarrow \nu_\tau + \text{had})}{\Gamma(\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e)} = 12\pi \int_0^{m_\tau^2} \frac{ds}{m_\tau^2} \left(1 - \frac{s}{m_\tau^2}\right)^2 \left[\left(1 + 2\frac{s}{m_\tau^2}\right) \operatorname{Im} \Pi^{(1)}(s) + \operatorname{Im} \Pi^{(0)}(s) \right]$$

$$\Pi^{(J)}(s) \equiv |V_{ud}|^2 \left[\Pi_{ud,V}^{(J)}(s) + \Pi_{ud,A}^{(J)}(s) \right] + |V_{us}|^2 \left[\Pi_{us,V}^{(J)}(s) + \Pi_{us,A}^{(J)}(s) \right]$$

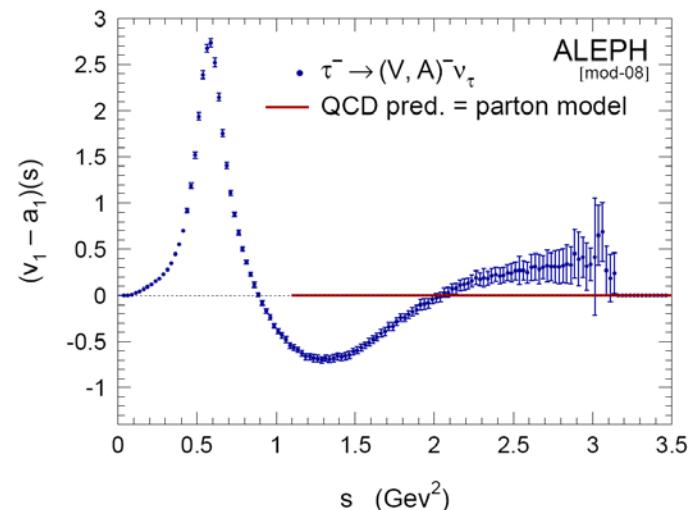
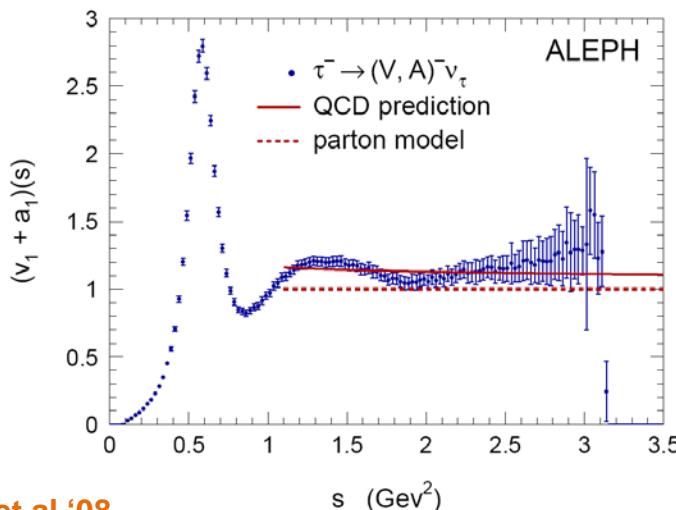
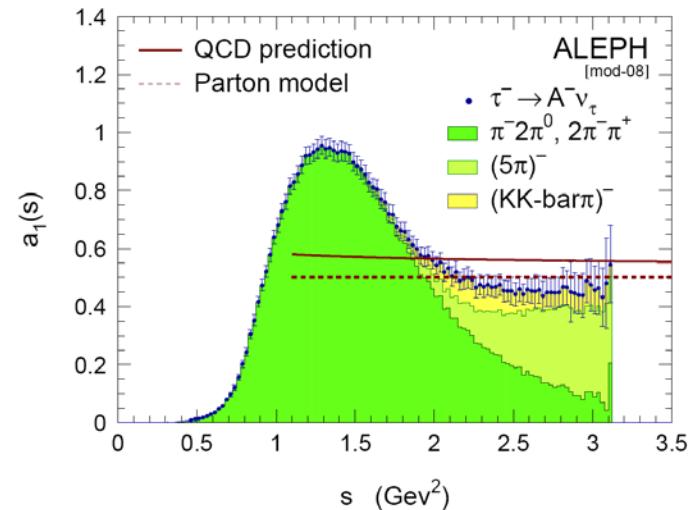
$$\Pi_{ij,J}^{\mu\nu}(q) \equiv i \int d^4x e^{iqx} \langle 0 | T[J_{ij}^\mu(x) J_{ij}^\nu(0)^\dagger] | 0 \rangle = (-g^{\mu\nu} q^2 + q^\mu q^\nu) \Pi_{ij,J}^{(1)}(q^2) + q^\mu q^\nu \Pi_{ij,J}^{(0)}(q^2)$$

