



Kinematical constraint revisited and nonlinear evolution at large values of coupling constant

Krzysztof Kutak



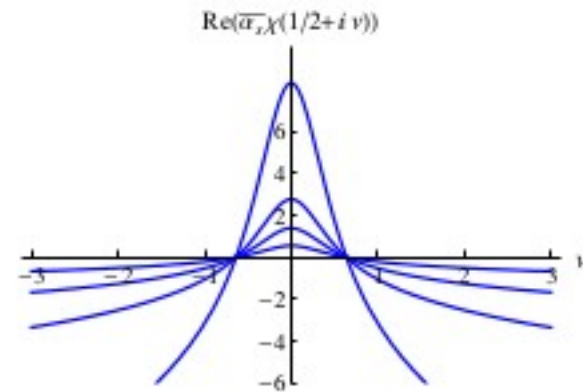
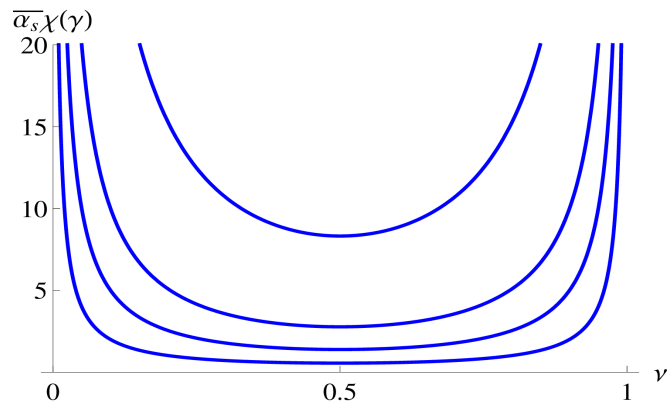
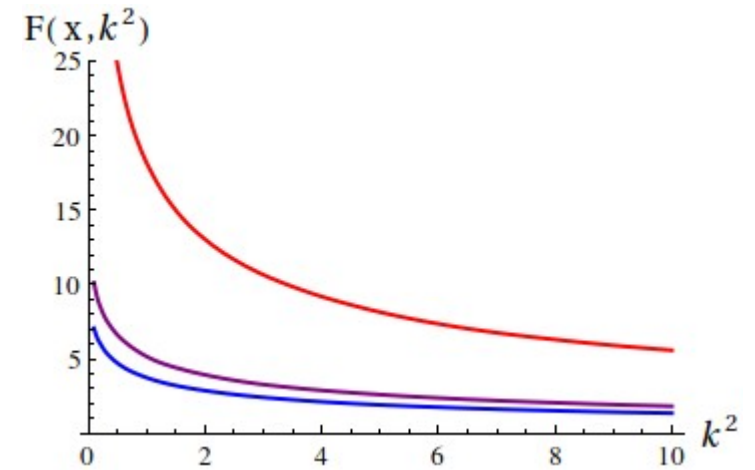
Supported by grant: LIDER/02/35/L-2/10/NCBiR/2011

The BFKL equation and its solution

$$f(x, k^2) = f_0(x, k^2) + \bar{\alpha}_s k^2 \int_{x/x_0}^1 \frac{dz}{z} \int_0^\infty \frac{dl^2}{l^2} \left[\frac{f(x/z, l^2) - f(x/z, k^2)}{|l^2 - k^2|} + \frac{f(x/z, k^2)}{\sqrt{4l^4 + k^4}} \right]$$

$$f(x, k^2) = \frac{1}{2\pi i} \int d\gamma (k^2)^\gamma f(x_0, \gamma) \left(\frac{x}{x_0} \right)^{-\bar{\alpha}_s \chi(\gamma)}, \quad \chi(\gamma) = 2\psi(1) - \psi(1-\gamma) - \psi(\gamma).$$

