



Higgs Pair Production via Vector Boson Fusion at the LHC

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Andrea Massironi
Northeastern University

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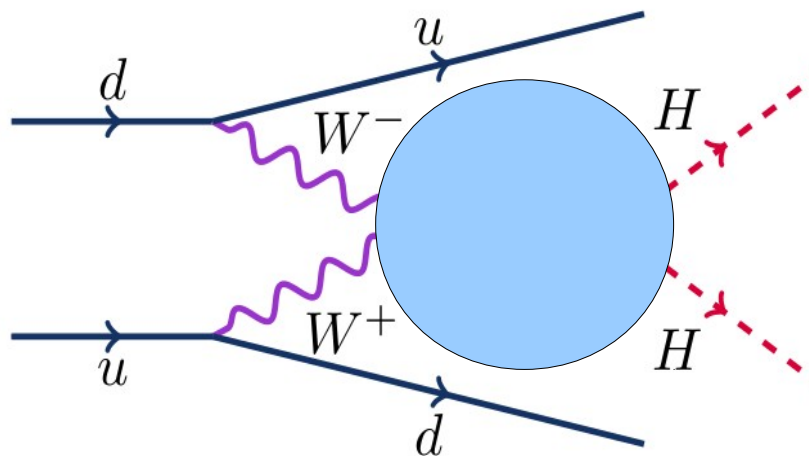
Layout

- **Vector Boson Fusion di-Higgs production**
 - Goal of the analysis
 - LHC @13 TeV
- **New physics**
 - **Direct search for new physics**
 - Resonant di-Higgs production
 - **Indirect search for new physics**
 - Strong Double Higgs production
 - Parton level analysis
 - Hadron level analysis
- Results



Vector Boson Fusion di-Higgs searches

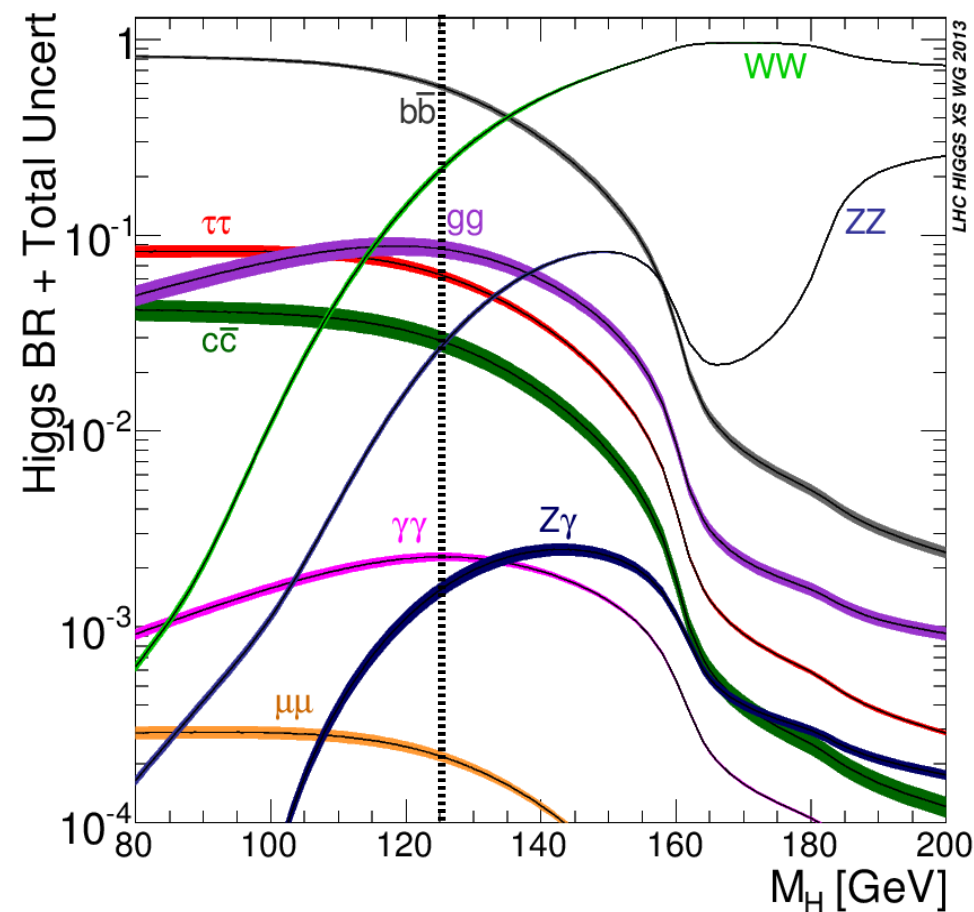
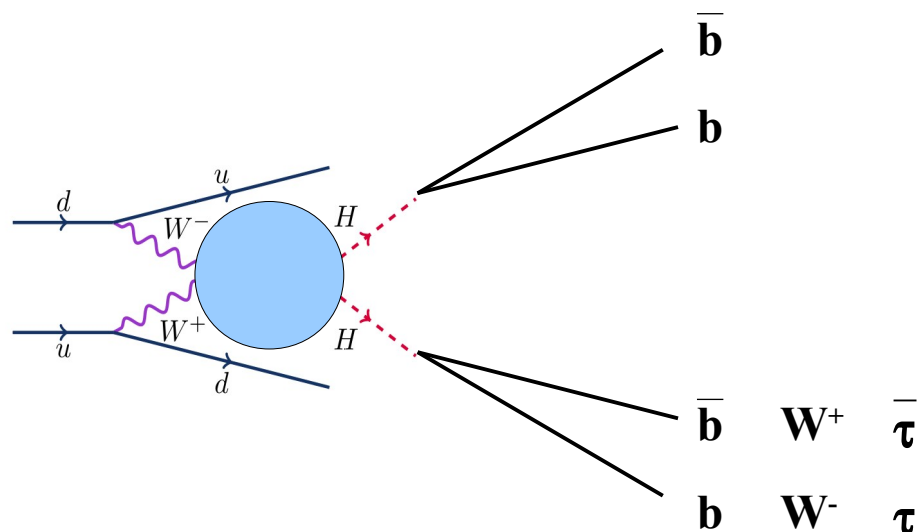
- **Higgs pair production** is one of the most crucial processes for future LHC runs
 - stringent tests of our understanding of electroweak symmetry breaking
- In the SM, the dominant process is **gluon fusion**, with 33 fb at 14 TeV
 - direct sensitivity to the **Higgs trilinear coupling** c_3
- Higgs pair production in **Vector-Boson Fusion** (VBF) is small in the SM, 2 fb at 14 TeV
 - provides direct information on the HVV and HHVV coupling
- Higgs pair production can be substantially enhanced in various BSM scenarios
 - Production of an on-shell **resonance** decaying into two Higgs bosons
 - Modification of the **couplings** and indirect search of new physics



di-Higgs decays

Final states considered

- Highest BR $H \rightarrow b\bar{b}$
- Clean and QCD-free final state $H \rightarrow WW \rightarrow l\nu l\nu$





