

# Exclusive photoproduction of charmonia in hadronic collisions

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# Outline

- 1 Introduction
- 2 Photoproduction in photon-proton collisions
- 3 Exclusive photoproduction in  $pp$  and  $p\bar{p}$  collisions
  - Formalism
  - Results
- 4 Conclusions

- Anna Cisek, Wolfgang Schäfer and Antoni Szczurek



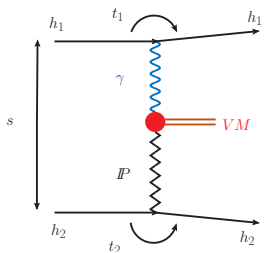
## Introduction

- Exclusive production of  $J/\Psi$  meson in **photon-proton** collisions has been studied in the energy range  $W \sim 20 - 300 \text{ GeV}$  (recently at HERA)
- This **energy range is relevant** for the exclusive photoproduction in proton-antiproton collisions at Tevatron energies for **not too large rapidities** of the meson
- For **Tevatron** we have only **one experimental point** for  $J/\Psi$  and  $\Psi'$  mesons at  $y = 0$
- New **experimental data** in proton-proton collisions for  $J/\Psi$  and  $\Psi'$  mesons in the rapidity range  $y \sim 2.0 - 4.5$  (LHCb)
- We **include absorption effects** in hadronic reaction



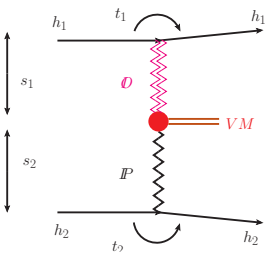
# The possible mechanism to production of vector meson in hadronic collisions

### Photoproduction



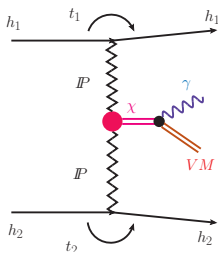
Khoze-Martin-Ryskin 2002  
Klein-Nystrand 2004

### Oderon-Pomeron fusion



Schäfer, Mankiewicz, Nachtmann 1991  
Bzdak, Motyka, Szymanowski, Cudell 2007

### Radiative Decay of $\chi_c$

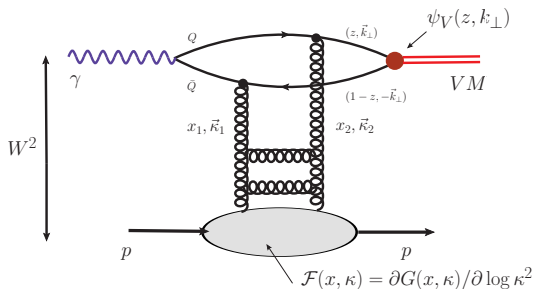


Pasechnik, Szczurek, Teryaev 2008

- In our analysis we restrict only to photon-Pomeron fusion mechanism



## Diagram for exclusive photoproduction $\gamma p \rightarrow J/\Psi p$



- $\psi_V(z, k^2) \rightarrow$  wave function of the vector meson
- $\mathcal{F}(x, \kappa^2) \rightarrow$  unintegrated gluon distribution function
- $x \sim (Q^2 + M_{J/\Psi}^2) / W^2$

## The production amplitude for $\gamma p \rightarrow J/\Psi p$

The imaginary part of the amplitude can be written as:

$$\Im \mathcal{M}_T(W, \Delta^2 = 0, Q^2 = 0) = W^2 \frac{c_v \sqrt{4\pi\alpha_{em}}}{4\pi^2} \int_0^1 \frac{dz}{z(1-z)} \int_0^\infty \pi dk^2 \psi_V(z, k^2) \int_0^\infty \frac{\pi dk^2}{\kappa^4} \alpha_S(q^2) \mathcal{F}(x_{eff}, \kappa^2) \left( A_0(z, k^2) W_0(k^2, \kappa^2) + A_1(z, k^2) W_1(k^2, \kappa^2) \right)$$

