

# High Energy Jets at the LHC

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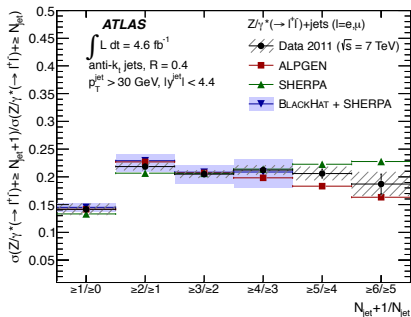
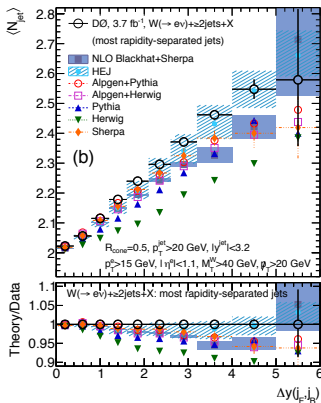


# Outline

- ▶ Motivation - Why study high energy jets?
- ▶ The Problem At Hand
- ▶ Traditional Approaches
- ▶ The Need For Something New: HEJ
  - ▶ The MRK Limit
  - ▶  $t$ -channel Dominance
  - ▶ Effective Vertices
- ▶ Extensions of the formalism
- ▶ Results

# Motivation - Why study jets at the LHC?

- As we open up the available phase space we observe an increase in jet activity (1302.6508, 1304.7098)



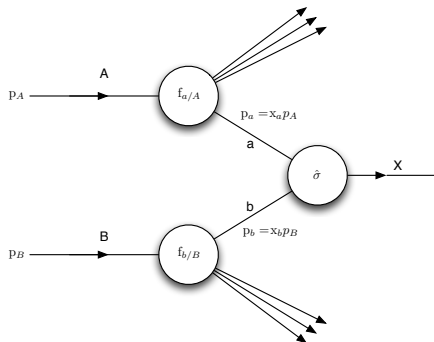
- and we have plenty of new phase space them at the LHC...

# Motivation - Why High Energy Jets?

- ▶ Jets at the LHC provide an insight in to QCD processes at scales previously unseen.
- ▶ Useful for constraining PDF's
- ▶ Jets events are an important background to understand for the study of
  - ▶ Top Physics,
  - ▶ Higgs physics,
  - ▶ BSM Physics.

# The Problem At Hand...

- QCD 'Factorisation' at the LHC



$$\sigma_{AB \rightarrow X} = \int dx_a dx_b f_{a/A}(x_a, \mu_F) f_{b/B}(x_b, \mu_F) \hat{\sigma}_{ab \rightarrow X}$$

- Here I will focus on the calculation of  $\hat{\sigma}_{ab \rightarrow X}$

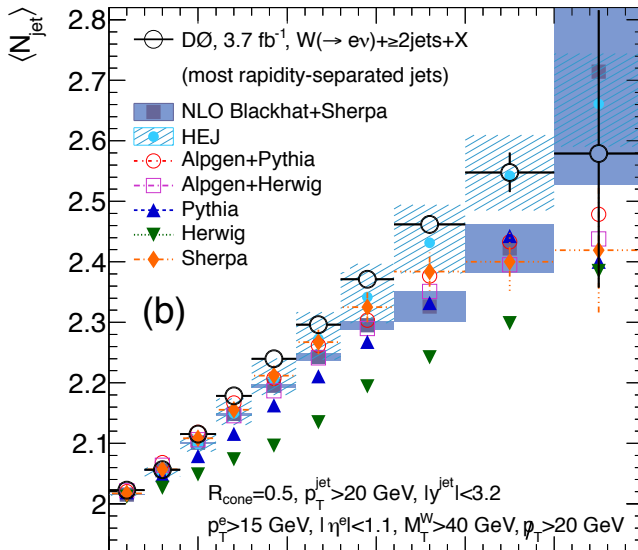
# Problems with traditional approaches

- ▶ 'Fixed Order' perturbation expansion:

$$\hat{\sigma}_{ab \rightarrow X} = \hat{\sigma}_{ab \rightarrow X}^{(0)} + \alpha_s^2(\mu_r^2) \hat{\sigma}_{ab \rightarrow X}^{(1)} + \dots$$

- ▶ The idea being that we may now truncate this series and calculate the terms in the series that remain.
- ▶ *Assumption:*  $\hat{\sigma}^{(i)}$  are assume to be  $\sim \mathcal{O}(1)$ . But these higher order terms are logarithmically enhanced in some regions of phase space.

