

# Radiative $B$ decays and new physics searches at $BABAR$

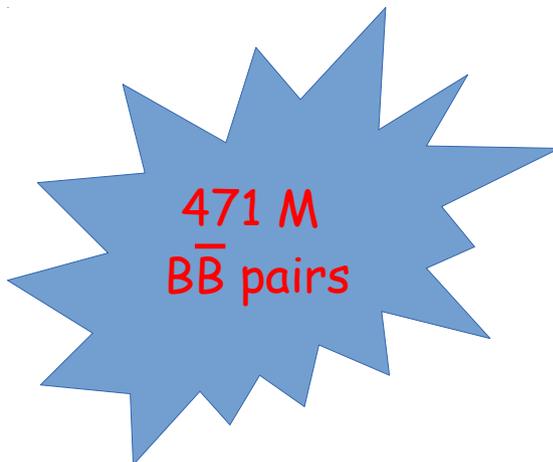


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(on behalf of the  $BABAR$  collaboration)

International Workshop on Deep-Inelastic Scattering and Related Subjects  
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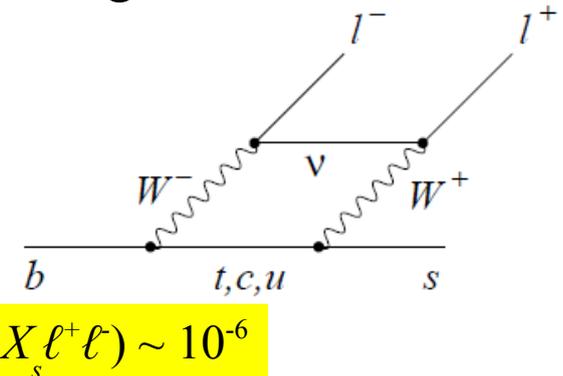
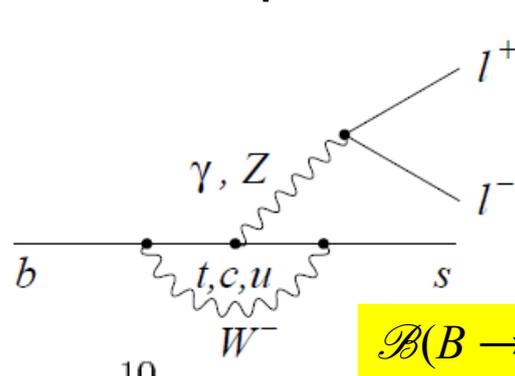
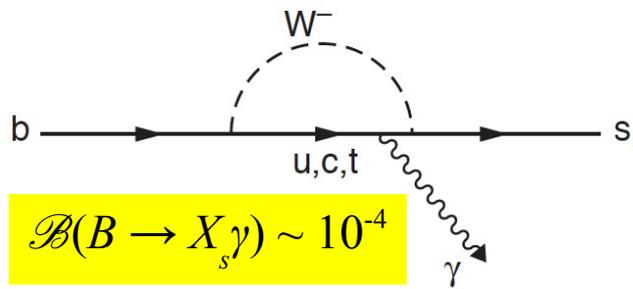
471 M  
 $B\bar{B}$  pairs

# Included BABAR results

- Search for direct CP violation in  $B \rightarrow X_s \gamma$  using a sum of exclusive final states
  - To be submitted to PRD
- Time-dependent analysis of  $B^0 \rightarrow K_s \pi^- \pi^+ \gamma$  and studies of the  $K^+ \pi^- \pi^+$  system in  $B^+ \rightarrow K^+ \pi^- \pi^+ \gamma$  decays
  - To be submitted to PRD
- Measurement of the  $B \rightarrow X_s \ell^+ \ell^-$  Branching Fraction (BF) and search for direct CP violation using a sum of exclusive final states
  - Submitted to PRL ( [arXiv:1312.5364](https://arxiv.org/abs/1312.5364) )

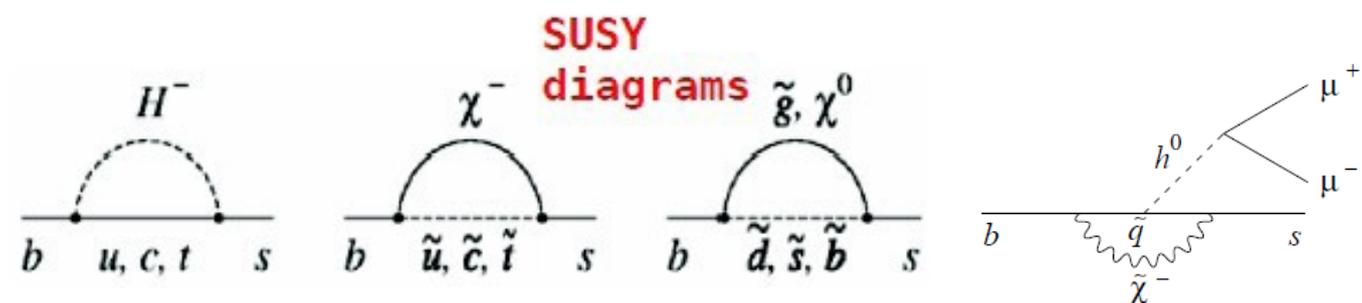
# Radiative B decays

- Forbidden at tree-level in SM, flavor-changing neutral-current (FCNC) processes proceed via loop and box diagrams such as:



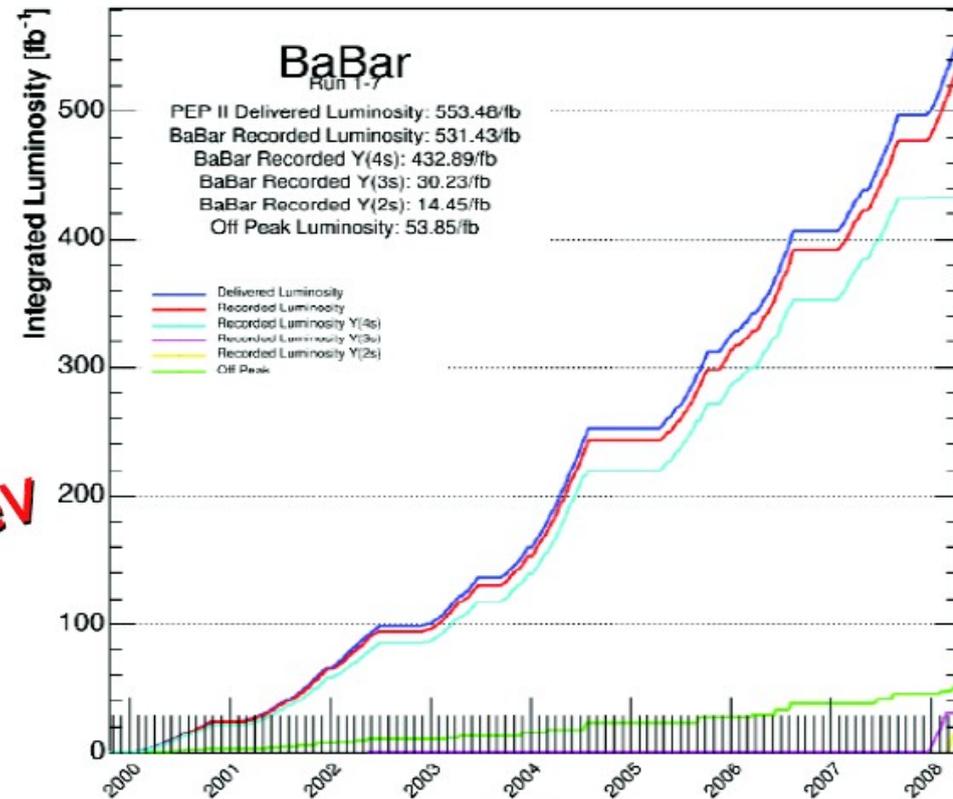
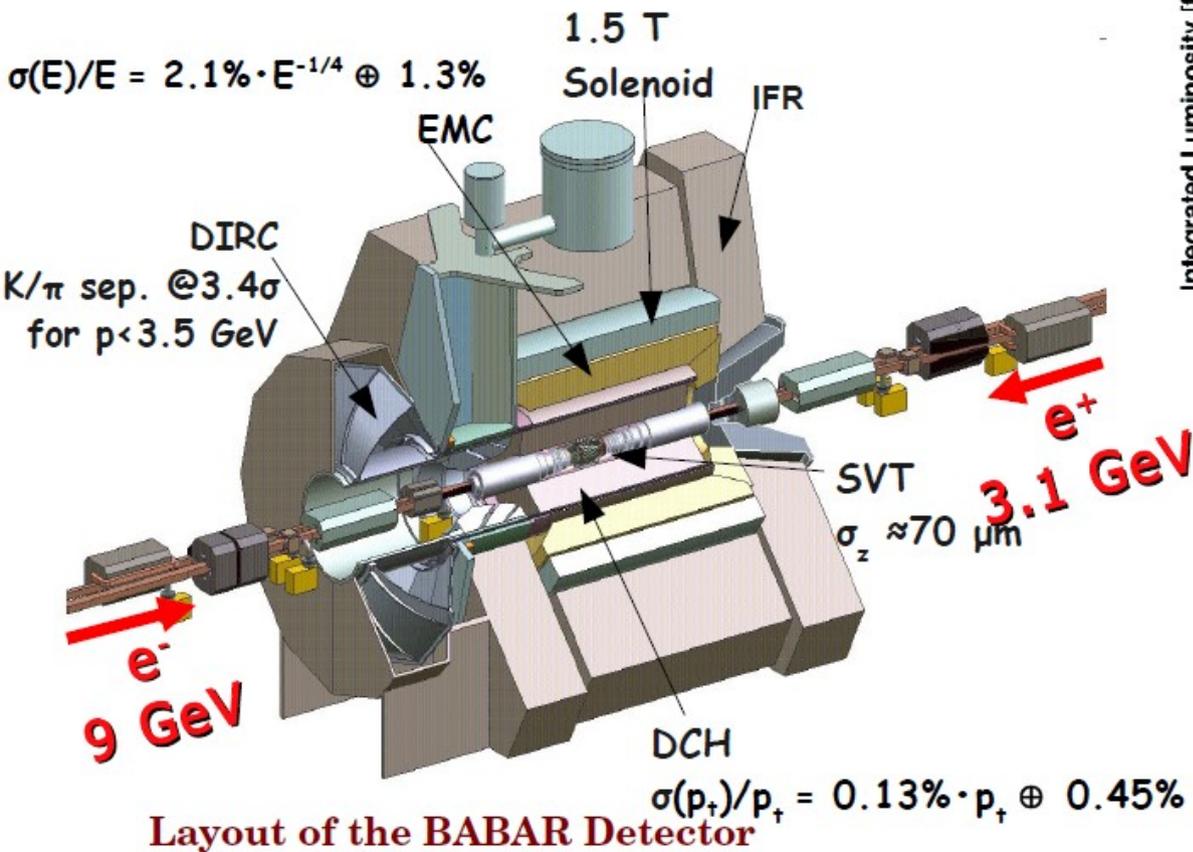
- Effective Hamiltonian  $H_{Eff} \propto \sum_{i=1}^{10} C_i \mathcal{O}_i$  factorizes short-distance Wilson coefficients  $C_i$  from long-distance effects