SPECTRAL FUNCTIONS

$$v_1(s) = 2\pi \operatorname{Im} \Pi_{ud,V}^{(0+1)}(s)$$



$$a_1(s) = 2\pi \operatorname{Im} \Pi_{ud,A}^{(0+1)}(s)$$



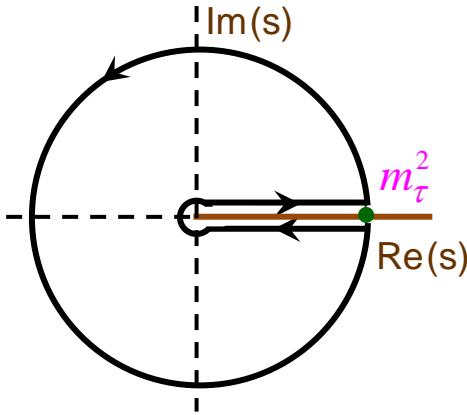
Better
data
needed

Davier et al '08

QCD Prediction of R_τ

Braaten-Narison-Pich'92

$$R_\tau \equiv \frac{\Gamma(\tau^- \rightarrow \nu_\tau + \text{had})}{\Gamma(\tau^- \rightarrow \nu_\tau e^- \bar{v}_e)} = 12\pi \int_0^1 dx (1-x)^2 \left[(1+2x) \text{Im} \Pi^{(1)}(xm_\tau^2) + \text{Im} \Pi^{(0)}(xm_\tau^2) \right]$$



$$R_\tau = 6\pi i \oint_{|x|=1} dx (1-x)^2 \left[(1+2x) \Pi^{(0+1)}(xm_\tau^2) - 2x \Pi^{(0)}(xm_\tau^2) \right]$$

$$\Pi^{(J)}(s) = \sum_{D=2n} \frac{C_D^{(J)}(s, \mu) \langle O_D(\mu) \rangle}{(-s)^{D/2}}$$

OPE

$$R_\tau = N_C S_{\text{EW}} (1 + \delta_{\text{P}} + \delta_{\text{NP}}) = R_{\tau,V} + R_{\tau,A} + R_{\tau,S}$$

$$S_{\text{EW}} = 1.0201 (3)$$

$$\delta_{\text{NP}} = -0.0059 \pm 0.0014$$

Marciano-Sirlin, Braaten-Li, Erler

Fitted from data (Davier et al)

$$\delta_{\text{P}} = a_\tau + 5.20 a_\tau^2 + 26 a_\tau^3 + 127 a_\tau^4 + \dots \approx 20\% \quad ; \quad a_\tau \equiv \alpha_s(m_\tau)/\pi$$

Baikov-Chetyrkin-Kühn

Perturbative ($m_q=0$)

$$-s \frac{d}{ds} \Pi^{(0+1)}(s) = \frac{1}{4\pi^2} \sum_{n=0} K_n \left(\frac{\alpha_s(-s)}{\pi} \right)^n$$

$$K_0 = K_1 = 1 \quad , \quad K_2 = 1.63982 \quad , \quad K_3 = 6.37101 \quad , \quad K_4 = 49.07570$$