# Traditional Approaches - Problems



## A New Approach...

- ▶ The High Energy Jets (HEJ) package provides a systematic all-order description of QCD emissions.
- ▶ Motivated by behaviour of MEs in Multi-Regge Kinematic region of phase space
- ▶ Key ingredients:
  - ▶ Large invariant mass:  $s_{ij} = 2p_i \cdot p_j \rightarrow \infty$
  - ▶ Effective vertices for extra emission,
  - ▶ The Lipatov ansatz to describe extra real emissions.



# The Multi-Regge Kinematic (MRK) Limit

- ▶ Infinite invariant mass:  $s_{ij} = 2p_i \cdot p_j \rightarrow \infty$

$$p_i = p_{\perp} \cdot (\cosh y, \cos \phi, \sin \phi, \sinh y)$$

$$\Rightarrow s_{ij} = 2p_{\perp i} p_{\perp j} (\cosh \Delta y - \underbrace{\cos \Delta \phi}_{\sim \mathcal{O}(1)})$$

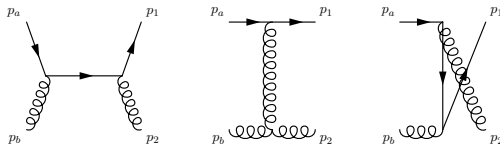
$$\Rightarrow s_{ij} \sim 2p_{\perp i} p_{\perp j} \cosh \Delta y$$

- ▶ So we can get large invariant mass having either:
  - ▶ Large perpendicular momentum or,
  - ▶ Large rapidity differences in final state jets.

# The HEJ Package

- We start from the naturally 'factorised' form for the  $2 \rightarrow 2$  scattering amplitude.

Naturally written as a contraction of a term depending only on  $p_a \sim p_1$  and a term depending on  $p_b \sim p_2$ .

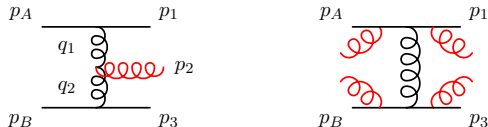


$$\mathcal{A}_{qg \rightarrow qg} \propto \frac{\langle a|\mu|1\rangle \cdot \langle b|\mu|2\rangle}{t}$$

- In the MRK limit the  $t$ -channel pole dominates

# Adding extra emissions

- ▶ We then include further hard emissions by looking at the possible ways in which we can emit an extra gluon:



- ▶ These emission sites can be combined and, in the relevant limit, expressed as an effective vertex:

$$V^\mu(q_1, q_2) = -(q_1 + q_2)^\mu + \frac{p_a^\mu}{2} \cdot \left( \frac{q_1^2}{p_a \cdot p_2} + \frac{p_b \cdot p_2}{p_a \cdot p_b} + \frac{p_2 \cdot p_3}{p_a \cdot p_3} \right) + (p_a \leftrightarrow p_1) \\ - \frac{p_b^\mu}{2} \cdot \left( \frac{q_2^2}{p_b \cdot p_2} + \frac{p_a \cdot p_2}{p_a \cdot p_b} + \frac{p_1 \cdot p_2}{p_1 \cdot p_b} \right) - (p_b \leftrightarrow p_3).$$

- ▶ Gauge invariant in all of phase-space!

## Recent Developments

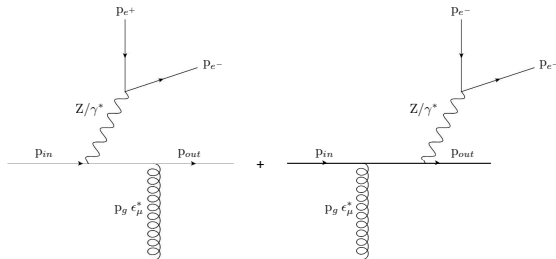
- ▶ Because of the simple structure of these amplitudes it is easy to extend the description to final states with EW bosons in.
- ▶ This has been done successfully for  $W+j$ 's and  $H+j$ 's (1206.6763,  $H$ 's paper pending) ( $Z+j$ 's and in progress).
- ▶ E.g. The inclusion of a  $Z^0$  in our final state can be included by modifying one of our contracted currents:

$$\mathcal{M}^{qg \rightarrow (Z^+ \rightarrow) e^+ e^- qg} = \frac{j_\mu^Z(p_a, p_1, p_e^+, p_e^-) \cdot \langle b|\mu|2 \rangle}{t} \quad (1)$$

- ▶ Where  $j_\mu^Z$  is the 'current' which encodes all possible emission sites of the  $Z^0$ .