$$\mathcal{F}(x, k^2) = \mathcal{F}(x_0, 1/2) \frac{1}{\sqrt{4\pi \ln(x_0/x) 1/2\lambda'}} e^{\lambda \ln(x_0/x) - 1/2 \ln(k^2/k_0^2)} e^{-\frac{\ln(k^2/k_0^2)^2}{41/2\lambda' \ln(x_0/x)}}$$



$$\partial_Y \mathcal{F}(Y, \rho) = \frac{1}{2} \lambda' \partial_\rho^2 \mathcal{F}(Y, \rho) + \frac{1}{2} \lambda' \partial_\rho \mathcal{F}(Y, \rho) + (\lambda + \lambda'/8) \mathcal{F}(Y, \rho).$$

The kinematical constraint effects

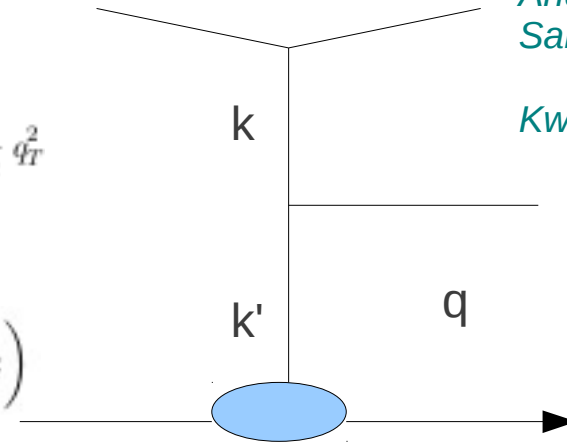
$$k^2 = k^+ k^- - k_T^2$$

$$k_T^2 > |k^+ k^-| \quad k^+ k^- \simeq -\frac{k^+}{q^+} q_T^2 = -\frac{k^+}{k^{'+} - k^+} q_T^2 = -\frac{z}{1-z} q_T^2$$

$$\theta\left(\frac{k_T^2}{q_T^2} - z\right) \quad \text{antilinear limit} \quad l \gg k \quad \theta\left(\frac{k_T^2}{l_T^2} - z\right)$$

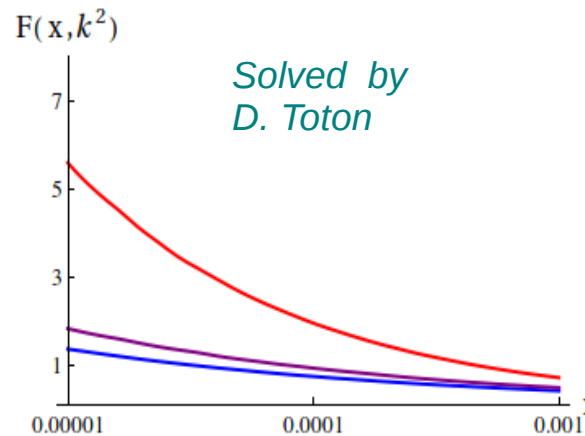
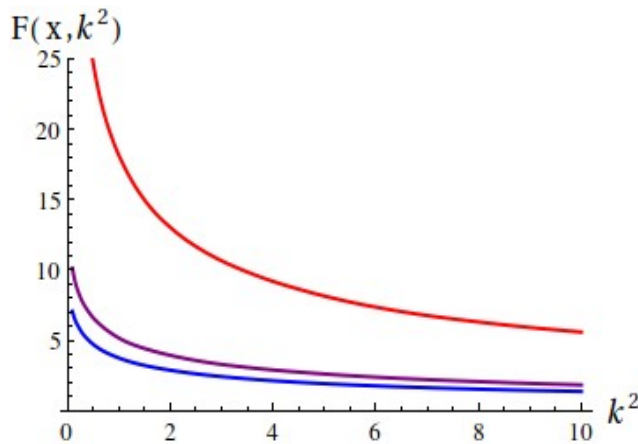
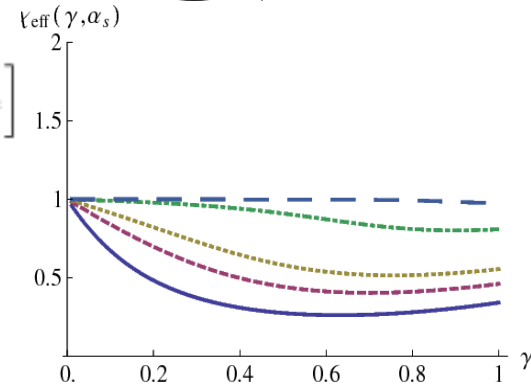
$$f(x, k^2) = f_0(x, k^2) + \bar{\alpha}_s \int_x^1 \frac{dz}{z} \int_0^\infty \frac{dl^2}{l^2} \left[\frac{\theta(k^2/l^2 - z) f(x/z, l^2) - f(x/z, k^2)}{|k^2 - l^2|} + \frac{f(x/z, k)}{\sqrt{4l^4 + k^4}} \right]$$