where

$$A_0(z, k^2) = m_c^2 + \frac{k^2 m_c}{M_{J/\Psi} + 2m_c},$$

$$A_1(z, k^2) = \left[ z^2 + (1-z)^2 - (2z-1)^2 \frac{m_c}{M_{J/\Psi} + 2m_c} \right] \frac{k^2}{k^2 + m_c^2},$$

$$W_0(k^2, \kappa^2) = \frac{1}{k^2 + m_c^2} - \frac{1}{\sqrt{(k^2 - m_c^2 - \kappa^2)^2 + 4m_c^2 k^2}},$$

$$W_1(k^2, \kappa^2) = 1 - \frac{k^2 + m_c^2}{2k^2} \left( 1 + \frac{k^2 - m_c^2 - \kappa^2}{\sqrt{(k^2 - m_c^2 - \kappa^2)^2 + 4m_c^2 k^2}} \right).$$



## Cross section for $\gamma p \rightarrow J/\Psi(\Psi') p$

The full amplitude:

$$\mathcal{M}_T(W, \Delta^2) = (i + \rho_T) \Im m \mathcal{M}_T(W, \Delta^2 = 0, Q^2 = 0) \exp(-B(W)\Delta^2).$$

where

$$\rho_T = \frac{\Re e \mathcal{M}_T}{\Im m \mathcal{M}_T} = \frac{\pi}{2} \frac{\partial \log \left( \Im m \mathcal{M}_T / W^2 \right)}{\partial \log W^2} = \frac{\pi}{2} \Delta_{\mathbf{P}},$$

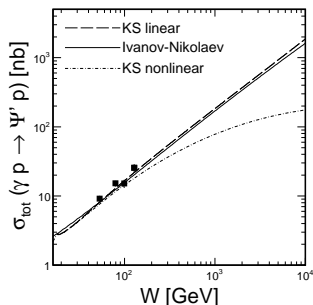
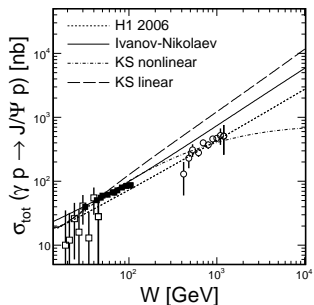
$$B(W) = B_0 + 2\alpha'_{eff} \log \left( \frac{W^2}{W_0^2} \right).$$

Total cross section can be written as:

$$\sigma_T(\gamma p \rightarrow J/\Psi p) = \frac{1 + \rho_T^2}{16\pi B(W)} \left| \frac{\Im m \mathcal{M}_T(W, \Delta^2 = 0, Q^2 = 0)}{W^2} \right|^2$$



# Total cross section for $\gamma p \rightarrow J/\Psi(\Psi') p$



HERA data and [extracted LHCb data](#)

- H1 Collaboration, Phys. Lett. B541 (2002) 251
- H1 Collaboration, Eur. Phys. J. C46 (2006) 585
- H1 Collaboration, Eur. Phys. J. C73 (2013) 2466





## Parameters of the vector meson wave functions

### How to choose parameters of the wave function

- $\Gamma(V \rightarrow e^+e^-) \Rightarrow g_V$
- $\psi_V(z, k) \Rightarrow g_V$
- normalization
- orthogonality

Decay electronic width: 
$$\Gamma(V \rightarrow e^+e^-) = \frac{4\pi\alpha_{em}^2 c_V^2}{3M_V^3} \cdot g_V^2 \cdot K_{NLO}$$

$g_V$  -leptonic decay constant: 
$$g_V = \frac{8N_c}{3} \int \frac{d^3\vec{p}}{(2\pi)^3} (M + m_q) \psi_V(z, k)$$

I.P.Ivanov, N.N.Nikolaev, A.A.Savin: Phys.Part.Nucl.37 (2006)



## Radial excitations

**Gauss:**  $\psi_{1S}(k^2) = C_1 \exp\left(-\frac{k^2 a_1^2}{2}\right)$

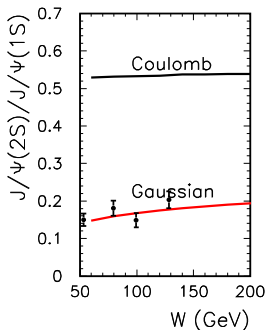
$$\psi_{2S}(k^2) = C_2 (\xi_0 - p^2 a_2^2) \exp\left(-\frac{k^2 a_2^2}{2}\right)$$

