

Beam-Beam studies for TLEP (and update on TMCI)

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Thanks to E. Metral, R. Tomas and F. Zimmermann

Outline

Model and benchmarking

*Parameters, layout,
simulation code, synchrotron radiation*

Dynamic effects

Working point, phase advances

Simulated luminosity

Preliminary performance estimates

Lifetime

*Preliminary estimates, number of IPs,
energy acceptance*

Update on TMCI

Model, synchro-betatron effects, thresholds

Beam parameters

- Latest beam parameters provided by F. Zimmermann

	TELP Z	TLEP W	TLEP H	TLEP t	TLEP t B
E [GeV]	45	80	120	175	175
N [e/bunch]	4.0e11	1.0e11	3.7e11	0.88e11	7.0e11
$\epsilon_{x/y}$ [nm]	29.2/0.06	3.3/0.017	7.5/0.015	2.0/0.002	16.0/0.016
$\beta_{x/y}$ [m]	0.5/1.0e-3	0.2/1.0e-3	0.5/1.0e-3	1.0/1.0e-3	1.0/1.0e-3
τ_{\parallel} [turns]	1319	242	72	23	23
$\xi_{x/y}/IP$	0.068	0.086	0.094	0.057	0.057
L/IP [cm ² .s ⁻¹]	5.9e35	1.6e35	5.1e34	1.3e34	1.0e34

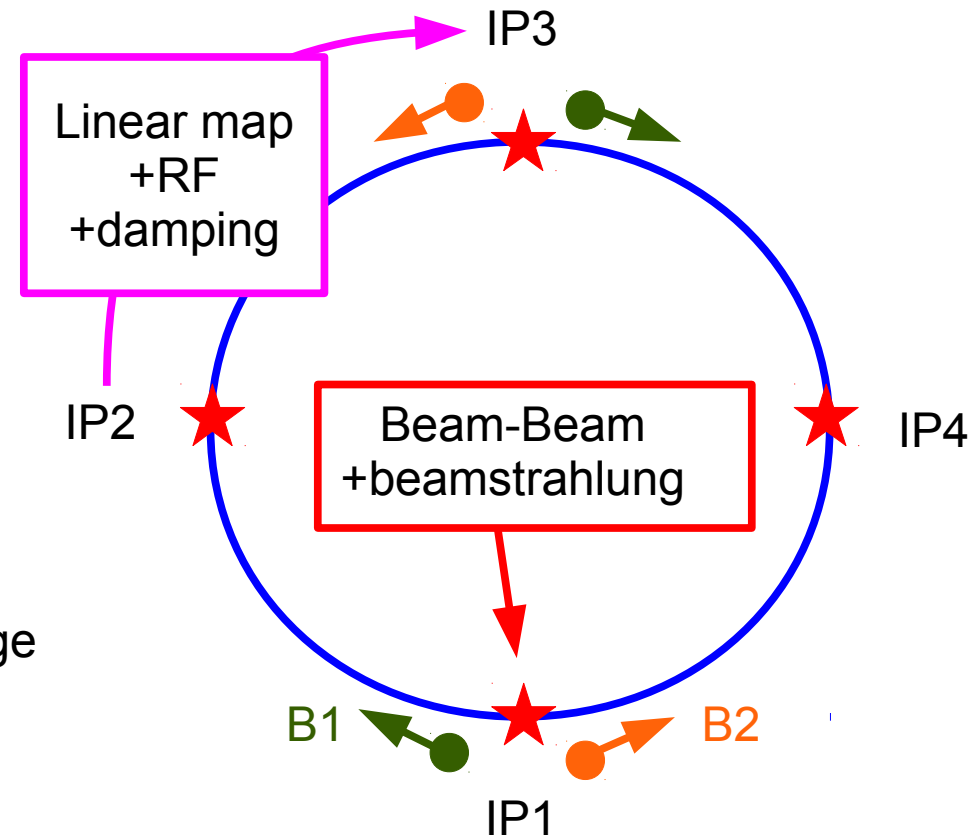
→ Beam parameters relevant for the beam-beam simulations

→ **Assume 100km circumference in all cases**

→ By design there are 4 IPs: the total beam-beam parameter scales accordingly

Model

- Strong-strong beam-beam model based on the code BeamBeam3D by J. Qiang
- Beam-beam module fully benchmarked against data and theory
- Allows to include impedance effects benchmarked with HEADTAIL and theory
- Track 2 bunches per beam: neglecting long-range interactions this covers the full picture
- Recently added synchrotron radiation (damping and beamstrahlung as relevant for TLEP)
- Benchmarking done against existing code and theory (will be presented in the following slides)
- Luminosity is computed in 3D taking into account the focusing effect of the beam-beam interaction



→ Only non linear elements are the beam-beam interactions

→ Arcs are modeled by linear transfer maps

→ Optics are the same in all IPs

Synchrotron radiation

- The average number of photons emitted for a given bending radius is:

$$n_y = \frac{5\sqrt{3}}{6} \frac{\alpha \gamma}{\rho} \Delta s$$

- Photon emission follows a Poisson distribution:

$$P(n) = \frac{n_y^n e^{-n_y}}{n!}$$

- Cumulative distribution of energy probability law:

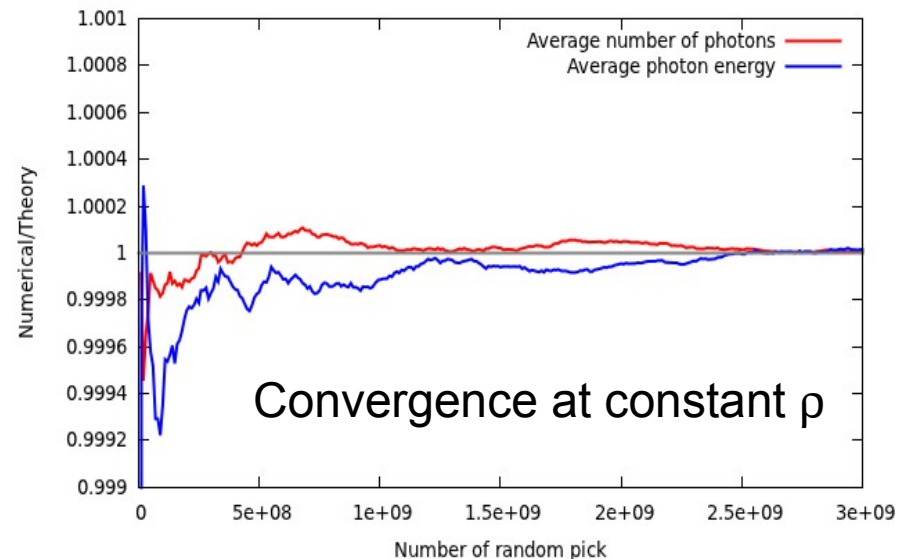
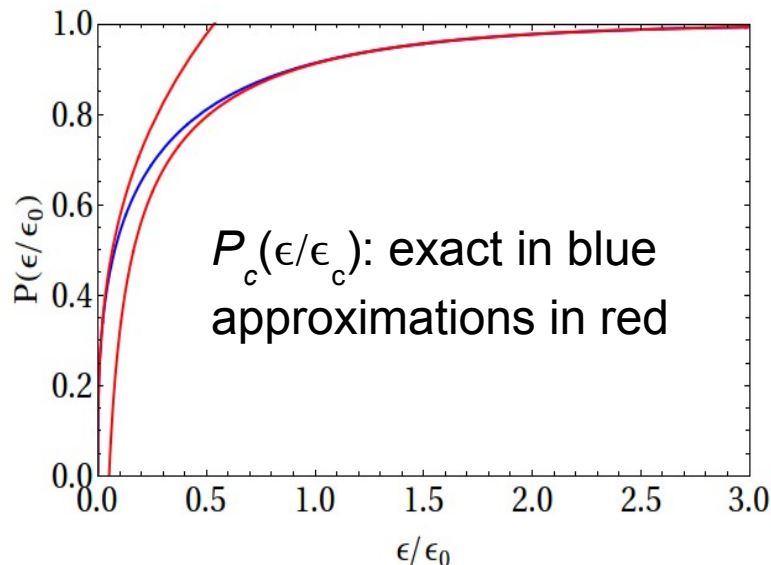
$$P_c(\epsilon/\epsilon_c) = \frac{3}{5\pi} \int_0^{\epsilon/\epsilon_c} \int_{\epsilon/\epsilon_c}^{\infty} K_{5/3}(x) dx$$

with $P_c(\epsilon/\epsilon_c \gg 1) \approx 1 - 0.24 \frac{e^{-\frac{\epsilon}{\epsilon_c}}}{\sqrt{\frac{\epsilon}{\epsilon_c}}}$ and $P_c(\epsilon/\epsilon_c \ll 1) \approx 1.23 \left(\frac{\epsilon}{\epsilon_c}\right)^{(1/3)}$

- Photon critical energy:
$$\epsilon_c = \frac{3\hbar \gamma^3 c}{2\rho}$$

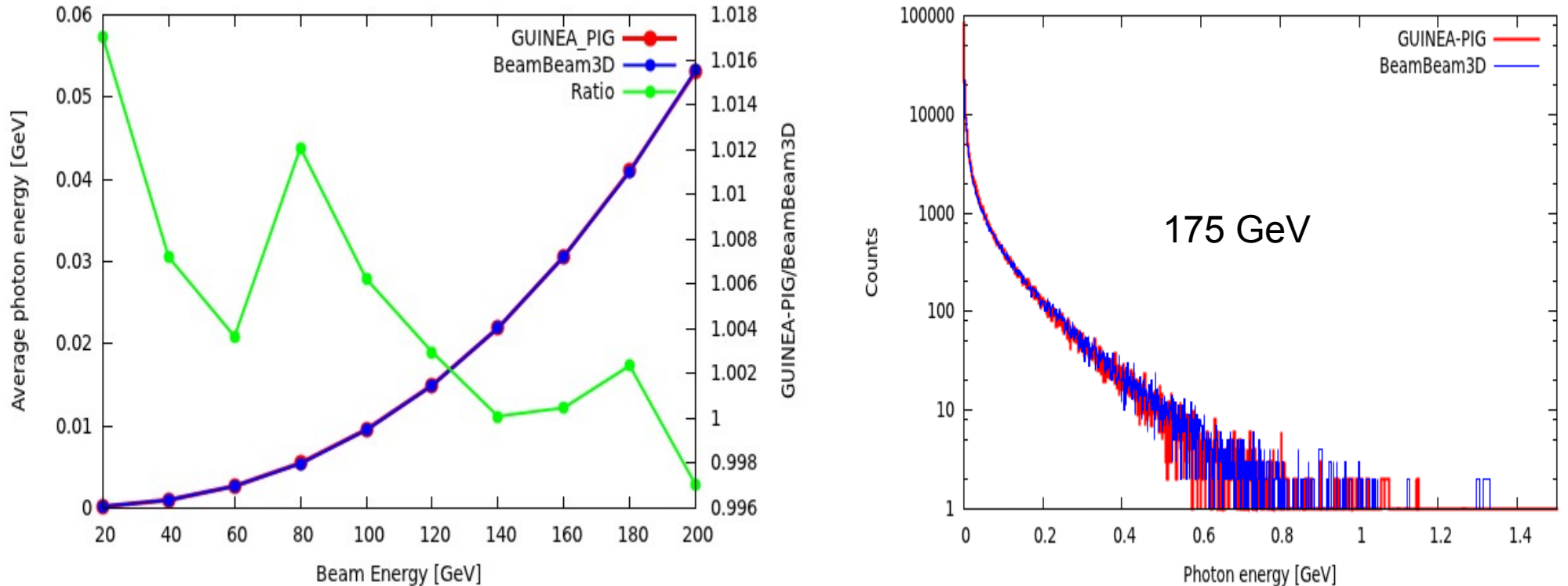
Photon emission

- Each time a kick is given to a particle ρ and ϵ_c are computed accordingly
- The number of photons emitted over the interaction length is selected at random following $P(n)$
- The energy of each photon is then computed by picking a random number between 0 and 1 and inverting $P_c(\epsilon/\epsilon_c)$
- For small and large energies analytical approximations are used. For intermediate energies a look-up table of the exact numerical integration is used: $0.02 < \epsilon/\epsilon_c < 5.0$



Comparison with GUINEA-PIG

- Test case with round beams at the IP. Photon energy distribution for 2M macro-particles, 13 slices and a grid of 128x128 after a single collision

