Space for New Physics in Neutral B mixing observables

Gilberto Tetlalmatzi

IPPP Durham University gilberto.tetlalmatzi-xolocotz@durham.ac.uk

April 30, 2014

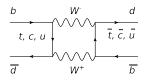
$$B_d = \{\bar{b}, d\}$$

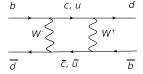
$$B_d = \{\bar{b}, d\}$$
 $\bar{B_d} = \{b, \bar{d}\}$

$$\begin{split} B_d &= \{\bar{b}, d\} \qquad \bar{B_d} = \{b, \bar{d}\} \\ i\frac{d}{dt} \begin{pmatrix} \left|B_d\right\rangle \\ \bar{B}_d \end{pmatrix} \end{pmatrix} &= \Sigma^d \begin{pmatrix} \left|B_d\right\rangle \\ \bar{B}_d \end{pmatrix} \end{split}$$

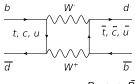
$$\begin{split} B_d &= \{\bar{b}, d\} \qquad \bar{B_d} = \{b, \bar{d}\} \\ i \frac{d}{dt} \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{pmatrix} &= \Sigma^d \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{split}$$

$$\begin{split} B_d &= \{\bar{b}, d\} \qquad \bar{B_d} = \{b, \bar{d}\} \\ i \frac{d}{dt} \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{pmatrix} &= \Sigma^d \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{pmatrix} \end{split}$$

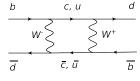




$$\begin{split} B_d &= \{\bar{b}, d\} \qquad \bar{B_d} = \{b, \bar{d}\} \\ i \frac{d}{dt} \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{pmatrix} &= \Sigma^d \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{split}$$



$$B_d \Longleftrightarrow \bar{B_d}$$



$$\begin{split} B_d &= \{\bar{b}, d\} \qquad \bar{B_d} = \{b, \bar{d}\} \\ i \frac{d}{dt} \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{pmatrix} &= \Sigma^d \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{pmatrix} \end{split}$$

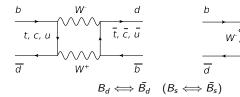
 $\Sigma^q = M^q - \frac{i}{2}\Gamma^q$; M^q and Γ^q are hermitian matrices.

c, u

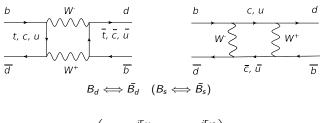
ō, ū

d

h



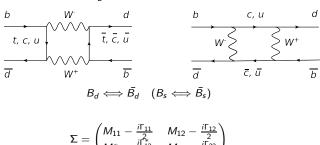
$$\begin{split} B_d &= \{\bar{b}, d\} \qquad \bar{B_d} = \{b, \bar{d}\} \\ i \frac{d}{dt} \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{pmatrix} &= \Sigma^d \begin{pmatrix} \left| B_d \right\rangle \\ \bar{B_d} \end{pmatrix} \end{split}$$



$$\Sigma = \begin{pmatrix} M_{11} - \frac{i\Gamma_{11}}{2} & M_{12} - \frac{i\Gamma_{12}}{2} \\ M_{12}^* - \frac{i\Gamma_{12}}{2} & M_{22} - \frac{i\Gamma_{22}}{2} \end{pmatrix}$$

$$\begin{split} B_d &= \{\bar{b}, d\} \qquad \bar{B}_d = \{b, \bar{d}\} \\ i \frac{d}{dt} \begin{pmatrix} \begin{vmatrix} B_d \\ \bar{B}_d \end{pmatrix} \end{pmatrix} &= \Sigma^d \begin{pmatrix} \begin{vmatrix} B_d \\ \bar{B}_d \end{pmatrix} \end{pmatrix} \end{split}$$

 $\Sigma^q = M^q - \frac{i}{2}\Gamma^q$; M^q and Γ^q are hermitian matrices.



$$\Sigma = \begin{pmatrix} M_{11} - \frac{i\Gamma_{11}}{2} & M_{12} - \frac{i\Gamma_{12}}{2} \\ M_{12}^* - \frac{i\Gamma_{12}}{2} & M_{22} - \frac{i\Gamma_{22}}{2} \end{pmatrix}$$

 Γ_{12} On-shell M_{12} Off-shell

$$\lambda_L = M_L - \frac{i}{2}\Gamma_L$$
 $\lambda_H = M_H - \frac{i}{2}\Gamma_H$

$$\lambda_L = M_L - \frac{i}{2}\Gamma_L$$
 $\lambda_H = M_H - \frac{i}{2}\Gamma_H$ $\Delta M = M_H - M_L$

$$\lambda_{L} = M_{L} - \frac{i}{2}\Gamma_{L} \qquad \lambda_{H} = M_{H} - \frac{i}{2}\Gamma_{H}$$

$$\Delta M = M_{H} - M_{L}$$

$$\Delta \Gamma = \Gamma_{L} - \Gamma_{H}$$

$$\lambda_{L} = M_{L} - \frac{i}{2}\Gamma_{L} \qquad \lambda_{H} = M_{H} - \frac{i}{2}\Gamma_{H}$$

$$\Delta M = M_{H} - M_{L}$$

$$\Delta \Gamma = \Gamma_{L} - \Gamma_{H}$$

$$\phi \equiv \arg\left(-\frac{M_{12}}{\Gamma_{12}}\right)$$

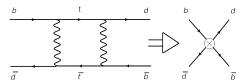
$$\lambda_{L} = M_{L} - \frac{i}{2}\Gamma_{L} \qquad \lambda_{H} = M_{H} - \frac{i}{2}\Gamma_{H}$$

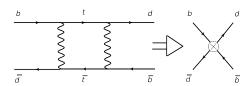
$$\Delta M = M_{H} - M_{L}$$

$$\Delta \Gamma = \Gamma_{L} - \Gamma_{H}$$

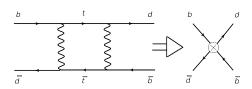
$$\phi \equiv \arg\left(-\frac{M_{12}}{\Gamma_{12}}\right)$$

$$\begin{array}{lcl} \Delta M & \approx & 2|M_{12}| \\ \Delta \Gamma & \approx & 2|\Gamma_{12}|\cos(\phi) \\ \\ a_{sl} & = & \left|\frac{\Gamma_{12}}{M_{12}}\right|\sin(\phi) \end{array}$$



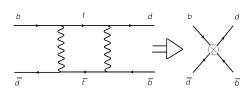


$$\hat{Q}^{|\Delta B|=2} \quad = \quad \left(\bar{d}b\right)_{V-A} \left(\bar{d}b\right)_{V-A}$$



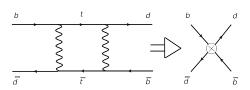
$$\hat{Q}^{|\Delta B|=2} = (\bar{d}b)_{V-A} (\bar{d}b)_{V-A}$$

$$\hat{H}^{|\Delta B|=2} = \frac{G_F^2}{4\pi^2} (V_{tb}V_{tq}^*)^2 C^{|\Delta B|=2}(m_t, M_W, \mu) \hat{Q}^{|\Delta B|=2}(\mu) + h.c.$$



$$\begin{split} \hat{Q}^{|\Delta B|=2} &= \left(\bar{d}b\right)_{V-A} \left(\bar{d}b\right)_{V-A} \\ \hat{H}^{|\Delta B|=2} &= \frac{G_F^2}{4\pi^2} \left(V_{tb}V_{tq}^*\right)^2 C^{|\Delta B|=2} (m_t, M_W, \mu) \hat{Q}^{|\Delta B|=2} (\mu) + h.c. \\ \Delta M_d &= \left|\frac{\langle B_d | \hat{H}^{|\Delta B|=2} | \bar{B}_d \rangle}{M_{B_d}} \right| = \frac{G_F^2}{6\pi^2} \eta_B B_{B_d} f_{B_d}^2 M_W^2 S_0(x_t) (V_{tb}V_{td}^*)^2 \end{split}$$

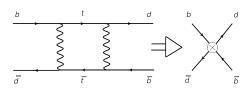
Calculation of ΔM_q



$$\begin{split} \hat{Q}^{|\Delta B|=2} &= \left(\bar{d}b\right)_{V-A} \left(\bar{d}b\right)_{V-A} \\ \hat{H}^{|\Delta B|=2} &= \frac{G_F^2}{4\pi^2} \left(V_{tb}V_{tq}^*\right)^2 C^{|\Delta B|=2} (m_t, M_W, \mu) \hat{Q}^{|\Delta B|=2} (\mu) + h.c. \\ \Delta M_d &= \left|\frac{\langle B_d | \hat{H}^{|\Delta B|=2} | \bar{B}_d \rangle}{M_{B_d}} \right| = \frac{G_F^2}{6\pi^2} \eta_B B_{B_d} f_{B_d}^2 M_W^2 S_0(\mathbf{x}_t) \left(V_{tb}V_{td}^*\right)^2 \end{split}$$

 $S_0(x_t)$: Inami-Lim function. Inami and Lim. Prog. Theor. Phys. 65 (1981) 297.