**Coulomb:**

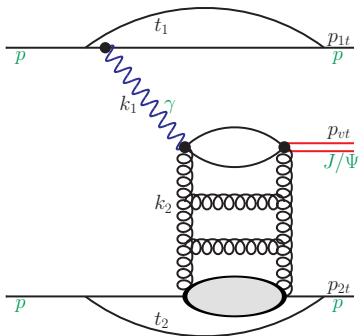
$$\psi_{1S}(k^2) = \frac{C_1}{\sqrt{M}} \frac{1}{(1 + a_1^2 k^2)^2}$$

$$\psi_{2S}(k^2) = \frac{C_2}{\sqrt{M}} \frac{\xi_0 - a_2^2 k^2}{(1 + a_2^2 k^2)^3}$$

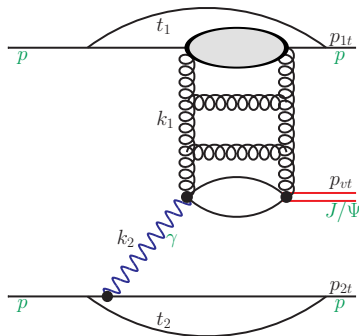
- strong dependence on the wave function
- H1 Collaboration, Phys. Lett. B541 (2002) 251



# Diagram for exclusive production of $J/\Psi$ ( $\Psi'$ ) meson in proton-proton collisions

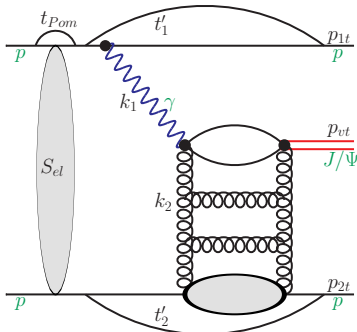


photon-Pomeron

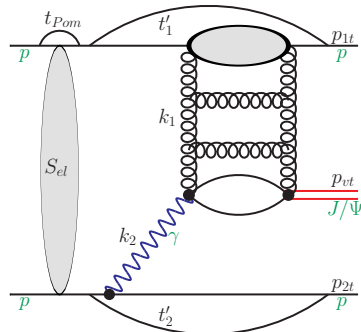


Pomeron-photon

# Diagram for $pp \rightarrow p J/\Psi(\Psi') p$ with absorptive corrections



photon-Pomeron



Pomeron-photon



## Amplitude for process $pp \rightarrow p J/\Psi(\Psi') p$

Amplitude without absorption:

$$\begin{aligned}
 \mathbf{M}^{(0)}(\mathbf{p}_1, \mathbf{p}_2) &= e_1 \frac{2}{z_1} \frac{\mathbf{p}_1}{t_1} \mathcal{F}_{\lambda'_1 \lambda_1}(\mathbf{p}_1, t_1) \mathcal{M}_{\gamma h_2 \rightarrow \nu h_2}(s_2, t_2, \mathbf{Q}_1^2) \\
 &+ e_2 \frac{2}{z_2} \frac{\mathbf{p}_2}{t_2} \mathcal{F}_{\lambda'_2 \lambda_2}(\mathbf{p}_2, t_2) \mathcal{M}_{\gamma h_1 \rightarrow \nu h_1}(s_1, t_1, \mathbf{Q}_2^2)
 \end{aligned}$$

Full amplitude for  $pp \rightarrow pVp$ :

$$\begin{aligned}
 \mathbf{M}(\mathbf{p}_1, \mathbf{p}_2) &= \int \frac{d^2 \mathbf{k}}{(2\pi)^2} S_{el}(\mathbf{k}) \mathbf{M}^{(0)}(\mathbf{p}_1 - \mathbf{k}, \mathbf{p}_2 + \mathbf{k}) \\
 &= \mathbf{M}^{(0)}(\mathbf{p}_1, \mathbf{p}_2) - \delta \mathbf{M}(\mathbf{p}_1, \mathbf{p}_2)
 \end{aligned}$$