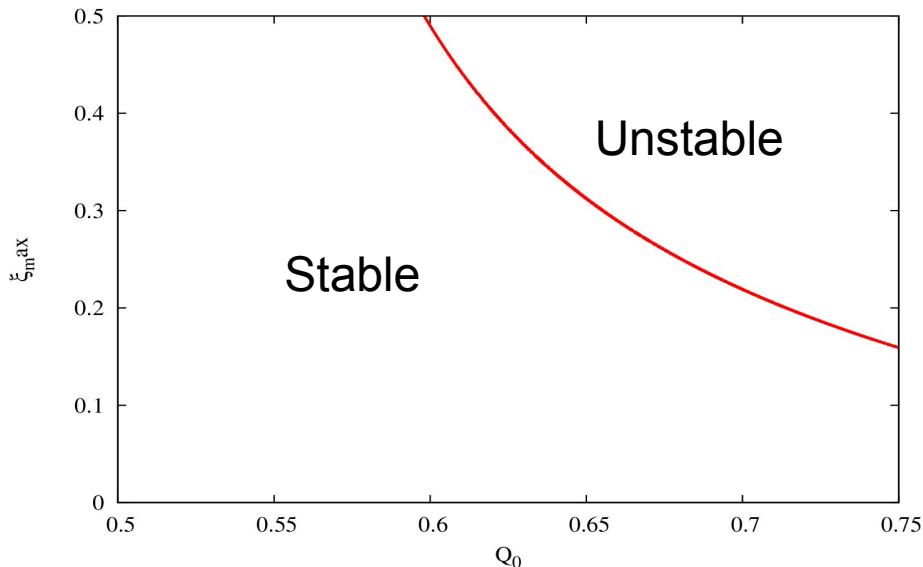
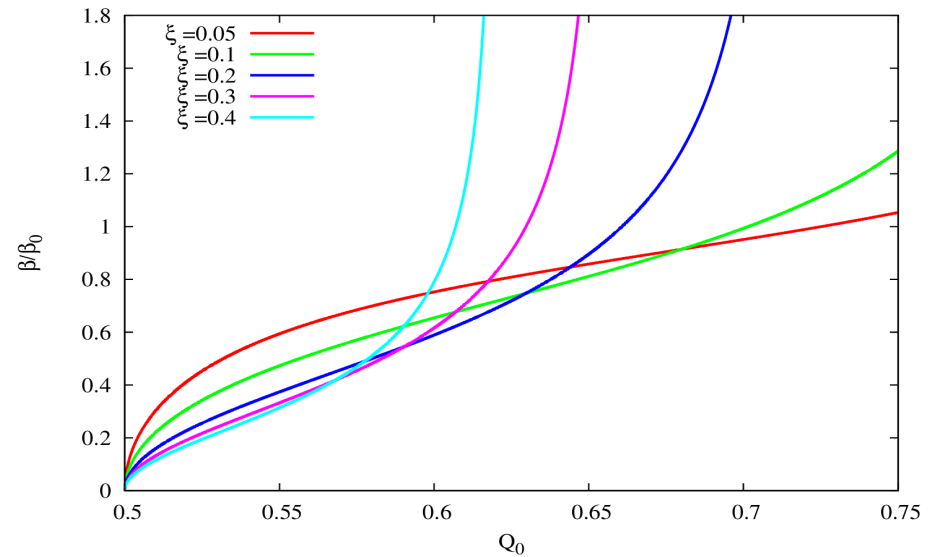
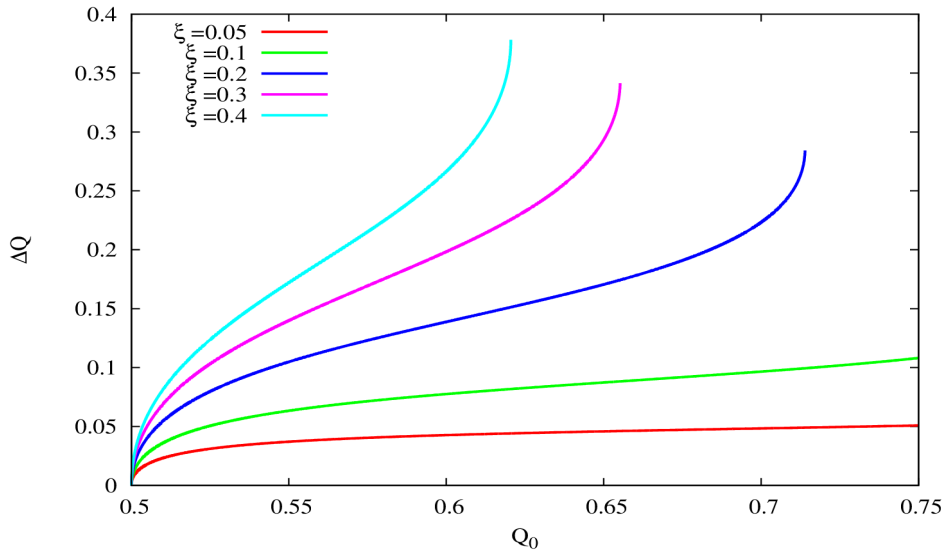


→ Photons energy distribution and average energy in good agreement

→ Some differences ($\sim 2\%$) at low energy – **agreement of 1% or better at higher energies**

Dynamic effects – single IP

- Depending on the working point strong beam-beam effects can distort the optics



→ Linear dynamic effects: see Chao's Handbook, condition for stability ($0 \rightarrow 0.5$):

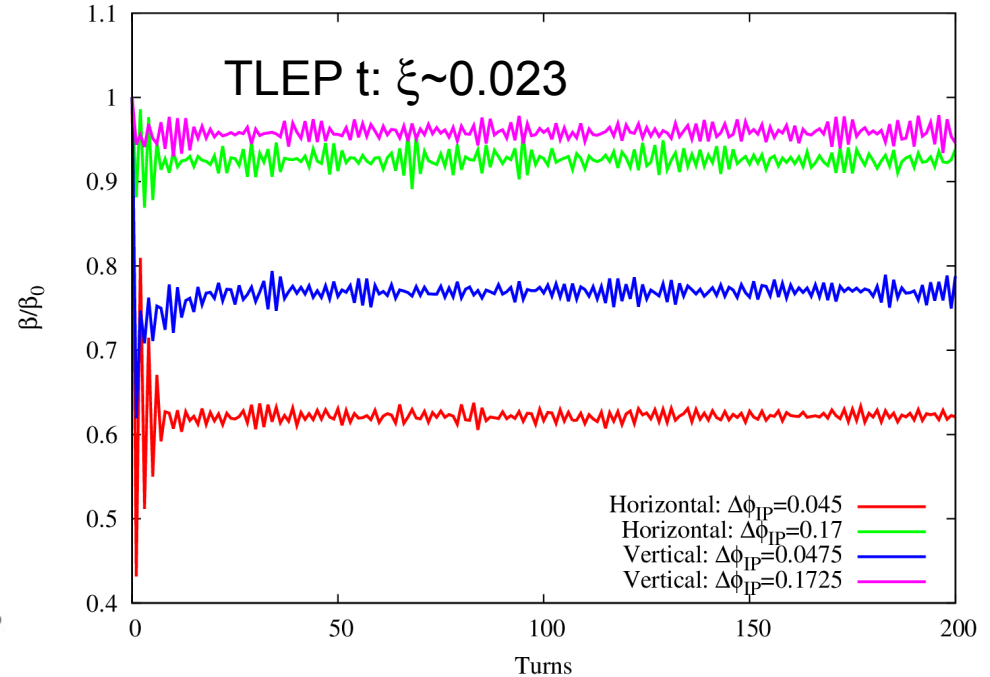
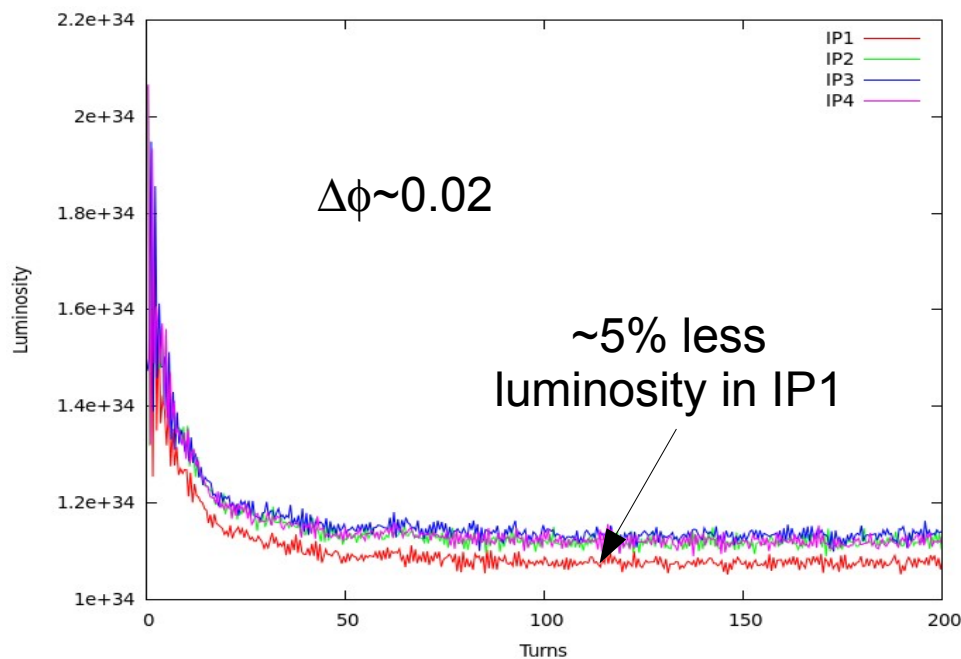
$$\xi_0 < \frac{1}{2\pi} \cot \frac{2\pi Q_0}{2}$$

→ Large beam-beam parameters requires to operate close to 0.5 or 1.0

→ **Very strong dynamic effects**

Multiple IPs

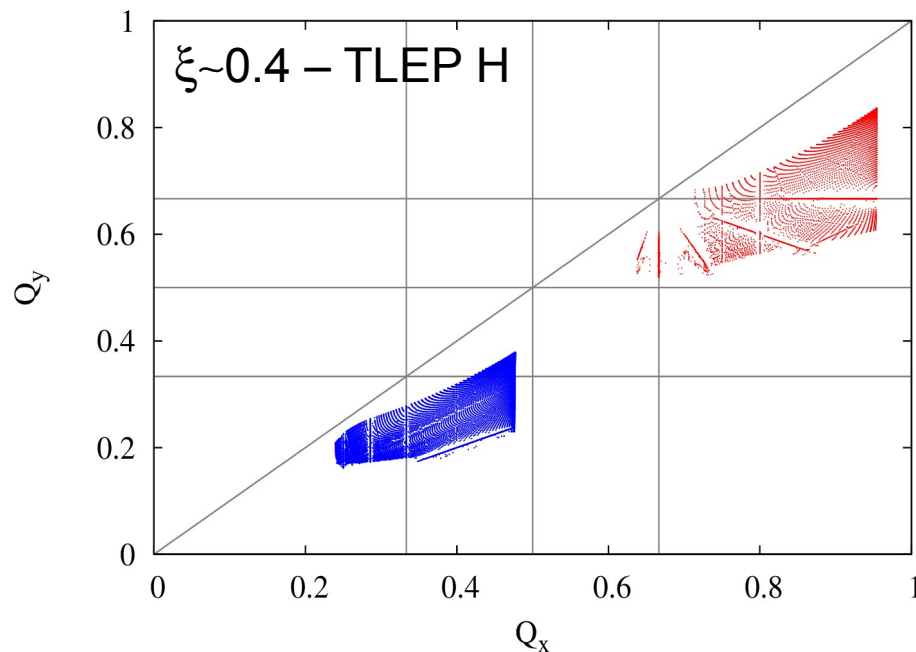
- Past studies have shown that breaking the collision symmetry can lead to excitation of additional resonances and different β -functions, i.e luminosities, at the different IPs
- For a symmetric case it is the phase advance between IPs that determines dynamic effects



- **Keeping collision symmetry is important: excellent optics control**
- Dynamic effects can be violent: in case of problems can be mitigated by a **proper choice of the working point** – here β -functions after the collision in IP1

Choice of the working point

- Assumptions – **final values should be determined by detailed study**:
 - The footprint should not overlap half integer and integer resonances
 - Resonances of order 3 may be crossed (to be checked case by case)
 - Phase advance between IPs: $0.0 < \Delta\phi < 0.25$: integer part multiple of 2
 - Dynamic effects should not be too strong: $\Delta\phi$ not too small
 - Avoid the coupling resonance $Q_x - Q_y = 0.0$
- As a start take 4x LEP as integer part (for TLEPH) – two possibilities for fractional part:



→ Below 0.5 smaller footprint but strong optics distortion

→ Above 0.5 larger footprint but very little optics distortion

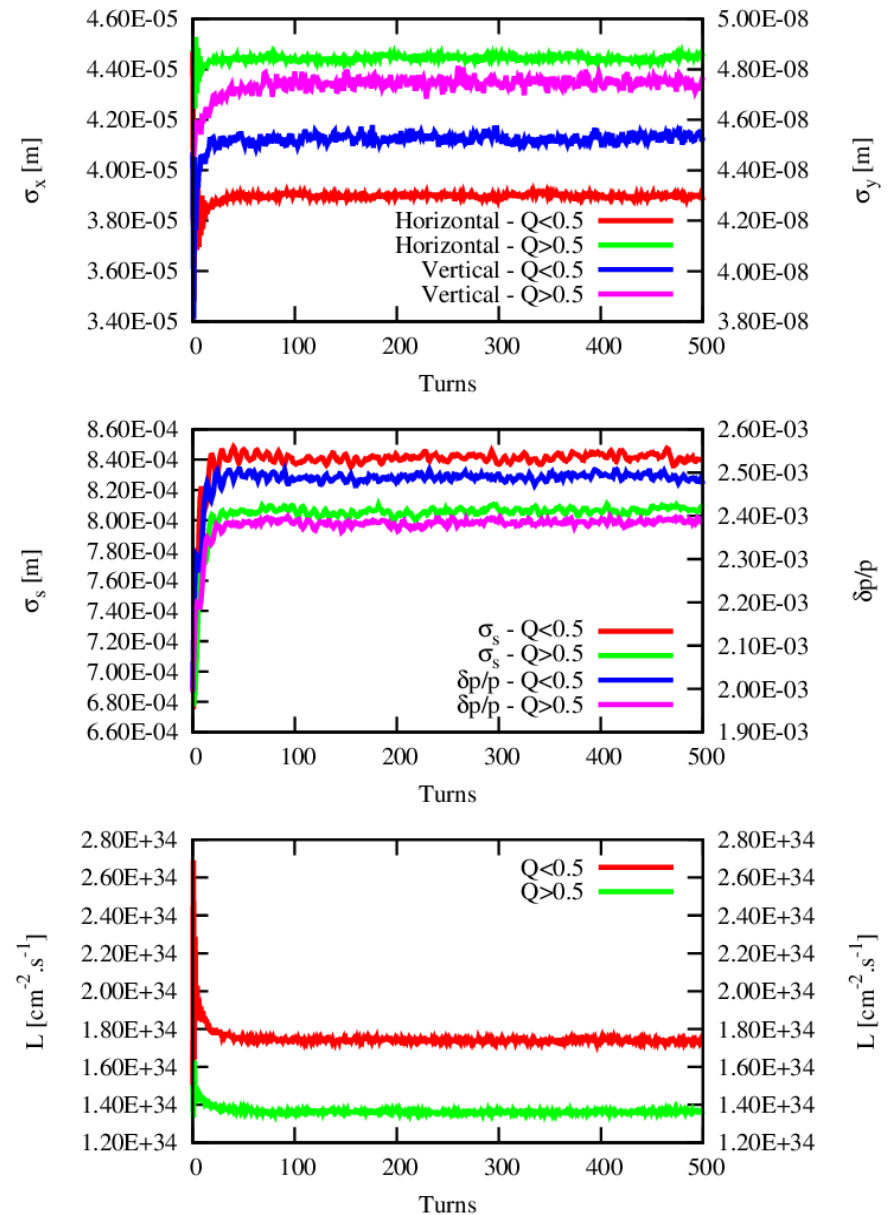
TLEP t

	Design	Simulated Q>0.5	Simulated Q<0.5
E [GeV]	175	175	175
N [e/bunch]	0.88e11	0.88e11	0.88e11
$\epsilon_{x/y}$ [nm]	2.0/0.002	2.1/0.0023	2.25/0.0025
$\beta_{x/y}$ [m]	1.0/1.0e-3	1.0/1.0e-3	1.0/1.0e-3
σ_s [mm] (BS)	0.77	0.8	0.84
$\sigma_{\delta p/p}$ [%] (BS)	0.23	0.24	0.25
$\tau_{ }$ [turns]	23	23	23
$\xi_{x/y}/IP$	0.057/0.057	0.056/0.047	0.051/0.053
L/IP [cm².s⁻¹]	1.3e34	1.38e34	1.75e34

→ The simulated cases correspond to the footprints previously shown

→ **Better than design luminosity achieved in both cases**

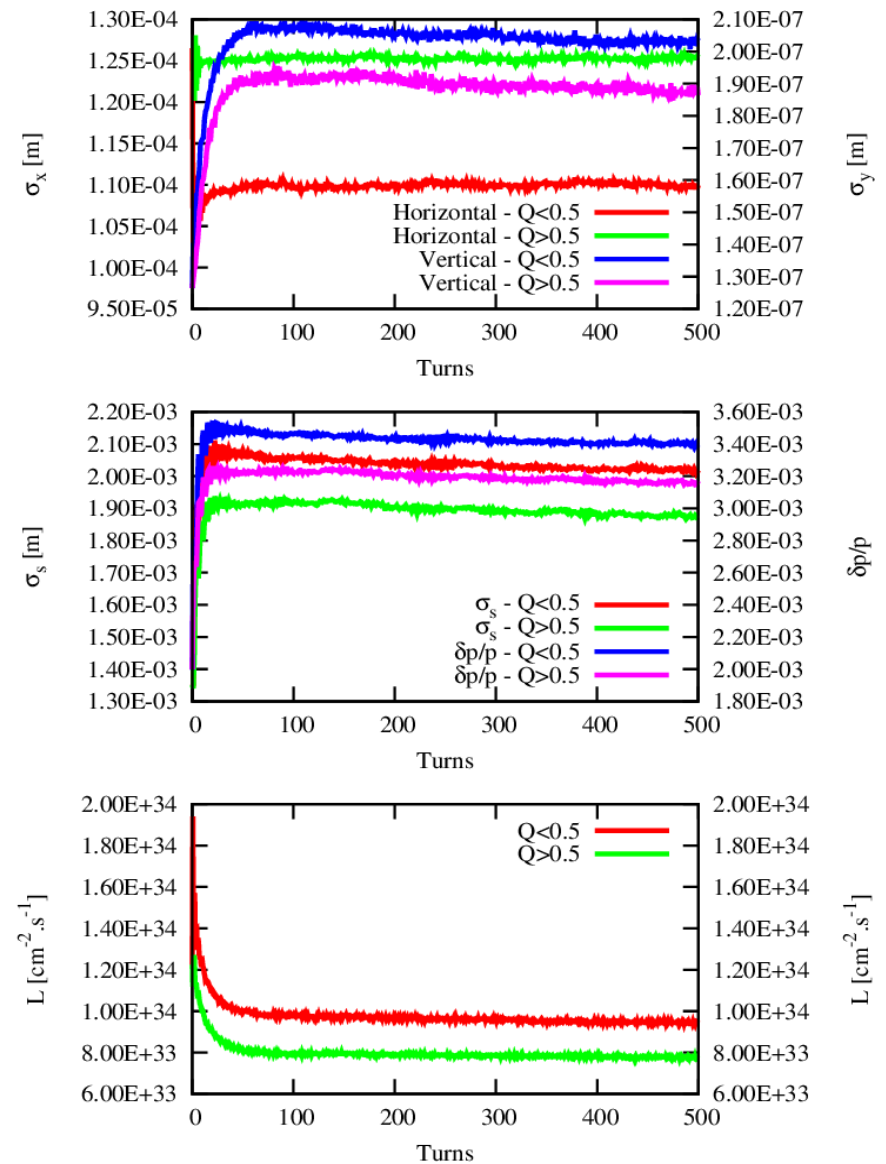
→ Dynamic effects boost the peak luminosity



TLEP t B

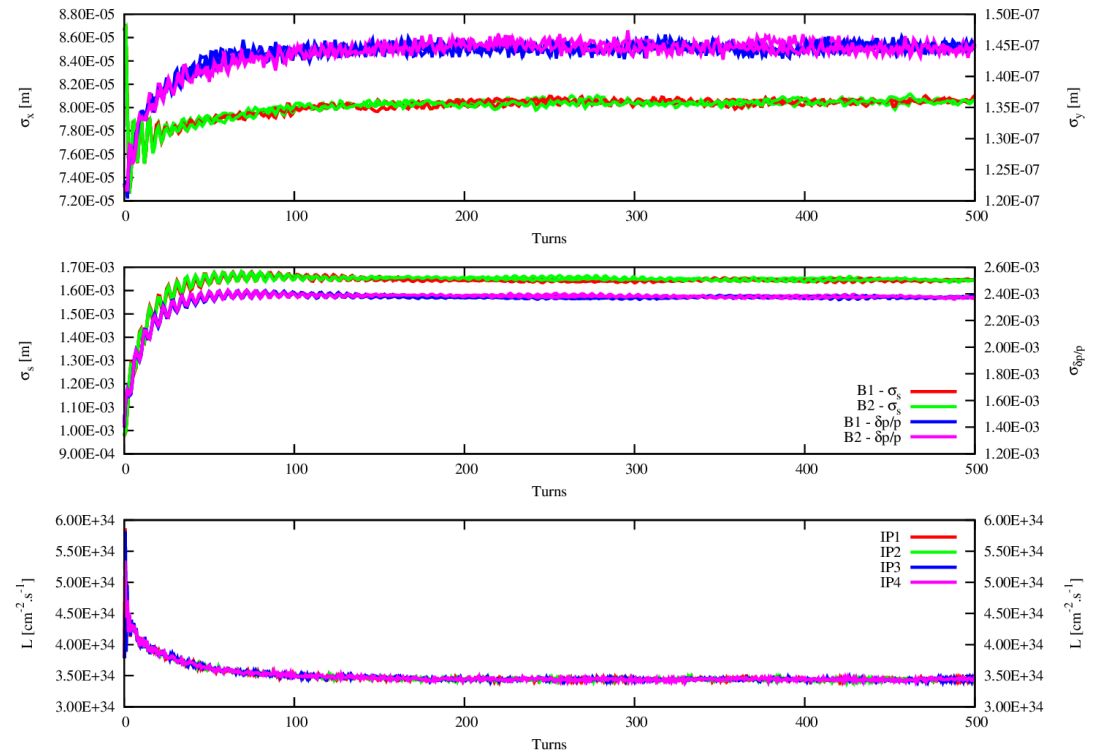
	Design	Simulated Q>0.5	Simulated Q<0.5
E [GeV]	175	175	175
N [e/bunch]	7.0e11	7.0e11	7.0e11
$\epsilon_{x/y}$ [nm]	16.0/0.016	16.7/0.035	18.0/0.035
$\beta_{x/y}$ [m]	1.0/1.0e-3	0.96/1.03e-3	0.67/1.2e-3
σ_s [mm] (BS)	1.95	1.9	2.05
$\sigma_{\delta p/p}$ [%] (BS)	0.29	0.32	0.34
$\tau_{ }$ [turns]	23	23	23
$\xi_{x/y}/IP$	0.057/0.057	0.055/0.04	0.051/0.050
L/IP [cm ² .s ⁻¹]	1.04e34	0.8e34	1.0e34

- Same tunes as for TLEP t simulations
- **Design luminosity barely reached for Q<0.5**
- This scenario looks less promising than TLEP t: vertical blow-up much stronger



TLEP H

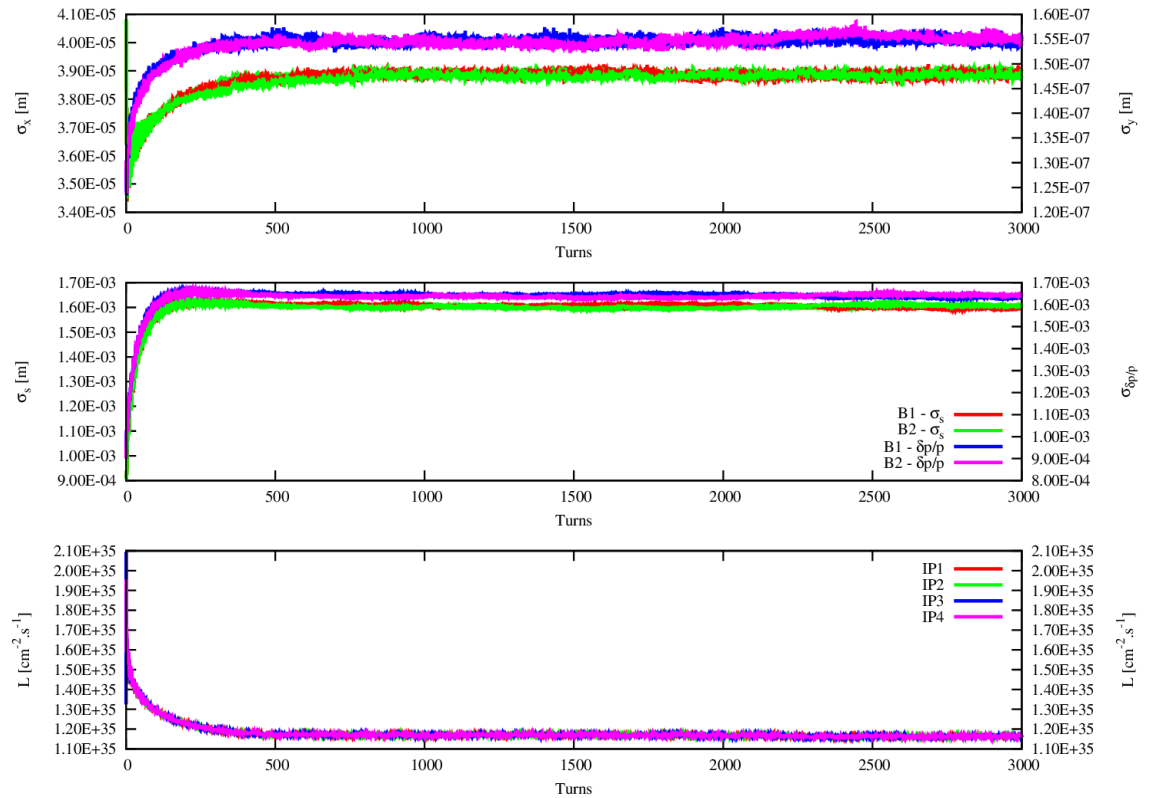
	Design	Simulated $Q < 0.5$
E [GeV]	120	120
N [e/bunch]	3.7e11	3.0e11
$\epsilon_{x/y}$ [nm]	7.5/0.015	8.8/0.019
$\beta_{x/y}$ [m]	0.5/1.0e-3	1.0/1.0e-3
σ_s [mm] (BS)	2.11	1.65
$\sigma_{\delta p/p}$ [%] (BS)	0.3	0.24
$\tau_{ }$ [turns]	72	72
$\xi_{x/y}/IP$	0.094/0.094	0.065/0.054
L/IP [cm ² .s ⁻¹]	5.08e34	3.5e34



- **TLEPH much more difficult:** try to achieve twice the beam-beam parameter with 3 times damping time: emittance blow-up of a factor 5 in few turns with nominal parameters
- Reduced beam-beam parameter: relaxed β_x and bunch intensity, keep current constant
- Making use of dynamic effects **~70% of design achieved:** could be pushed further but then lifetime becomes an issue

TLEP W

	Design	Simulated Q<0.5
E [GeV]	80	80
N [e/bunch]	1.0e11	1.0e11
$\epsilon_{x/y}$ [nm]	3.3/0.017	4.1/0.022
$\beta_{x/y}$ [m]	0.2/1.0e-3	0.5/1.0e-3
σ_s [mm] (BS)	1.98	1.6
$\sigma_{\delta p/p}$ [%] (BS)	0.2	0.165
$\tau_{ }$ [turns]	242	242
$\xi_{x/y}/IP$	0.086/0.086	0.068/0.051
L/IP [cm².s⁻¹]	1.64e35	1.15e35