Baikov-Chetyrkin-Kühn '08

$\rightarrow \delta_P = \sum_{n=1} K_n A^{(n)}(\alpha_s) = a_\tau + 5.20 a_\tau^2 + 26 a_\tau^3 + 127 a_\tau^4 + \dots$

Le Diberder- Pich '92

CIPT

FOPT

$$A^{(n)}(\alpha_s) \equiv \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} (1 - 2x + 2x^3 - x^4) \left(\frac{\alpha_s(-s)}{\pi} \right)^n = a_\tau^n + \dots ; \quad a_\tau \equiv \alpha_s(m_\tau)/\pi$$

Power Corrections

Braaten-Narison-Pich '92

$$\Pi_{\text{OPE}}^{(0+1)}(s) \approx \frac{1}{4\pi^2} \sum_{n \geq 2} \frac{C_{2n} \langle O_{2n} \rangle}{(-s)^n}$$

$$C_4 \langle O_4 \rangle \approx \frac{2\pi}{3} \langle 0 | \alpha_s G^{\mu\nu} G_{\mu\nu} | 0 \rangle$$

$$\delta_{\text{NP}} \approx \frac{-1}{2\pi i} \oint_{|x|=1} dx (1 - 3x^2 + 2x^3) \sum_{n \geq 2} \frac{C_{2n} \langle O_{2n} \rangle}{(-xm_\tau^2)^n} = -3 \frac{C_6 \langle O_6 \rangle}{m_\tau^6} - 2 \frac{C_8 \langle O_8 \rangle}{m_\tau^8}$$

Suppressed by m_τ^6

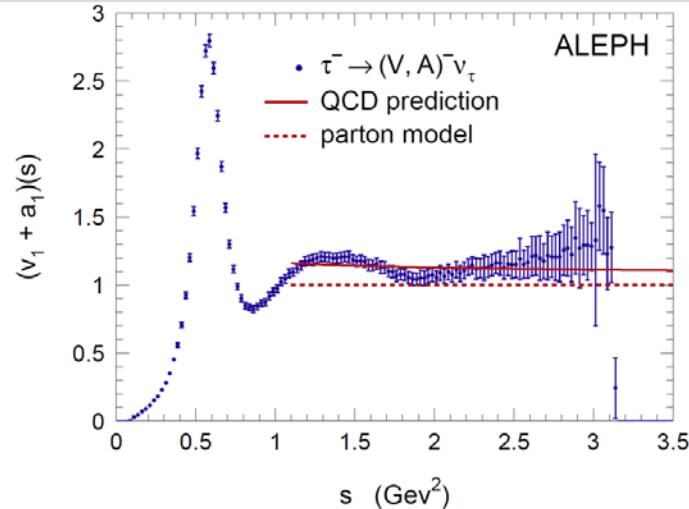
[additional chiral suppression in $C_6 \langle O_6 \rangle^{V+A}$]

Spectral Function Distribution

■ Moments:

$$R_\tau^{kl}(s_0) \equiv \int_0^{s_0} ds \left(1 - \frac{s}{s_0}\right)^k \left(\frac{s}{m_\tau^2}\right)^l \frac{dR_\tau}{ds}$$

Sensitivity to power corrections (k, l)



The non-perturbative contribution to R_τ can be obtained from the invariant-mass distribution of the final hadrons

$$\delta_{NP} = -0.0059 \pm 0.0014$$

Davier et al. (ALEPH data)

■ Duality Violations: (non-pinched moments)

$$\delta_{NP} = -0.003 \pm 0.012$$

$$\oint_{x=1} dx \Pi(x m_\tau^2)$$

$$\text{Im } \Pi(s) \sim \exp(-\delta - \gamma s) \sin(\alpha + \beta s)$$

Boito et al. (OPAL data)