- New physics brings in new loops, and may change the SM values of Wilson coefficients / introduce new terms, e.g.  $C_S$  &  $C_P$



# BABAR experiment

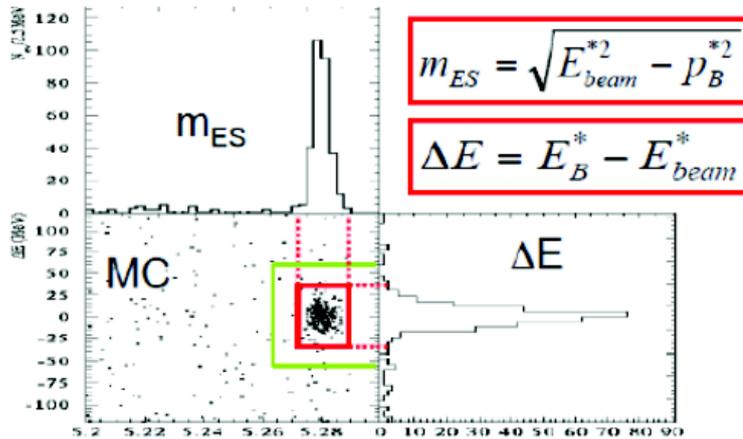
- The experimental data were collected between 1999 and 2008 with the BABAR detector at the PEP-II asymmetric energy  $e^+e^-$  collider located at SLAC



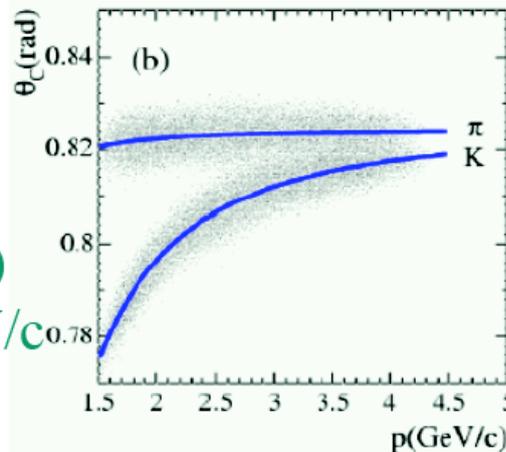
**$471 \times 10^6$   $B\bar{B}$  pairs collected at  $Y(4S)$  resonance**

# Common experimental techniques

## Kinematics of fully reconstructed B



Good charged particle ID (in particular K/ $\pi$ ) up to few GeV/c



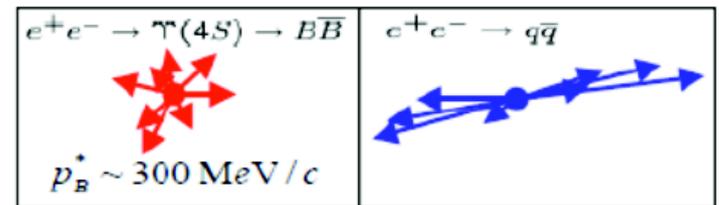
## Background characterization:

→ Mainly continuum:  $e^+e^- \rightarrow q\bar{q}$  ( $q = u, d, s, c$ ).

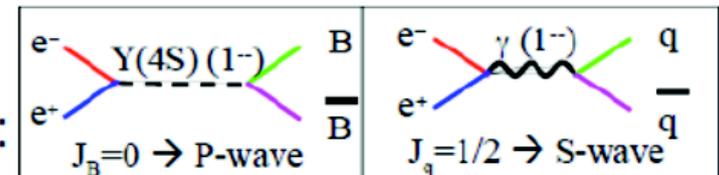
Suppression by multi-variable classifiers based on event-shape variables:

Fisher discriminant, Boosted decision trees (BDT)

Topology:

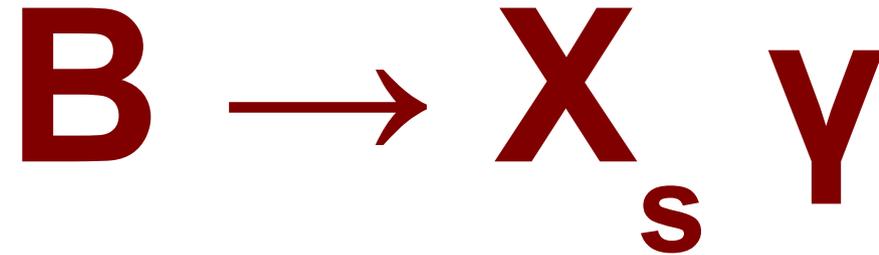


Angular distribution:

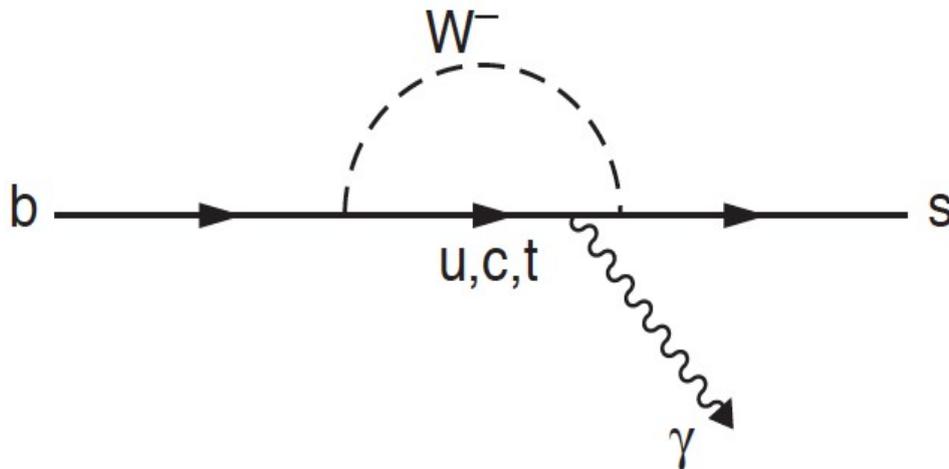


→ Background from B decays: classified by kinematic and topological properties

Variables are often combined to a likelihood function, used in a maximum likelihood fit for signal/background separation and to measure parameters of interest



- Measurements of direct CP asymmetries ( $A_{CP}$ ) from a sum of exclusive final states



# Direct $A_{CP}$ in $B \rightarrow X_s \gamma$

- Direct  $A_{CP}$  for the sum of exclusive decays is measured by:

$$A_{CP} = \frac{\Gamma_{\bar{B}^0/B^- \rightarrow X_s \gamma} - \Gamma_{B^0/B^+ \rightarrow X_{\bar{s}} \gamma}}{\Gamma_{\bar{B}^0/B^- \rightarrow X_s \gamma} + \Gamma_{B^0/B^+ \rightarrow X_{\bar{s}} \gamma}}$$

HFAG avg  $A_{CP} = (-0.8 \pm 2.9)\%$

- $A_{CP}$  in the SM is expected to be small and within  $(-0.6, +2.8)\%$  (1)

- Difference in direct  $A_{CP}$  in charged and neutral B mesons:

$$\Delta A_{X_s \gamma} = A_{B^\pm \rightarrow X_s \gamma} - A_{B^0/\bar{B}^0 \rightarrow X_s \gamma}$$

interference amplitude

$$17 \text{ MeV} < \bar{\Lambda}_{78} < 190 \text{ MeV}$$

- Depends on two Wilson coefficients:  $\Delta A_{X_s \gamma} = 4\pi^2 \alpha_s \frac{\bar{\Lambda}_{78}}{m_b} \Im m \left( \frac{C_{8g}}{C_{7\gamma}} \right)$  (1)

- In the SM,  $\Delta A_{X_s \gamma} = 0$  as the two Wilson coefficients are real

- NP may modify the two Wilson coefficients

- $C_{7\gamma}$  is constrained from  $b \rightarrow s \ell^+ \ell^-$  and  $\mathcal{B}(b \rightarrow s \gamma)$

- $C_{8g}$  can thus be constrained by  $\Delta A_{X_s \gamma}$ , **for the first time**