Calculation of ΔM_a



$$\hat{Q}^{|\Delta B|=2} \quad = \quad \left(\bar{d}b\right)_{V-A} \left(\bar{d}b\right)_{V-A}$$

$$\hat{H}^{|\Delta B|=2} = \frac{G_F^2}{4\pi^2} \left(V_{tb} V_{tq}^* \right)^2 C^{|\Delta B|=2} (m_t, M_W, \mu) \hat{Q}^{|\Delta B|=2} (\mu) + h.c.$$

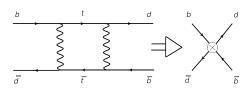
$$\Delta M_d = \left| \frac{\langle B_d | \hat{H}^{|\Delta B|=2} | \bar{B}_d \rangle}{M_{B_d}} \right| = \frac{G_F^2}{6\pi^2} \eta_B B_{B_d} f_{B_d}^2 M_W^2 S_0(x_t) (V_{tb} V_{td}^*)^2$$

$$\Delta M_d = |\frac{\langle B_d | H^{|\Delta B|=2} | B_d \rangle}{M_{B_d}}| = \frac{G_F^2}{6\pi^2} \eta_B B_{B_d} f_{B_d}^2 M_W^2 S_0(x_t) (V_{tb} V_{td}^*)^2$$

 $S_0(x_t)$: Inami-Lim function. Inami and Lim. Prog. Theor. Phys. 65 (1981) 297.

QCD corrections to the box diagram. Buras et al., Nucl. Phys. B 347 (1990) 491. η_{R}

Calculation of ΔM_a



$$\hat{Q}^{|\Delta B|=2} = (\bar{d}b)_{V-A} (\bar{d}b)_{V-A}$$

$$\begin{split} \hat{H}^{|\Delta B|=2} &= \frac{G_F^2}{4\pi^2} \left(V_{tb}V_{tq}^*\right)^2 C^{|\Delta B|=2}(m_t, M_W, \mu) \hat{Q}^{|\Delta B|=2}(\mu) + h.c. \\ \Delta M_d &= |\frac{\langle B_d | \hat{H}^{|\Delta B|=2} | \bar{B}_d \rangle}{M_{B_d}} | = \frac{G_F^2}{6\pi^2} \eta_B B_{B_d} f_{B_d}^2 M_W^2 S_0(x_t) (V_{tb}V_{td}^*)^2 \end{split}$$

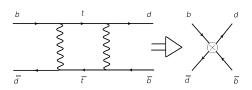
$$\Delta M_d = |\frac{\langle B_d | H^{|\Delta B|=2} | B_d \rangle}{M_{B_d}}| = \frac{G_F^2}{6\pi^2} \eta_B B_{B_d} f_{B_d}^2 M_W^2 S_0(x_t) (V_{tb} V_{td}^*)^2$$

 $S_0(x_t)$ Inami-Lim function. Inami and Lim. Prog. Theor. Phys. 65 (1981) 297.

QCD corrections to the box diagram. Buras et al., Nucl. Phys. B 347 (1990) 491. η_{R}

 f_{B_d} (190.5 ± 4.2) MeV. Aoki et al. ArXiv:1310.8555.

Calculation of ΔM_a



$$\hat{Q}^{|\Delta B|=2} \quad = \quad \left(\bar{d}b\right)_{V-A} \left(\bar{d}b\right)_{V-A}$$

$$\begin{split} \hat{H}^{|\Delta B|=2} &= \frac{G_F^2}{4\pi^2} \left(V_{tb}V_{tq}^*\right)^2 C^{|\Delta B|=2}(m_t, M_W, \mu) \hat{Q}^{|\Delta B|=2}(\mu) + h.c. \\ \Delta M_d &= |\frac{\langle B_d | \hat{H}^{|\Delta B|=2} | \bar{B}_d \rangle}{M_{B_d}} | = \frac{G_F^2}{6\pi^2} \eta_B B_{B_d} f_{B_d}^2 M_W^2 S_0(x_t) (V_{tb}V_{td}^*)^2 \end{split}$$

$$\Delta M_d = |\frac{\langle B_d | H^{|\Delta B|=2} | B_d \rangle}{M_{B_d}}| = \frac{G_F^2}{6\pi^2} \eta_B B_{B_d} f_{B_d}^2 M_W^2 S_0(x_t) (V_{tb} V_{td}^*)^2$$

 $S_0(x_t)$ Inami-Lim function. Inami and Lim. Prog. Theor. Phys. 65 (1981) 297.

QCD corrections to the box diagram. Buras et al., Nucl. Phys. B 347 (1990) 491. η_{R}

 f_{B_d} = (190.5 ± 4.2) MeV. Aoki et al. ArXiv:1310.8555.

 (0.84 ± 0.07) GeV. Aoki et al. ArXiv:1310.8555. B_{B_d}

Theoretical predictions vs experimental results for the mass splittings:

Theoretical predictions vs experimental results for the mass splittings:

$$\Delta M_d^{Theo} = (0.543 \pm 0.091) \, ps^{-1}$$
. Lenz and Nierste, arXiv:1102.4274. $\Delta M_d^{Exp} = (0.510 \pm 0.004) \, ps^{-1}$. Online update of HFAG.

$$\Delta M_d^{\text{Exp}} = (0.510 \pm 0.004) \, ps^{-1}$$
. Online update of HFAG

Theoretical predictions vs experimental results for the mass splittings:

$$\Delta M_d^{Theo} = (0.543 \pm 0.091) \, ps^{-1}$$
. Lenz and Nierste, arXiv:1102.4274. $\Delta M_d^{Exp} = (0.510 \pm 0.004) \, ps^{-1}$. Online update of HFAG.

$$\Delta M_s^{Theo} = (17.3 \pm 2.6) ps^{-1}$$
. Lenz and Nierste, arXiv:1102.4274.
 $\Delta M_s^{Exp} = (17.69 \pm 0.08) ps^{-1}$. Online update of HFAG.

Theoretical predictions vs experimental results for the mass splittings:

$$\Delta M_d^{Theo} = (0.543 \pm 0.091) \, ps^{-1}$$
. Lenz and Nierste, arXiv:1102.4274. $\Delta M_d^{Exp} = (0.510 \pm 0.004) \, ps^{-1}$. Online update of HFAG.

$$\Delta M_s^{Theo} = (17.3 \pm 2.6) \, ps^{-1}$$
. Lenz and Nierste, arXiv:1102.4274. $\Delta M_s^{Exp} = (17.69 \pm 0.08) \, ps^{-1}$. Online update of HFAG.

The main source of the theoretical uncertainty comes from the hadronic matrix elements.

Theoretical predictions vs experimental results for the mass splittings:

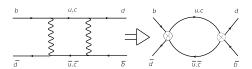
$$\Delta M_d^{Theo} = (0.543 \pm 0.091) \, ps^{-1}$$
. Lenz and Nierste, arXiv:1102.4274. $\Delta M_d^{Exp} = (0.510 \pm 0.004) \, ps^{-1}$. Online update of HFAG.

$$\Delta M_s^{Theo} = (17.3 \pm 2.6) \, ps^{-1}$$
. Lenz and Nierste, arXiv:1102.4274. $\Delta M_s^{Exp} = (17.69 \pm 0.08) \, ps^{-1}$. Online update of HFAG.