Three analyses

Resonant VBF di-Higgs search

Non-resonant VBF di-Higgs search at parton level

Non-resonant VBF di-Higgs search at hadron level

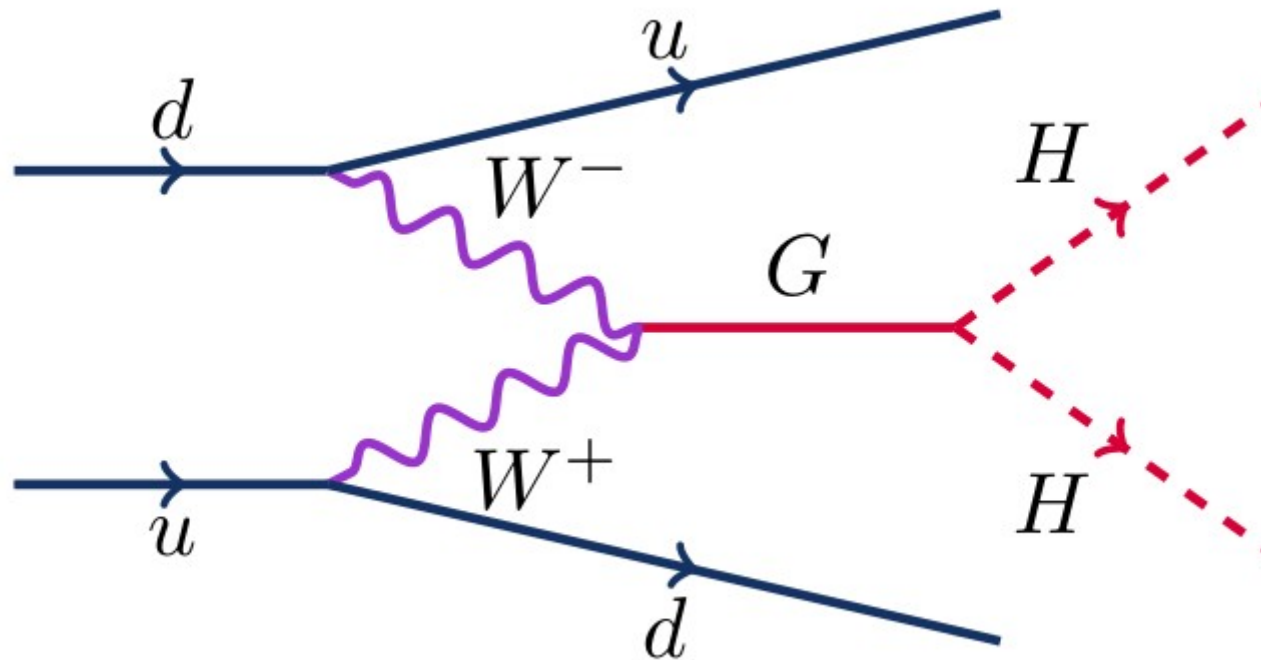


Resonant Higgs pair production

- Direct search [0]

- If new particles have masses within LHC reach
- On-shell decay into a Higgs boson pair
- Couplings to gauge bosons play an important role
- Benchmark: KK-graviton ($\tilde{k} = 0.1$) [1,2]

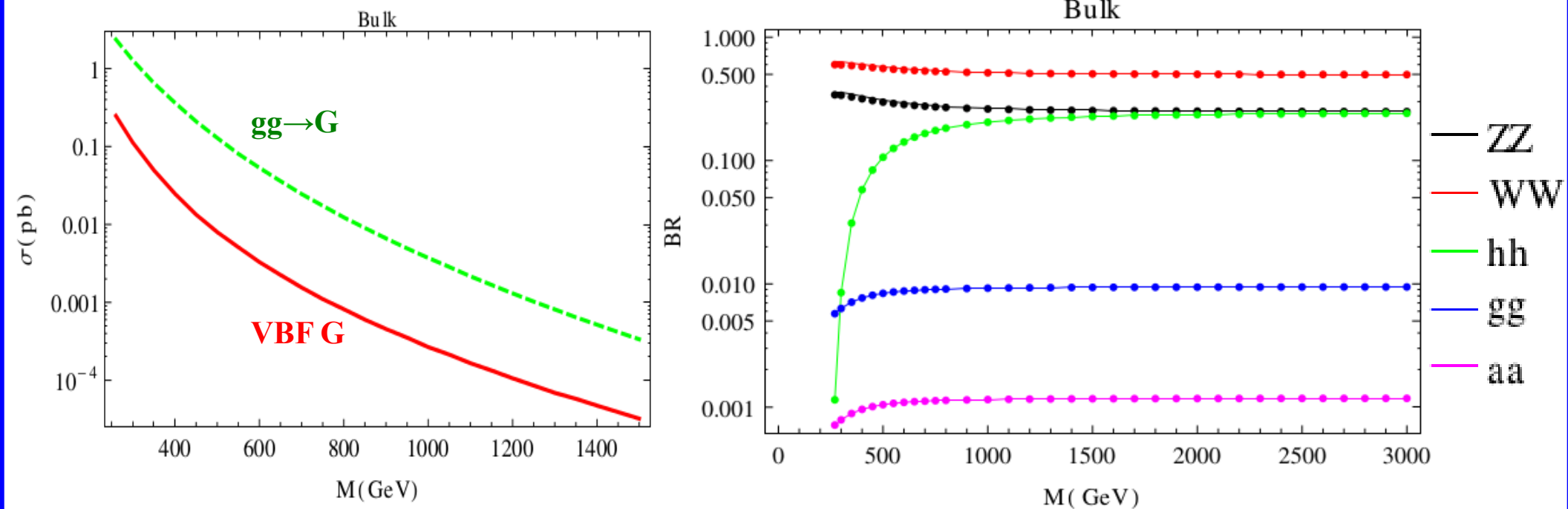
Alexander Belayaev,
Olivier Bondu,
Andrea Massironi,
Alexandra Oliveira,
Rogerio Rosenfeld,
Veronica Sanz





$G \rightarrow HH$

- Production cross-section: VBF production $\sim 1/10$ of gluon fusion production
- Branching ratio: di-Higgs is one of the highest BR ($\sim 25\%$)





Selections

- Parton level analysis, with jets algorithm run on top of partons
- Exploit boosted topology from $H \rightarrow b\bar{b}$ decays

$G \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ analysis

- Objects:

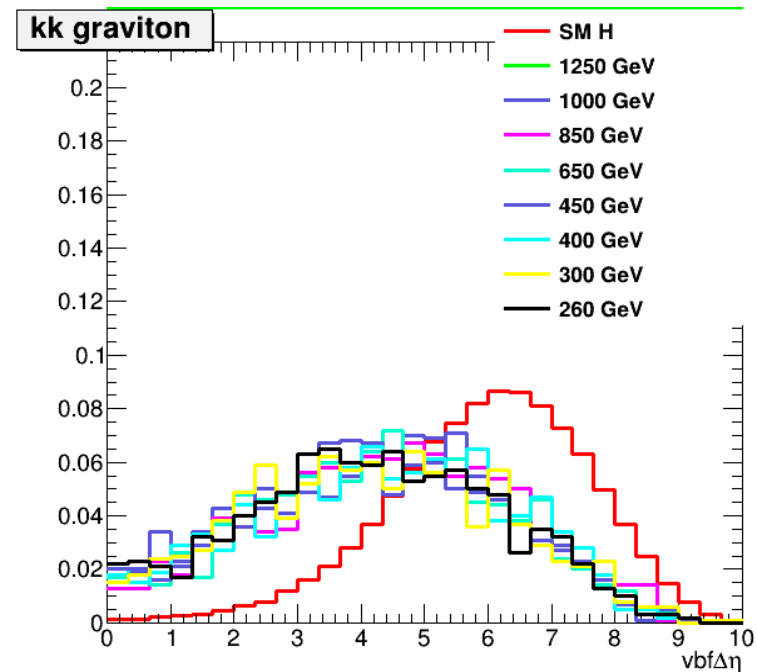
- Jet $p_T > 30$ GeV
- Jet $|\eta| < 4.7$

- VBF topology:

- Invariant mass of 2 jets: $m_{jj} > 400$ GeV
- No $\Delta\eta_{jj}$ cut is applied, due to p_T -dependence of the KK-graviton vertices with vector bosons

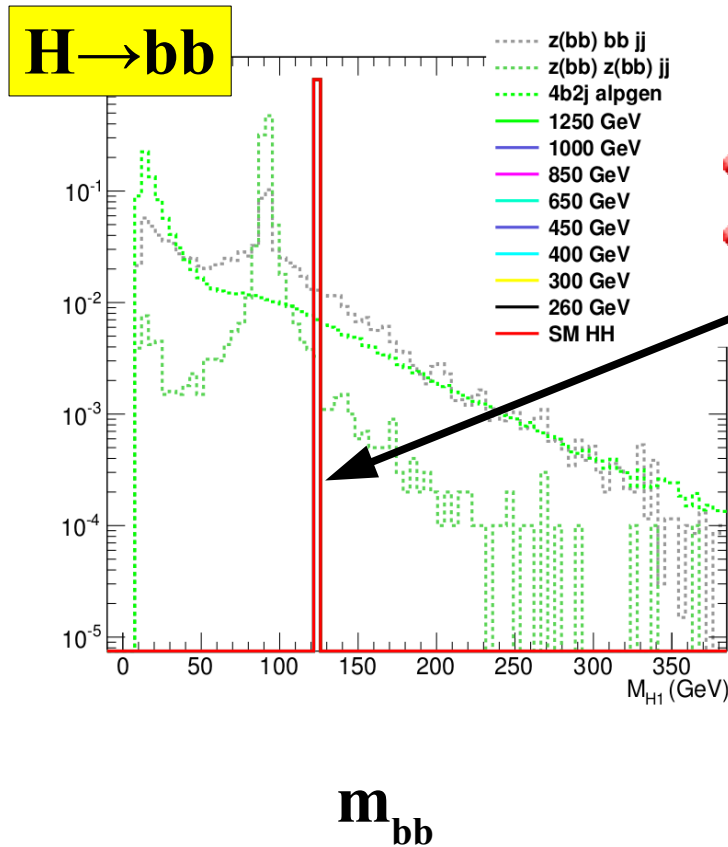
- Boosted di-Higgs system:

- $M_{HH} > 250$ GeV
- $p_T^{HH} > 60$ GeV
- $\Delta\eta_{HH} < 2$

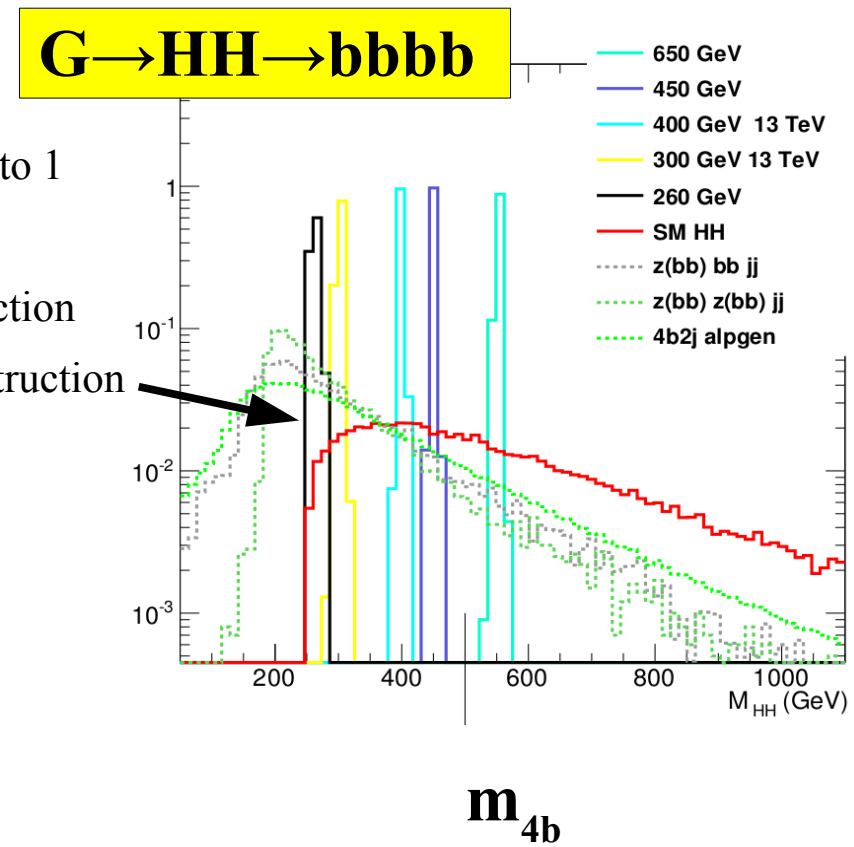


Caveat

- No selections concerning $H \rightarrow bb$ and $G \rightarrow HH \rightarrow bbbb$ invariant mass are applied
 - To be addressed in a dedicated study after showering and hadronization (ongoing study...)
 - Very conservative result



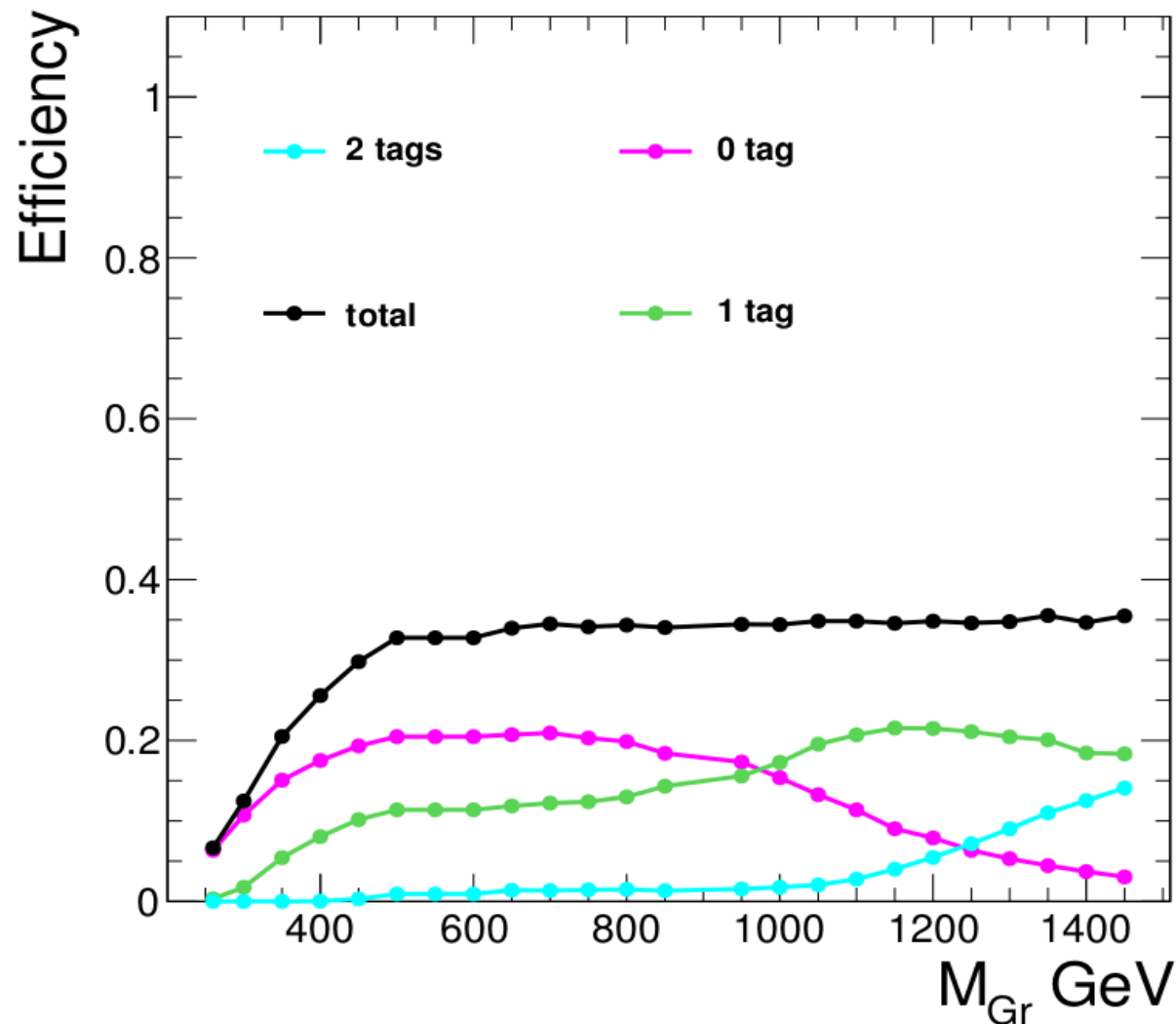
- Distributions normalized to 1
- Improvements from:
 - Higgs mass reconstruction
 - Graviton mass reconstruction





Boosted jets

- The higher the G mass the higher is the probability for the two Higgs to be boosted and the two b-jets to be merged
- We recover an uniform analysis efficiency combining different categories





Disclaimer

- Feasibility study to establish more competitive channels and search strategies
 - Analysis performed at parton level, no showering and hadronization effects are taken into account
 - Perfect b-tag efficiency is considered withing experimental acceptance volume
 - Simple cut based analysis performed
-
- Complete studies, with showering effects, reconstruction effects and possible additional final states are ongoing





Yields

- Yields with 3 ab⁻¹ at 13 TeV
- High signal yields
- Still big background contamination, but it can be reduced by means of m_{bb} and m_{4b} selections