# Event selections

- 16 self-tagging final states used to reconstruct B meson decays:
  - Other modes rejected for being not flavor-specific or low in yield
- Requiring in each event at least one high energy photon ( $1.6 < E_\gamma^* < 3.0$  GeV in the  $Y(4S)$  rest frame)
- Requiring  $0.6 < m_{X_s} < 3.2$  GeV,  $m_{ES} > 5.24$  GeV, and  $|\Delta E| < 0.15$  GeV
- Primary (non-peaking) background of continuum events is suppressed by requiring the ratio of angular moments  $L_{12}/L_{10} < 0.46$ . Further background suppression with 2 BDTs

## Charged Modes:

$K_S^0 \pi^+ \gamma$   
 $K^+ \pi^0 \gamma$   
 $K^+ \pi^+ \pi^- \gamma$   
 $K_S^0 \pi^+ \pi^0 \gamma$   
 $K^+ \pi^0 \pi^0 \gamma$   
 $K_S^0 \pi^+ \pi^- \pi^+ \gamma$   
 $K^+ \pi^+ \pi^- \pi^0 \gamma$   
 $K_S^0 \pi^+ \pi^0 \pi^0 \gamma$   
 $K^+ \eta \gamma$   
 $K^+ K^- K^+ \gamma$

## Neutral Modes:

$K^+ \pi^- \gamma$   
 $K^+ \pi^- \pi^0 \gamma$   
 $K^+ \pi^+ \pi^- \pi^- \gamma$   
 $K^+ \pi^- \pi^0 \pi^0 \gamma$   
 $K^+ \eta \pi^- \gamma$   
 $K^+ K^- K^+ \pi^- \gamma$

Requiring  $K_s$  from  $K_s \rightarrow \pi^+ \pi^-$ ,  
 $\pi^0/\eta$  from  $\pi^0/\eta \rightarrow \gamma\gamma$

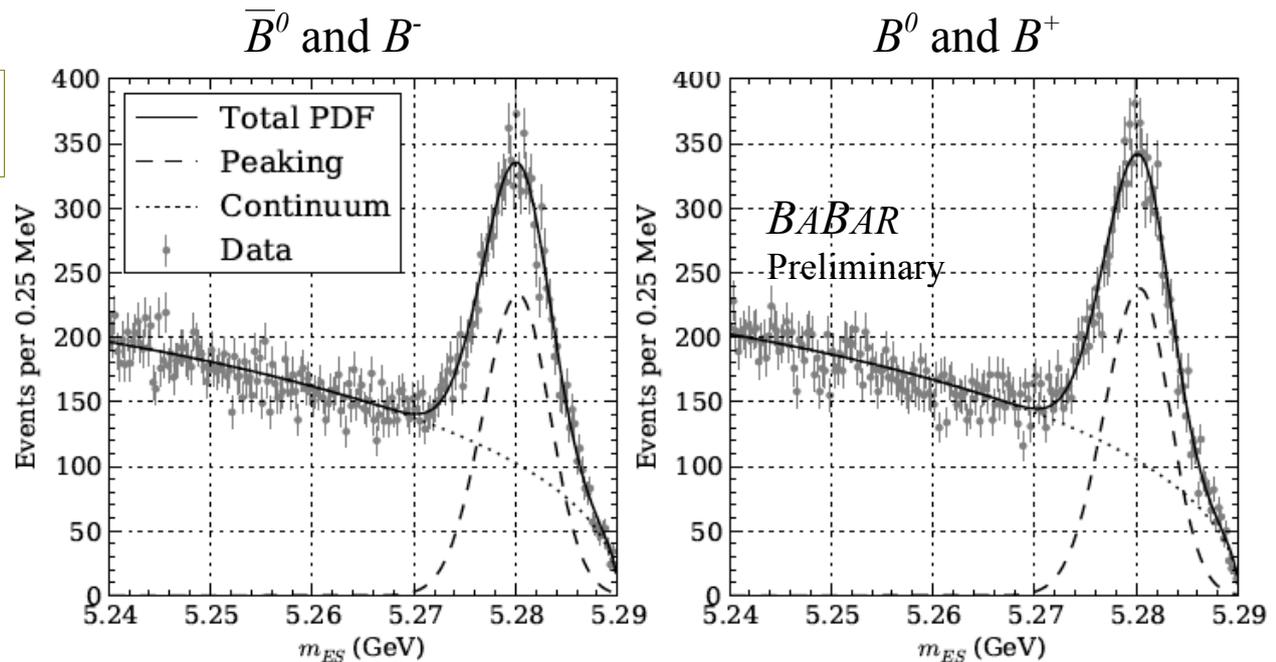
# Fitting for asymmetry

- Raw asymmetry is extracted via fitting simultaneously for B and  $\bar{B}$  tagged samples
- $A_{CP}$  is measured by correcting the raw asymmetry with detector asymmetry, which is determined from the  $m_{ES} < 5.27$  GeV sideband
- In the entire  $m_{Xs}$  range, we have:  $A_{CP} = +(1.7 \pm 1.9_{\text{stat}} \pm 1.0_{\text{syst}})\%$

– Consistent with SM

$$-0.6\% < A_{CP}^{\text{SM}} < 2.8\%$$

– The most precise measurement to date



# $\Delta A_{Xsy}$ and implication on $\text{Im}(C_{8g}/C_{7\gamma})$

- From the simultaneous fits to charged and neutral  $B$  samples measure

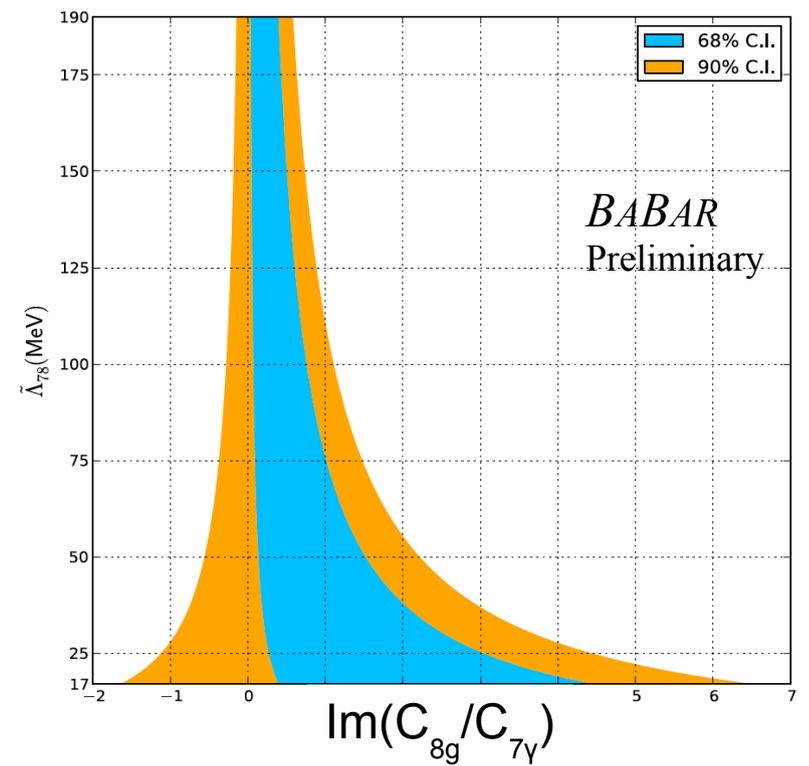
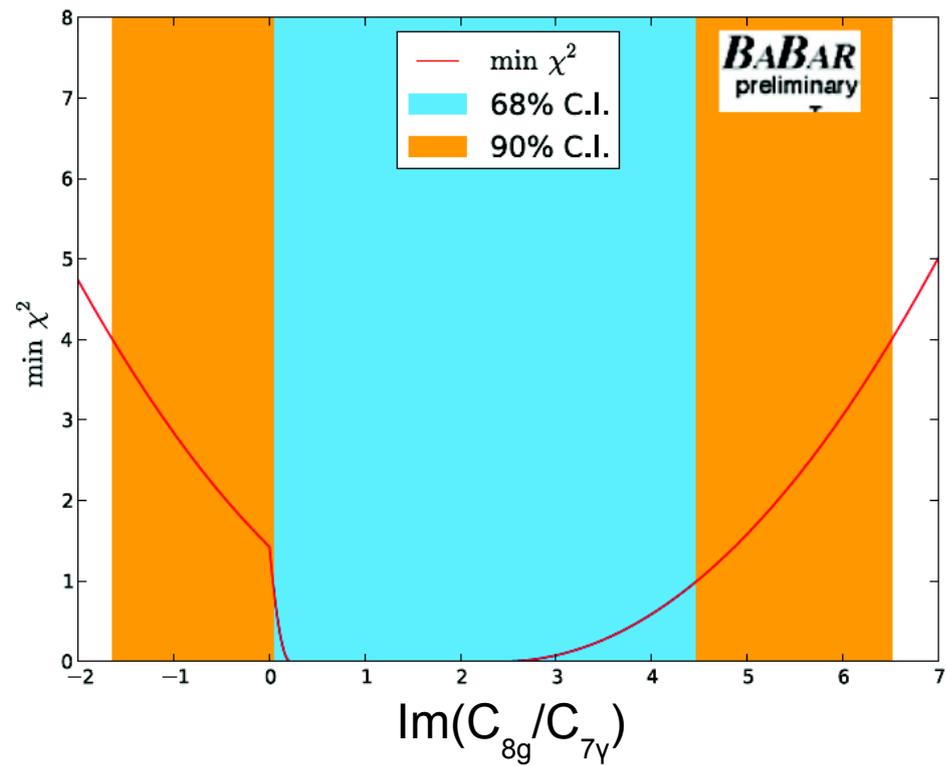
$$\Delta A_{Xsy} = +(5.0 \pm 3.9_{\text{stat}} \pm 1.5_{\text{syst}})\%$$