Recent $\alpha_s(m_\tau)$ Analyses

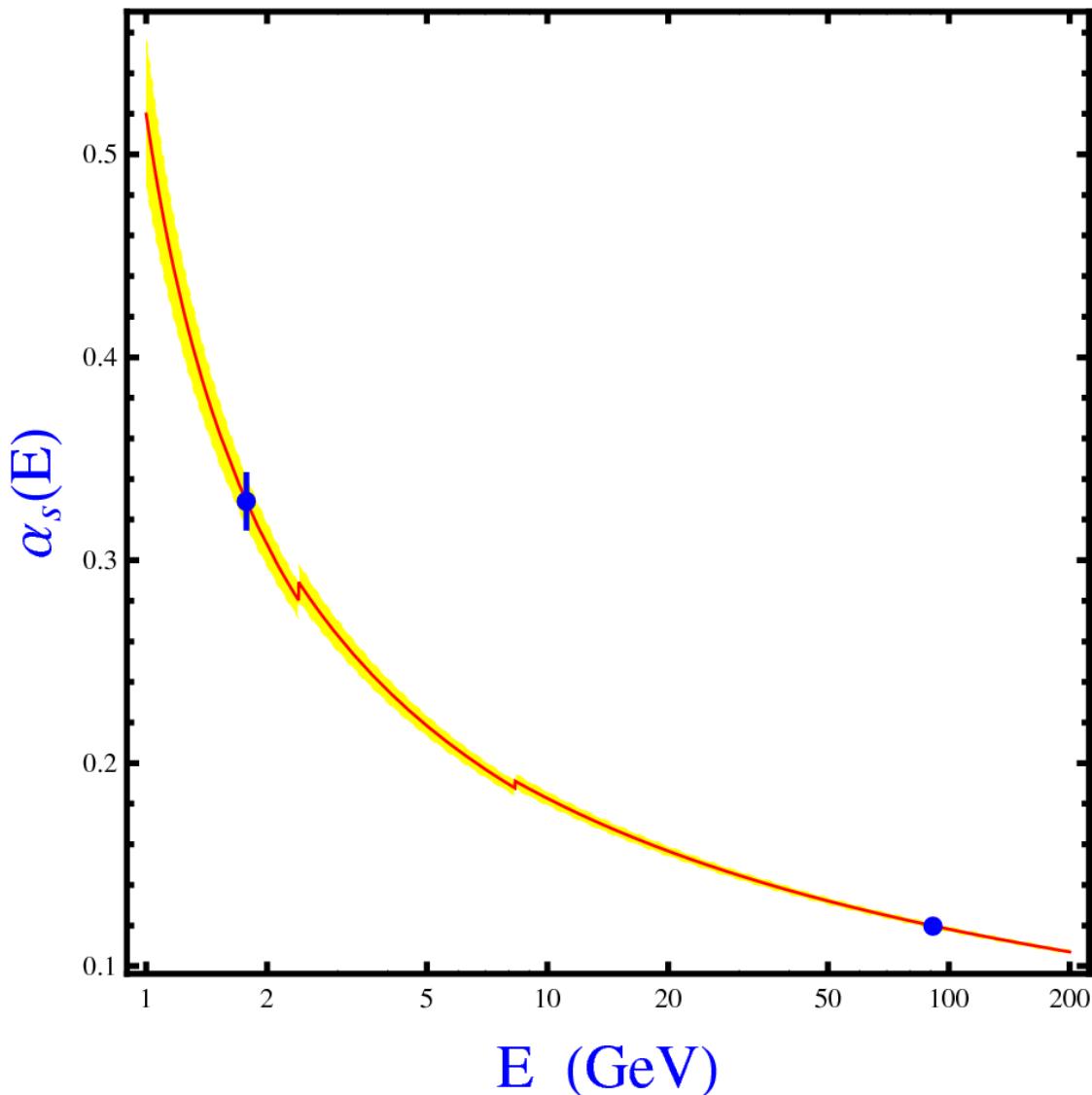
Reference	Method	δ_{NP}	δ_P	$\alpha_s(m_\tau)$	$\alpha_s(m_Z)$
Baikov et al	CIPT, FOPT		0.1998 (43)	0.332 (16)	0.1202 (19)
Davier et al	CIPT	- 0.0059 (14)	0.2066 (70)	0.344 (09)	0.1212 (11)
Beneke-Jamin	BSR + FOPT	- 0.007 (3)	0.2042 (50)	0.316 (06)	0.1180 (08)
Maltman-Yavin	PWM + CIPT	+ 0.012 (18)	-	0.321 (13)	0.1187 (16)
Menke	CIPT, FOPT		0.2042 (50)	0.342 (11)	0.1213 (12)
Narison	CIPT, FOPT		-	0.324 (08)	0.1192 (10)
Caprini-Fischer	BSR + CIPT		0.2037 (54)	0.322 (16)	-
Abbas et al	IFOPT		0.2037 (54)	0.338 (10)	
Cvetič et al	$\beta_{exp} + CIPT$		0.2040 (40)	0.341 (08)	0.1211 (10)
Boito et al 2012	CIPT, DV	- 0.002 (12)	-	0.347 (25)	0.1216 (27)
	FOPT, DV	- 0.004 (12)		0.325 (18)	0.1191 (22)
Pich	CIPT	- 0.0059 (14)	0.1995 (33)	0.339 (13)	0.1210 (15)
	FOPT			0.318 (14)	0.1185 (18)
My Average	CIPT, FOPT		0.1995 (33)	0.329 (13)	0.1198 (15)