	$pp \rightarrow Gr(HH)jj$ σ (pb)	σ^*_{eff} (pb)	Nevents (3000/fb)
1450 GeV	6.91E-06	8.17E-07	2.45E+00
1250 GeV	1.86E-05	2.14E-06	6.41E+00
1050 GeV	7.87E-05	9.13E-06	2.74E+01
850 GeV	2.63E-04	3.00E-05	9.01E+01
650 GeV	1.28E-03	1.40E-04	4.20E+02
450 GeV	8.87E-03	8.80E-04	2.64E+03
400 GeV	1.41E-02	1.20E-03	3.60E+03
300 GeV	1.55E-02	6.42E-04	1.93E+03
260 GeV	5.72E-03	1.27E-04	3.80E+02



Non-resonant VBF di-Higgs search at parton level

- Indirect search [0]

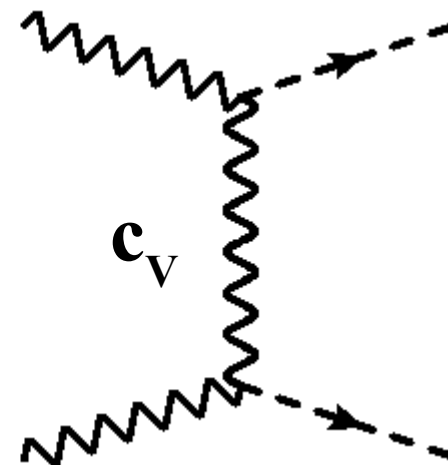
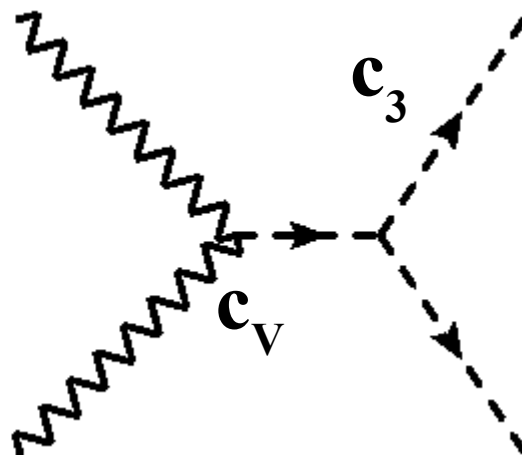
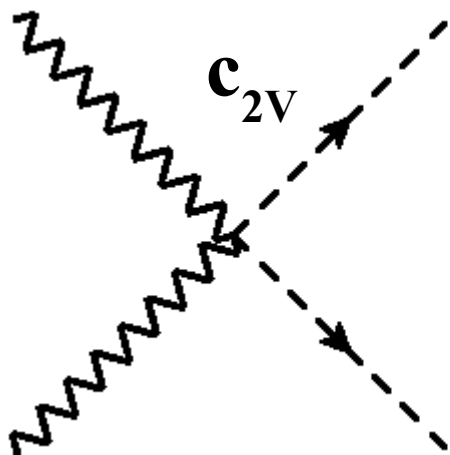
- Effective lagrangian up to dimension 6 operators [3,4]
- Modification of coupling parameters
 - $c_V g_{HVV}^{\text{SM}} HVV \rightarrow$ partially constrained through single Higgs searches
 - $c_{2V} g_{HHVV}^{\text{SM}} HHVV \rightarrow$ the only process to measure it is via VBF HH
 - $c_3 g_{HHH}^{\text{SM}} HHH \rightarrow$ accessible also via gluon fusion HH
- SM: $c_V = c_{2V} = c_3 = 1$

Olivier Bondu,
Roberto Contino,
Maxime Gouzevitch,
Andrea Massironi,
Alexandra Oliveira,
Juan Rojo

$$0.5 < c_V < 1.5$$

$$0.0 < c_{2V} < 2.0$$

$$0.0 < c_3 < 2.0$$





Selections

• Objects:

- Light and b jet and $p_T > 25$ GeV
- Light (b) jet $|\eta| < 4.5$ (2.5) and lepton $|\eta| < 2.5$

• VBF topology:

- Invariant mass of 2 jets: $m_{jj} > 500$ GeV
- High separation: $\Delta R_{jj} > 4.0$

Two analyses:

• **4b + 2j**

- Tighter VBF cuts to suppress QCD background
 - $m_{jj} > 800$ GeV
- $m_{bb} \sim 125$ GeV $\pm 15\%$ for both the bb pairs

Main backgrounds:

• **QCD 4b + 2j**

• **Z(bb) + 2b + 2j**

• **2b + 2l2v + 2j**

- Exploit $H \rightarrow WW$ kinematic
 - Invariant mass of di-lepton system $m_{ll} < 70$ GeV
 - Transverse mass of di-W system $m_T < 125$ GeV
- $m_{bb} \sim 125$ GeV $\pm 15\%$

• **WW + 2b + 2j**

- Mainly dominated by **tt + 2j**



High energy regime

• For $(c_V - c_{2V}) \neq 0$ the $VV \rightarrow HH$ cross section grows with the partonic energy