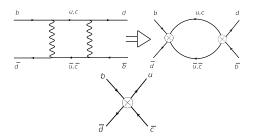
The main source of the theoretical uncertainty comes from the hadronic matrix elements.

The hadronic uncertainties allow space for new physics.

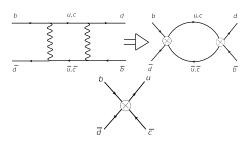
$\Delta \Gamma_{d,s}$



$\Delta\Gamma_{d,s}$

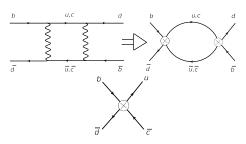


$\Delta\Gamma_{d,s}$



$$H_{eff}^{\Delta B=1} = \frac{4G_F}{\sqrt{2}} \sum_{q,q'=u,c} \lambda_{qq'} \sum_{i=1,2} C_i^{q,q'} (M_W,\mu) Q_i^{qq'} + h.c.$$

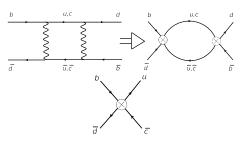
$\Delta\Gamma_{d,s}$



$$H_{\mathrm{eff}}^{\Delta B=1} = \frac{4G_F}{\sqrt{2}} \sum_{q,q'=u,c} \lambda_{qq'} \sum_{i=1,2} C_i^{q,q'} (M_W,\mu) Q_i^{qq'} + h.c. \label{eq:eff_def}$$

$$\Gamma^{d}_{12} \quad = \quad \frac{1}{2M_{B_d}} < \bar{B_d} | \text{Im} \left(i \int d^4 x \, \hat{T} \left[H_{\text{eff}}^{\Delta B=1}(x) H_{\text{eff}}^{\Delta B=1}(0) \right] \right) | B_d >$$

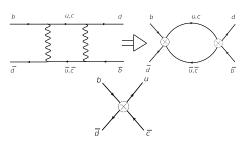
$\Delta \Gamma_{d,s}$



$$H_{eff}^{\Delta B=1} = \frac{4G_F}{\sqrt{2}} \sum_{q,q'=u,c} \lambda_{qq'} \sum_{i=1,2} C_i^{q,q'} (M_W,\mu) Q_i^{qq'} + h.c.$$

$$\begin{array}{rcl} \Gamma^d_{12} & = & \frac{1}{2M_{B_d}} < \bar{B_d} | \mbox{Im} \left(i \int d^4 x \, \hat{T} \left[H_{\rm eff}^{\Delta B=1}(x) H_{\rm eff}^{\Delta B=1}(0) \right] \right) | B_d > \\ \\ \phi_d & \equiv & \mbox{arg} \left(- \frac{M_{12}}{\Gamma_{12}} \right) \end{array}$$

Calculation of $\Delta\Gamma_{d,s}$

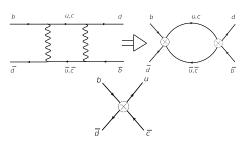


$$H_{eff}^{\Delta B=1} = rac{4G_F}{\sqrt{2}} \sum_{q,q'=u,c} \lambda_{qq'} \sum_{i=1,2} C_i^{q,q'} (M_W,\mu) Q_i^{qq'} + h.c.$$

$$\begin{array}{rcl} \Gamma^d_{12} & = & \frac{1}{2M_{B_d}} < \bar{B_d} | \mathit{Im} \left(i \int d^4x \, \hat{T} \left[H_{eff}^{\Delta B=1}(x) H_{eff}^{\Delta B=1}(0) \right] \right) | B_d > \\ \\ \phi_d & \equiv & \mathit{arg} \left(- \frac{M_{12}}{\Gamma_{12}} \right) \end{array}$$

$$\Delta\Gamma_d \approx 2|\Gamma_{12}^d|cos(\phi_d)$$

Calculation of $\Delta\Gamma_{d,s}$



$$H_{eff}^{\Delta B=1} = rac{4G_F}{\sqrt{2}} \sum_{q,q'=u,c} \lambda_{qq'} \sum_{i=1,2} C_i^{q,q'} (M_W,\mu) Q_i^{qq'} + h.c.$$

$$\begin{array}{rcl} \Gamma^d_{12} & = & \frac{1}{2M_{B_d}} < \bar{B_d} | \mathit{Im} \left(i \int d^4x \, \hat{T} \left[H_{eff}^{\Delta B=1}(x) H_{eff}^{\Delta B=1}(0) \right] \right) | B_d > \\ \\ \phi_d & \equiv & \mathit{arg} \left(- \frac{M_{12}}{\Gamma_{12}} \right) \end{array}$$

$$\Delta\Gamma_d \approx 2|\Gamma_{12}^d|cos(\phi_d)$$

Experimental results vs theoretical prediction for $\Delta\Gamma_s$:

Experimental results vs theoretical prediction for $\Delta\Gamma_s$:

$$\Delta\Gamma_s^{HFAG} = (0.081 \pm 0.011) \, ps^{-1} (\text{LHCb}(2013), \, \text{ATLAS}(2012), \, \text{CDF} \, (2012) \, \text{and D0} \, (2012)).$$

Experimental results vs theoretical prediction for $\Delta\Gamma_s$:

$$\Delta\Gamma_s^{HFAG} = (0.081 \pm 0.011) \, ps^{-1} (\text{LHCb}(2013), \, \text{ATLAS}(2012), \, \text{CDF} \, (2012) \, \text{and D0} \, (2012)).$$

$$\Delta\Gamma_s^{Theo} = (0.087 \pm 0.021) \, ps^{-1} (\text{A. Lenz and U. Nierste, arXiv:1102.4274}).$$

Experimental results vs theoretical prediction for $\Delta\Gamma_s$:

```
\Delta\Gamma_s^{HFAG} = (0.081 \pm 0.011) \, ps^{-1} (\text{LHCb}(2013), \, \text{ATLAS}(2012), \, \text{CDF} \, (2012) \, \text{and D0} \, (2012)). \Delta\Gamma_s^{Theo} = (0.087 \pm 0.021) \, ps^{-1} (\text{A. Lenz and U. Nierste, arXiv:1102.4274}).
```

Experimental results vs theoretical prediction for $\Delta\Gamma_d$:

Experimental results vs theoretical prediction for $\Delta\Gamma_s$:

$$\Delta \Gamma_s^{HFAG} = (0.081 \pm 0.011) \, ps^{-1} (\text{LHCb}(2013), \, \text{ATLAS}(2012), \, \text{CDF} \, (2012) \, \text{and D0} \, (2012)).$$

$$\Delta \Gamma_s^{Theo} = (0.087 \pm 0.021) \, ps^{-1} (\text{A. Lenz and U. Nierste, arXiv:} 1102.4274).$$

Experimental results vs theoretical prediction for $\Delta\Gamma_d$:

$$\begin{array}{lll} \frac{\Delta \Gamma_d^{\textit{HFAG}}}{\Gamma_d} & = & (1.5 \pm 1.8)\% (\text{BABAR}(2006) \text{ and Belle}(2012)). \\ \\ \frac{\Delta \Gamma_d^{D0}}{\Gamma_d} & = & (0.50 \pm 1.38)\% (2014). \\ \\ \frac{\Delta \Gamma_d^{\textit{LHCb}}}{\Gamma_d} & = & (-4.4 \pm 2.7)\% (2014). \end{array}$$

Experimental results vs theoretical prediction for $\Delta\Gamma_s$:

$$\Delta\Gamma_s^{HFAG} = (0.081 \pm 0.011) \, ps^{-1} (\text{LHCb}(2013), \, \text{ATLAS}(2012), \, \text{CDF} \, (2012) \, \text{and D0} \, (2012)).$$

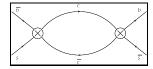
$$\Delta\Gamma_s^{Theo} = (0.087 \pm 0.021) \, ps^{-1} (\text{A. Lenz and U. Nierste, arXiv:} 1102.4274).$$