First measurement

- Confidence limits are set on  $\text{Im}(C_{8g}/C_{7\gamma})$ :  
First constraint of this kind for NP

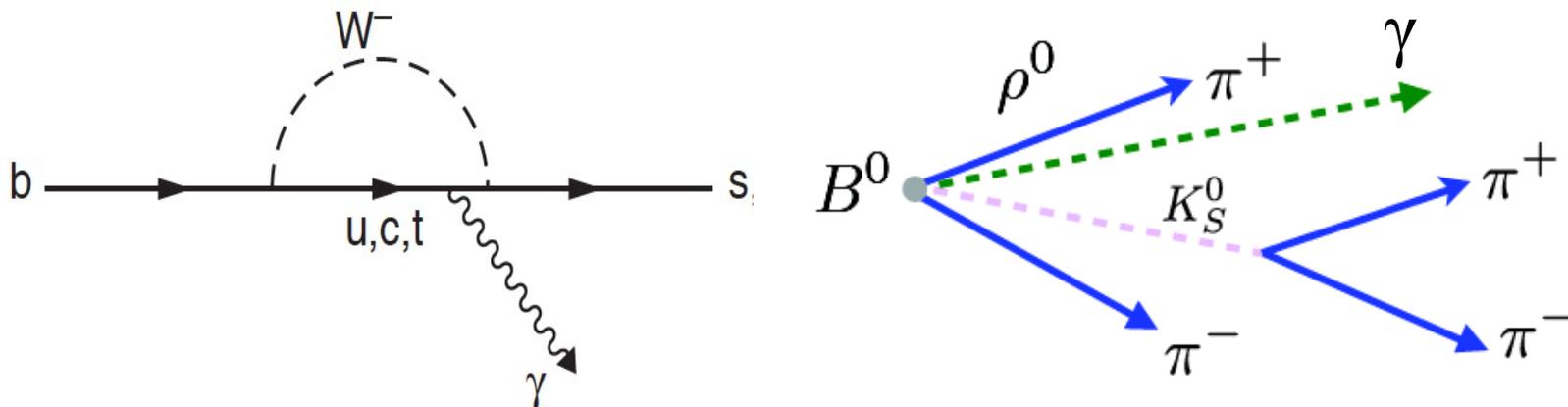
$$0.07 \leq \text{Im} \frac{C_{8g}}{C_{7\gamma}} \leq 4.48, \quad 68\% \text{ CL},$$

$$-1.64 \leq \text{Im} \frac{C_{8g}}{C_{7\gamma}} \leq 6.52, \quad 90\% \text{ CL}.$$

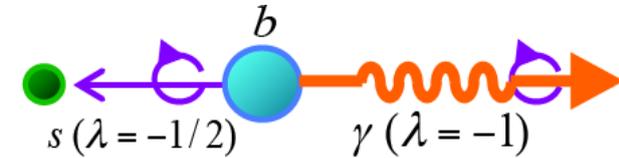


# $B \rightarrow K \pi^+ \pi^- \gamma$

- Time dependent analysis of  $B^0 \rightarrow K_S \pi^+ \pi^- \gamma$  and studies of the  $K^+ \pi^- \pi^+$  system in  $B^+ \rightarrow K^+ \pi^- \pi^+ \gamma$  decays



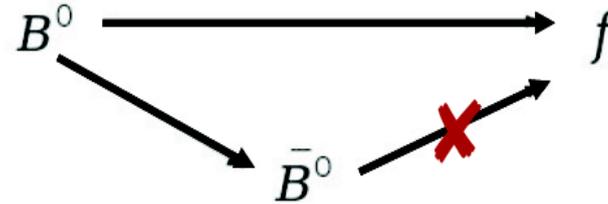
# Introduction



- In the SM, left-handed quarks and right-handed anti-quarks

$$\text{SM} \Rightarrow b \rightarrow s\gamma_L \text{ OR } \bar{b} \rightarrow \bar{s}\gamma_R \Rightarrow$$

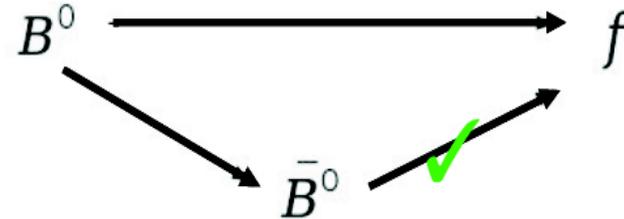
**CP asymmetry parameters  $\approx 0$**



- In the presence of NP, opposite helicity photons may get involved

$$\text{NP} \Rightarrow b \rightarrow s\gamma_{L,R} \text{ OR } \bar{b} \rightarrow \bar{s}\gamma_{R,L} \Rightarrow$$

**CP asymmetry parameters  $\neq 0$**



- Time-dependent CP asymmetry:

$$\begin{aligned} \mathcal{A}_{CP}(\Delta t) &= \frac{\Gamma(\bar{B}^0(\Delta t) \rightarrow f_{CP}\gamma) - \Gamma(B^0(\Delta t) \rightarrow f_{CP}\gamma)}{\Gamma(\bar{B}^0(\Delta t) \rightarrow f_{CP}\gamma) + \Gamma(B^0(\Delta t) \rightarrow f_{CP}\gamma)} \\ &= \mathcal{S}_{f_{CP}} \sin(\Delta m_d \Delta t) - \mathcal{C}_{f_{CP}} \cos(\Delta m_d \Delta t) \end{aligned}$$

Observable

$$\mathcal{S}_{f_{CP}} \stackrel{\text{SM}}{\propto} \frac{m_s}{m_b} \simeq 0.02$$

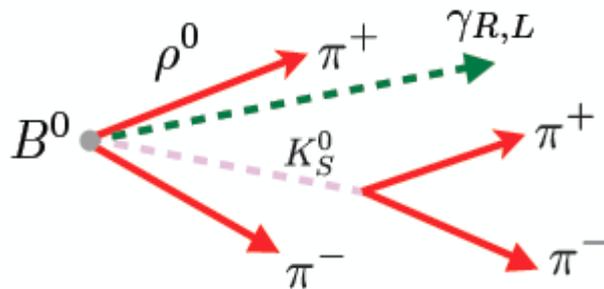
$\Delta m_d$ :  $B^0\bar{B}^0$  oscillation frequency, fixed to the measurement (1)

**Objective: measurement of  $\mathcal{S}$  in  $B^0 \rightarrow K_S \rho \gamma$  decays**

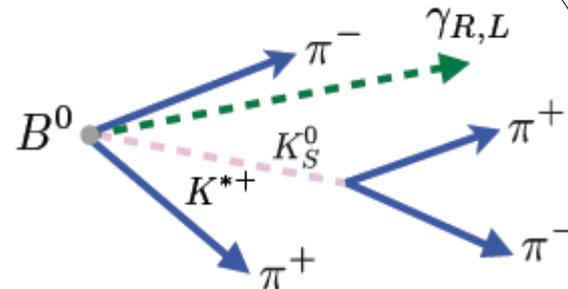
# Dilution factor

- Irreducible backgrounds from non-CP eigenstates ( $K^{*\pm}\pi^\mp$ ) which lie underneath the  $\rho^0(770)$  resonance and dilute the value of  $\mathcal{S}_{K_S^0\rho\gamma}$

$B^0 \rightarrow K_S \rho^0 (\rightarrow \pi^- \pi^+) \gamma :$



$B^0 \rightarrow K^{*\pm} (\rightarrow K_S \pi^\pm) \pi^\mp \gamma :$



- An amplitude study is needed to extract the **dilution factor**:

- Not possible for  $K_S \pi^+ \pi^-$  events due to low statistics

$$\mathcal{D}_{K_S^0\rho\gamma} \equiv \frac{\mathcal{S}_{K_S^0\pi^+\pi^-\gamma}}{\mathcal{S}_{K_S^0\rho\gamma}}$$

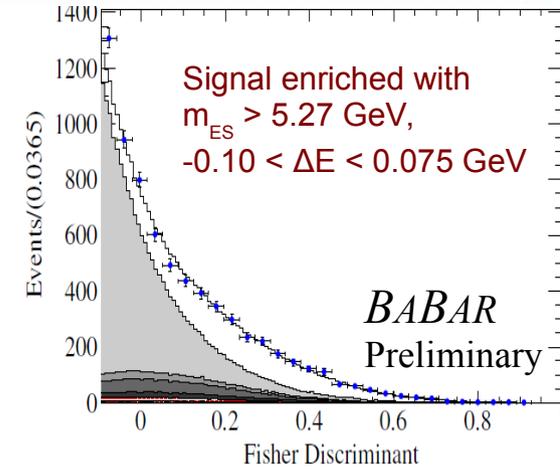
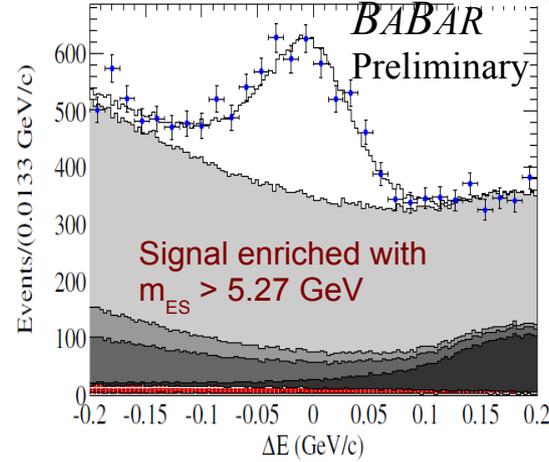
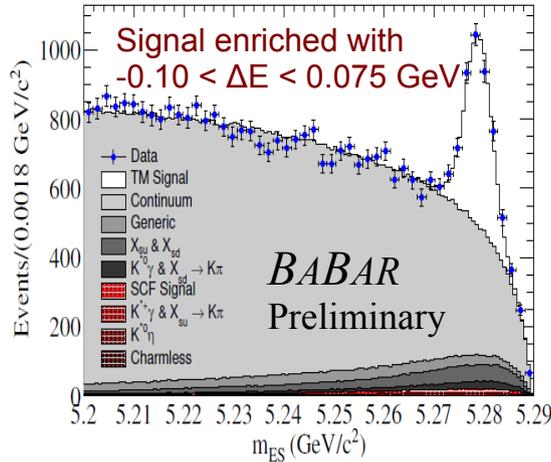
- By assuming isospin symmetry, we determine the dilution factor from the amplitude study in the  $B^+ \rightarrow K^+ \pi^+ \pi^- \gamma$  decay
  - Full 2D  $m_{K\pi} - m_{\pi\pi}$  amplitude study is still not possible given the sample size