$$f(x, k^2) = \frac{1}{2\pi i} \int d\gamma (k^2)^\gamma \frac{1}{2\pi i} \int d\omega x^{-\omega} \frac{\omega \bar{f}_0(\omega, \gamma)}{\omega - \bar{\alpha}_s \chi_{k,c}(\gamma, \omega)}$$

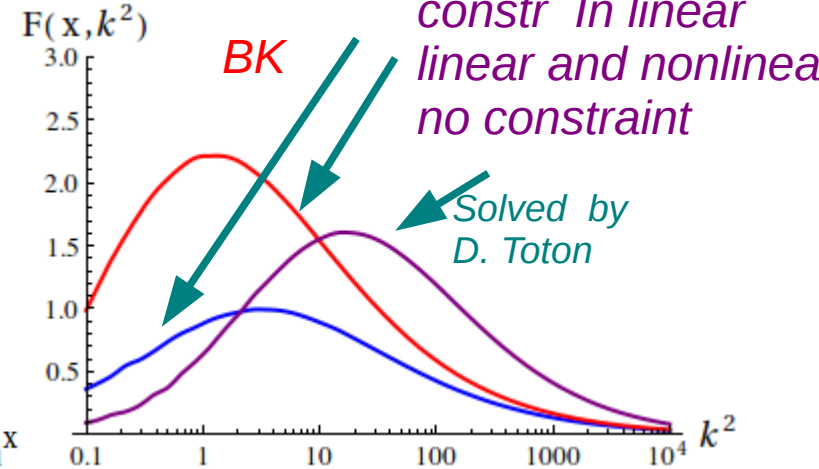


Andersson, Gustafson
Samuelsson, Kharraziha '96

Kwiecinski, Martin, Sutton '97



Solved by
D. Toton



constr In linear
linear and nonlinear
no constraint

Solved by
D. Toton

Stasto model for resummed BFKL with kinematical constraint and DGLAP effects

Stasto, '07

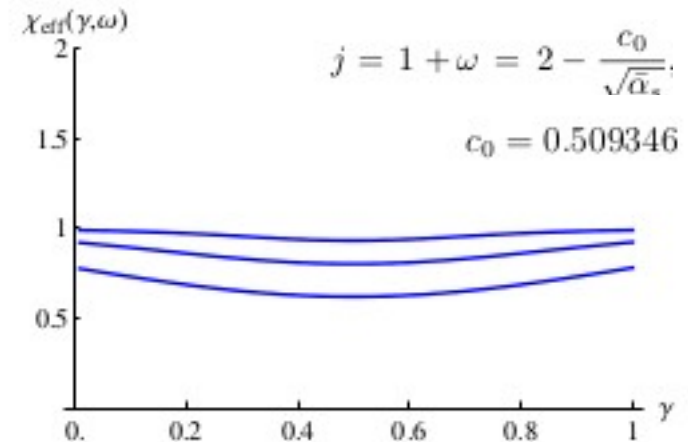
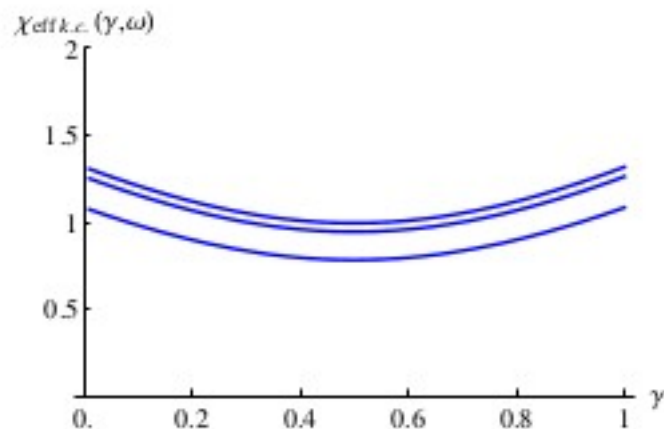
$$f(x, k^2) = f_0(x, k^2) + \bar{\alpha}_s k^2 \int_x^1 \frac{dz}{z} \int_0^\infty \frac{dl^2}{l^2} \left[\frac{f(x/z, l^2) \theta(l - kz) \theta(k/z - l) - f(x/z, k^2)}{|l^2 - k^2|} + \frac{f(x/z, k^2)}{\sqrt{4l^4 + k^4}} \right]$$

$$\chi_{k.c.}(\gamma, \omega) = 2\psi(1) - \psi(1 - \gamma + \omega/2) - \psi(\gamma + \omega/2).$$