CIPT: Contour-improved perturbation theory
 FOPT: Fixed-order perturbation theory
 BSR: Borel summation of renormalon series
 IFOPT: Improved FOPT

β_{exp} : Expansion in derivatives of α_s (β function)
 PWM: Pinched-weight moments
 CIPTm: Modified CIPT (conformal mapping)
 DV: Duality violation (OPAL only)

Present Status

A. Pich, arXiv:1303.2262



$$\alpha_s(m_\tau^2) = 0.329 \pm 0.013$$



$$\alpha_s(M_Z^2) = 0.1198 \pm 0.0015$$

$$\alpha_s(M_Z^2)_{Z\text{ width}} = 0.1197 \pm 0.0028$$

The most precise test of
Asymptotic Freedom

$$\alpha_s^\tau(M_Z^2) - \alpha_s^Z(M_Z^2) = 0.0001 \pm 0.0015 \tau \pm 0.0028_Z$$

V_{us} Determination

Gámiz-Jamin-Pich-Prades-Schwab

$$|V_{us}|^2 = \frac{R_{\tau,S}^{00}}{\frac{R_{\tau,V+A}^{00}}{|V_{ud}|^2} - \delta R_{\tau,\text{th}}^{00}}$$

$$\delta R_{\tau}^{kl} \approx 24 \frac{m_s^2(m_{\tau}^2)}{m_{\tau}^2} \Delta_{kl}(\alpha_s)$$

$$\delta R_{\tau,\text{th}}^{00} \equiv \underbrace{0.1544(37)}_{J=0} + \underbrace{0.086(32)}_{m_s(2 \text{ GeV}) = 94(6) \text{ MeV}} = 0.240(32)$$