The absorptive corrections for amplitude:

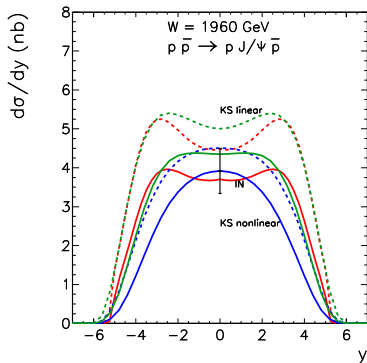
$$\delta \mathbf{M}(\mathbf{p}_1, \mathbf{p}_2) = \int \frac{d^2 \mathbf{k}}{2(2\pi)^2} T(\mathbf{k}) \mathbf{M}^{(0)}(\mathbf{p}_1 - \mathbf{k}, \mathbf{p}_2 + \mathbf{k})$$

where

$$T(\mathbf{k}) = \sigma_{tot}^{pp}(s) \exp\left(-\frac{1}{2} B_{el} \mathbf{k}^2\right)$$



# Old results for Tevatron



- old results
- CDF Collaboration, T.Aaltonen et al. Phys. Rev. Lett. 102 (2009)
- A.Cisek PhD thesis (2012)
- different UGDFs
- Gauss wave function

## Helicity conserving and helicity flip amplitudes

The full amplitude for the  $pp \rightarrow pVp$  process can be written as:

$$\begin{aligned}
 \mathcal{M}_{h_1 h_2 \rightarrow h_1 h_2 V}^{\lambda_1 \lambda_2 \rightarrow \lambda_1' \lambda_2' \lambda_V}(s, s_1, s_2, t_1, t_2) &= \mathcal{M}_{\gamma \mathbf{P}} + \mathcal{M}_{\mathbf{P} \gamma} \\
 &= \langle p_1', \lambda_1' | J_\mu | p_1, \lambda_1 \rangle \epsilon_\mu^*(q_1, \lambda_V) \frac{\sqrt{4\pi\alpha_{em}}}{t_1} \mathcal{M}_{\gamma^* h_2 \rightarrow V h_2}^{\lambda_\gamma^* \lambda_2 \rightarrow \lambda_V \lambda_2}(s_2, t_2, Q_1^2) \\
 &+ \langle p_2', \lambda_2' | J_\mu | p_2, \lambda_2 \rangle \epsilon_\mu^*(q_2, \lambda_V) \frac{\sqrt{4\pi\alpha_{em}}}{t_2} \mathcal{M}_{\gamma^* h_1 \rightarrow V h_1}^{\lambda_\gamma^* \lambda_1 \rightarrow \lambda_V \lambda_1}(s_1, t_1, Q_2^2)
 \end{aligned}$$

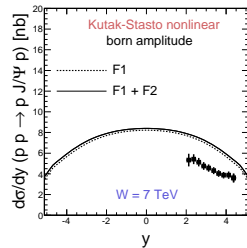
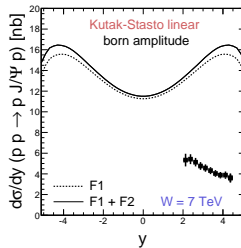
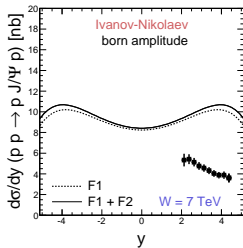
Simple structure:

$$\begin{aligned}
 \langle p_1', \lambda_1' | J_\mu | p_1, \lambda_1 \rangle \epsilon_\mu^*(q_1, \lambda_V) &= \frac{(\mathbf{e}^{*(\lambda_V)} \mathbf{q}_1)}{\sqrt{1-z_1}} \frac{2}{z_1} \cdot \\
 \cdot \chi_{\lambda'}^\dagger \left\{ F_1(Q_1^2) - \frac{i\kappa_p F_2(Q_1^2)}{2m_p} (\boldsymbol{\sigma}_1 \cdot [\mathbf{q}_1, \mathbf{n}]) \right\} \chi_\lambda
 \end{aligned}$$