$$\chi_{eff}(\gamma, \omega) = \bar{\alpha}_s \chi_{k.c.}(\gamma, \omega) (1 + A\omega)$$

Crucial behaviour vanishing eigenvalue when $\omega \rightarrow 1$

Contains DGLAP anomalous Dimension at LO in $\ln Q^2$



The leading order in AdS/CFT *Brower, Polchinski, Strassler*

$$j = 1 + \omega = 2 - \frac{c_0}{\sqrt{\bar{\alpha}_s}}, \quad c_0 = 1/\pi$$

Higher orders:
Costa, Goncalves, Penedones '12
Kotikov, Lipatov '13
Janik '14

Stasto model for BFKL with kinematical constraint and DGLAP effects

Kutak, Surowka, '13

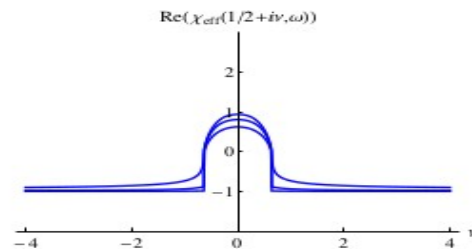
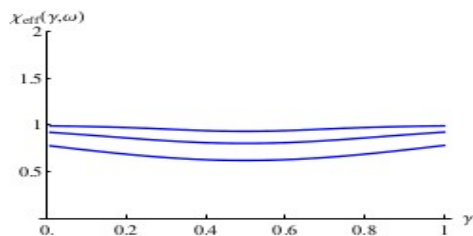
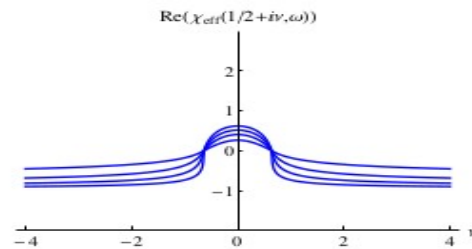
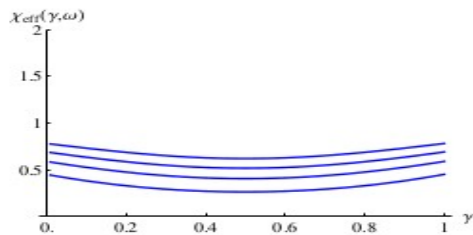
$$f(x, k^2) = f_0(x, k^2) + \bar{\alpha}_s k^2 \int_x^1 \frac{dz}{z} \int_0^\infty \frac{dl^2}{l^2} \left[\frac{f(x/z, l^2) \theta(l - kz) \theta(k/z - l) - f(x/z, k^2)}{|l^2 - k^2|} + \frac{f(x/z, k^2)}{\sqrt{4l^4 + k^4}} \right]$$

we choose the symmetric form of k.c.