Experimental results vs theoretical prediction for $\Delta\Gamma_d$:

$$\begin{array}{lll} \frac{\Delta \Gamma_d^{HFAG}}{\Gamma_d} & = & (1.5 \pm 1.8)\% (\text{BABAR}(2006) \text{ and Belle}(2012)). \\ \frac{\Delta \Gamma_d^{D0}}{\Gamma_d} & = & (0.50 \pm 1.38)\% (2014). \\ \frac{\Delta \Gamma_d^{LHCb}}{\Gamma_d} & = & (-4.4 \pm 2.7)\% (2014). \\ \frac{\Delta \Gamma_d^{Theo}}{\Gamma_d} & = & (0.42 \pm 0.08)\% (\text{A. Lenz and U. Nierste, arXiv:} 1102.4274). \end{array}$$

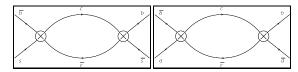
New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:



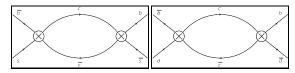
 $\Delta \Gamma_s$ triggered by b o c ar c s

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:



 $\Delta \Gamma_s$ triggered by b o c ar c s $\Delta \Gamma_d$ triggered by b o c ar c d

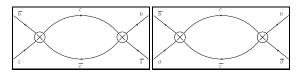
New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:



 $\Delta\Gamma_s$ triggered by $b o c\bar{c}s$ $\Delta\Gamma_d$ triggered by $b o c\bar{c}d$

$$Br(b \to c\bar{c}s) = (23.7 \pm 1.3)\% \mid\mid Br(b \to c\bar{c}d) = (1.31 \pm 0.07)\%$$

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:

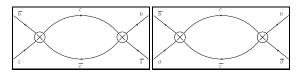


 $\Delta\Gamma_s$ triggered by $b \to c\bar{c}s$ $\Delta\Gamma_d$ triggered by $b \to c\bar{c}d$

$$Br(b \to c\bar{c}s) = (23.7 \pm 1.3)\% \mid\mid Br(b \to c\bar{c}d) = (1.31 \pm 0.07)\%$$

 \implies 100% enhancement on $\Gamma(b \to c \bar c s)$ leads to sizable effect on Γ_{tot}

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:

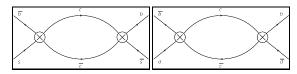


 $\Delta\Gamma_s$ triggered by $b o c\bar{c}s$ $\Delta\Gamma_d$ triggered by $b o c\bar{c}d$

$$Br(b \to c\bar{c}s) = (23.7 \pm 1.3)\% \mid\mid Br(b \to c\bar{c}d) = (1.31 \pm 0.07)\%$$

- \implies 100% enhancement on $\Gamma(b \to c \bar c s)$ leads to sizable effect on Γ_{tot}
- \implies 100% enhancement on $\Gamma(b \to c\bar{c}d)$ hidden in the hadronic uncertainties.

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:



 $\Delta\Gamma_s$ triggered by $b \to c\bar{c}s$ $\Delta\Gamma_d$ triggered by $b \to c\bar{c}d$

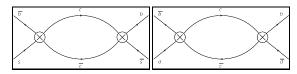
$$Br(b \to c\bar{c}s) = (23.7 \pm 1.3)\% \mid\mid Br(b \to c\bar{c}d) = (1.31 \pm 0.07)\%$$

 \Rightarrow 100% enhancement on $\Gamma(b \to c\bar{c}s)$ leads to sizable effect on Γ_{tot}

100% enhancement on $\Gamma(b \to c \bar c d)$ hidden in the hadronic uncertainties.

Enhancements in $\Delta\Gamma_d$ arise from:

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:



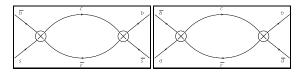
 $\Delta\Gamma_s$ triggered by $b \to c\bar{c}s$ $\Delta\Gamma_d$ triggered by $b \to c\bar{c}d$

$$Br(b \to c\bar{c}s) = (23.7 \pm 1.3)\% \mid\mid Br(b \to c\bar{c}d) = (1.31 \pm 0.07)\%$$

- \implies 100% enhancement on $\Gamma(b \to c\bar{c}s)$ leads to sizable effect on Γ_{tot}
- \implies 100% enhancement on $\Gamma(b \to c\bar{c}d)$ hidden in the hadronic uncertainties.

Enhancements in $\Delta\Gamma_d$ arise from:

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:

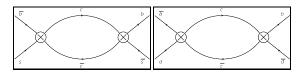


 $\Delta\Gamma_s$ triggered by $b \to c\bar{c}s$ $\Delta\Gamma_d$ triggered by $b \to c\bar{c}d$

$$Br(b \to c\bar{c}s) = (23.7 \pm 1.3)\% \mid\mid Br(b \to c\bar{c}d) = (1.31 \pm 0.07)\%$$

- \implies 100% enhancement on $\Gamma(b \to c\bar{c}s)$ leads to sizable effect on Γ_{tot} \implies 100% enhancement on $\Gamma(b \to c\bar{c}d)$ hidden in the hadronic uncertainties.
 - Enhancements in $\Delta\Gamma_d$ arise from:
- CKM Unitarity violations.
- New Physics at SM tree level decays.

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:

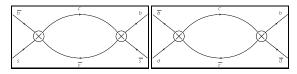


 $\Delta\Gamma_s$ triggered by $b \to c\bar{c}s$ $\Delta\Gamma_d$ triggered by $b \to c\bar{c}d$

$$Br(b \to c\bar{c}s) = (23.7 \pm 1.3)\% \mid\mid Br(b \to c\bar{c}d) = (1.31 \pm 0.07)\%$$

- \implies 100% enhancement on $\Gamma(b \to c\bar{c}s)$ leads to sizable effect on Γ_{tot} \implies 100% enhancement on $\Gamma(b \to c\bar{c}d)$ hidden in the hadronic uncertainties.
 - Enhancements in $\Delta\Gamma_d$ arise from:
- CKM Unitarity violations.
- New Physics at SM tree level decays.
- $(\bar{d}b)(\bar{\tau}\tau)$ operators.

New Physics (NP) effects in $\Delta\Gamma_s$ are strongly constrained in comparison with $\Delta\Gamma_d$ because:



 $\Delta\Gamma_s$ triggered by $b \to c\bar{c}s$ $\Delta\Gamma_d$ triggered by $b \to c\bar{c}d$

$$Br(b \to c\bar{c}s) = (23.7 \pm 1.3)\% \mid\mid Br(b \to c\bar{c}d) = (1.31 \pm 0.07)\%$$

- \implies 100% enhancement on $\Gamma(b \to c\bar{c}s)$ leads to sizable effect on Γ_{tot} \implies 100% enhancement on $\Gamma(b \to c\bar{c}d)$ hidden in the hadronic uncertainties.
 - Enhancements in $\Delta\Gamma_d$ arise from:
- CKM Unitarity violations.
- New Physics at SM tree level decays.
- $(\bar{d}b)(\bar{\tau}\tau)$ operators.

Bobeth, Haisch, Lenz, Pecjak and Tetlalmatzi-Xolocotzi, arXiv:1404.2531.

Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.

Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.
SM: $\lambda_u + \lambda_c + \lambda_t = 0$

Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.
SM: $\lambda_u + \lambda_c + \lambda_t = 0 \Longrightarrow NP$: $\lambda_u + \lambda_c + \lambda_t + \delta_{CKM} = 0$

Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.
SM: $\lambda_u + \lambda_c + \lambda_t = 0 \Longrightarrow \text{NP}$: $\lambda_u + \lambda_c + \lambda_t + \delta_{CKM} = 0$
As a very rough estimate (4th family studies): $\frac{\delta_{CKM}^d}{\lambda_c^d} = \mathcal{O}(1)$, $\frac{\delta_{CKM}^s}{\lambda_c^s} = \mathcal{O}(\lambda)$ with $\lambda \approx 0.23$

Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.
SM: $\lambda_u + \lambda_c + \lambda_t = 0 \Longrightarrow \text{NP}$: $\lambda_u + \lambda_c + \lambda_t + \delta_{CKM} = 0$
As a very rough estimate (4th family studies): $\frac{\delta_{CKM}^d}{\lambda_t^d} = \mathcal{O}(1)$, $\frac{\delta_{CKM}^s}{\lambda_t^s} = \mathcal{O}(\lambda)$ with $\lambda \approx 0.23$
Enhancement factors: 4 for $\Delta \Gamma_d$ and 1.4 for $\Delta \Gamma_s$.