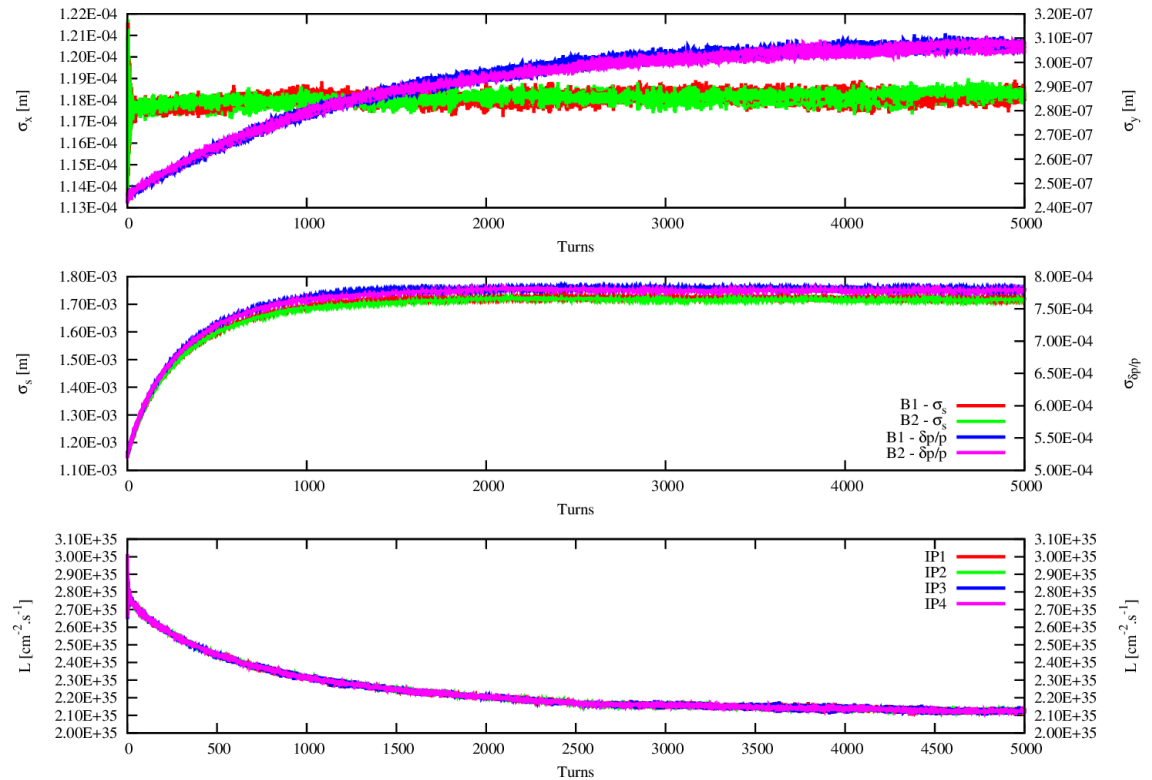
→ Again with design parameters strong emittance blow-up: reduced beam-beam parameter, relaxed β_x

→ Making use of dynamic effects **~70% of design achieved**

→ Q>0.5: luminosity well below design

TLEP Z

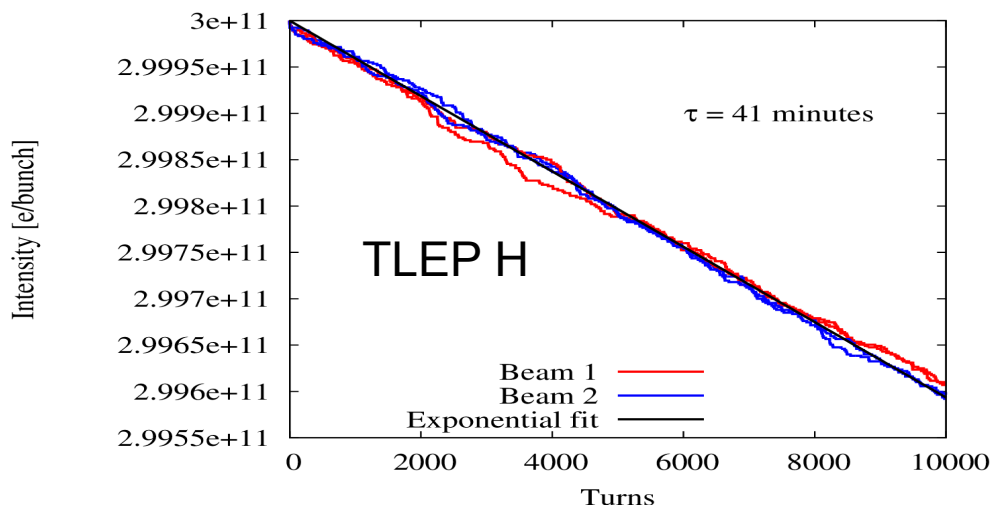
	Design	Simulated Q>0.5
E [GeV]	45	45
N [e/bunch]	4.0e11	1.5e11
$\epsilon_{x/y}$ [nm]	29.2/0.06	29.7/0.09
$\beta_{x/y}$ [m]	0.5/1.0e-3	0.5/1.0e-3
σ_s [mm] (BS)	2.93	1.7
$\sigma_{\delta p/p}$ [%] (BS)	0.13	0.075
$\tau_{ }$ [turns]	1319	1319
$\xi_{x/y}/IP$	0.068/0.068	0.029/0.024
L/IP [cm ² .s ⁻¹]	5.86e35	2.1e35



- LEP at 45.6 GeV achieved $\xi_{x/y} \sim 0.03/0.044$ per IP with a damping time of ~ 360 turns
- Can TLEP do better with $\sim 4x$ longer damping time?
- Had to significantly reduce the beam-beam parameter to achieve reasonable emittance blow-up
- In this case **$\sim 36\%$ of design reached but requires 20000 bunches (probably not realistic – electron cloud?)**

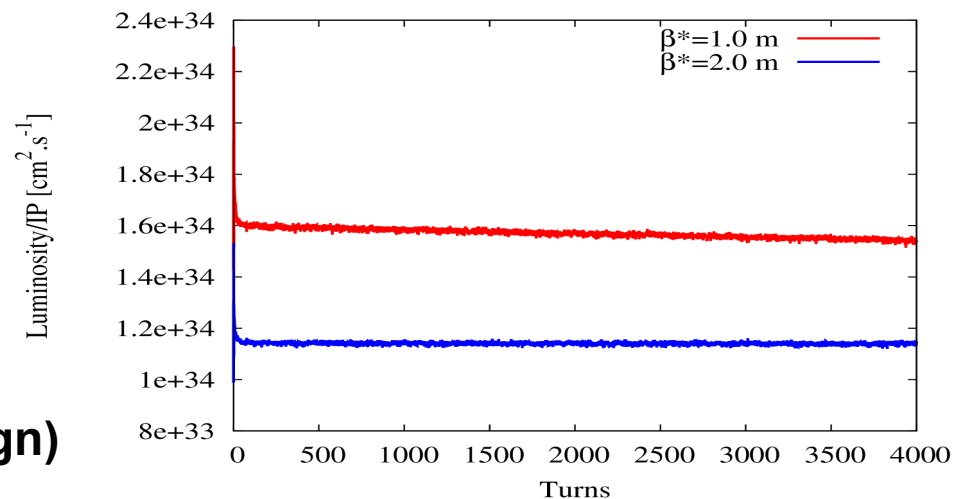
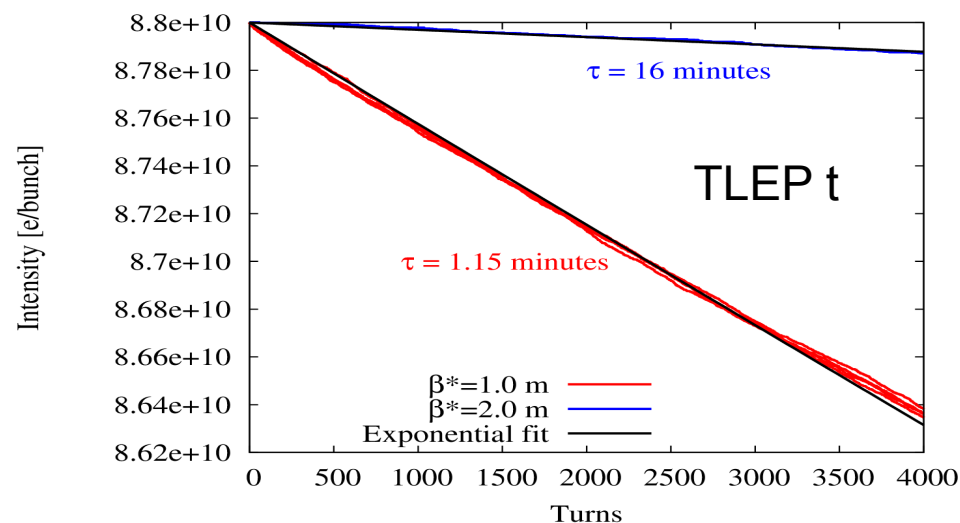
Lifetimes

- Cases with strong beamstrahlung: TLEP t and TLEP H. Does not include burn off. Collimation performed at each IP with 14σ transverse and $\eta=0.02$



→ **TLEP H:** bunch intensity and β_x were already relaxed to reduce emittance blow-up. **simulated lifetime 41 minutes**

→ **TLEP t:** with nominal parameters lifetime barely above 1 minute. Relaxing β_x to 2.0m allows to increase it to **15 minutes with a luminosity reduce to 1.17e34 (~90% of design)**

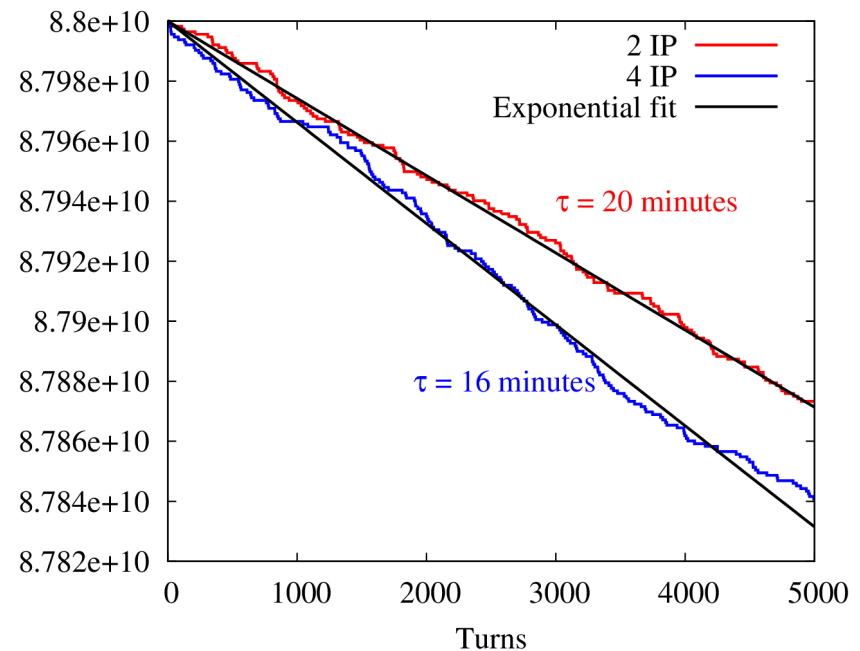
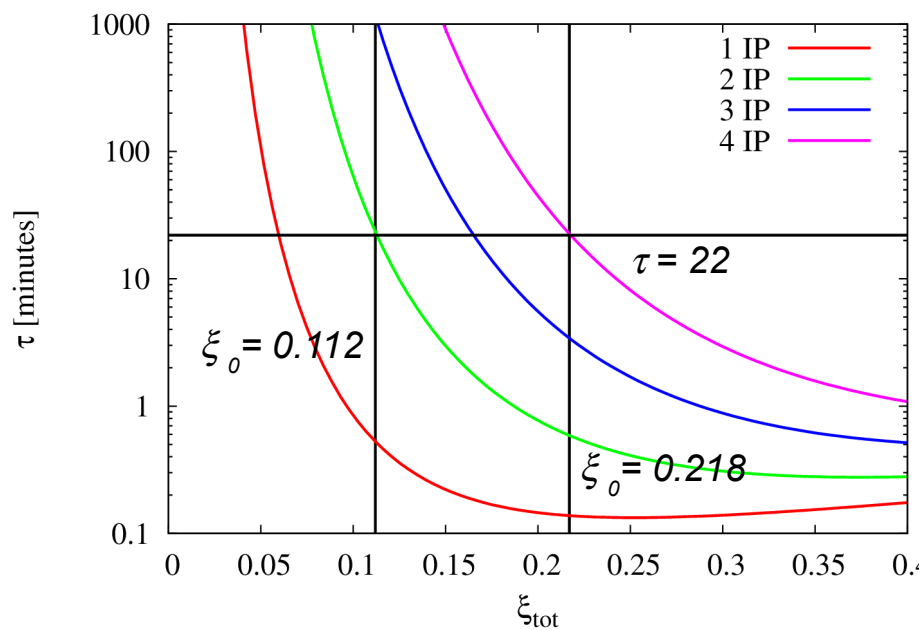


Performance with 2IPs

- Can we scale the beam-beam parameter, i.e luminosity per IP, with the number of IPs?
Analytical lifetime estimate by Telnov:

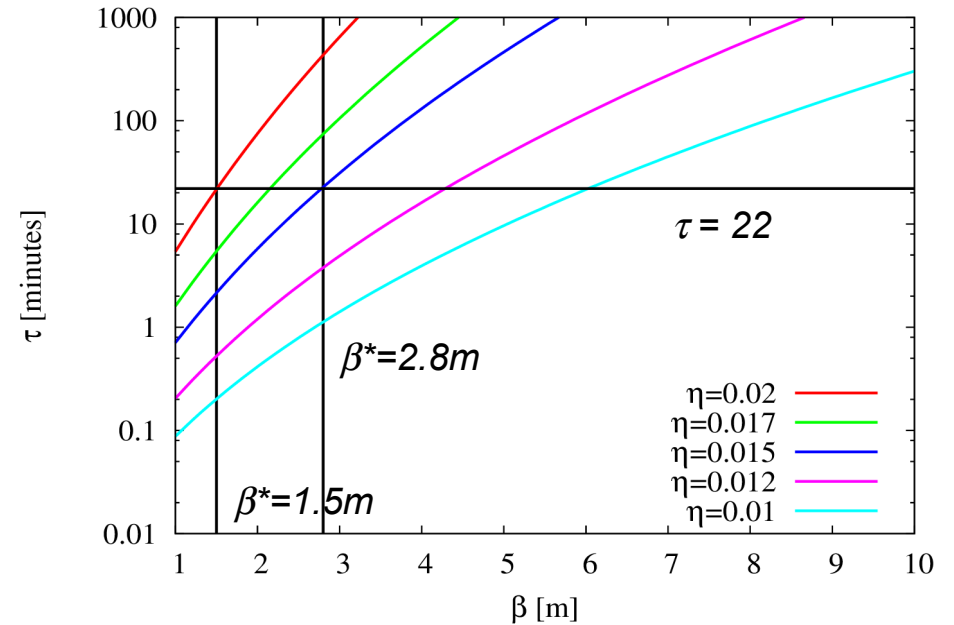
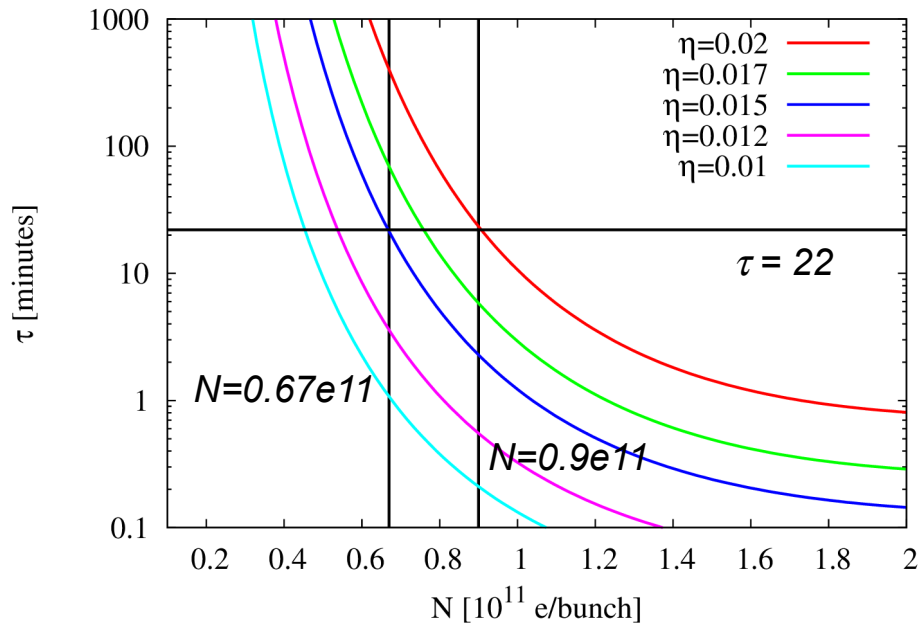
$$\tau = \frac{20}{n_{IP}} \frac{\sqrt{6} \pi r_e \gamma}{\alpha^2 \eta \sigma_s} \frac{2 \pi R}{c} u^{3/2} e^u \quad \text{with} \quad u = \eta \frac{\alpha \sigma_x \sigma_s}{3 \gamma r_e^2 N_p}$$

→ Much stronger dependency on the bunch intensity than number of IPs. The equilibrium bunch length and β also vary with ξ . Example of TLEP t ($\Delta\sigma_s$ included assume no dynamic β):



→ **Reducing the number of IPs will degrade the total luminosity performance**

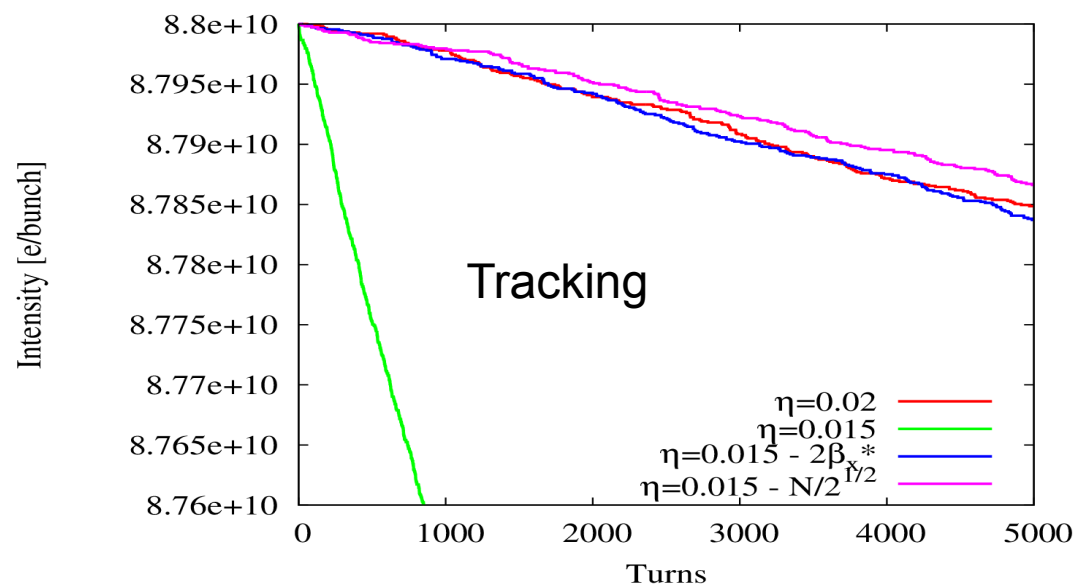
Reduced energy acceptance



→ Reducing the energy acceptance will significantly degrade the beam lifetime

→ **Lifetime can be recovered by either increasing β_x or decreasing the bunch intensity**

→ For luminosity performance increasing β_x is clearly better



Summary

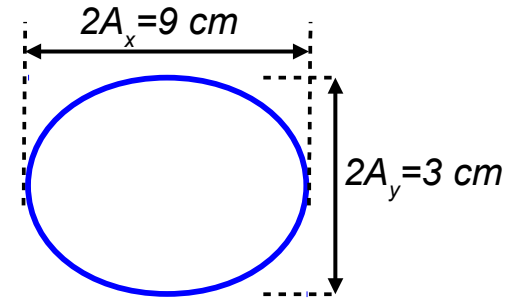
	TELP Z	TLEP W	TLEP H	TLEP t
E [GeV]	45	80	120	175
N [e/bunch]	1.5e11	1.0e11	3.0e11	0.88e11
# bunches	20000	3200	206	160
$\epsilon_{x/y}$ [nm]	29.7/0.09	4.1/0.022	8.8/0.019	2.17/0.00217
$\beta_{x/y}$ [m] init.	0.5/1.0e-3	0.5/1.0e-3	1.0/1.0e-3	2.0/1.0e-3
$\beta_{x/y}$ [m] dist.	0.47/1.0e-3	0.36/1.1e-3	0.75/1.1e-3	1.6/0.85e-3
$\xi_{x/y}$ /IP dist.	0.029/0.024	0.068/0.051	0.065/0.054	0.055/0.036
L/IP [cm ² .s ⁻¹]	2.1e35	1.15e35	3.5e34	1.17e34
L/IP [% design]	36	70	69	89
τ_{BS} [min], $\eta=0.02$	-	-	~41	~15

- **Good performance achieved in all cases except TLEP Z** (only 36% of design)
- **Dynamic effects are essential** to reach these performance: choice of working point
- **TLEP H and TLEP t are limited by lifetime not beam-beam**
- Reducing the number of IPs would result in significant loss in total luminosity
- **Reducing the energy acceptance is possible if β_x is increased**: lower luminosity
- In case of strong hourglass (TLEP Z) we could try traveling focus to recover some luminosity

Impedance model: resistive wall

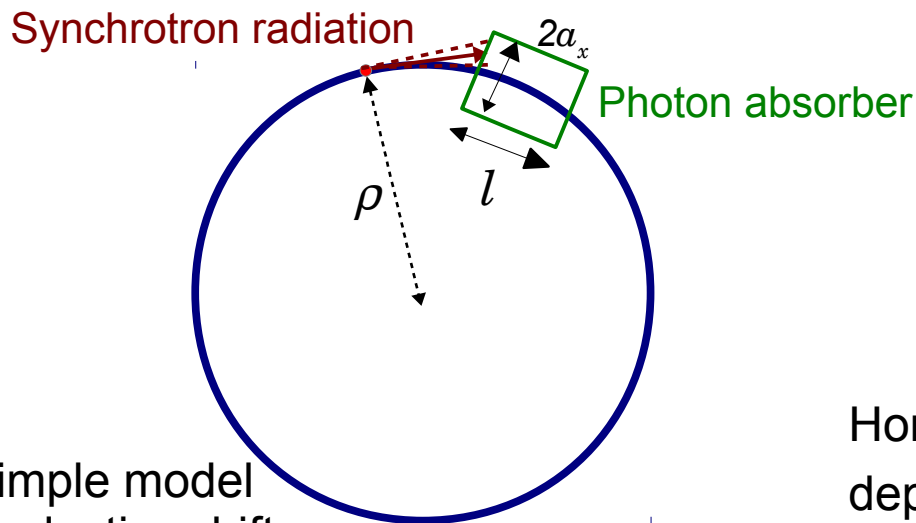
- Resistive-wall impedances considered:

- Aluminum ($\sigma = 3.7 \cdot 10^7 \text{ S/m}$) main vacuum pipe



Elliptic shape

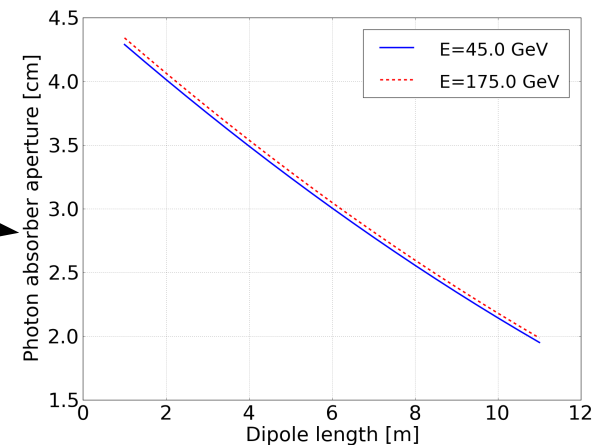
- Tungsten ($\sigma = 1.85 \cdot 10^7 \text{ S/m}$) photon absorbers, intercepting synchrotron radiation from dipoles (elliptic with reduced aperture in horizontal). Transitions between vacuum pipe and photon absorbers cross-sections: imp. approximated using K. Yokoya's formula for round tapers [CERN SL/90-88]:



Simple model neglecting drift and straight sections

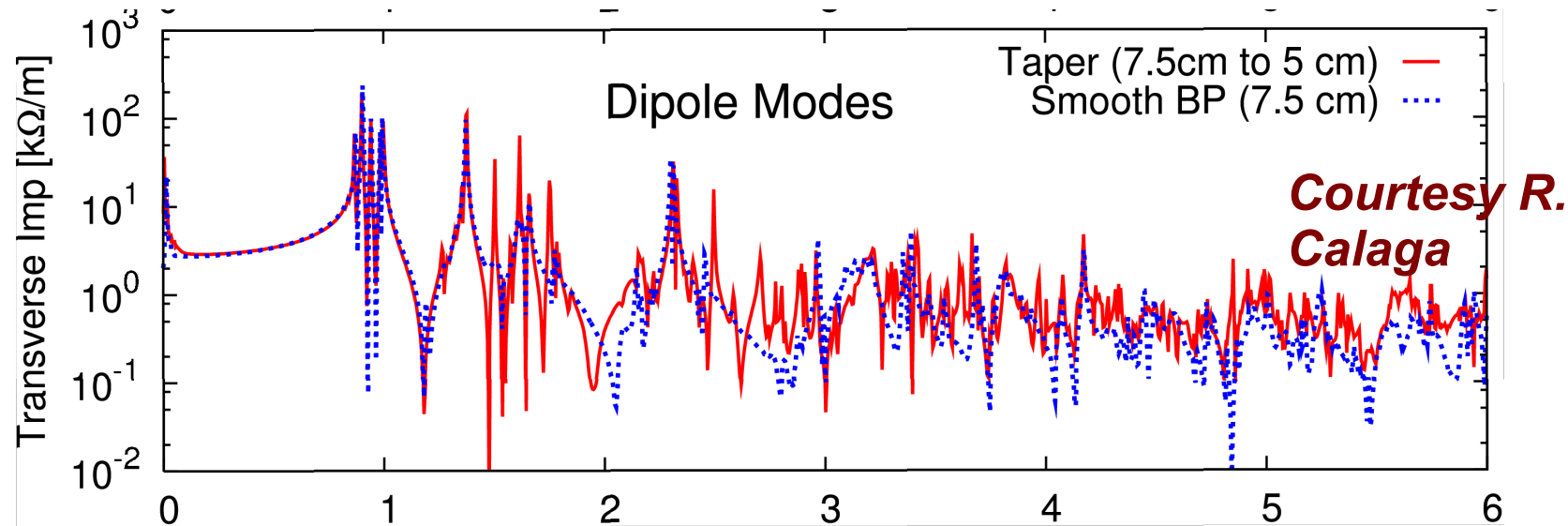
$$a_x = \rho \left(\frac{\cos\left(\frac{1}{\gamma}\right)}{\cos\left[\arccos\left(\frac{\rho \cos\left(\frac{1}{\gamma}\right)}{\rho + A_x}\right) + \frac{l - l_d}{\rho} - \frac{2}{\gamma}\right]} - 1 \right)$$