# Recent Developments

- Diagrammatically:

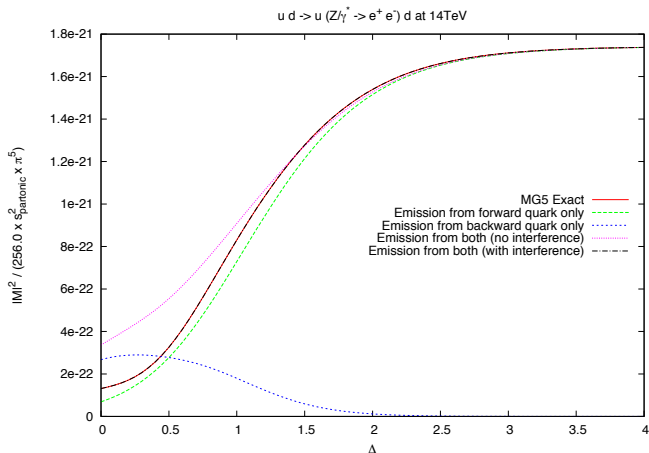


- The *only* change to our matrix element is that we must replace the  $\langle 1|\mu|a\rangle$  current with a more complicated form:

$$j_{\mu}^Z(h_q, h_l) = \left( \frac{2p_1^{\sigma} \langle 1^{hq} | \gamma^{\mu} | a^{hq} \rangle + \langle 1^{hq} | \gamma^{\sigma} | e_{hq}^{+} \rangle \langle e_{hq}^{+} | \gamma^{\mu} | a^{hq} \rangle + \langle 1^{hq} | \gamma^{\sigma} | e_{hq}^{-} \rangle \langle e_{hq}^{-} | \gamma^{\mu} | a^{hq} \rangle}{t_a} + \dots \right. \\ \left. \dots + \frac{2p_a^{\sigma} \langle 1^{hq} | \gamma^{\mu} | a^{hq} \rangle - \langle 1^{hq} | \gamma^{\mu} | e_{hq}^{+} \rangle \langle e_{hq}^{+} | \gamma^{\sigma} | a^{hq} \rangle - \langle 1^{hq} | \gamma^{\mu} | e_{hq}^{-} \rangle \langle e_{hq}^{-} | \gamma^{\sigma} | a^{hq} \rangle}{t_b} \right) \langle e_{h_l}^{+} | \gamma_{\sigma} | e_{h_l}^{-} \rangle$$

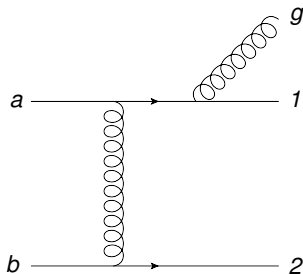
# Recent Developments - $Z$ +jets

- ▶ Other complexities arise due to the possibility of the  $\gamma^*$  channel interference
- ▶ Need an improved regularisation to include all possible emission sites for the boson.



## Recent Developments - $H$ +jets and 'unordered' emissions

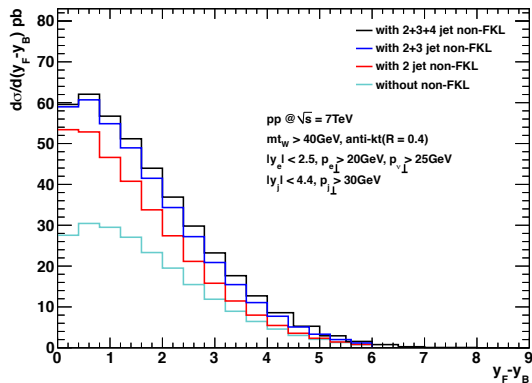
- ▶ HEJ requires a strong rapidity ordering of all final state partons.
- ▶ We have recently widened the region of applicability of our resummation to include diagrams which have one 'unordered' emission' e.g.



- ▶ This has been successfully implemented in the  $H$ +j's code.