CKM Unitarity violations

Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.
SM: $\lambda_u + \lambda_c + \lambda_t = 0 \Longrightarrow \text{NP}$: $\lambda_u + \lambda_c + \lambda_t + \delta_{CKM} = 0$
As a very rough estimate (4th family studies): $\frac{\delta_{CKM}^d}{\lambda_t^d} = \mathcal{O}(1)$, $\frac{\delta_{CKM}^s}{\lambda_t^s} = \mathcal{O}(\lambda)$ with $\lambda \approx 0.23$
Enhancement factors: 4 for $\Delta \Gamma_d$ and 1.4 for $\Delta \Gamma_s$.

CKM Unitarity violations

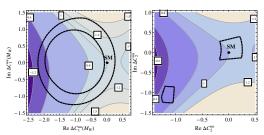
Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.
SM: $\lambda_u + \lambda_c + \lambda_t = 0 \Longrightarrow \text{NP}$: $\lambda_u + \lambda_c + \lambda_t + \delta_{CKM} = 0$
As a very rough estimate (4th family studies): $\frac{\delta_{CKM}^d}{\lambda_t^d} = \mathcal{O}(1)$, $\frac{\delta_{CKM}^s}{\lambda_t^s} = \mathcal{O}(\lambda)$ with $\lambda \approx 0.23$
Enhancement factors: 4 for $\Delta \Gamma_d$ and 1.4 for $\Delta \Gamma_s$.

Calculation of the effects of
$$C_{1,2}^{qq^{'},SM}\Longrightarrow C_{1,2}^{qq^{'},SM}+\Delta C_{1,2}^{qq^{'}}$$
 on $\Delta\Gamma_d$

CKM Unitarity violations

Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.
SM: $\lambda_u + \lambda_c + \lambda_t = 0 \Longrightarrow \text{NP}$: $\lambda_u + \lambda_c + \lambda_t + \delta_{CKM} = 0$
As a very rough estimate (4th family studies): $\frac{\delta_{CKM}^d}{\lambda_t^d} = \mathcal{O}(1)$, $\frac{\delta_{CKM}^s}{\lambda_t^s} = \mathcal{O}(\lambda)$ with $\lambda \approx 0.23$
Enhancement factors: 4 for $\Delta \Gamma_d$ and 1.4 for $\Delta \Gamma_s$.

Calculation of the effects of
$$C_{1,2}^{qq^{'},SM}\Longrightarrow C_{1,2}^{qq^{'},SM}+\Delta C_{1,2}^{qq^{'}}$$
 on $\Delta\Gamma_d$

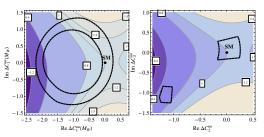


The contours show the ratio $\Delta\Gamma_d^{SM+NP}/\Delta\Gamma_d^{SM}$ for qq'=uu

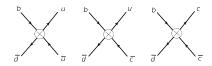
CKM Unitarity violations

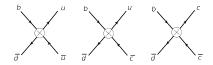
Let
$$\lambda_u = V_{ud}^* V_{ub}$$
, $\lambda_c = V_{cd}^* V_{cb}$, $\lambda_t = V_{td}^* V_{tb}$.
SM: $\lambda_u + \lambda_c + \lambda_t = 0 \Longrightarrow \text{NP}$: $\lambda_u + \lambda_c + \lambda_t + \delta_{CKM} = 0$
As a very rough estimate (4th family studies): $\frac{\delta_{CKM}^d}{\lambda_t^d} = \mathcal{O}(1)$, $\frac{\delta_{CKM}^s}{\lambda_t^s} = \mathcal{O}(\lambda)$ with $\lambda \approx 0.23$
Enhancement factors: 4 for $\Delta \Gamma_d$ and 1.4 for $\Delta \Gamma_s$.

Calculation of the effects of
$$C_{1,2}^{qq^{'},SM}\Longrightarrow C_{1,2}^{qq^{'},SM}+\Delta C_{1,2}^{qq^{'}}$$
 on $\Delta\Gamma_d$



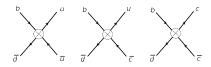
The contours show the ratio $\Delta\Gamma_d^{SM+NP}/\Delta\Gamma_d^{SM}$ for qq'=uu





Cuu bounds from:

$$\begin{array}{cccc} B^{0} & \rightarrow & \pi^{+}\pi^{-} \\ B^{-} & \rightarrow & \pi^{-}\pi^{0} \\ B^{0} & \rightarrow & \rho^{+}\rho^{-} \\ B^{0} & \rightarrow & \rho^{+/-}\pi^{-/+} \end{array}$$



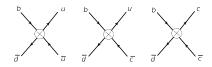
Cuu bounds from:

$$\begin{array}{cccc} B^0 & \rightarrow & \pi^+\pi^- \\ B^- & \rightarrow & \pi^-\pi^0 \\ B^0 & \rightarrow & \rho^+\rho^- \\ B^0 & \rightarrow & \rho^{+/-}\pi^{-/+} \end{array}$$

Cuc bounds from:

$$ar{\mathcal{B}}_d \rightarrow \mathcal{D}^{*+}\pi^-$$

 $\Gamma(\mathcal{B}_d)$



Cuu bounds from:

$$\begin{array}{cccc} B^0 & \rightarrow & \pi^+\pi^- \\ B^- & \rightarrow & \pi^-\pi^0 \\ B^0 & \rightarrow & \rho^+\rho^- \\ B^0 & \rightarrow & \rho^{+/-}\pi^{-/+} \end{array}$$

Cuc bounds from:

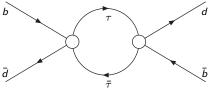
$$ar{B_d}
ightarrow D^{*+}\pi^- \Gamma(B_d)$$

Ccc bounds from:

$$egin{array}{ccc} B &
ightarrow & X_d \gamma \ sin(2eta) & & & \ a_{sl}^d & & & \end{array}$$

$(\bar{d}b)(\bar{ au} au)$ operators

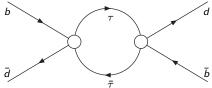
$$\begin{array}{lcl} Q_{S,AB} & = & \left(\bar{d} \, P_A \, b \right) \left(\bar{\tau} \, P_B \, \tau \right) \, , \\ Q_{V,AB} & = & \left(\bar{d} \, \gamma^\mu P_A \, b \right) \left(\bar{\tau} \, \gamma_\mu P_B \, \tau \right) \, , \\ Q_{T,A} & = & \left(\bar{d} \, \sigma^{\mu\nu} P_A \, b \right) \left(\bar{\tau} \, \sigma_{\mu\nu} P_A \, \tau \right) \, , \end{array}$$



Bounds from the following channels

$(\bar{d}b)(\bar{ au} au)$ operators

$$\begin{array}{lcl} Q_{S,AB} & = & \left(\bar{d} \, P_A \, b \right) \left(\bar{\tau} \, P_B \, \tau \right) \, , \\ Q_{V,AB} & = & \left(\bar{d} \, \gamma^\mu P_A \, b \right) \left(\bar{\tau} \, \gamma_\mu P_B \, \tau \right) \, , \\ Q_{T,A} & = & \left(\bar{d} \, \sigma^{\mu\nu} P_A \, b \right) \left(\bar{\tau} \, \sigma_{\mu\nu} P_A \, \tau \right) \, , \end{array}$$



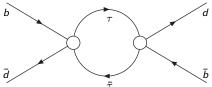
Bounds from the following channels

$$\bullet$$
 $B_d \rightarrow \tau^+\tau^-$

Room for new physics in $\Delta\Gamma_d$

$(\bar{d}b)(\bar{ au} au)$ operators

$$\begin{array}{lcl} Q_{S,AB} & = & \left(\bar{d} \, P_A \, b \right) \left(\bar{\tau} \, P_B \, \tau \right) \, , \\ Q_{V,AB} & = & \left(\bar{d} \, \gamma^\mu P_A \, b \right) \left(\bar{\tau} \, \gamma_\mu P_B \, \tau \right) \, , \\ Q_{T,A} & = & \left(\bar{d} \, \sigma^{\mu\nu} P_A \, b \right) \left(\bar{\tau} \, \sigma_{\mu\nu} P_A \, \tau \right) \, , \end{array}$$