# Fitting to $K^+\pi^-\pi^+$ events

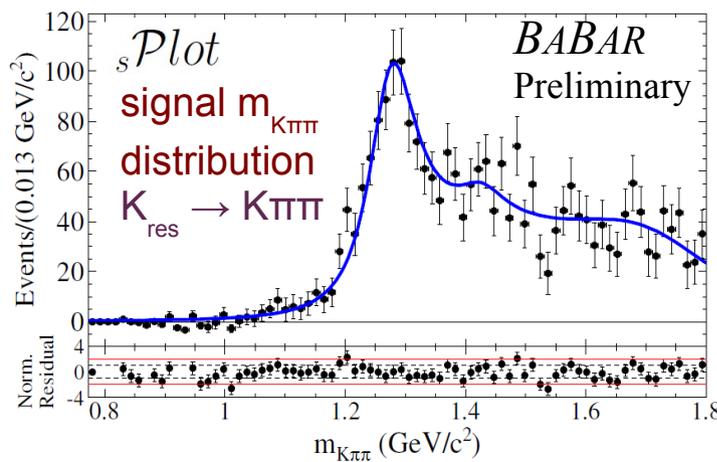
- Three stages of ML fits to find the dilution factor:

$$\mathcal{D}_{K_S^0\rho\gamma} = 0.549^{+0.096}_{-0.094}$$

1) 3D fits to extract  $K^+\pi^-\pi^+$  signals

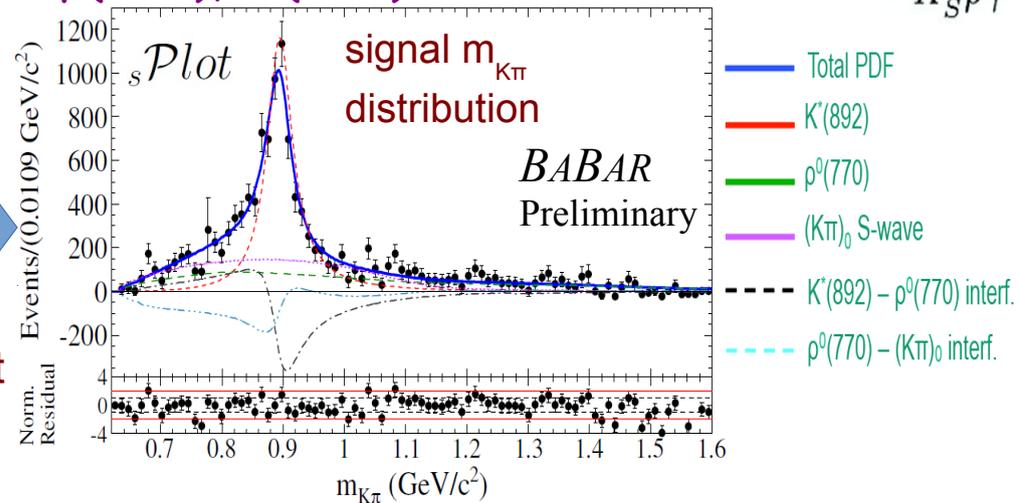


2) Fit to determine BF's of different kaonic resonances ( $K_{res}$ )



$K_{res}$  info as input

3) Fit to determine amplitudes of  $\rho(770)$ ,  $K^*(892)$  ...  $\Rightarrow$  Dilution factor  $\mathcal{D}_{K_S^0\rho\gamma}$



# BFs in the $K^+\pi^-\pi^+$ system

BFs for  $K_{\text{res}} \rightarrow K^+\pi^-\pi^+$

Mode	$\mathcal{B}(B^+ \rightarrow \text{Mode}) \times \mathcal{B}(K_{\text{res}} \rightarrow K^+\pi^+\pi^-) \times 10^{-6}$	$\mathcal{B}(B^+ \rightarrow \text{Mode}) \times 10^{-6}$	PDG values ( $\times 10^{-6}$ )
Inclusive $B^+ \rightarrow K^+\pi^+\pi^-\gamma$	...	$27.2 \pm 1.0^{+1.1}_{-1.3}$	$27.6 \pm 2.2$
$K_1(1270)^+\gamma$	$14.5^{+2.0+1.1}_{-1.3-1.2}$	$44.0^{+6.0+3.5}_{-4.0-3.7} \pm 4.6$	$43 \pm 13$
$K_1(1400)^+\gamma$	$4.1^{+1.9+1.3}_{-1.2-0.8}$	$9.7^{+4.6+3.1}_{-2.9-1.8} \pm 0.6$	$< 15 \text{ CL} = 90\%$
$K^*(1410)^+\gamma$	$9.7^{+2.1+2.4}_{-1.9-0.7}$	$23.8^{+5.2+5.9}_{-4.6-1.4} \pm 2.4$	$\emptyset$
$K_2^*(1430)^+\gamma$	$1.5^{+1.2+0.9}_{-1.0-1.4}$	$10.4^{+8.7+6.3}_{-7.0-9.9} \pm 0.5$	$14 \pm 4$
$K^*(1680)^+\gamma$	$17.0^{+1.7+3.5}_{-1.4-3.0}$	$71.7^{+7.2+15}_{-5.7-13} \pm 5.8$	$< 1900 \text{ CL} = 90\%$

*BABAR*

Preliminary

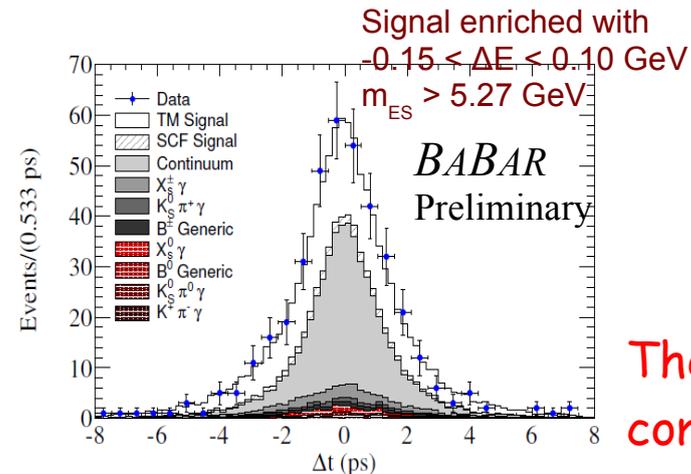
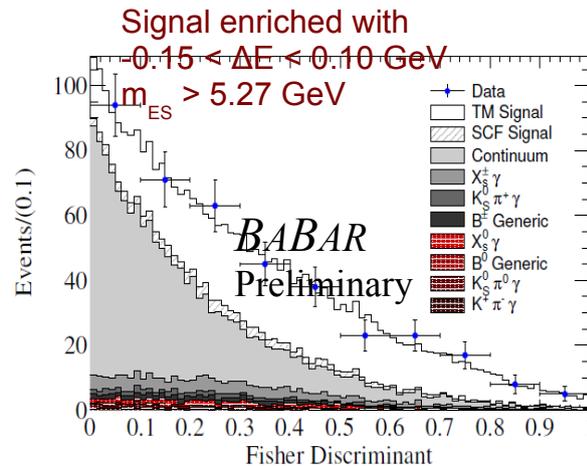
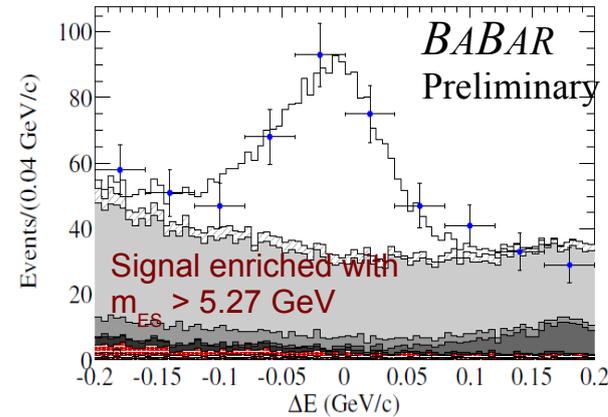
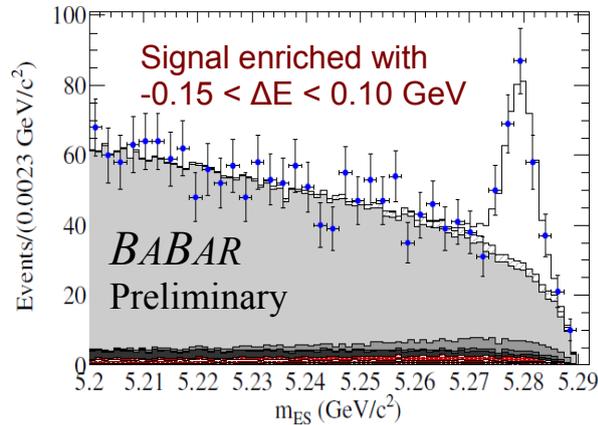
BFs of different resonances decaying to  $K^+\pi^-$  and  $\pi^+\pi^-$