• Modification from SM prediction \rightarrow high energies

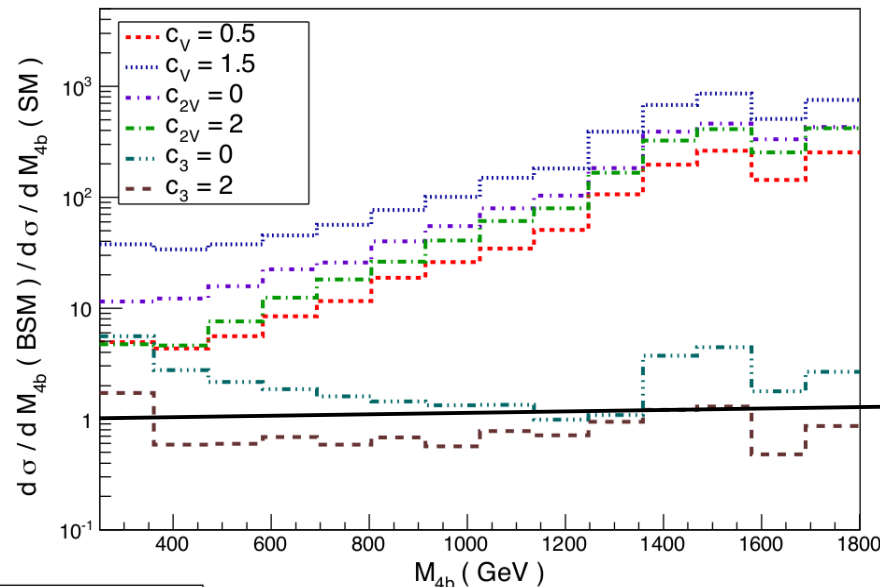
• **4b + 2j**

• $m_{4b} > 1000$ GeV

• **2b + 2l2v + 2j**

• $m_{llbb} > 500$ GeV

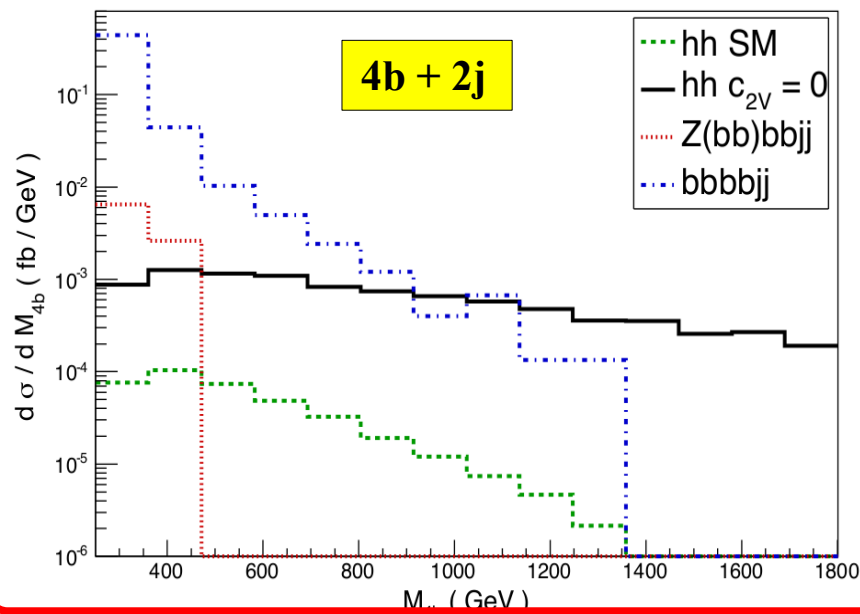
• $p_T^{bb} > 200$ GeV



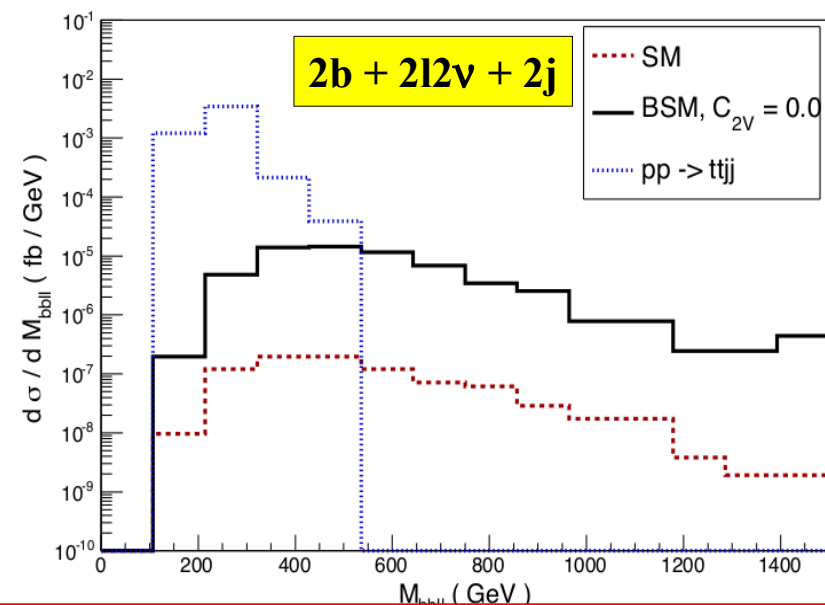
• Normalized to SM HH

• Enhancement up to x100 !

SM HH



• Suppression of most of background





Yields

- After all the selections applied
 - Selections efficiency quite high $\sim 25\%$ for non-SM VBF HH signal
 - Good sensitivity for non-SM VBF HH signal and coupling
 - HL-LHC is needed: $L = 3 \text{ ab}^{-1}$

4b + 2j

Sample	$N_{\text{ev}} = \sigma_{IV} \cdot \mathcal{L} (3 \text{ ab}^{-1})$
SM	6.788
$c_V = 0.5$	587.5
$c_V = 1.5$	2039
$c_{2V} = 0$	1160
$c_{2V} = 2$	982.5
$c_3 = 0$	10.21
$c_3 = 2$	5.385
$4bjj$	355.5
$Zbbjj \rightarrow 4bjj$	< 7

2b + 2l2v + 2j

Sample	$N_{\text{ev}} = \sigma_{III} \cdot \mathcal{L} (3 \text{ ab}^{-1})$
SM	0.1
$c_V = 0.5$	1.4
$c_V = 1.5$	45.2
$c_{2V} = 0.0$	10.3
$c_{2V} = 2.0$	8.1
$c_3 = 0.0$	0.2
$c_3 = 2.0$	0.1
$WWbbjj$	–
$t\bar{t}jj \rightarrow WWbbjj$	< 6.2



Non-resonant VBF di-Higgs search at hadron level

Olivier Bondu,
Roberto Contino,
Andrea Massironi,
Juan Rojo

$$0.0 < c_{2V} < 2.0$$

$$-4.0 < c_3 < 6.0$$

- Parton level events **showered** and **hadronized** with **Pythia8**
- Jet clustering using **FastJet** with the anti-kT algorithm with **R=0.4**
- Final states:
 - **4b2j**
 - **2b2τ2j**
 - **2b2W2j** (in progress)
- Realistic b-tagging and τ-tagging, including mistag rate
- Only hadronic decays of τ used in the analysis.
- **14 TeV** and **100 TeV** hadron colliders are used as benchmarks
- c_{2V} and c_3 estimated by means of a likelihood fit [5]:

$$\sigma = c_V^4 \sigma_{\text{SM}} (1 + A\delta_{c_{2V}} + B\delta_{c_3} + C\delta_{c_{2V}}\delta_{c_3} + D\delta_{c_{2V}}^2 + E\delta_{c_3}^2)$$

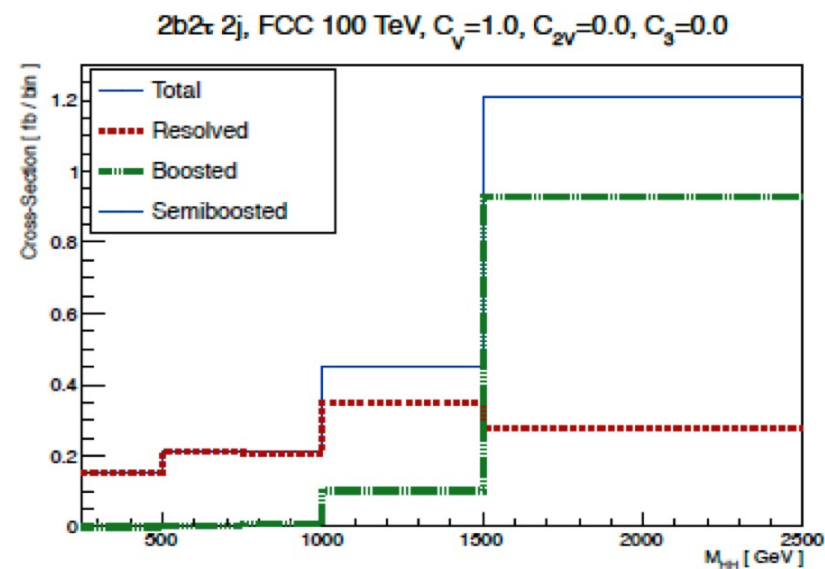
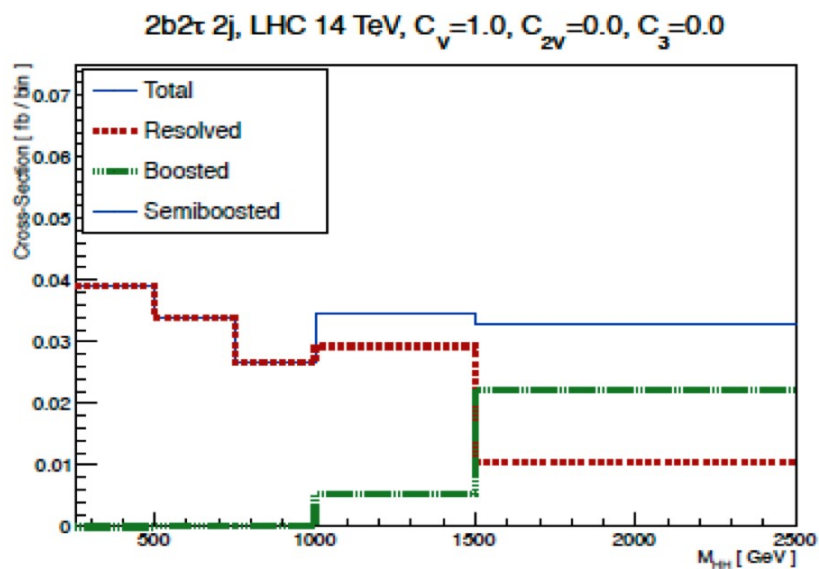
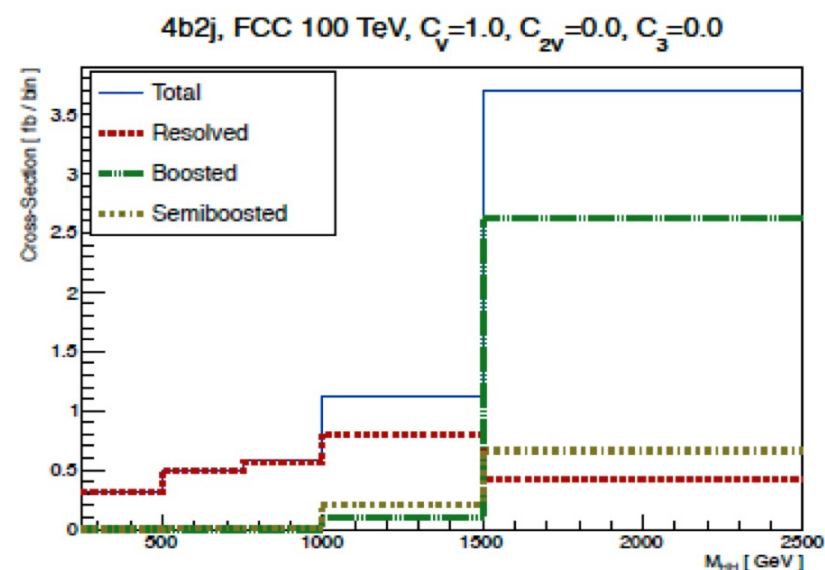
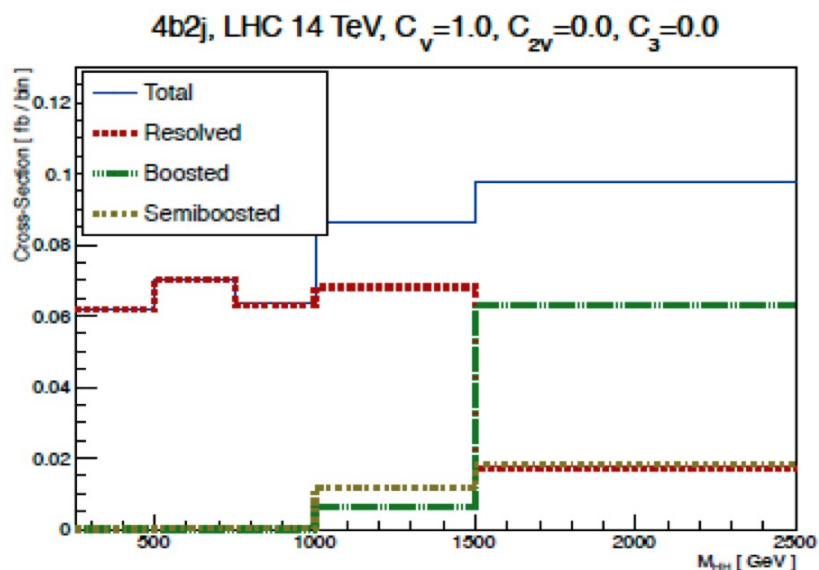
$$\delta_{c_{2V}} \equiv 1 - \frac{c_{2V}}{c_V^2}, \quad \delta_{c_3} \equiv 1 - \frac{c_3}{c_V}$$