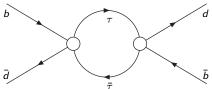
Bounds from the following channels

- \bullet $B_d \rightarrow \tau^+\tau^-$
- $B \to X_d \tau^+ \tau^-$ and $B^+ \to \pi^+ \tau^+ \tau^-$

Room for new physics in $\Delta\Gamma_d$

$(\bar{d}b)(\bar{ au} au)$ operators

$$\begin{array}{lcl} Q_{S,AB} & = & \left(\bar{d} \, P_A \, b \right) \left(\bar{\tau} \, P_B \, \tau \right) \, , \\ Q_{V,AB} & = & \left(\bar{d} \, \gamma^\mu P_A \, b \right) \left(\bar{\tau} \, \gamma_\mu P_B \, \tau \right) \, , \\ Q_{T,A} & = & \left(\bar{d} \, \sigma^{\mu\nu} P_A \, b \right) \left(\bar{\tau} \, \sigma_{\mu\nu} P_A \, \tau \right) \, , \end{array}$$



Bounds from the following channels

- \bullet $B_d \rightarrow \tau^+\tau^-$
- $B \to X_d \tau^+ \tau^-$ and $B^+ \to \pi^+ \tau^+ \tau^-$
- \bullet $B^+ o \pi^+ \mu^+ \mu^-$

Enhancement factor: 3.7 for $\Delta\Gamma_d$.

$$a_{sl}^q = \frac{\Gamma(\bar{B}(t)_q \to f) - \Gamma(B(t)_q \to \bar{f})}{\Gamma(\bar{B}(t)_q) \to f) + \Gamma(B(t)_q \to \bar{f})} = |\frac{\Gamma_{12}^q}{M_{12}^q}|sin(\phi_q)$$

$$a_{sl}^q = \frac{\Gamma(\bar{B}(t)_q o f) - \Gamma(B(t)_q o \bar{f})}{\Gamma(\bar{B}(t)_q) o f) + \Gamma(B(t)_q o \bar{f})} = |\frac{\Gamma_{12}^q}{M_{12}^q}|sin(\phi_q)$$

f: semileptonic state.

$$a_{sl}^q = \frac{\Gamma(\bar{B}(t)_q \to f) - \Gamma(B(t)_q \to \bar{f})}{\Gamma(\bar{B}(t)_q) \to f) + \Gamma(B(t)_q \to \bar{f})} = |\frac{\Gamma_{12}^q}{M_{12}^q}|sin(\phi_q)$$

f: semileptonic state.

$$Prob(B_q o ar{B}_q)
eq Prob(ar{B}_q o B_q)$$

$$a_{sl}^q = \frac{\Gamma(\bar{B}(t)_q \to f) - \Gamma(B(t)_q \to \bar{f})}{\Gamma(\bar{B}(t)_q) \to f) + \Gamma(B(t)_q \to \bar{f})} = |\frac{\Gamma_{12}^q}{M_{12}^q}|sin(\phi_q)$$

f: semileptonic state.

$$Prob(B_q o \bar{B}_q) \neq Prob(\bar{B}_q o B_q)$$

SM

$$a_{sl}^q = \frac{\Gamma(\bar{B}(t)_q \to f) - \Gamma(B(t)_q \to \bar{f})}{\Gamma(\bar{B}(t)_q) \to f) + \Gamma(B(t)_q \to \bar{f})} = |\frac{\Gamma_{12}^q}{M_{12}^q}| sin(\phi_q)$$

f: semileptonic state.

$$Prob(B_q o \bar{B}_q) \neq Prob(\bar{B}_q o B_q)$$

SM

$$a_{sl}^{d,HQE} = (-4.1 \pm 0.6) \times 10^{-4} \quad a_{sl}^{s,HQE} = (1.9 \pm 0.3) \times 10^{-5} \text{Lenz \& Nierste, arxiv:} \\ 1102.4274 \times 10^{-5} \text{Lenz & Nierste, arxiv:} \\ 1102.4274 \times 10^{-5} \times 10^{-$$

$$a_{sl}^q \quad = \quad \frac{\Gamma(\bar{B}(t)_q \to f) - \Gamma(B(t)_q \to \bar{f})}{\Gamma(\bar{B}(t)_q) \to f) + \Gamma(B(t)_q \to \bar{f})} = |\frac{\Gamma_{12}^q}{M_{12}^q}| sin(\phi_q)$$

f: semileptonic state.

$$Prob(B_q o \bar{B}_q) \neq Prob(\bar{B}_q o B_q)$$

SM

$$a_{sl}^{d,HQE} = (-4.1 \pm 0.6) \times 10^{-4} \quad a_{sl}^{s,HQE} = (1.9 \pm 0.3) \times 10^{-5} \text{Lenz \& Nierste, arxiv:} \\ 1102.4274 \times 10^{-5} \text{Lenz & Nierste, arxiv:} \\ 1102.4274 \times 10^{-5} \times 10^{-$$

Experimental results

$$a_{sl}^q \quad = \quad \frac{\Gamma(\bar{B}(t)_q \to f) - \Gamma(B(t)_q \to \bar{f})}{\Gamma(\bar{B}(t)_q) \to f) + \Gamma(B(t)_q \to \bar{f})} = |\frac{\Gamma_{12}^q}{M_{12}^q}| sin(\phi_q)$$

f: semileptonic state.

$$Prob(B_q o ar{B}_q)
eq Prob(ar{B}_q o B_q)$$

SM

$$a_{sl}^{d,HQE} = (-4.1 \pm 0.6) \times 10^{-4} \quad a_{sl}^{s,HQE} = (1.9 \pm 0.3) \times 10^{-5} \text{Lenz \& Nierste, arxiv:} \\ 1102.4274 \times 10^{-5} \text{Lenz & Nierste, arxiv:} \\ 1102.4274 \times 10^{-5} \times 10^{-$$

Experimental results

$$a_{sl}^{d,D0} = (0.68 \pm 0.45 \pm 0.14)\%(2012)$$
 $a_{sl}^{s,D0} = (-1.12 \pm 0.74 \pm 0.17)\%(2013)$ $a_{sl}^{d,BaBar} = (0.06 \pm 0.17 \pm 0.17^{+0.38}_{-0.23})\%(2013)$ $a_{sl}^{s,LHCb} = (-0.06 \pm 0.050 \pm 0.36)\%(2014)$

Like-sign dimuon asymmetry

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

 $A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$ $N^{++/--} : \# \text{ of events with two } +/\text{- muons from B hadron decays}$

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

$$N^{++/--} : \# \text{ of events with two } +/\text{- muons from B hadron decays}$$

$$A = A_{CP} + A_{bkg}$$

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

$$N^{++/--} : \# \text{ of events with two } +/- \text{ muons from B hadron decays}$$

$$A = A_{CP} + A_{bkg}$$

Standard interpretation $A_{CP} \propto A_{sl}^b = C_d a_{sl}^d + C_s a_{sl}^s$

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

$$N^{++/--} : \# \text{ of events with two } +/\text{- muons from B hadron decays}$$

$$A = A_{CP} + A_{bkg}$$

Standard interpretation
$$A_{CP} \propto A_{sl}^b = C_d a_{sl}^d + C_s a_{sl}^s$$

$$A_{sl}^{b,D0} = (-0.787 \pm 0.172 \pm 0.093)\%(2011)$$
 3.9 σ deviation from the SM

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

$$N^{++/--} : \# \text{ of events with two } +/- \text{ muons from B hadron decays}$$

$$A = A_{CP} + A_{bkg}$$

Standard interpretation $A_{CP} \propto A_{sl}^b = C_d a_{sl}^d + C_s a_{sl}^s$

$$A_{sl}^{b,D0} = (-0.787 \pm 0.172 \pm 0.093)\%(2011)$$
 3.9 σ deviation from the SM