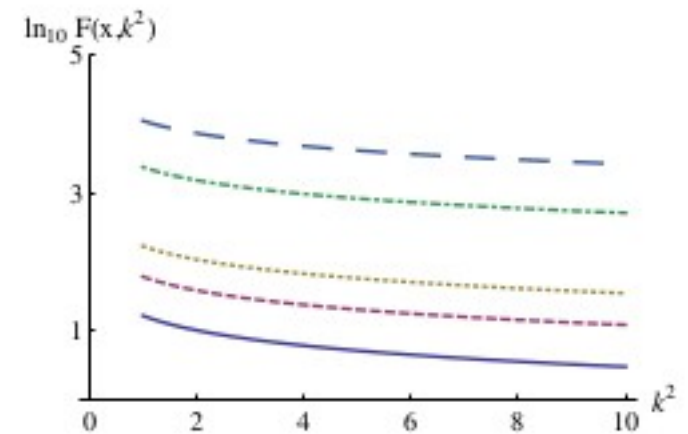
$$\chi_{k.c.}(\gamma, \omega) = 2\psi(1) - \psi(1 - \gamma + \omega/2) - \psi(\gamma + \omega/2).$$

$$f(x, k^2) = \frac{1}{2\pi i} \int d\gamma (k^2)^\gamma \frac{1}{2\pi i} \int d\omega x^{-\omega} \frac{\omega \bar{f}_0(\omega, \gamma)}{\omega - \bar{\alpha}_s \chi_{k.c.}(\gamma, \omega)}$$

$$\omega = \text{Re}(\chi_{eff k.c.}(1/2 + i\nu, \omega)) \quad \chi_{eff}(\gamma, \omega) = \bar{\alpha}_s \chi_{k.c.}(\gamma, \omega) (1 + A\omega)$$



$$f(x, k^2) = \frac{1}{2\pi i} \int d\gamma (k^2)^\gamma f(x_0, \gamma) \left(\frac{x}{x_0}\right)^{-\bar{\alpha}_s \chi(\gamma)}$$

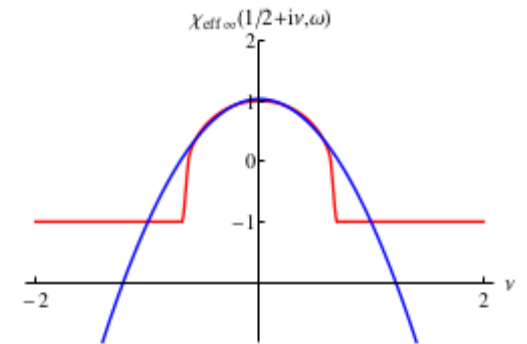


BFKL equation at strong coupling

Kutak, Surowka, '13

Fit function $\chi_{eff\infty}(\omega, 1/2 + i\nu) = \sum_{n=-M}^N A_n \nu^n$

$$\chi_{eff\infty}(\omega, 1/2 + i\nu) = P_{10}(\nu)\theta(\nu + 0.683)\theta(0.683 - \nu) - \theta(-\nu - 0.683) - \theta(\nu - 0.683)$$