Some of these measurements are the world's best

Mode	$\mathcal{B}(B^+ \rightarrow \text{Mode}) \times \mathcal{B}(R \rightarrow hh) \times 10^{-6}$	$\mathcal{B}(B^+ \rightarrow \text{Mode}) \times 10^{-6}$	PDG values ( $\times 10^{-6}$ )
Inclusive $B^+ \rightarrow K^+\pi^+\pi^-\gamma$	...	$27.2 \pm 1.0^{+1.1}_{-1.3}$	$27.6 \pm 2.2$
$K^{*0}(892)\pi^+\gamma$	$17.3 \pm 0.9^{+1.2}_{-1.1}$	$26.0^{+1.4}_{-1.3} \pm 1.8$	$20^{+7}_{-6}$
$K^+\rho(770)^0\gamma$	$9.1^{+0.8}_{-0.7} \pm 1.3$	$9.2^{+0.8}_{-0.7} \pm 1.3 \pm 0.02$	$< 20 \text{ CL} = 90\%$
$(K\pi)_0^{*0}\pi^+\gamma$	$11.3 \pm 1.5^{+2.0}_{-2.6}$	...	$\emptyset$
$(K\pi)_0^0\pi^+\gamma$ (NR)	...	$10.8^{+1.4+1.9}_{-1.5-2.5}$	$< 9.2 \text{ CL} = 90\%$
$K_0^*(1430)^0\pi^+\gamma$	$0.51 \pm 0.07^{+0.09}_{-0.12}$	$0.82 \pm 0.11^{+0.15}_{-0.19} \pm 0.08$	$\emptyset$

# TDCP analysis in $K_S\pi^-\pi^+$

- 4D ML fit to extract CP asymmetry parameters



The result is compatible with SM value

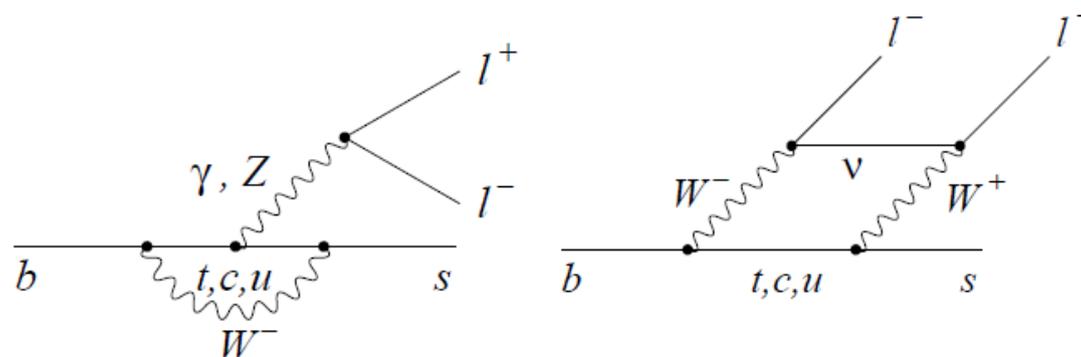
$$S_{K_S^0 \rho \gamma} = \frac{S_{K_S^0 \pi^+ \pi^- \gamma}}{D_{K_S^0 \rho \gamma}} = 0.249 \pm 0.455^{+0.076}_{-0.060} \quad (\text{stat.}) \quad (\text{syst.})$$

$$S_{K_S^0 \rho^0 \gamma} \stackrel{\text{SM}}{\sim} 0.02$$

$$B \rightarrow X_s e^+ e^-$$

- Measurement of the branching fraction and search for direct CP violation using a sum of exclusive final states

– [arXiv:1312.5364](https://arxiv.org/abs/1312.5364)

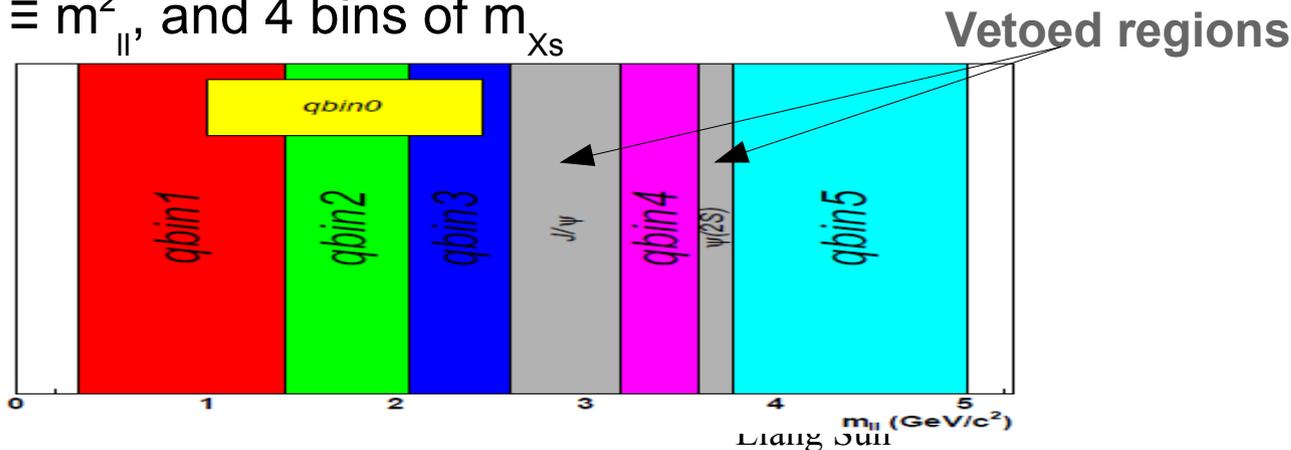


# Event selection

$$B \rightarrow X_s \ell^+ \ell^-$$

$$m_{X_s} < 1.8 \text{ GeV}$$

- 20 exclusive final states with  $\leq 2$  pions and  $\leq 1 \pi^0$ :
- $m_{ES} > 5.225 \text{ GeV}$
- $-0.1 (-0.05) < \Delta E < 0.05 \text{ GeV}$  for  $X_s e^+ e^-$  ( $X_s \mu^+ \mu^-$ )
- Selection represent  $\sim 70\%$  of the inclusive rate with  $m_{X_s} < 1.8 \text{ GeV}$ , accounting for  $K_L^0$  modes,  $K_s^0 \rightarrow \pi^0 \pi^0$  and  $\pi^0$  Dalitz decays
- We extrapolate for the missing modes, and those with  $m_{X_s} > 1.8 \text{ GeV}$ , using JETSET fragmentation and theory prediction
- We measure the total BF and partial BFs in 6 bins of  $q^2 \equiv m_{ll}^2$ , and 4 bins of  $m_{X_s}$



$$B^0 \rightarrow K_s^0 \mu^+ \mu^-$$

$$B^+ \rightarrow K^+ \mu^+ \mu^-$$

$$B^0 \rightarrow K_s^0 e^+ e^-$$

$$B^+ \rightarrow K^+ e^+ e^-$$

$$B^0 \rightarrow K^{*0} (K_s^0 \pi^0) \mu^+ \mu^-$$

$$B^+ \rightarrow K^{*+} (K^+ \pi^0) \mu^+ \mu^-$$

$$B^+ \rightarrow K^{*+} (K_s^0 \pi^+) \mu^+ \mu^-$$

$$B^0 \rightarrow K^{*0} (K^+ \pi^-) \mu^+ \mu^-$$

$$B^0 \rightarrow K^{*0} (K_s^0 \pi^0) e^+ e^-$$

$$B^+ \rightarrow K^{*+} (K^+ \pi^0) e^+ e^-$$

$$B^+ \rightarrow K^{*+} (K_s^0 \pi^+) e^+ e^-$$

$$B^0 \rightarrow K^{*0} (K^+ \pi^-) e^+ e^-$$

$$B^+ \rightarrow K_s^0 \pi^+ \pi^0 \mu^+ \mu^-$$

$$B^0 \rightarrow K^+ \pi^- \pi^0 \mu^+ \mu^-$$

$$B^0 \rightarrow K_s^0 \pi^+ \pi^- \mu^+ \mu^-$$

$$B^+ \rightarrow K^+ \pi^+ \pi^- \mu^+ \mu^-$$

$$B^+ \rightarrow K_s^0 \pi^+ \pi^0 e^+ e^-$$

$$B^0 \rightarrow K^+ \pi^- \pi^0 e^+ e^-$$

$$B^0 \rightarrow K_s^0 \pi^+ \pi^- e^+ e^-$$

$$B^+ \rightarrow K^+ \pi^+ \pi^- e^+ e^-$$

Requiring  $K_s$  from  $K_s \rightarrow \pi^+ \pi^-$ ,  $\pi^0$  from  $\pi^0 \rightarrow \gamma \gamma$