Hor. aperture a_x depends on dipole length l_d



Impedance model: RF cavities

- Broad-band resonator estimates ($Q=1$, $f_r=6$ GHz):
 - RF cavities: from R. Calaga's PhD (BNL-SERL cavity, 700 MHz)



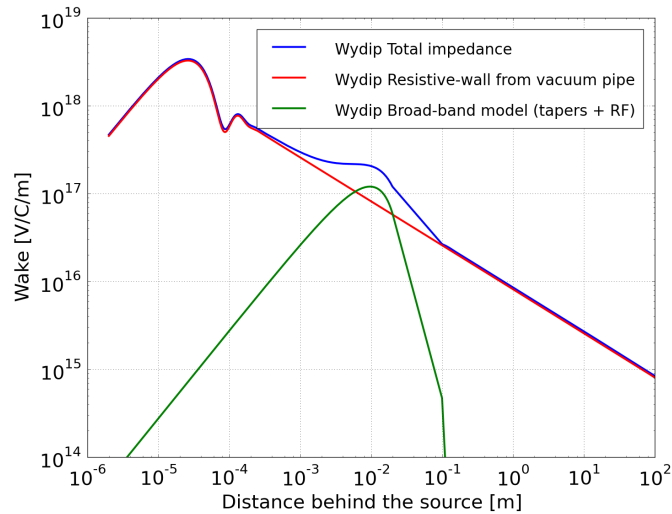
→ "fit" low-frequency part by constant inductive impedance of 3 kΩ/m, multiplied by 600 (number of cavities).

- **All impedances weighted by approximative average beta function $\sim 2R/Q$:**
 - Assume 4xLEP for TLEPH (312,392)
 - The tunes range from (104,130) for TLEPZ to (624,784) for TLEPt due to the different FODO cells length: **strong impact on the effect of impedance**

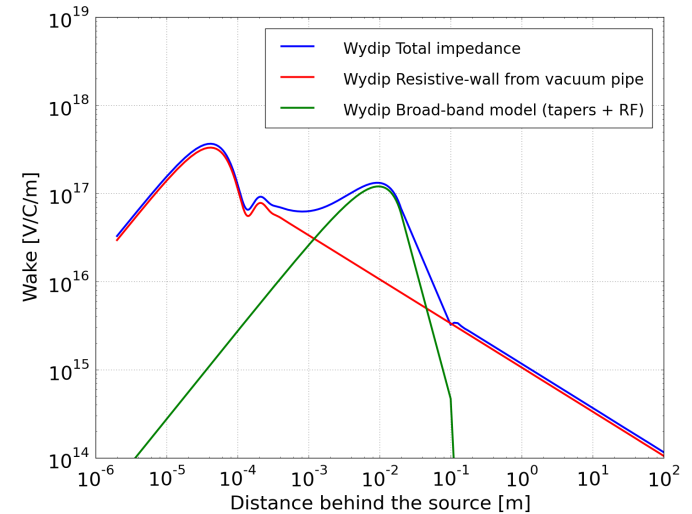
Dipolar wake function

- Transverse dipolar wake function in vertical: relative contributions for different vertical apertures of the vacuum pipe and different dipole lengths

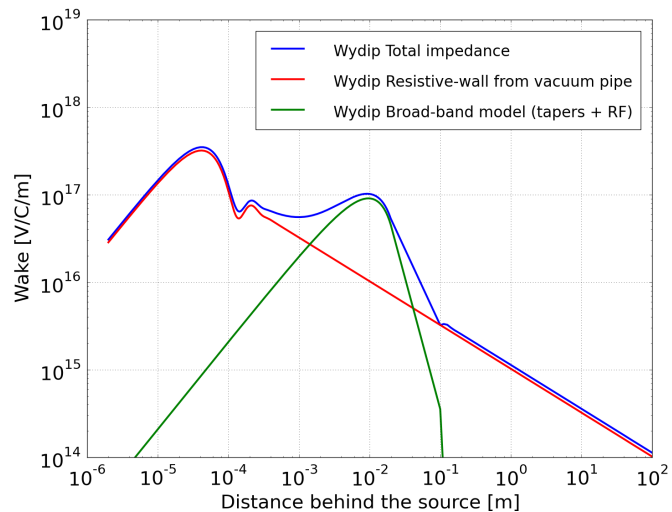
$A_y = 1.5$
cm,
 $l_d = 11$ m



$A_y = 3$ cm,
 $l_d = 11$ m



$A_y = 3$ cm,
 $l_d = 5.5$ m



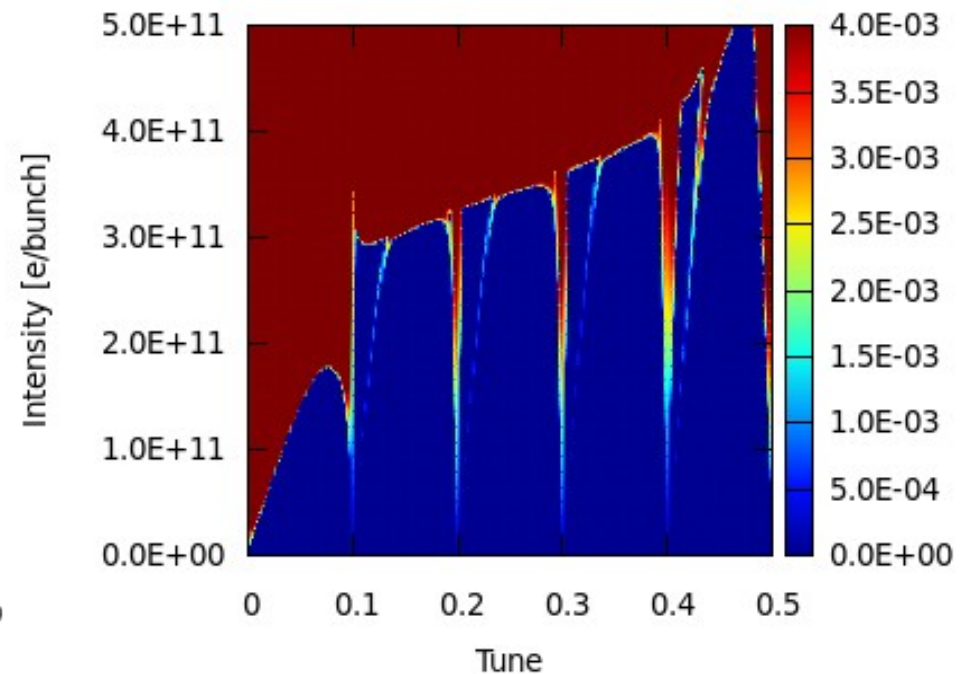
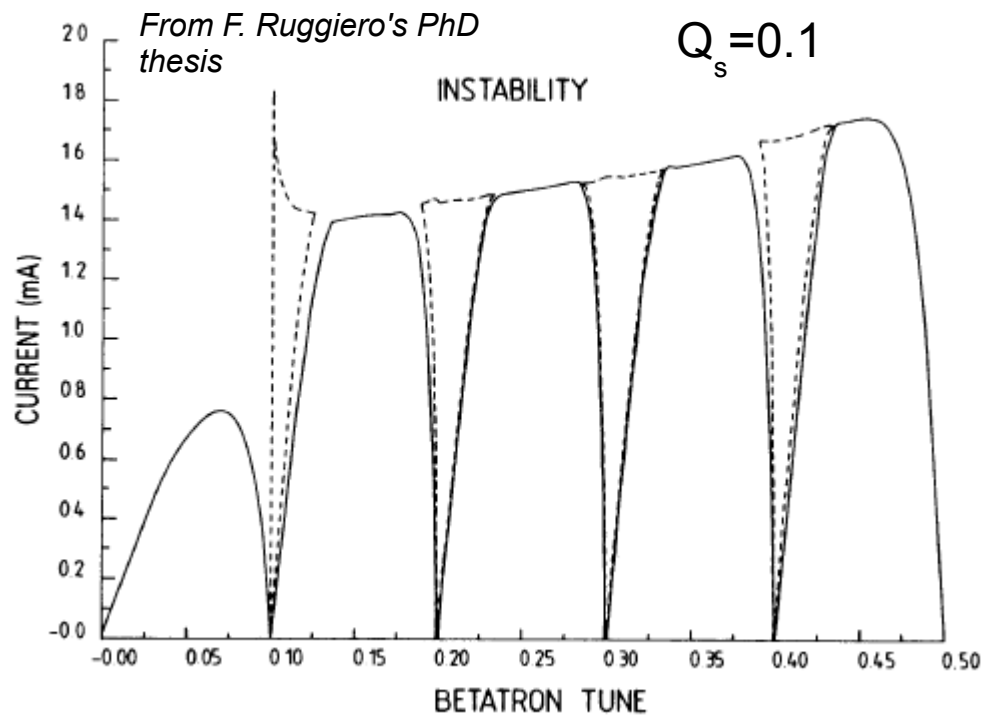
→ Increase of vertical aperture has a clear beneficial effect

→ Length of dipoles has a smaller impact

→ TMCI will be estimated for these three cases

Synchro-betatron effects

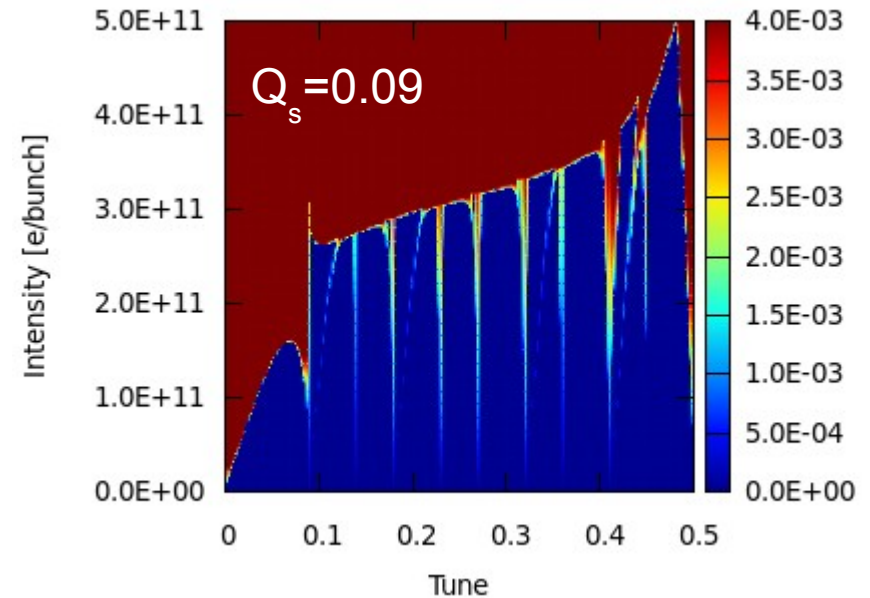
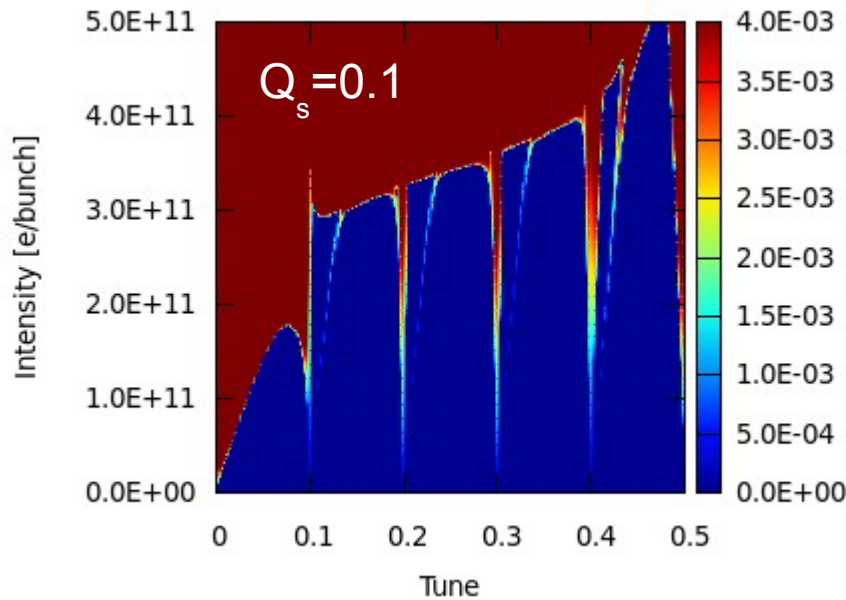
- It was shown by F. Ruggiero that in case of large synchrotron tune one has to consider synchro-betatron resonances ($Q_\beta = Q_s$) and eventual coupling between modes 0 or -1 and reflexions of the synchrotron sidebands:



- **Good qualitative agreement**, difference may be explained by impedance model
- Consider enough sidebands to cover a full integer