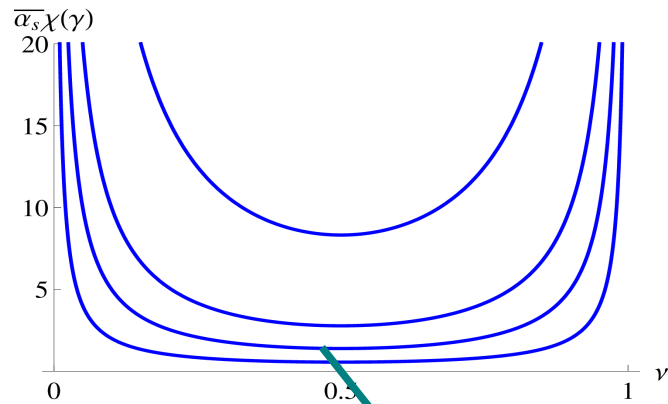


Works very well: $\chi_{eff\infty}(\omega, 1/2 + i\nu) = 1.02795 - 2.04635\nu^2 \equiv \lambda_{st} - \frac{1}{2}\lambda'_{st}\nu^2$

$$\partial_Y \Phi(Y, \rho) = \frac{1}{2}\lambda'_{st}\partial_\rho^2 \Phi(Y, \rho) + \frac{1}{2}\lambda'_{st}\partial_\rho \Phi(Y, \rho) + (\lambda_{st} + \lambda'_{st}/8)\Phi(Y, \rho) \quad \lambda'_{st} = 4.08, \lambda_{st} = 1.02$$