$$R_{\tau,S}^{00} = 0.1612(28)$$

$$R_{\tau,V+A}^{00} = 3.4671(84)$$

$$|V_{ud}| = 0.97425(22)$$

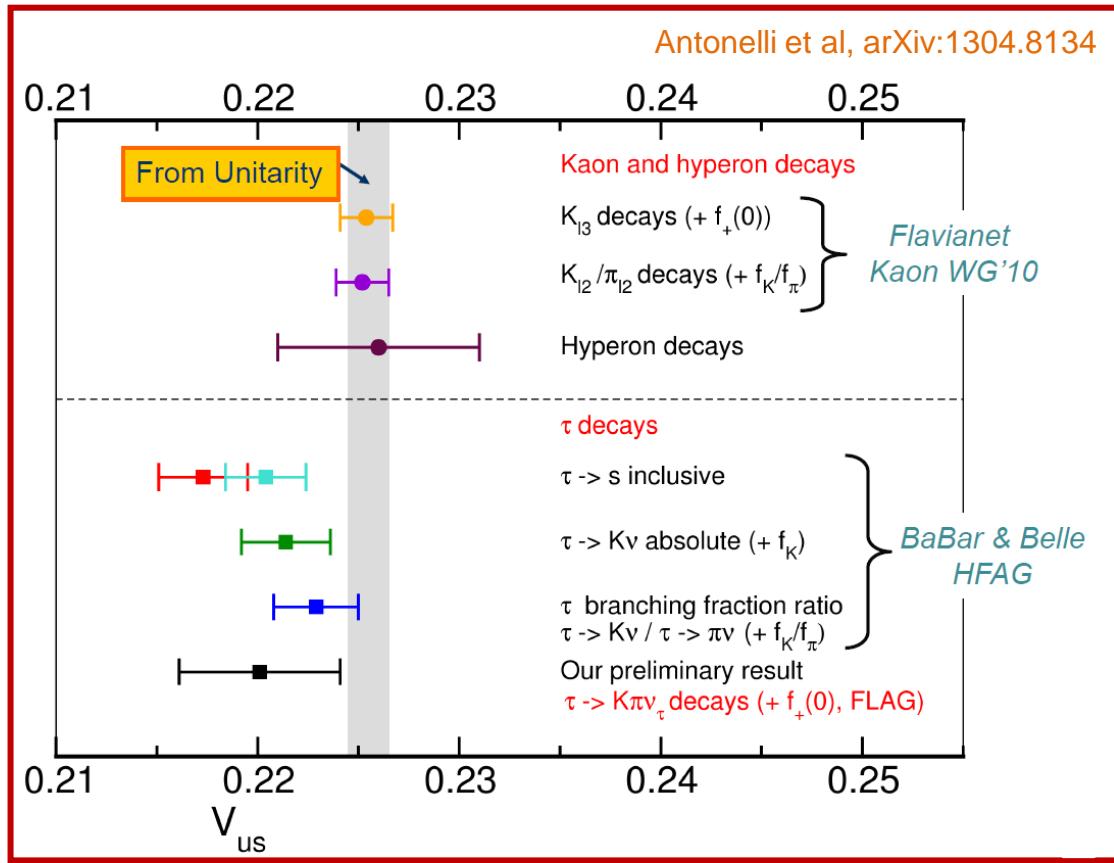


$$|V_{us}| = 0.2173 \pm 0.0020_{\text{exp}} \pm 0.0010_{\text{th}}$$

$$\mathbf{K_{I3}:} \quad |V_{us}| = 0.2238 \pm 0.0011 \quad [f_+(0) = 0.967 \pm 0.004]$$

The τ could give the most precise V_{us} determination

V_{us} from τ



Branching fraction	HFAG Winter 2012 fit	
$\Gamma_{10} = K^- \nu_\tau$	$(0.6955 \pm 0.0096) \cdot 10^{-2}$	$\Rightarrow (0.713 \pm 0.003) \cdot 10^{-2}$
$\Gamma_{16} = K^- \pi^0 \nu_\tau$	$(0.4322 \pm 0.0149) \cdot 10^{-2}$	$\Rightarrow (0.4473 \pm 0.0244) \cdot 10^{-2}$
$\Gamma_{23} = K^- 2\pi^0 \nu_\tau$ (ex. K^0)	$(0.0630 \pm 0.0222) \cdot 10^{-2}$	
$\Gamma_{28} = K^- 3\pi^0 \nu_\tau$ (ex. K^0, η)	$(0.0419 \pm 0.0218) \cdot 10^{-2}$	
$\Gamma_{35} = \pi^- \bar{K}^0 \nu_\tau$	$(0.8206 \pm 0.0182) \cdot 10^{-2}$	$\Rightarrow (0.8627 \pm 0.0353) \cdot 10^{-2}$
$\Gamma_{110} = X_s^- \nu_\tau$	$(2.8746 \pm 0.0498) \cdot 10^{-2}$	$\Rightarrow (2.9496 \pm 0.0571) \cdot 10^{-2}$

Slight underestimate of
 τ branching ratios ?