Non-resonant VBF di-Higgs search at hadron level

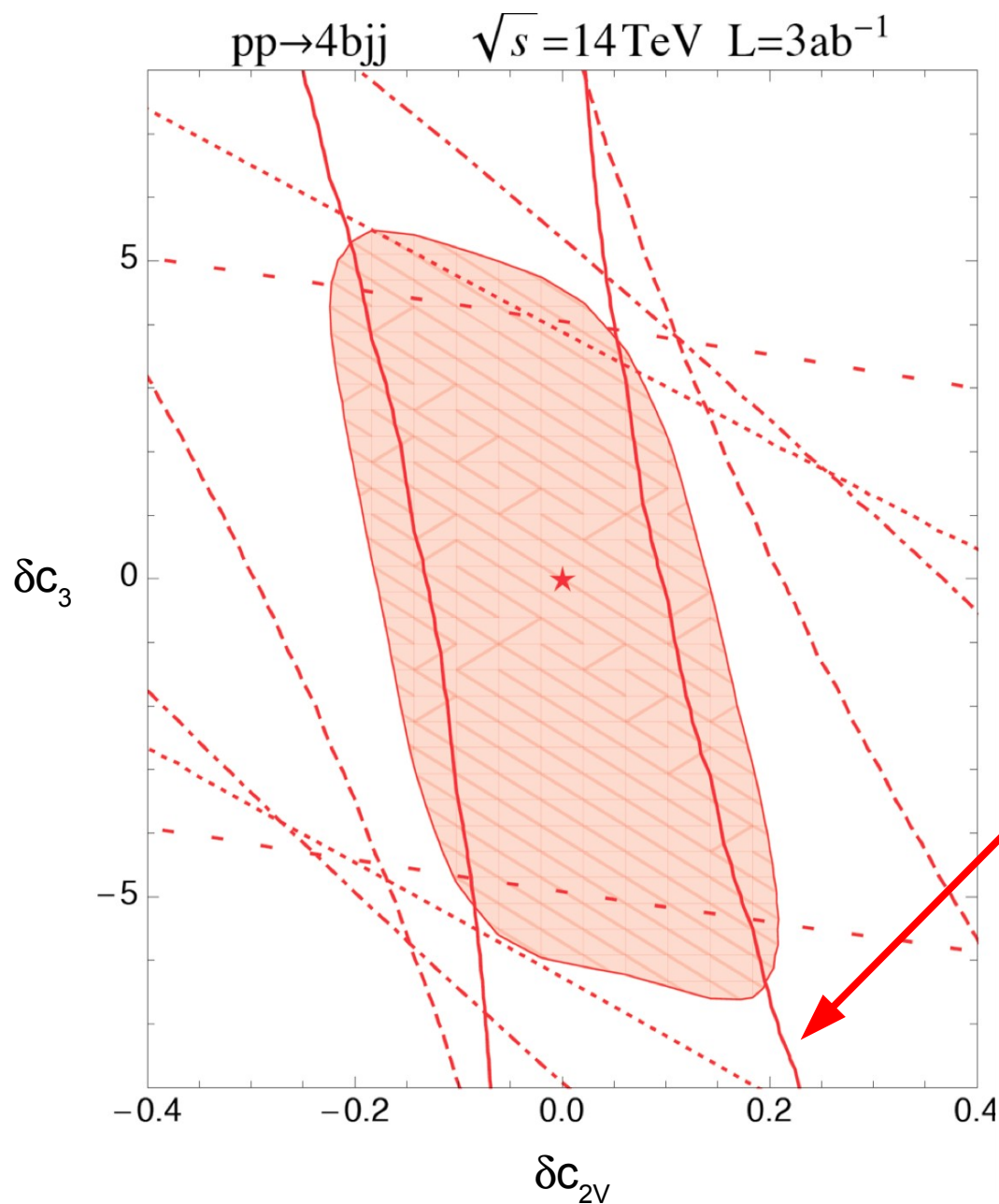
- Scale-invariant resonance tagging, which provides a smooth matching between boosted and resolved kinematics [6]