Borissov and Hoeneisen $A_{CP} \propto C_d a_{sl}^d + C_s a_{sl}^s + C_{\Gamma_d} \frac{\Delta \Gamma_d}{\Gamma_d} + C_{\Gamma_s} \frac{\Delta \Gamma_s}{\Gamma_s}$ Phys. Rev. D 87, 074020

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

$$N^{++/--} : \# \text{ of events with two } +/- \text{ muons from B hadron decays}$$

$$A = A_{CP} + A_{bkg}$$

Standard interpretation
$$A_{CP} \propto A_{sl}^b = C_d a_{sl}^d + C_s a_{sl}^s$$

$$A_{sl}^{b,D0} = (-0.787 \pm 0.172 \pm 0.093)\%(2011) 3.9 \sigma$$
 deviation from the SM

Borissov and Hoeneisen
$$A_{CP} \propto C_d a_{sl}^d + C_s a_{sl}^s + C_{\Gamma_d} \frac{\Delta \Gamma_d}{\Gamma_d} + C_{\Gamma_s} \frac{\Delta \Gamma_s}{\Gamma_s}$$
 Phys. Rev. D 87, 074020

$$a_{sl}^d = (-0.62 \pm 0.43)\%$$
 $a_{sl}^s = (-0.82 \pm 0.99)\%$ $\frac{\Delta \Gamma_d}{\Gamma_d} = (0.50 \pm 1.38)\% D0 (2014)$

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

$$N^{++/--} : \# \text{ of events with two } +/- \text{ muons from B hadron decays}$$

$$A = A_{CP} + A_{bkg}$$

Standard interpretation
$$A_{CP} \propto A_{sl}^b = C_d a_{sl}^d + C_s a_{sl}^s$$

$$A_{sl}^{b,D0} = (-0.787 \pm 0.172 \pm 0.093)\%(2011)$$
 3.9 σ deviation from the SM

Borissov and Hoeneisen
$$A_{CP} \propto C_d a_{sl}^d + C_s a_{sl}^s + C_{\Gamma_d} \frac{\Delta \Gamma_d}{\Gamma_d} + C_{\Gamma_s} \frac{\Delta \Gamma_s}{\Gamma_s}$$
 Phys. Rev. D 87, 074020

$$a_{sl}^d = (-0.62 \pm 0.43)\%$$
 $a_{sl}^s = (-0.82 \pm 0.99)\%$ $\frac{\Delta \Gamma_d}{\Gamma_d} = (0.50 \pm 1.38)\% D0 (2014)$

Phys. Rev. D 89, 012002 (2014)

Like-sign dimuon asymmetry

$$A = \frac{N^{++} - N^{--}}{N^{++} + N^{--}}$$

$$N^{++/--} : \# \text{ of events with two } +/- \text{ muons from B hadron decays}$$

$$A = A_{CP} + A_{bkg}$$

Standard interpretation $A_{CP} \propto A_{sl}^b = C_d a_{sl}^d + C_s a_{sl}^s$

$$A_{sl}^{b,D0} = (-0.787 \pm 0.172 \pm 0.093)\%(2011) 3.9 \sigma$$
 deviation from the SM

Borissov and Hoeneisen $A_{CP} \propto C_d a_{sl}^d + C_s a_{sl}^s + C_{\Gamma_d} \frac{\Delta \Gamma_d}{\Gamma_d} + C_{\Gamma_s} \frac{\Delta \Gamma_s}{\Gamma_s}$ Phys. Rev. D 87, 074020

$$a_{sl}^d = (-0.62 \pm 0.43)\%$$
 $a_{sl}^s = (-0.82 \pm 0.99)\%$ $\frac{\Delta \Gamma_d}{\Gamma_d} = (0.50 \pm 1.38)\% D0 (2014)$

Phys. Rev. D 89, 012002 (2014)

3.0 σ deviation from the SM

$$\beta = \arg \left(-\frac{V_{cd} V_{cb}^*}{V_{td} V_{tb}^*} \right)$$

$$\beta = \arg\left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*}\right)$$

$$Prob(\bar{B}_d \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\beta = \arg\left(-\frac{V_{cd}\,V_{cb}^*}{V_{td}\,V_{tb}^*}\right)$$

$$Prob(\bar{B_d} \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\frac{\Gamma[\bar{B}_0(t) \to J/\psi K_S^0] - \Gamma[B_d(t) \to J/\psi K_S^0]}{\Gamma[\bar{B}_d(t) \to J/\psi K_S^0] + \Gamma[B_d(t) \to J/\psi K_S^0]} = C_{J/\psi K_S^0} cos(\Delta M_d t) + S_{J/\psi K_S^0} sin(\Delta M_d t)$$

$$\beta = \arg\left(-\frac{V_{cd}\,V_{cb}^*}{V_{td}\,V_{tb}^*}\right)$$

$$Prob(\bar{B}_d \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\frac{\Gamma[\bar{B}_0(t) \to J/\psi K_S^0] - \Gamma[B_d(t) \to J/\psi K_S^0]}{\Gamma[\bar{B}_d(t) \to J/\psi K_S^0]} = C_{J/\psi K_S^0} cos(\Delta M_d t) + S_{J/\psi K_S^0} sin(\Delta M_d t)$$

$$\frac{\bar{B}_0^0}{\bar{B}_0^0} + \frac{\bar{B}_0^0}{\bar{B}_0^0} + \frac{\bar{B}_$$

$$\beta = \arg\left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*}\right)$$

$$Prob(\bar{B_d} \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\frac{\Gamma[\bar{B}_0(t) \to J/\psi K_S^0] - \Gamma[B_d(t) \to J/\psi K_S^0]}{\Gamma[\bar{B_d}(t) \to J/\psi K_S^0] + \Gamma[B_d(t) \to J/\psi K_S^0]} = C_{J/\psi K_S^0} cos(\Delta M_d t) + S_{J/\psi K_S^0} sin(\Delta M_d t)$$

$$\frac{c}{\bar{B_0}} \xrightarrow{c} J/\psi \qquad b \xrightarrow{c} J/\psi K_S^0 = C_{J/\psi K_S^0} cos(\Delta M_d t) + C_{J/\psi K_S^0} sin(\Delta M_d t)$$

$$C_{J/\psi K_S^0} \approx 0$$

$$\beta = \arg\left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*}\right)$$

$$Prob(\bar{B}_d \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\frac{\Gamma[\bar{B}_0(t) \to J/\psi K_S^0] - \Gamma[B_d(t) \to J/\psi K_S^0]}{\Gamma[\bar{B}_d(t) \to J/\psi K_S^0]} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$C_{J/\psi K_S^0} \approx 0$$

$$S_{J/\psi K_S^0} \approx 0$$

$$S_{J/\psi K_S^0} \approx \sin(2\beta)$$

$$\beta = \arg\left(-\frac{V_{cd}\,V_{cb}^*}{V_{td}\,V_{tb}^*}\right)$$

$$Prob(\bar{B}_d \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\frac{\Gamma[\bar{B}_0(t) \to J/\psi K_S^0] - \Gamma[B_d(t) \to J/\psi K_S^0]}{\Gamma[\bar{B}_d(t) \to J/\psi K_S^0]} = C_{J/\psi K_S^0} cos(\Delta M_d t) + S_{J/\psi K_S^0} sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = \frac{\bar{B}_0}{\bar{B}_0} + \frac{\bar{B}_$$

$$\beta = \arg\left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*}\right)$$

$$Prob(\bar{B}_d \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\frac{\Gamma[\bar{B}_0(t) \to J/\psi K_S^0] - \Gamma[B_d(t) \to J/\psi K_S^0]}{\Gamma[\bar{B}_d(t) \to J/\psi K_S^0]} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} = 0$$

$$C_{J/\psi K_S^0} \approx 0$$

$$S_{J/\psi K_S^0} \approx \sin(2\beta)$$

$$\sin(2\beta)^{HFAG} = 0.680 \pm 0.020 \qquad \sin(2\beta)^{CKMfitter} = 0.775_{-0.049}^{+0.020}$$