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- $-0.1 (-0.05) < \Delta E < 0.05 \text{ GeV}$  for  $X_s e^+ e^-$  ( $X_s \mu^+ \mu^-$ )
- Selection represent  $\sim 70\%$  of the inclusive rate with  $m_{X_s} < 1.8 \text{ GeV}$ , accounting for  $K_L^0$  modes,  $K_s^0 \rightarrow \pi^0 \pi^0$  and  $\pi^0$  Dalitz decays
- We extrapolate for the missing modes, and those with  $m_{X_s} > 1.8 \text{ GeV}$ , using JETSET fragmentation and theory prediction
- We separate B and  $\bar{B}$  samples in the 14 **self-tagging modes** to measure direct CP asymmetry

$$A_{CP} \equiv \frac{\Gamma(\bar{b} \rightarrow \bar{s} \ell \ell) - \Gamma(b \rightarrow s \ell \ell)}{\Gamma(\bar{b} \rightarrow \bar{s} \ell \ell) + \Gamma(b \rightarrow s \ell \ell)}$$

SM expectation well below 1% (1)

(1) Phys. Rev. D 54, 882 (1996); Eur. Phys. J. C 8, 619 (1999); JHEP 0807, 106 (2008); JHEP 0901, 019 (2009)

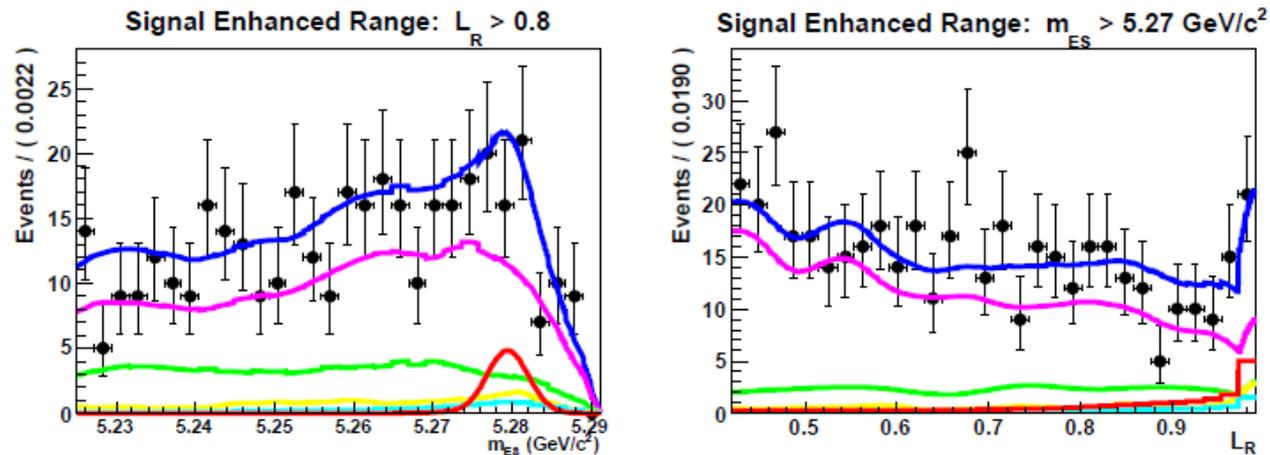
$B^+ \rightarrow K^+ \mu^+ \mu^-$
$B^+ \rightarrow K^+ e^+ e^-$
$B^+ \rightarrow K^{*+} (K^+ \pi^0) \mu^+ \mu^-$
$B^+ \rightarrow K^{*+} (K_s^0 \pi^+) \mu^+ \mu^-$
$B^0 \rightarrow K^{*0} (K^+ \pi^-) \mu^+ \mu^-$
$B^+ \rightarrow K^{*+} (K^+ \pi^0) e^+ e^-$
$B^+ \rightarrow K^{*+} (K_s^0 \pi^+) e^+ e^-$
$B^0 \rightarrow K^{*0} (K^+ \pi^-) e^+ e^-$
$B^+ \rightarrow K_s^0 \pi^+ \pi^0 \mu^+ \mu^-$
$B^0 \rightarrow K^+ \pi^- \pi^0 \mu^+ \mu^-$
$B^+ \rightarrow K^+ \pi^+ \pi^- \mu^+ \mu^-$
$B^+ \rightarrow K_s^0 \pi^+ \pi^0 e^+ e^-$
$B^0 \rightarrow K^+ \pi^- \pi^0 e^+ e^-$
$B^+ \rightarrow K^+ \pi^+ \pi^- e^+ e^-$

Requiring  $K_s$  from  $K_s \rightarrow \pi^+ \pi^-$ ,  $\pi^0$  from  $\pi^0 \rightarrow \gamma \gamma$

# Fitting strategy

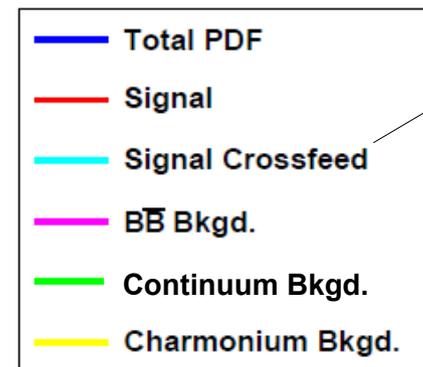
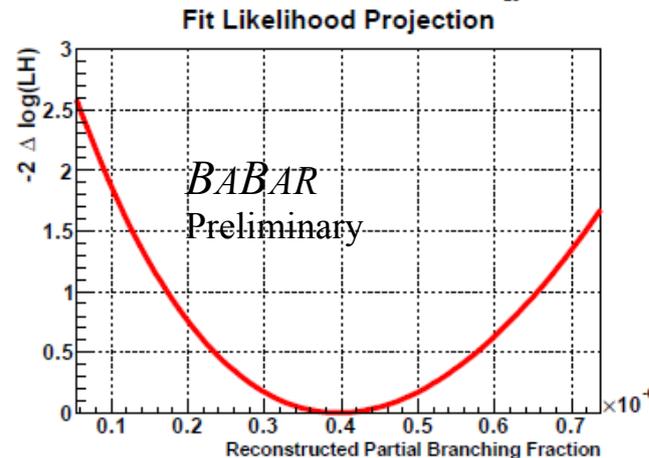
- Signal extraction with 2D ML fit to  $m_{ES}$  and likelihood ratio  $L_R$  which is based on boosted decision trees and peaks  $\sim 1$  for signals

Example fit for  $X_s e^+ e^-$  events with  $2.0 < q^2 < 4.3 \text{ GeV}^2$



$$L_R = \frac{\mathcal{P}_S}{\mathcal{P}_S + \mathcal{P}_B}$$

$\mathcal{P}_{S/B}$ : signal/bkg. probability from  $\overline{BB}$  BDT output



Mis-reconstructed signal events

Different background categories

# Branching fraction results

$$B \rightarrow X_s \ell^+ \ell^-$$

- Total branching fraction ( $q^2 > 0.1 \text{ GeV}^2$ )

$$B(B \rightarrow X_s \ell^+ \ell^-) = (6.73^{+0.70+0.34}_{-0.63-0.25} \pm 0.5) \times 10^{-6}$$

$$B_{SM}(B \rightarrow X_s \ell^+ \ell^-) = (4.6 \pm 0.8) \times 10^{-6} \quad (1)$$

- $< 2\sigma$  higher than the SM expectation

- In the low mass range ( $1 < q^2 < 6 \text{ GeV}^2$ ):

$$B(B \rightarrow X_s \ell^+ \ell^-) = (1.60^{+0.41+0.17}_{-0.39-0.13} \pm 0.18) \times 10^{-6}$$

$$B_{SM}(B \rightarrow X_s e^+ e^-) = (1.64 \pm 0.11) \times 10^{-6} \quad (2)$$

- In good agreement with the SM

- In the high mass range ( $q^2 > 14.2 \text{ GeV}^2$ )

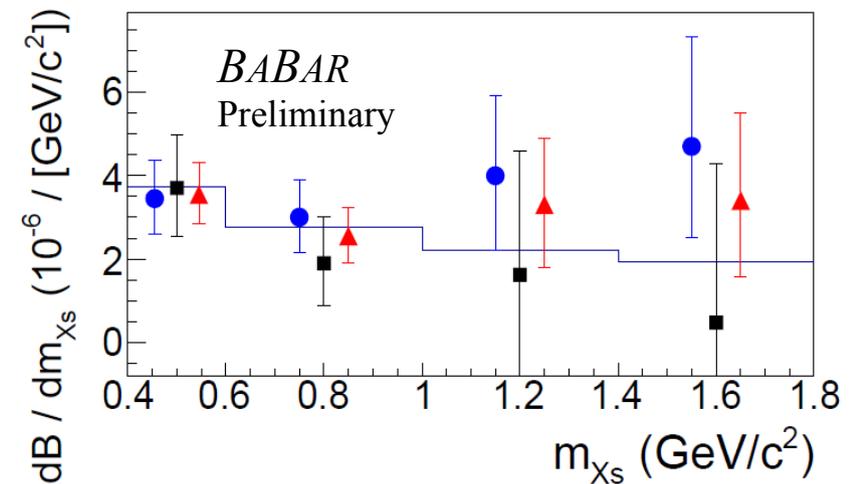
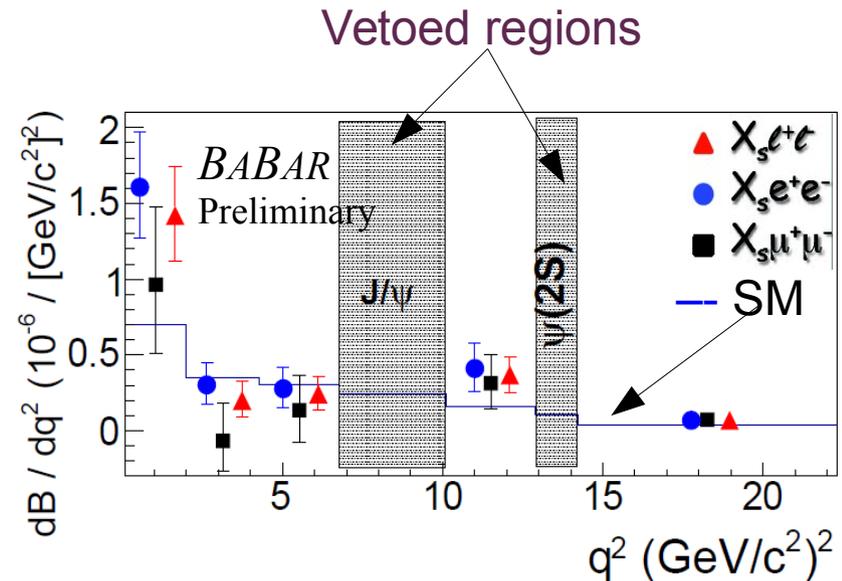
$$B(B \rightarrow X_s \ell^+ \ell^-) = (0.57^{+0.16+0.03}_{-0.15-0.02} \pm 0.0) \times 10^{-6}$$

$$B_{SM}(B \rightarrow X_s e^+ e^-) = (0.21 \pm 0.07) \times 10^{-6} \quad (2)$$

- $\sim 2\sigma$  higher than the SM expectation

(1) Nucl.Phys.B 685, 351 (2004)  
 (2) Nucl.Phys.B 802, 40 (2008)