variable
 M_{HH}





Non-resonant VBF di-Higgs search at hadron level



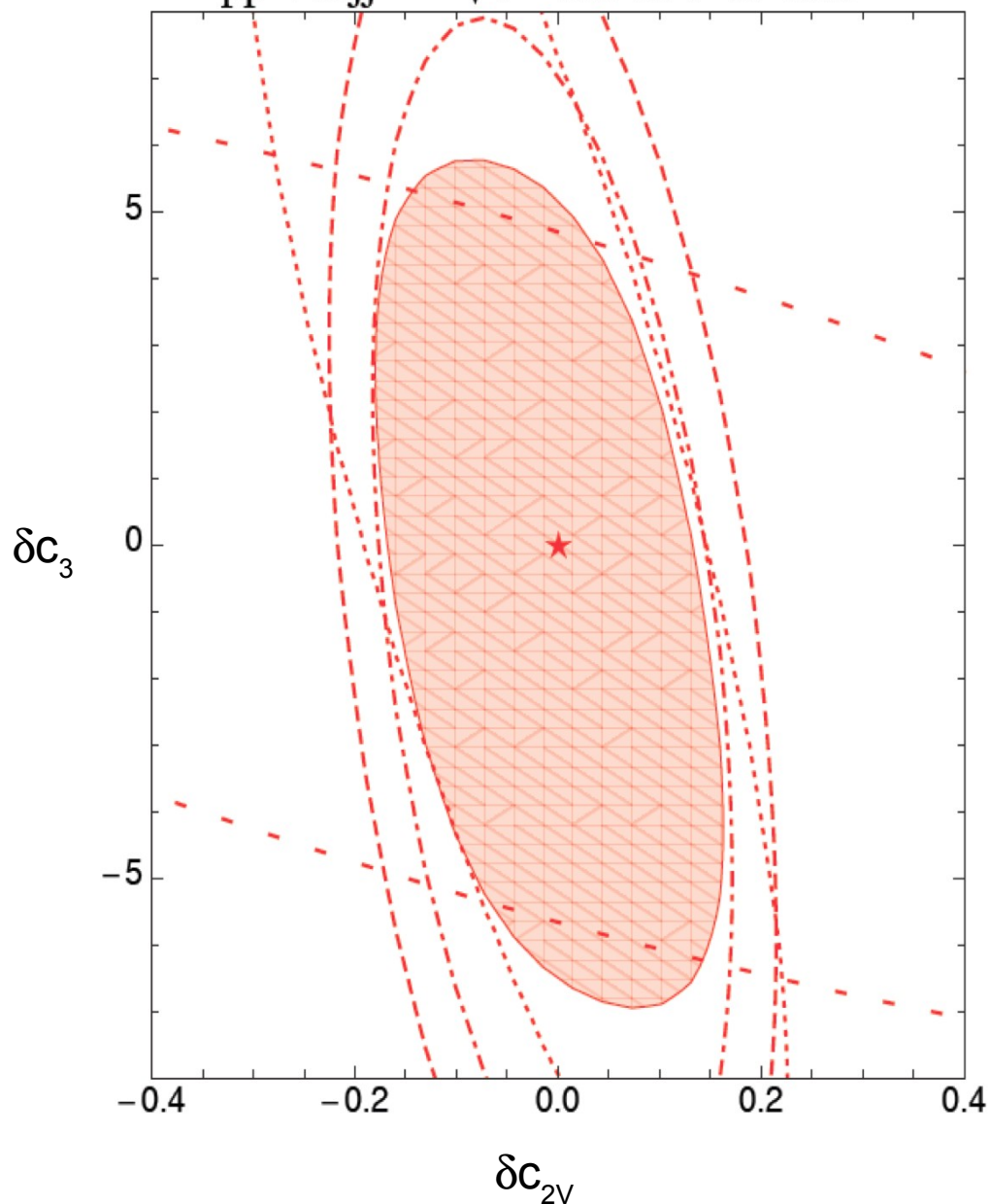
4b + 2j

- At HL-LHC deviations from the SM value of c_{2V} can be probed at the level of 20%
- The most stringent constraints come from the boosted region with $M_{HH} > 1.5\text{ TeV}$, where jet substructure is crucial
- The sensitivity to c_3 is worse than in $gg\rightarrow HH$ since in the threshold region the backgrounds are much larger than signal



Non-resonant VBF di-Higgs search at hadron level

$pp \rightarrow 4bjj$ $\sqrt{s} = 100 \text{ TeV}$ $L = 3 \text{ ab}^{-1}$



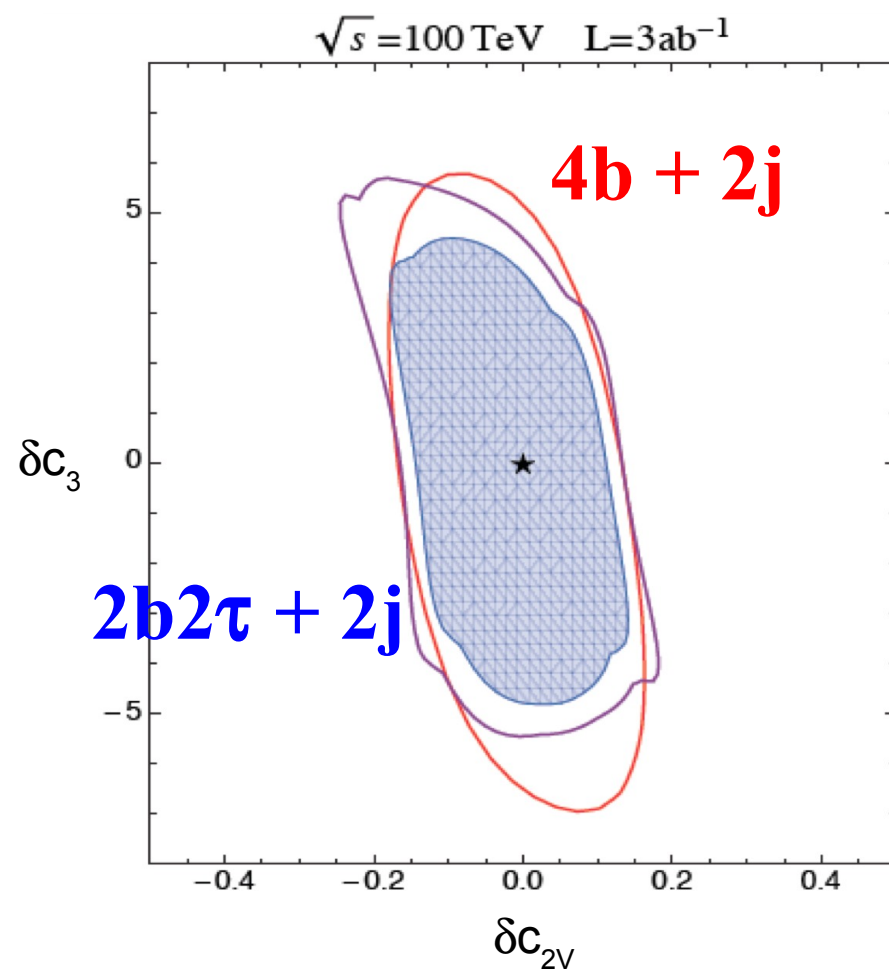
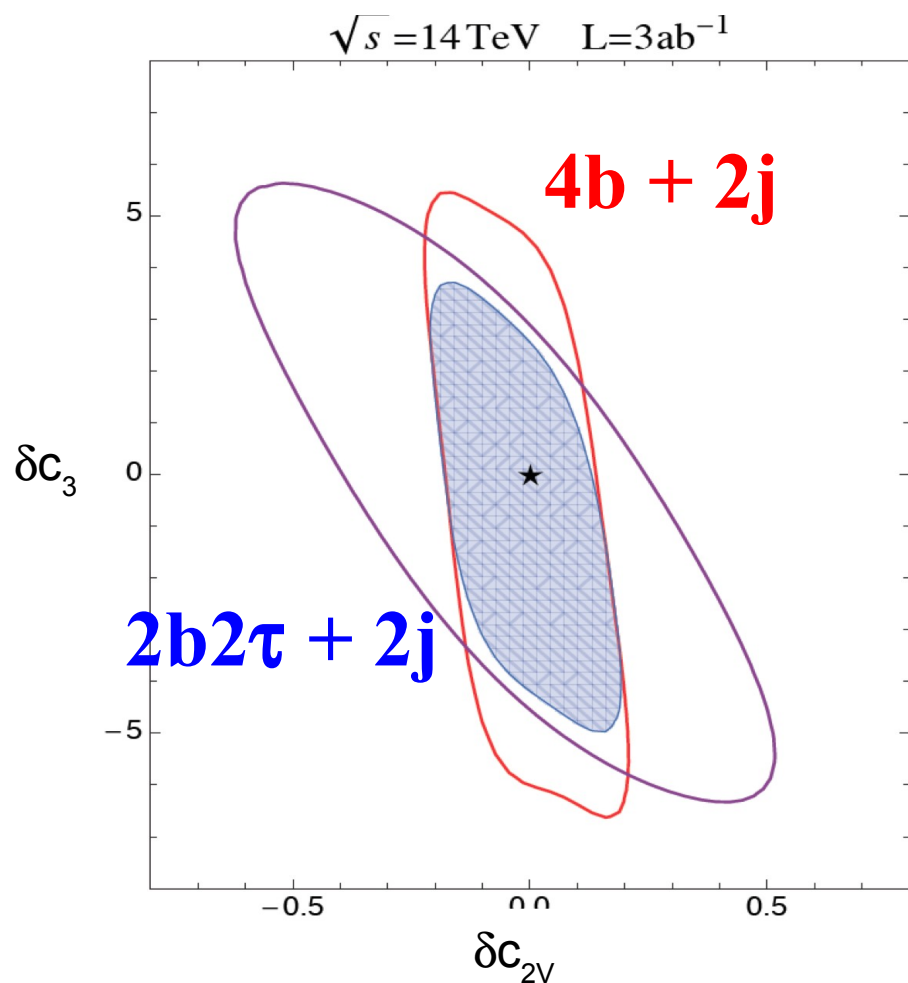
4b + 2j