Use modes
measured
in K decays

Larger R_s

$$R_s = 0.1663 (34)$$



Larger V_{us}

$$|V_{us}| = 0.2207 (25)$$

Better data needed!

Remaining Problems

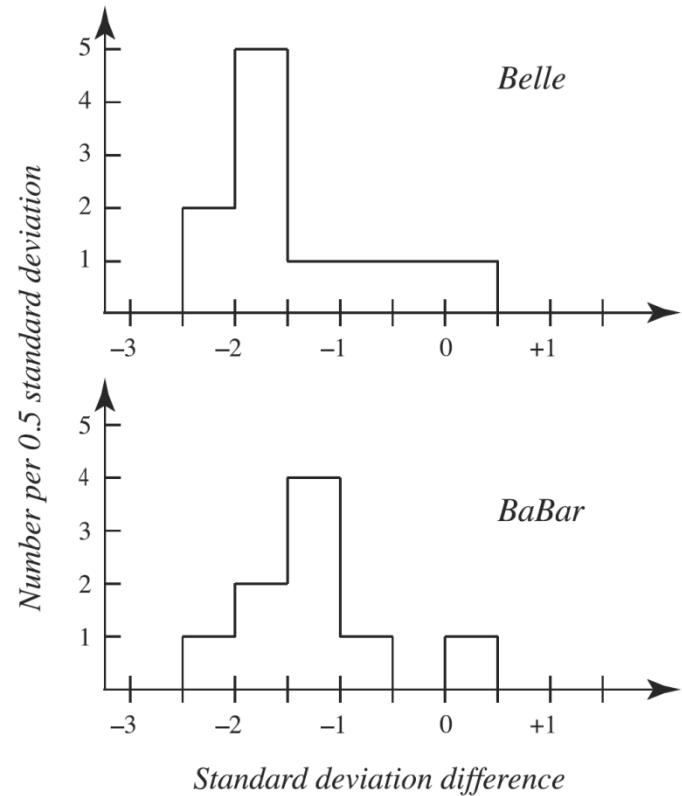
PDG 2012:

“Eighteen of the 20 *B*-factory branching fraction measurements are smaller than the non-*B*-factory values. The average normalized difference between the two sets of measurements is -1.30 (-1.41 for the 11 Belle measurements and -1.24 for the 9 BaBar measurements)”

Missing modes ?

Modes measured by the two BF experiments

Mode	BaBar – Belle Normalized Difference (# σ)
$\pi^-\pi^+\pi^-\nu_\tau$ (ex. K^0)	+1.4
$K^-\pi^+\pi^-\nu_\tau$ (ex. K^0)	-2.9
$K^-K^+\pi^-\nu_\tau$	-2.9
$K^-K^+K^-\nu_\tau$	-5.4
$\eta K^-\nu_\tau$	-1.0
$\phi K^-\nu_\tau$	-1.3



PHYSICS OUTLOOK



Many more interesting topics

- Tests of QCD and the Electroweak Theory
- Looking for Signals of New Phenomena
- Superb Tool for New Physics Searches

**High-statistics τ data samples at high energies
would provide many useful informations**

Systematic errors & backgrounds need to be studied

Backup Slides

Perturbative Uncertainty on $\alpha_s(m_\tau)$

$$-s \frac{d}{ds} \Pi^{(0+1)}(s) = \frac{1}{4\pi^2} \sum_{n=0} K_n a(-s)^n$$

$$\delta_P = \underbrace{\sum_{n=1} K_n A^{(n)}(\alpha_s)}_{\text{CIPT}} = \underbrace{\sum_{n=0} r_n a_\tau^n}_{\text{FOPT}}$$

$$r_n = K_n + g_n$$

$$A^{(n)}(\alpha_s) \equiv \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} (1 - 2x + 2x^3 - x^4) a_\tau (-x m_\tau^2)^n = a_\tau^n + \dots ; \quad a_\tau \equiv \alpha_s(m_\tau)/\pi$$