$$\beta_S = -\arg(-V_{CS}V_{cb}^*/V_{IS}V_{tb}^*)$$

$$\beta = \arg\left(-\frac{V_{cd}\,V_{cb}^*}{V_{td}\,V_{tb}^*}\right)$$

$$Prob(\bar{B_d} \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\frac{\Gamma[\bar{B}_0(t) \to J/\psi K_S^0] - \Gamma[B_d(t) \to J/\psi K_S^0]}{\Gamma[\bar{B_d}(t) \to J/\psi K_S^0]} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{b}{\bar{B_0}} \xrightarrow{c} J/\psi$$

$$\bar{B_0} \xrightarrow{c} J/\psi$$

$$\bar{B_0} \xrightarrow{c} J/\psi$$

$$\bar{B_0} \xrightarrow{c} J/\psi$$

$$C_{J/\psi K_S^0} \approx 0$$

$$S_{J/\psi K_S^0} \approx \sin(2\beta)$$

$$\sin(2\beta)^{HFAG} = 0.680 \pm 0.020$$

$$\sin(2\beta)^{CKMfitter} = 0.775^{+0.020}_{-0.049}$$

$$\beta_S = -\arg(-V_{cs}V_{cb}^*/V_{ts}V_{tb}^*)$$

$$B_S \to J/\psi K^+K^-$$

$$\beta = \arg\left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*}\right)$$

$$Prob(\bar{B}_d \to J/\psi K_S^0) \neq Prob(B_d \to J/\psi K_S^0)$$

$$\frac{\Gamma[\bar{B}_0(t) \to J/\psi K_S^0] - \Gamma[B_d(t) \to J/\psi K_S^0]}{\Gamma[\bar{B}_d(t) \to J/\psi K_S^0]} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \sin(\Delta M_d t)$$

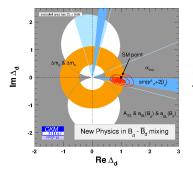
$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta M_d t) + S_{J/\psi K_S^0} \cos(\Delta M_d t)$$

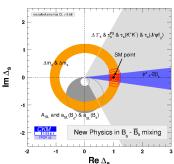
$$\frac{\bar{B}_0}{\bar{B}_0} = C_{J/\psi K_S^0} \cos(\Delta$$

$$M_{12}^q = M_{12}^{q,SM} \Delta_q = M_{12}^{q,SM} |\Delta_q| e^{i\phi_q^{\Delta}}$$

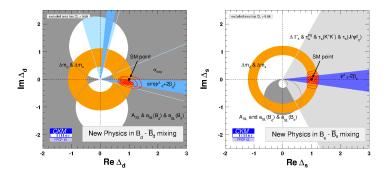
$$M_{12}^q = M_{12}^{q,SM} \Delta_q = M_{12}^{q,SM} |\Delta_q| \mathrm{e}^{i\phi_q^\Delta}$$
 $sin(2eta) o sin(2eta + \phi_d^\Delta)$

$$egin{align} \emph{M}_{12}^q = \emph{M}_{12}^{q,SM} \Delta_q = \emph{M}_{12}^{q,SM} |\Delta_q| e^{i\phi_q^\Delta} \ & sin(2eta)
ightarrow sin(2eta + \phi_d^\Delta) \ \end{split}$$



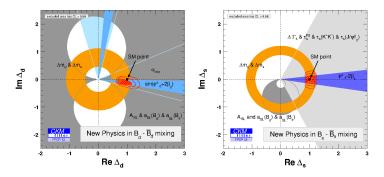


$$M_{12}^q = M_{12}^{q,SM} \Delta_q = M_{12}^{q,SM} |\Delta_q| e^{i\phi_q^\Delta}$$
 $sin(2eta)
ightarrow sin(2eta + \phi_d^\Delta)$



 Δ_d : Deviation from the SM less than 2σ , $\phi_d^{\Delta} < 10^{\circ}$.

$$M_{12}^q = M_{12}^{q,SM} \Delta_q = M_{12}^{q,SM} |\Delta_q| e^{i\phi_q^\Delta}$$
 $sin(2eta) o sin(2eta + \phi_d^\Delta)$



 Δ_d : Deviation from the SM less than $2\sigma,\ \phi_d^\Delta < 10^\circ.$

 Δ_s : Consistent with the SM, $-20^\circ < \phi_s^\Delta < 20^\circ$.

• Good agreement between SM and experiment for: ΔM_d , ΔM_s , $\Delta \Gamma_s$.

- Good agreement between SM and experiment for: ΔM_d , ΔM_s , $\Delta \Gamma_s$.
- lacktriangle Theoretical hadronic uncertainties in ΔM_d and ΔM_s allow space for new physics.

- Good agreement between SM and experiment for: ΔM_d , ΔM_s , $\Delta \Gamma_s$.
- Theoretical hadronic uncertainties in ΔM_d and ΔM_s allow space for new physics.
- There is room for new physics in $\Delta\Gamma_d$.

- Good agreement between SM and experiment for: ΔM_d , ΔM_s , $\Delta \Gamma_s$.
- Theoretical hadronic uncertainties in ΔM_d and ΔM_s allow space for new physics.
- There is room for new physics in $\Delta\Gamma_d$.
- Current experimental constraints allow the following enhacements:

- Good agreement between SM and experiment for: ΔM_d , ΔM_s , $\Delta \Gamma_s$.
- Theoretical hadronic uncertainties in ΔM_d and ΔM_s allow space for new physics.
- There is room for new physics in $\Delta\Gamma_d$.
- Current experimental constraints allow the following enhacements:

$$\frac{\Delta \Gamma_d}{\Delta \Gamma_d^{SM}} \leq \left\{ \begin{array}{ll} 4 & \text{CKM unitarity violations.} \\ 16 & \text{Current-current operators.} \\ 3.7 & (bd)(\tau \tau) \text{ operators.} \end{array} \right.$$

- Good agreement between SM and experiment for: ΔM_d , ΔM_s , $\Delta \Gamma_s$.
- Theoretical hadronic uncertainties in ΔM_d and ΔM_s allow space for new physics.
- There is room for new physics in $\Delta\Gamma_d$.
- Current experimental constraints allow the following enhacements:

$$\frac{\Delta \Gamma_d}{\Delta \Gamma_d^{SM}} \leq \left\{ \begin{array}{ll} 4 & \text{CKM unitarity violations.} \\ 16 & \text{Current-current operators.} \\ 3.7 & (bd)(\tau\tau) \text{ operators.} \end{array} \right.$$

 Direct measurements for sin(2β) show a small deviations with respect to indirect determinations.

- Good agreement between SM and experiment for: ΔM_d , ΔM_s , $\Delta \Gamma_s$.
- Theoretical hadronic uncertainties in ΔM_d and ΔM_s allow space for new physics.
- There is room for new physics in $\Delta\Gamma_d$.
- Current experimental constraints allow the following enhacements:

$$\frac{\Delta \Gamma_d}{\Delta \Gamma_d^{SM}} \leq \left\{ \begin{array}{ll} 4 & \text{CKM unitarity violations.} \\ 16 & \text{Current-current operators.} \\ 3.7 & (bd)(\tau\tau) \text{ operators.} \end{array} \right.$$

- Direct measurements for $sin(2\beta)$ show a small deviations with respect to indirect determinations.
- β_s is in good agreement with the SM.

- Good agreement between SM and experiment for: ΔM_d , ΔM_s , $\Delta \Gamma_s$.
- Theoretical hadronic uncertainties in ΔM_d and ΔM_s allow space for new physics.
- There is room for new physics in $\Delta\Gamma_d$.
- Current experimental constraints allow the following enhacements:

$$\frac{\Delta \Gamma_d}{\Delta \Gamma_d^{SM}} \leq \left\{ \begin{array}{ll} 4 & \text{CKM unitarity violations.} \\ 16 & \text{Current-current operators.} \\ 3.7 & (bd)(\tau\tau) \text{ operators.} \end{array} \right.$$

- Direct measurements for $sin(2\beta)$ show a small deviations with respect to indirect determinations.
- β_s is in good agreement with the SM.
- Further studies required to reduce uncertainties in hadronic parameters.