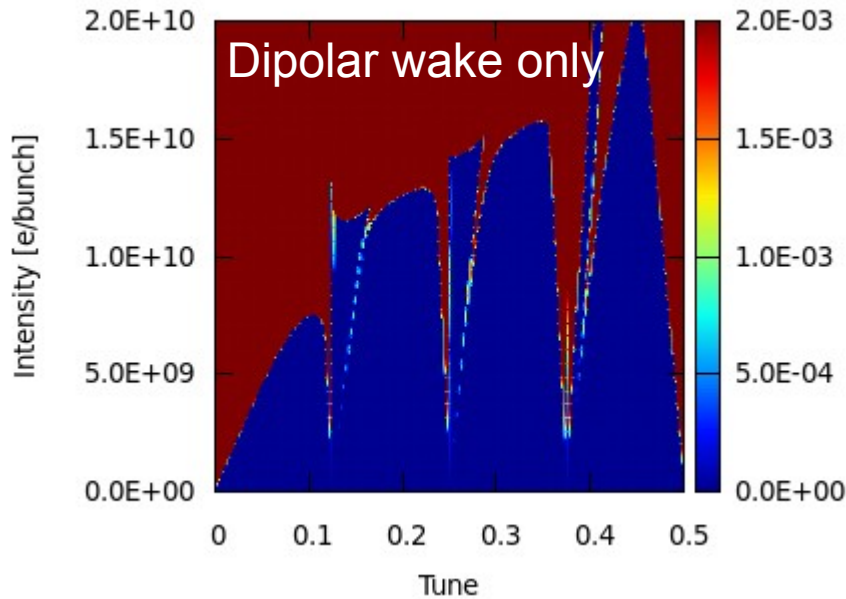
Synchrotron tune

- When $0.5/Q_s = \text{integer}$ all the reflexions line up and we get the clean picture shown in the previous slide. Example for $Q_s \sim 0.1$:

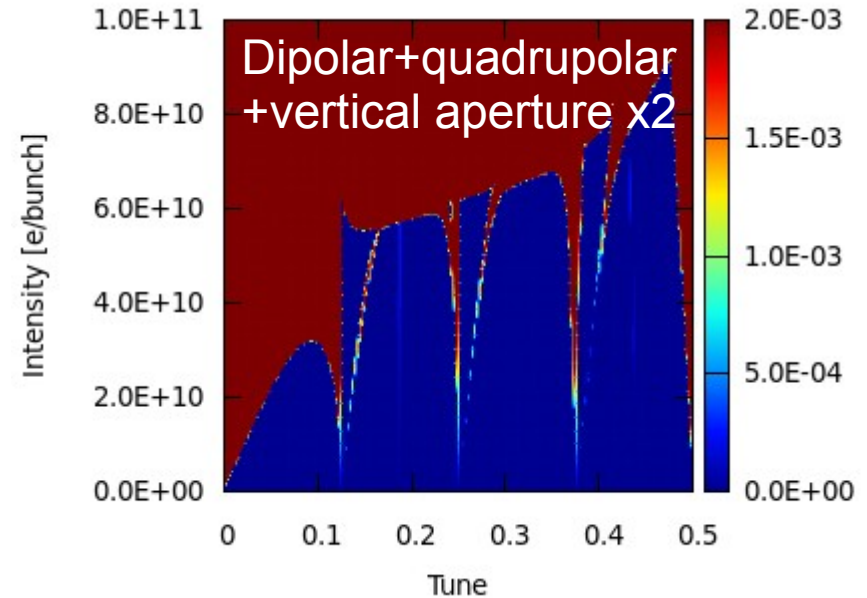
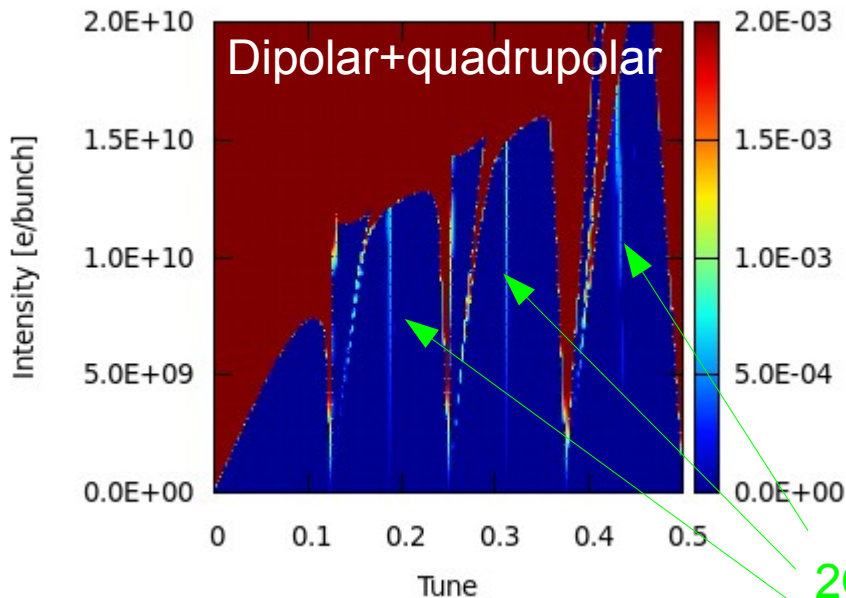


- The TMCI threshold approximately scales with the synchrotron tune
- When $0.5/Q_s$ is not an integer more synchro-betatron resonances are observed
- It is not clear up to which order these resonances are harmful: would require dedicated study, for now stick with design values

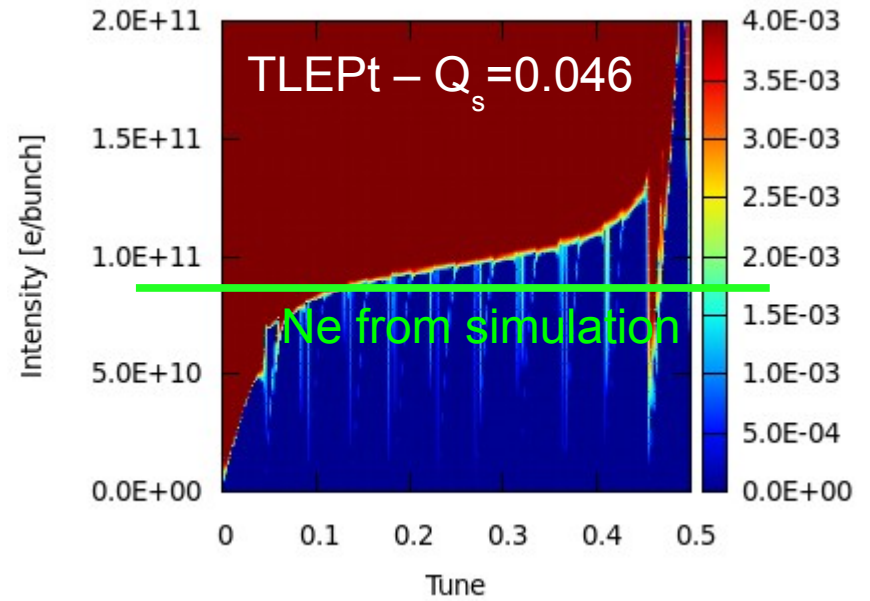
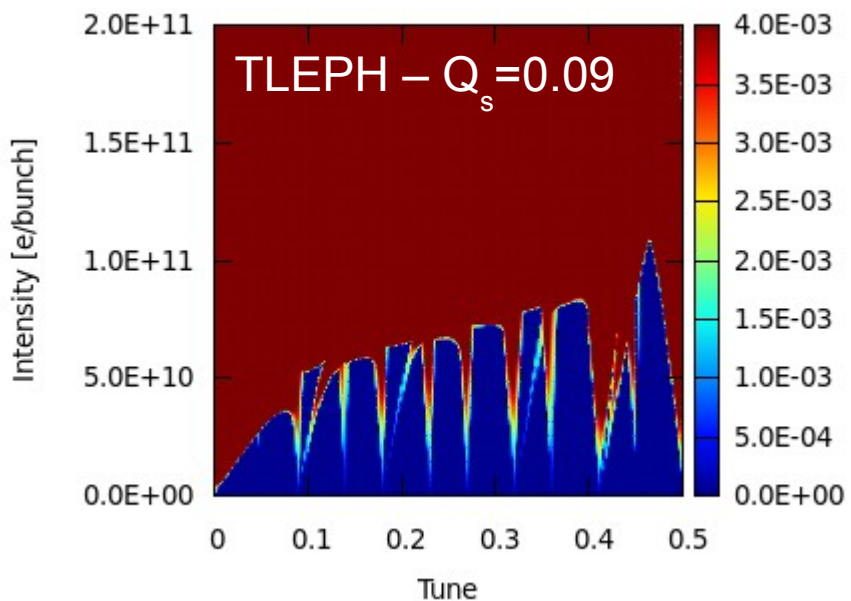
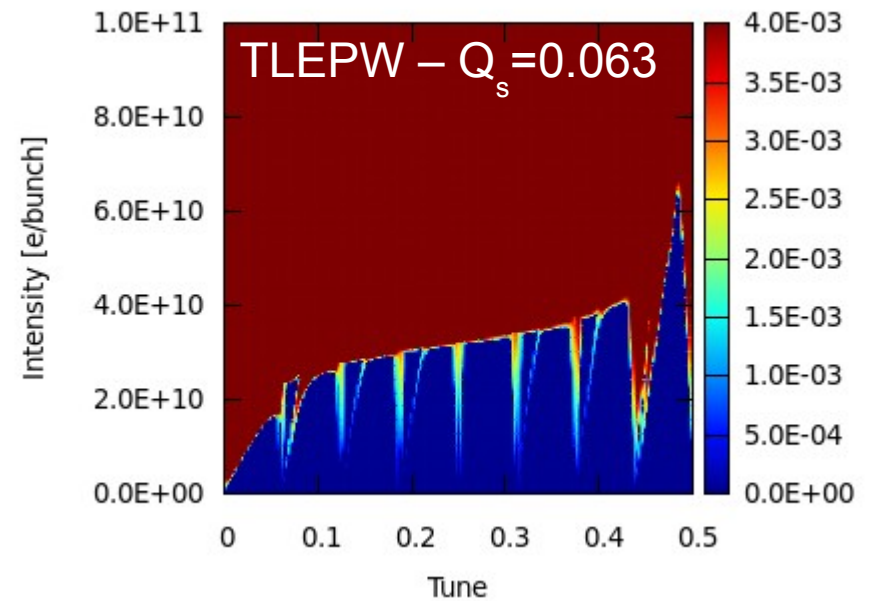
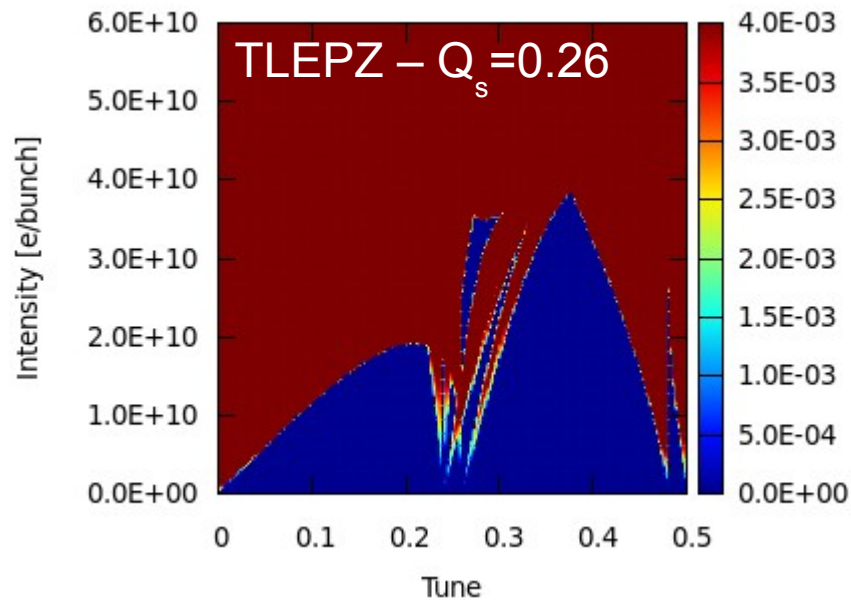
Quadrupolar wake