- Sensitivity at 100 TeV is only slightly better than at 14 TeV: increase in signal rates compensated by stronger growth of the QCD background



Non-resonant VBF di-Higgs search at hadron level

- HL-LHC constraints from the $2b2\tau2j$ final state are less severe on c_{2V} than those from $4b$:
 - deviations from the SM down to 40-50%
 - reduced sensitivity comes from the small number of signal events
 - At 100 TeV the sensitivity on the $2b2\tau2j$ final state is improved thanks to the increase in signal cross-section





Conclusions

- **Vector Boson Fusion di-Higgs production**

- Goal of the analysis → new physics at LHC @ 13 TeV

- **Direct search for new physics**

- Resonant di-Higgs production
- KK-graviton as a benchmark
- VBF $G \rightarrow HH \rightarrow bbbb$ final state considered
 - High signal yields

- **Indirect search for new physics**

- Strong Double Higgs production
- Sensitivity to c_{2V} (HHVV vertex) → only accessible via VBF HH

- **Parton level analysis:**

- VBF $HH \rightarrow bbbb / WW(l\nu l\nu)bb$ final state considered

- **Hadron level analysis:**

- VBF $HH \rightarrow bbbb / \tau\tau bb$ final state considered
- ~ 20% sensitivity to c_{2V} with HL-LHC



backup



Strong double Higgs production

- Cross section for signal with different coupling modification and different final states

Model	Final state	Cross section [fb]	$N_{\text{ev}} (\mathcal{L} = 3 \text{ ab}^{-1})$
SM (no cut)	$hhjj$	0.83	2500
SM	$hhjj$	0.12	360
SM	$hhjj \rightarrow 4bjj$	0.049	150
$c_V = 0.5$		0.54	1600
$c_V = 1.5$		2.72	8100
$c_{2V} = 0$		1.23	3700
$c_{2V} = 2$		0.78	2300
$c_3 = 0$		0.14	420
$c_3 = 2$		0.042	130
SM	$hhjj \rightarrow l^+l^- \cancel{E}_T 2bjj$	$8.6 \cdot 10^{-4}$	2.6
$c_V = 0.5$		$2.0 \cdot 10^{-3}$	6
$c_V = 1.5$		$9.8 \cdot 10^{-2}$	290
$c_{2V} = 0$		$1.9 \cdot 10^{-2}$	54
$c_{2V} = 2$		$1.1 \cdot 10^{-2}$	33
$c_3 = 0$		$2.4 \cdot 10^{-3}$	7
$c_3 = 2$		$7.4 \cdot 10^{-4}$	2.2



Transverse mass $H \rightarrow WW \rightarrow l\nu l\nu$

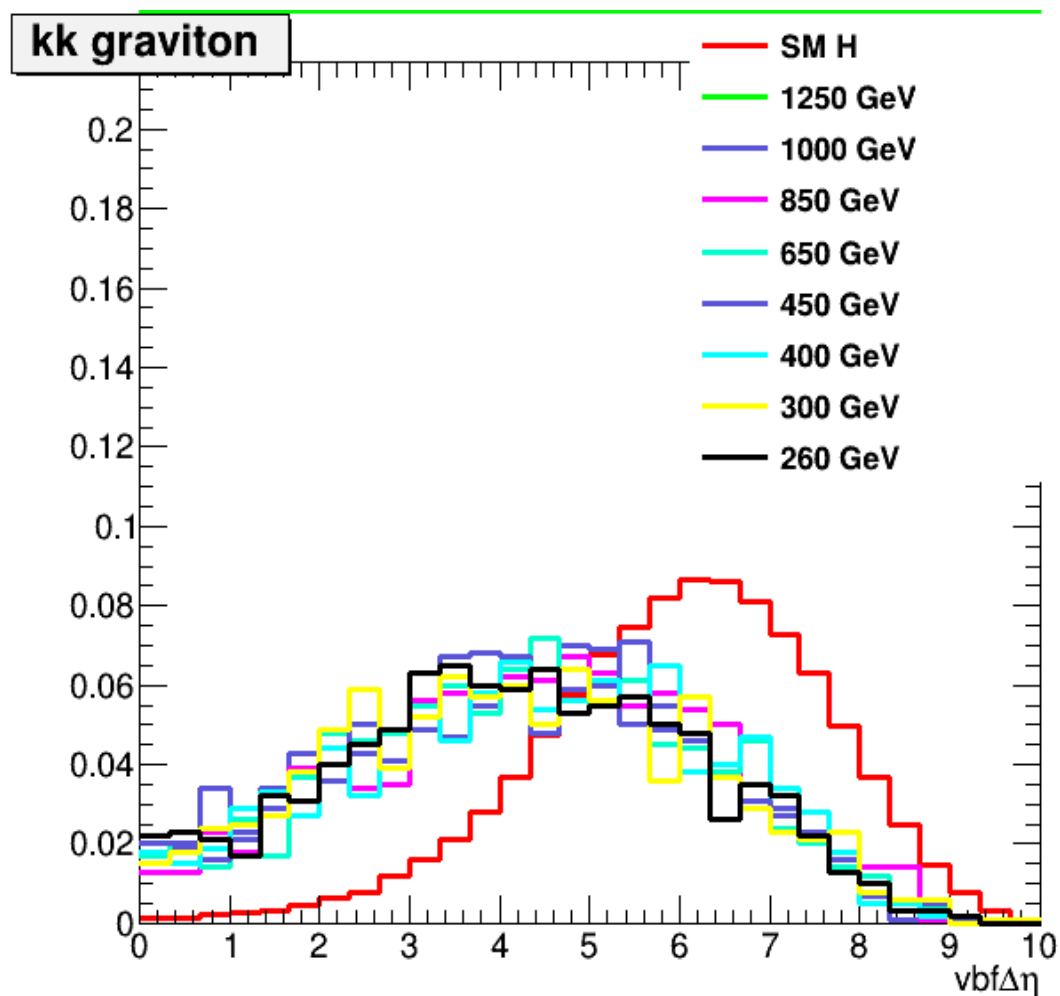
- Transverse mass used in $H \rightarrow WW \rightarrow l\nu l\nu$ analysis
 - Due to lack of information from neutrinos, invariant mass cannot be reconstructed
 - Transverse mass under the hypothesis $m_{ll} \sim m_{\nu\nu}$

$$m_T(WW) \equiv \left(\left(\sqrt{m_{ll}^2 + |\vec{p}_{Tl}|^2} + \sqrt{m_{ll}^2 + |\vec{p}_{Tmiss}|^2} \right)^2 - |\vec{p}_{Tl} + \vec{p}_{Tmiss}|^2 \right)^{1/2}$$



KK-Graviton

- KK-graviton $\Delta\eta_{jj}$ distribution
- KK-graviton vertices with vector bosons implies no high- $\Delta\eta_{jj}$ distribution
- Related spin 2 resonances





KK-Graviton efficiencies

• Cut flow efficiencies

Sample	basic cuts (eqs. 7,6)	jet merging (akt5)	$M_{jj} > 400 \text{ GeV}$	$M_{HH} > 250 \text{ GeV}$	$p_T^{HH} > 60 \text{ GeV}$	$\Delta\eta_{HH} < 2$
1450 GeV	0.53	0.49	0.44	0.44	0.41	0.35
1250 GeV	0.52	0.47	0.43	0.43	0.40	0.35
1050 GeV	0.52	0.48	0.44	0.44	0.40	0.35
850 GeV	0.51	0.47	0.42	0.42	0.39	0.34
650 GeV	0.51	0.43	0.39	0.39	0.36	0.33
450 GeV	0.54	0.38	0.34	0.34	0.32	0.30
400 GeV	0.55	0.33	0.29	0.29	0.28	0.26
300 GeV	0.58	0.17	0.14	0.14	0.13	0.12
260 GeV	0.60	0.09	0.07	0.07	0.07	0.07
SM $H(b\bar{b})H(b\bar{b})jj$	0.41	0.38	0.36	0.36	0.30	0.12
$Z(b\bar{b})b\bar{b}jj$	0.50	0.36	0.14	0.10	7.91E-02	4.55E-02
$Z(b\bar{b})Z(b\bar{b})jj$	0.62	0.51	0.17	0.12	9.66E-02	6.61E-02
$b\bar{b}b\bar{b}jj$	0.70	0.20	0.11	6.73E-02	5.49E-02	4.55E-02



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