- The coupling with  $F_1$  - proton helicity conserving,  $F_2$  - proton helicity flip



## Dirac vs Pauli form factors (Born)

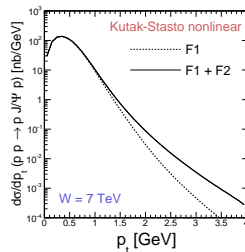
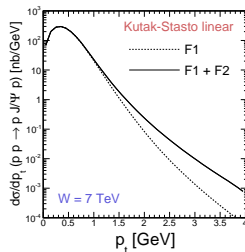
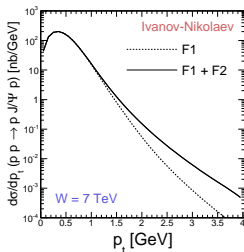


- R. Aaij et al. (LHCb collaboration), J. Phys. **G40** (2013) 045001
- R. Aaij et al. (LHCb collaboration), arXiv:1401.3288 [hep-ex]
- $F_2$  - only small change of normalization
- $F_1$  and  $F_2$  contributions do not interfere
- Absorption must be included





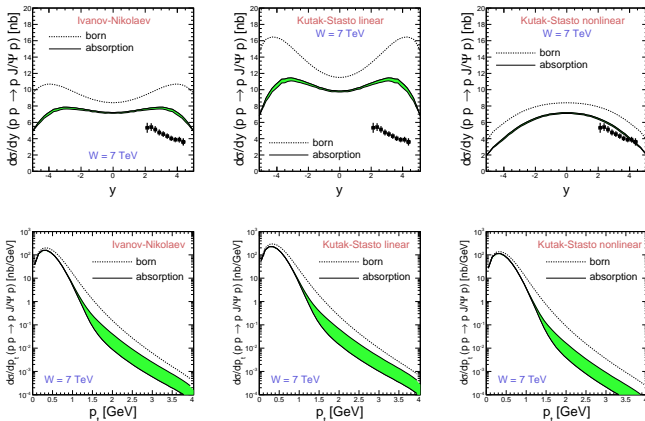
## Dirac vs Pauli form factors (Born)



- At large  $p_t$  we get an enhancement factor of the cross section of order of 10
- Pauli form factor changes the  $p_t$ -shape of elastic contribution at larger  $p_t$



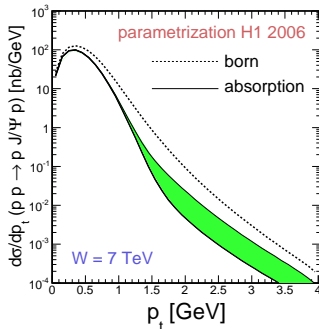
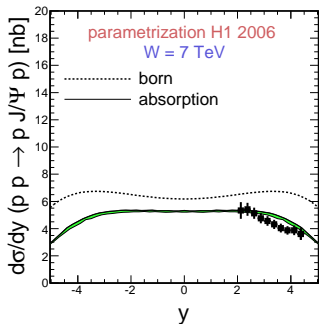
## Absorption effect



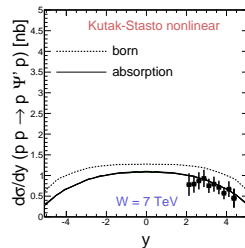
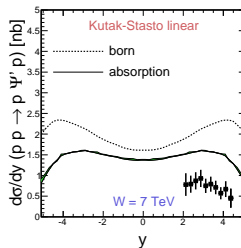
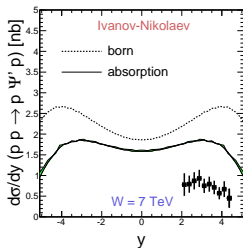
- R. Aaij et al. (LHCb collaboration), arXiv:1401.3288 [hep-ex]
- Absorption may be bigger