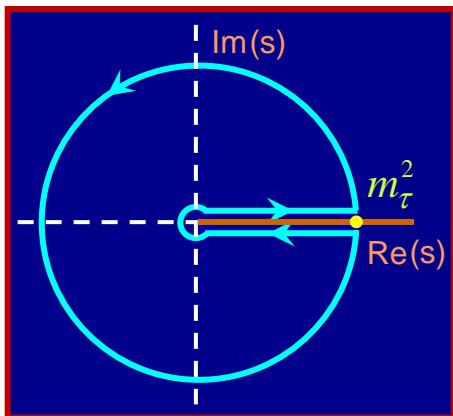
n	1	2	3	4	5
K_n	1	1.6398	6.3710	49.0757	
g_n	0	3.5625	19.9949	78.0029	307.78
r_n	1	5.2023	26.3659	127.079	

The dominant corrections come from the contour integration

Le Diberder- Pich 1992

Large running of a_s along the circle $s = m_\tau^2 e^{i\phi}$, $\phi \in [0, 2\pi]$

$$A^{(n)}(a_\tau) \equiv \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} (1 - 2x + 2x^3 - x^4) a_\tau (-x m_\tau^2)^n = a_\tau^n + \dots ; \quad a_\tau \equiv \alpha_s(m_\tau)/\pi$$



$$A^{(1)}(a_\tau) = a_\tau - \frac{19}{24} \beta_1 a_\tau^2 + \left[\beta_1^2 \left(\frac{265}{288} - \frac{\pi^2}{12} \right) - \frac{19}{24} \beta_2 \right] a_\tau^3 + \dots$$

$$a(-s) \simeq \frac{a_\tau}{1 - \frac{\beta_1}{2} a_\tau \log(-s/m_\tau^2)} = \frac{a_\tau}{1 - i \frac{\beta_1}{2} a_\tau \phi} = a_\tau \sum_n \left(i \frac{\beta_1}{2} a_\tau \phi \right)^n ; \quad \phi \in [0, 2\pi]$$

FOPT expansion only convergent if $a_\tau < 0.14$ (0.11) [at 1 (3) loops]

Experimentally $\alpha_\tau \approx 0.11$



FOPT should not be used
(divergent series)

FOPT suffers a large renormalization-scale dependence (Le Diberder- Pich , Menke)

The difference between FOPT and CIPT grows at higher orders

Renormalons

$$D(s) \equiv -s \frac{d}{ds} \Pi^{(0+1)}(s) = \frac{1}{4\pi^2} \sum_{n=0} K_n a(-s)^n$$

Asymptotic series

Borel Summation:

$$B(t) \equiv \sum_{n=0} K_{n+1} \frac{t^n}{n!} \quad \longrightarrow \quad D(s) = \frac{1}{4\pi^2} \left\{ 1 + \int_0^\infty dt e^{-t/a(-s)} B(t) \right\}$$

However, $B(t)$ has pole singularities at

- $u \equiv -\beta_1 t/2 = +n \quad (n \geq 2)$
- $u \equiv -\beta_1 t/2 = -n \quad (n \geq 1)$

Infrared Renormalons

Ultraviolet Renormalons

IR - n Renormalon



Ambiguity: $\delta D(s) \sim \left(\frac{\Lambda^2}{-s}\right)^n$

Renormalon Hypothesis: Asymptotics already reached

Modelling a better behaved FOPT

(Beneke – Jamin)

- Large higher-order K_n corrections could cancel the g_n ones
Happens in the “large- β_0 ” approximation (UV renormalon chain)
- $D = 4$ corrections very suppressed in R_τ
→ **n = 2 IR renormalons can do the job** ($K_n \approx -g_n$)
- No sign of renormalon behaviour in known coefficients
→ **n = -1,2,3 renormalons + linear polynomial**
5 unknown constants fitted to K_n ($2 \leq n \leq 5$). $K_5 = 283$ assumed
- **Borel summation:** large renormalon contributions. Smaller α

Nice model of higher orders. But too many different possibilities ...

(Descotes-Genon – Malaescu)