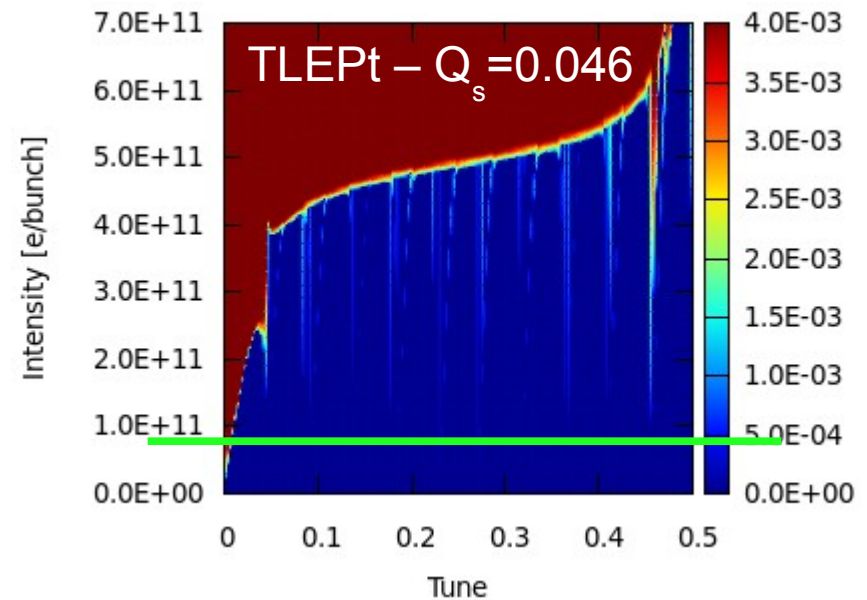
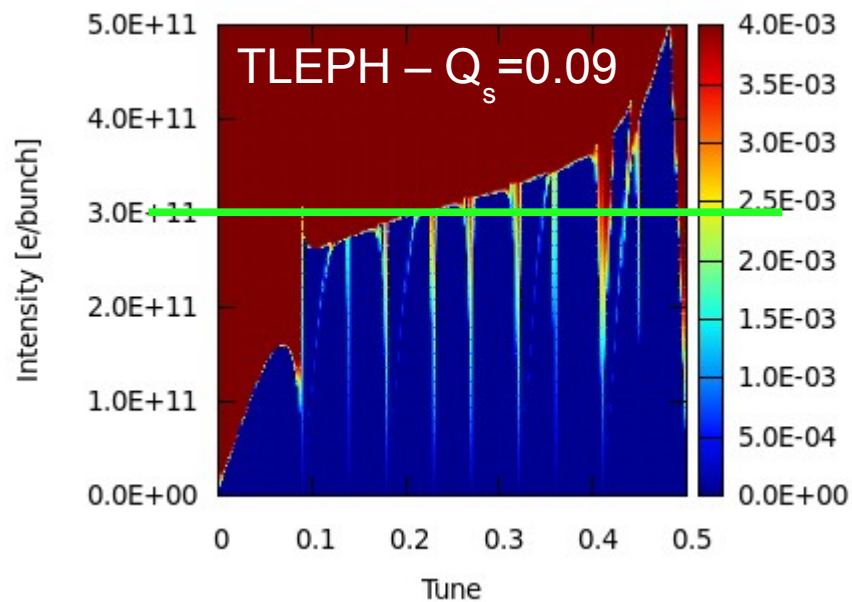
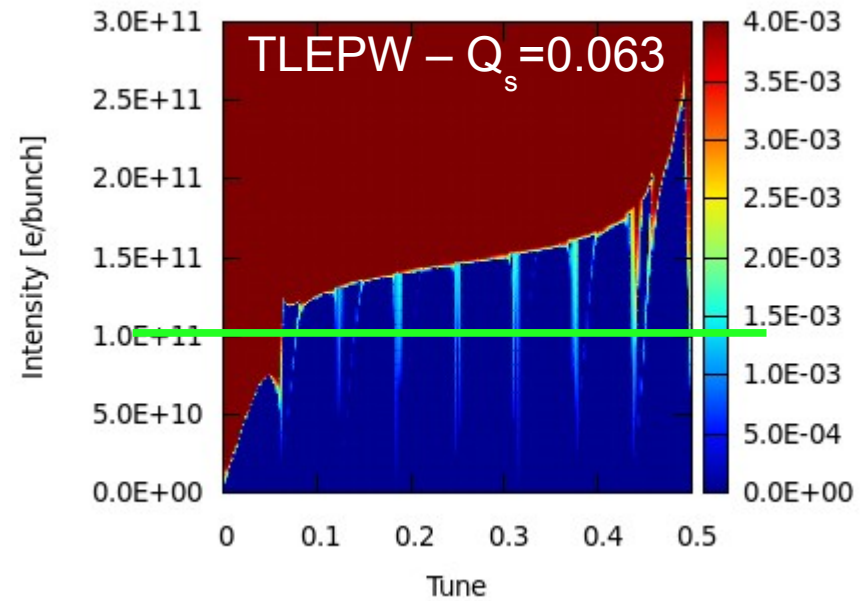
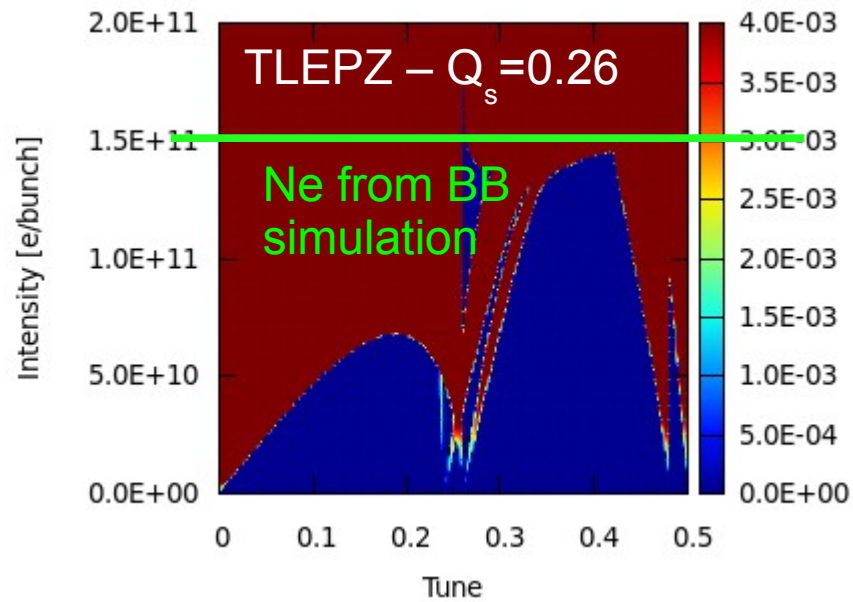
- Use an exemple with $Q_s = 0.125$
- Dipolar wake only: resonance of type $Q_y = nQ_s$ observed
- Dipolar and quadrupolar wake: additional resonances of type $2Q_y = nQ_s$
- Increasing the vertical beam pipe by a factor 2 strongly mitigates the quadrupolar resonances
- **No apparent effect from quadrupolar wake on TMCI threshold**



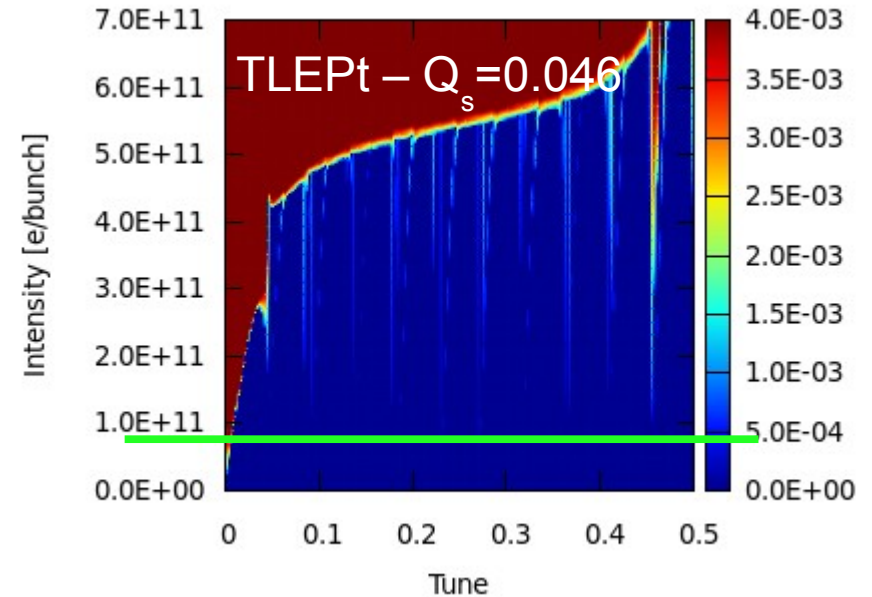
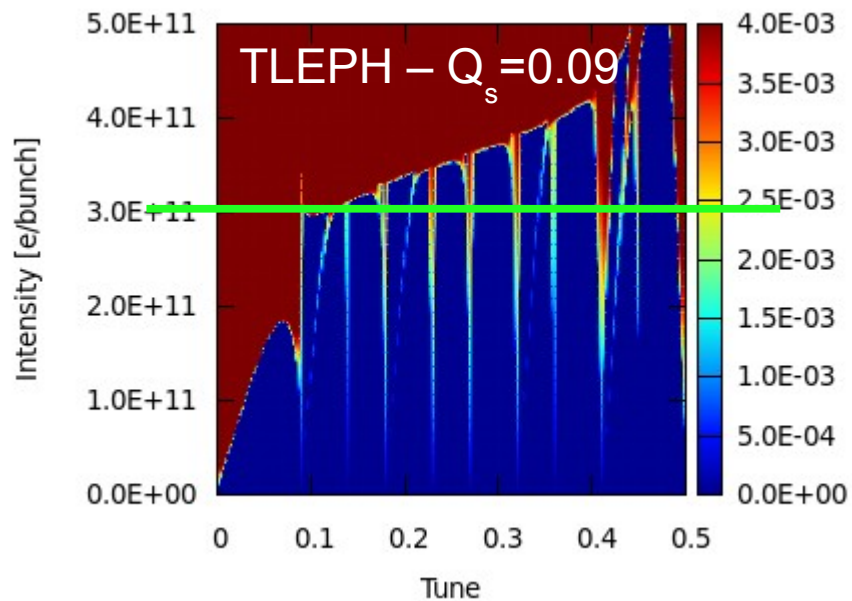
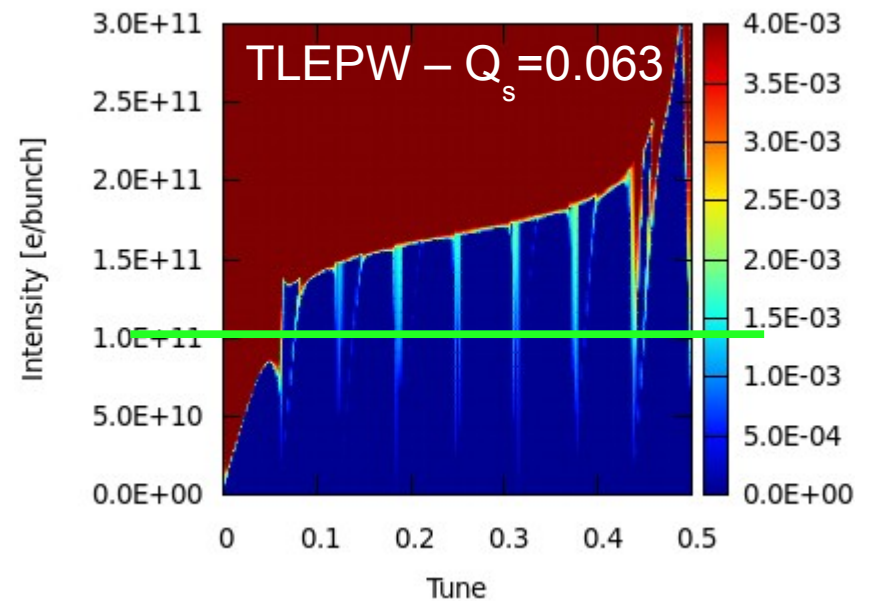
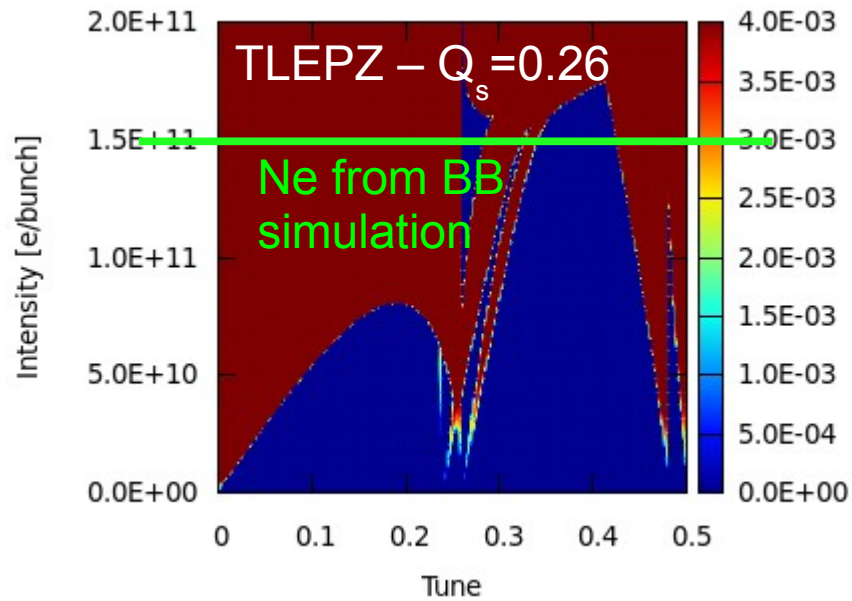
1.5cm half vertical aperture – 11m dipoles



3.0cm half vertical aperture – 11m dipoles



3.0cm half vertical aperture – 5.5m dipoles



Summary

- The impedance model was extended to include photo-absorber: impedance however dominated by the resistive wall from the beam pipe. **The results presented here are preliminary, more detailed studies (tracking) are required for validation.**
- Preliminary observations:
 - **Synchro-betatron resonances are important** and should be avoided: additional constraint on the choice of the working point
 - With a half vertical aperture of **1.5cm only TLEPt is below threshold**
 - Increasing the half vertical aperture to 3.0cm significantly improves the situation, reducing the dipole length and retracting the photo-absorbers further increases the threshold
 - With all this taken into account only **TELPZ remains as an issue** (already an issue for beam-beam, what about electron cloud? → review parameters?)
 - In the previous calculations the bunch lengthening from beamstrahlung was not taken into account: this will slightly increase the thresholds
- Next steps:
 - Validation of the results with tracking
 - Possible benefits of beam-beam tune spread, chromaticity, transverse damper, radiation damping
 - Multi-bunch resistive wall instability