Gluon density at the large coupling values

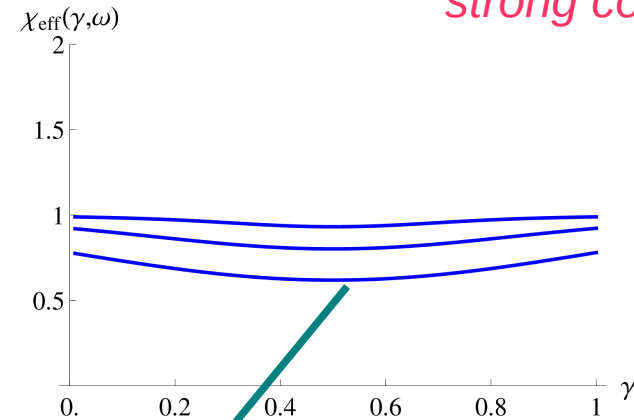
weak coupling



$$\chi(\gamma) = 2\psi(1) - \psi(1-\gamma) - \psi(\gamma)$$

weak coupling

strong coupling



Stasto '07

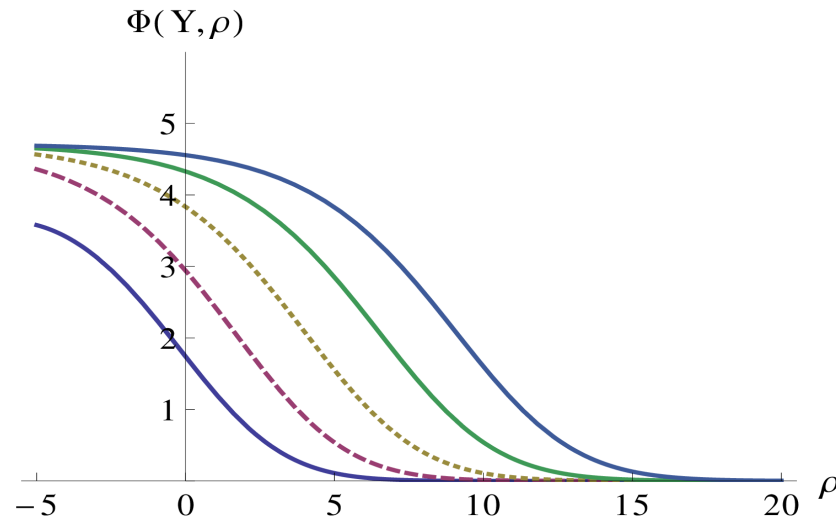
strong coupling

Kutak, Surowka '13

critical point dominates
at large coupling

WW density at the large coupling values

Kutak, Surowka, '13



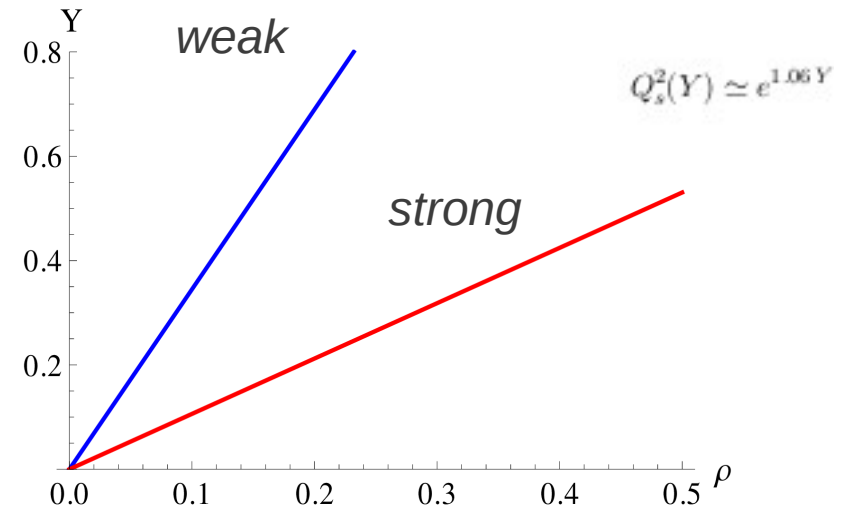
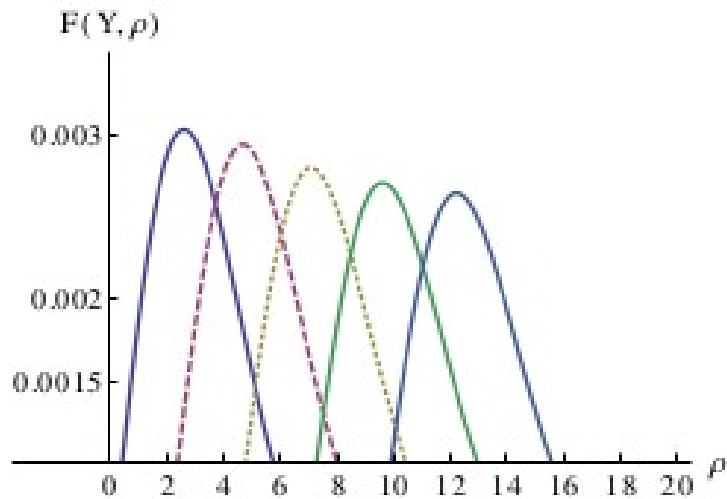
Nonlinear nonlinear equation valid at strong coupling limit

$$\partial_Y \Phi(Y, \rho) = \frac{1}{2} \lambda'_{st} \partial_\rho^2 \Phi(Y, \rho) + \frac{1}{2} \lambda'_{st} \partial_\rho \Phi(Y, \rho) + (\lambda_{st} + \lambda'_{st}/8) \Phi(Y, \rho) - \frac{\bar{\alpha}_s}{\pi R^2} \Phi^2(Y, \rho)$$

Saturation scale at large values of coupling constant

$$\mathcal{F}_{BK}(Y, \rho) = \frac{N_c}{4\pi\alpha_s} \partial_\rho^2 \Phi(Y, \rho)$$

$$\partial_\rho \mathcal{F}_{BK}(Y, \rho) \Big|_{\rho = \ln Q_s^2(Y)} = 0$$



Similar behaviour as in
Mueller, Shoshi, Xiao '10

Outlook

- *Entropy at large coupling*
- *Full range in running coupling effect*
- *Just for curiosity check the cross section for inclusive production*
- *Perhaps formulate directly in momentum space*