## Extrapolation of the HERA data

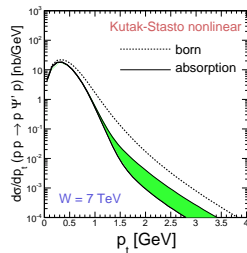
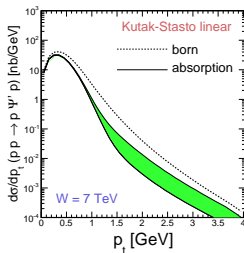
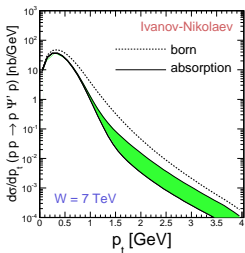


Cross section for  $\gamma p \rightarrow J/\psi p$  parametrized in the power-like form fitted to HERA data

Excited state  $\Psi'$ 

- R. Aaij et al. (LHCb collaboration), arXiv:1401.3288 [hep-ex]
- Absorption may be bigger

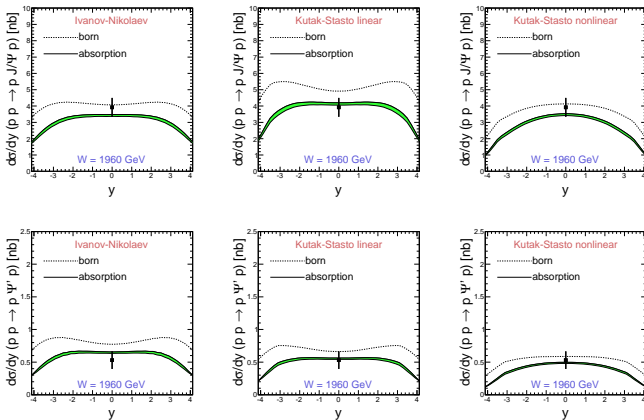


Excited state  $\Psi'$ 

- The same shape for  $\Psi'$  and  $J/\psi$



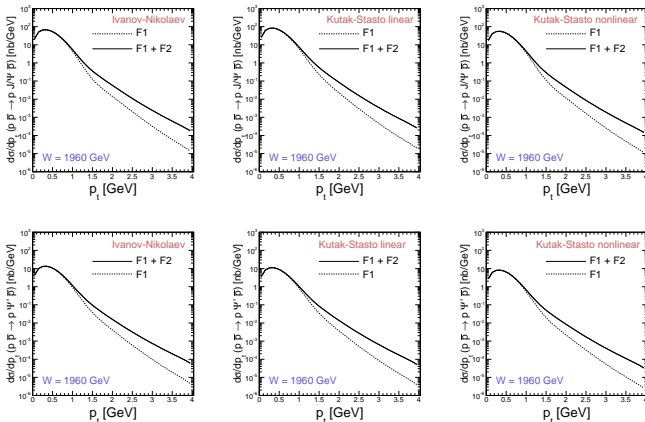
# Absorption effect for $J/\Psi$ and $\Psi'$ at the Tevatron



- CDF Collaboration, T.Aaltonen et al., Phys. Rev. Lett. 102 (2009)



# Dirac and Pauli form factor (Born) for $J/\Psi$ and $\Psi'$ at the Tevatron



## Conclusions

- We have compared our results with recent **HERA** ( $\gamma p \rightarrow J/\Psi(\Psi') p$ ) and **LHCb** ( $pp \rightarrow p J/\Psi(\Psi') p$ ) data.
- $d\sigma/dy = d\sigma^{\gamma\mathbf{P}}/dy + d\sigma^{\mathbf{P}\gamma}/dy$  only in the **Born approximation**.
- Sensitivity to the **quarkonium wave function** and **testing UGDF**.
- $d\sigma/d\phi$  is due to **interference of  $\gamma\mathbf{P}$  and  $\mathbf{P}\gamma$**  amplitudes.
- $d\sigma/dp_t$  is interesting (**spin flip**, **Pomeron-Odderon fusion**) but difficult to measure.
- **Absorptive corrections** have been included. Their **effect depends on  $p_t$  and  $y$** .