# Results - Non-FKL Matching

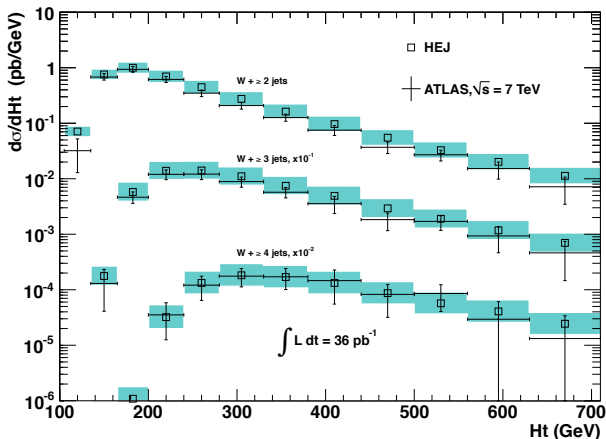
- ▶ We include the extra configurations which aren't resummed, up to 4 jets:



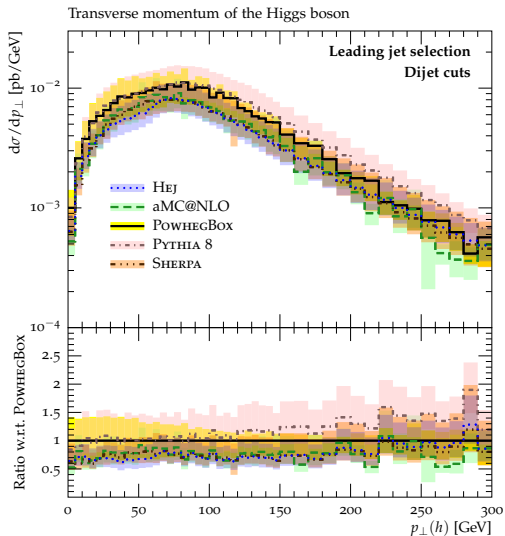


## Results - $W$ +jets

- ▶ An ATLAS study showing various results differential cross-section of  $W + j$ 's at different  $H_T$ . Even at quite low  $H_T$  we have a good description.

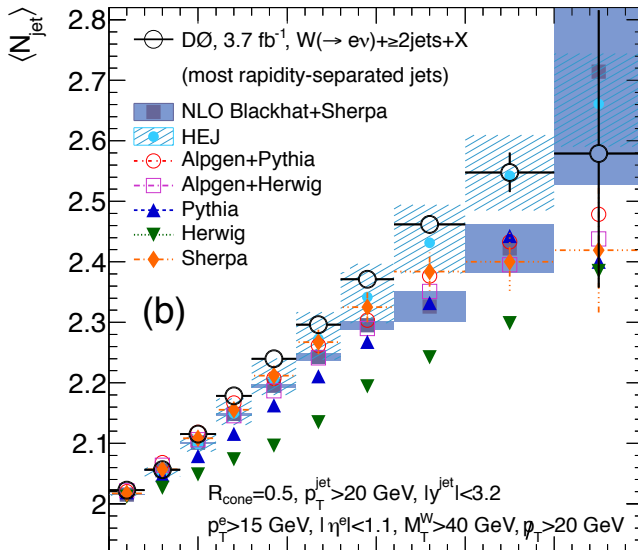


# Results - $H$ +jets



► Les Houches paper coming soon!

# Results - $W$ +jets



## In Summary

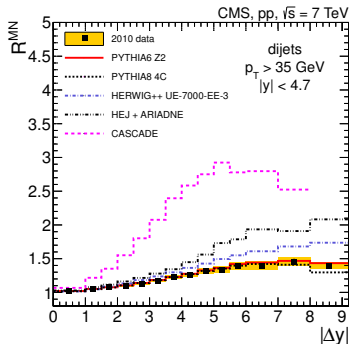
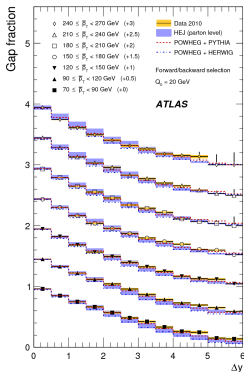
- ▶ The LHC has given us access to new regions of phase space where we need to rethink our approach to QCD calculations,
- ▶ Jets at the LHC provide us with a good check of SM processes as well as a look at what is coming next,
- ▶ The HEJ approach provides an elegant way to describe physics which was traditionally tricky to calculate,
- ▶ Recent extensions of our formalism mean we can describe a wide range of important LHC processes,
- ▶ Good description of data seen so far.

[www.cern.ch/HEJ/](http://www.cern.ch/HEJ/)

Thanks for listening!

# Backup - Mueller-Navalet Jets at the LHC

- ▶ Both CMS and ATLAS have published studies.
- ▶ CMS shows NLO matches with the data and HEJ overestimates jet activity.



- ▶ Why do NLO undershoot in ATLAS? Are we really probing the hard scatter here or testing MPI, UE tuning?