In the listed results, the first uncertainty is statistical, the second experimental systematics and the third model-dependent systematics



# Branching fraction results

$$B \rightarrow X_s \ell^+ \ell^-$$

- Total branching fraction ( $q^2 > 0.1 \text{ GeV}^2$ )

$$B(B \rightarrow X_s \ell^+ \ell^-) = (6.73^{+0.70+0.34}_{-0.63-0.25} \pm 0.5) \times 10^{-6}$$

$$B_{SM}(B \rightarrow X_s \ell^+ \ell^-) = (4.6 \pm 0.8) \times 10^{-6} \quad (1)$$

- $< 2\sigma$  higher than the SM expectation

The BF measurement for  $q^2 > 14.2 \text{ GeV}^2$  disfavors by  $\sim 2.5\text{-}3\sigma$  the  $\delta C_9$  NP interpretation of the recent LHCb "anomaly" for an angular observable from  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  (3)

- > Expected pBF lower than SM in this interpretation

- In the high mass range ( $q^2 > 14.2 \text{ GeV}^2$ )

$$B(B \rightarrow X_s \ell^+ \ell^-) = (0.57^{+0.16+0.03}_{-0.15-0.02} \pm 0.0) \times 10^{-6}$$

$$B_{SM}(B \rightarrow X_s e^+ e^-) = (0.21 \pm 0.07) \times 10^{-6} \quad (2)$$

- $\sim 2\sigma$  higher than the SM expectation

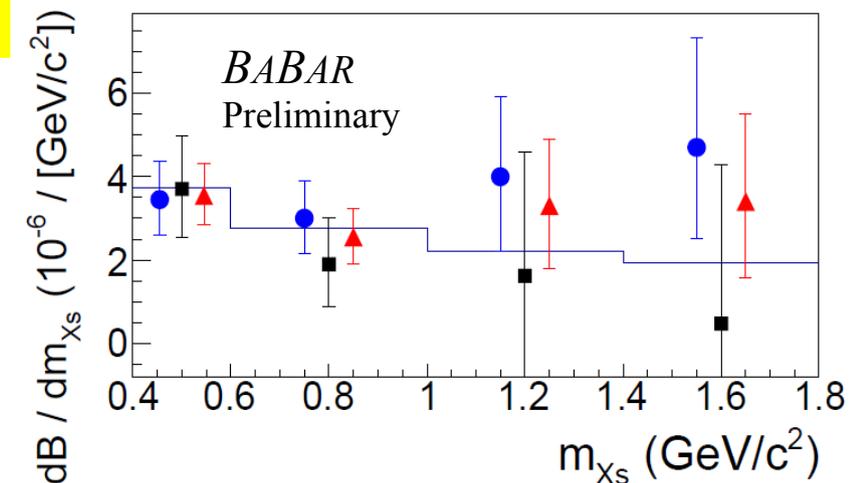
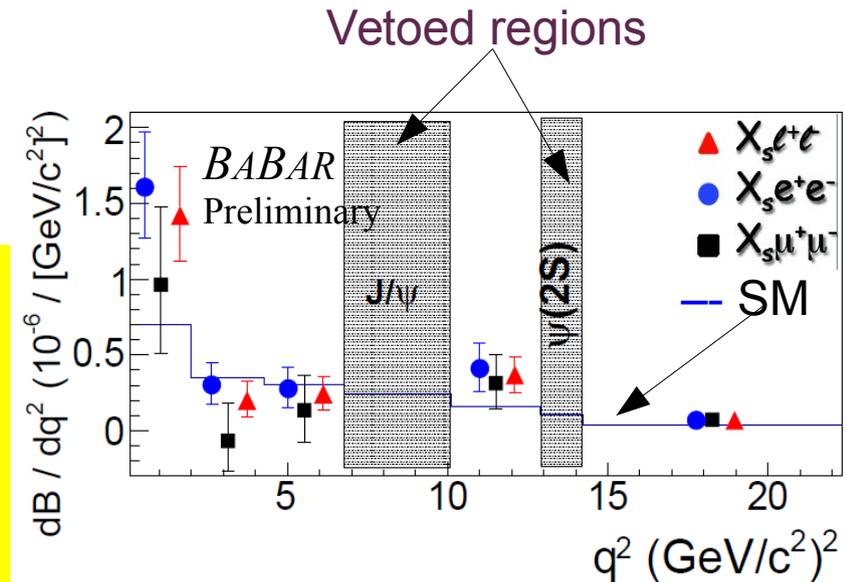
(1) Nucl.Phys.B 685, 351 (2004)

(2) Nucl.Phys.B 802, 40 (2008)

(3) PRL 111, 191801 (2013)

Liang Sun

In the listed results, the first uncertainty is statistical, the second experimental systematics and the third model-dependent systematics



# Results on direct $A_{CP}$

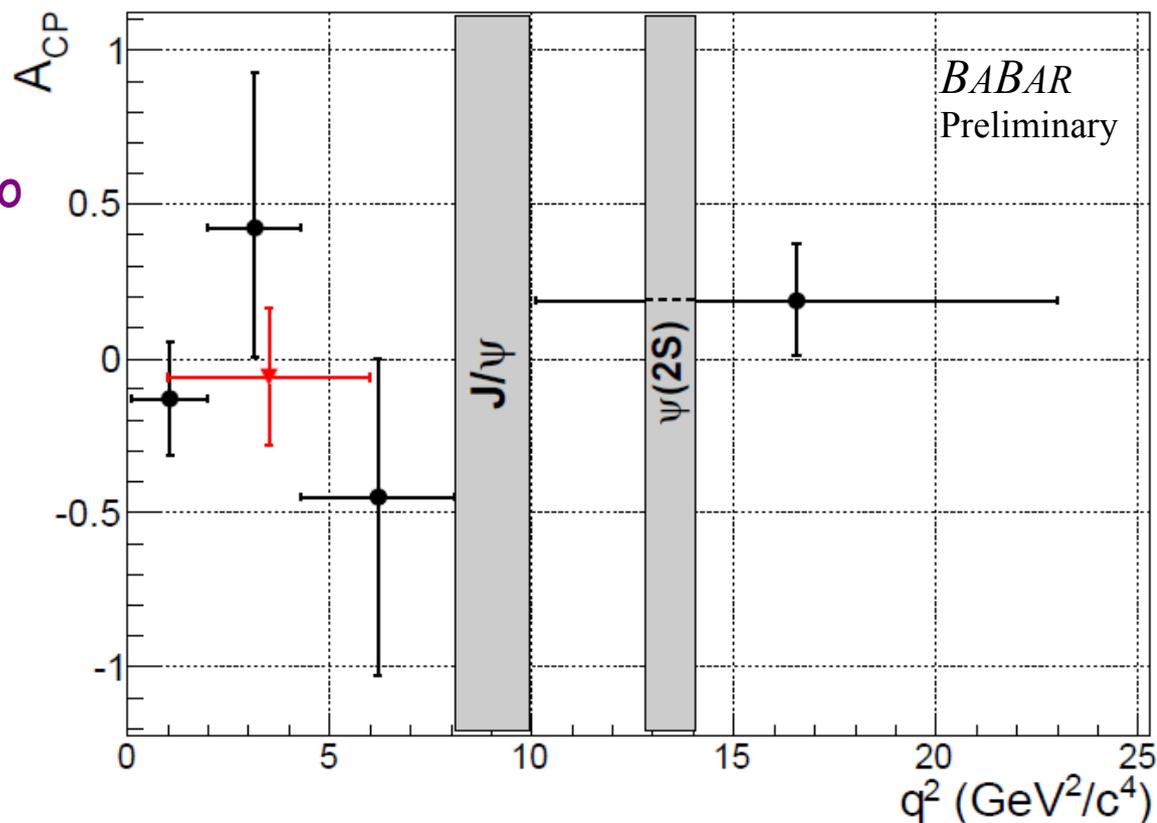
- Over the full dilepton mass range, we have:

*BABAR Preliminary*

$$A_{CP}(B \rightarrow X_s \ell^+ \ell^-) = 0.04 \pm 0.11[\text{stat}] \pm 0.01[\text{syst}]$$

- Direct  $A_{CP}$  is measured in different regions of  $q^2$ :

➤ *CP asymmetry is found to be consistent with zero everywhere*

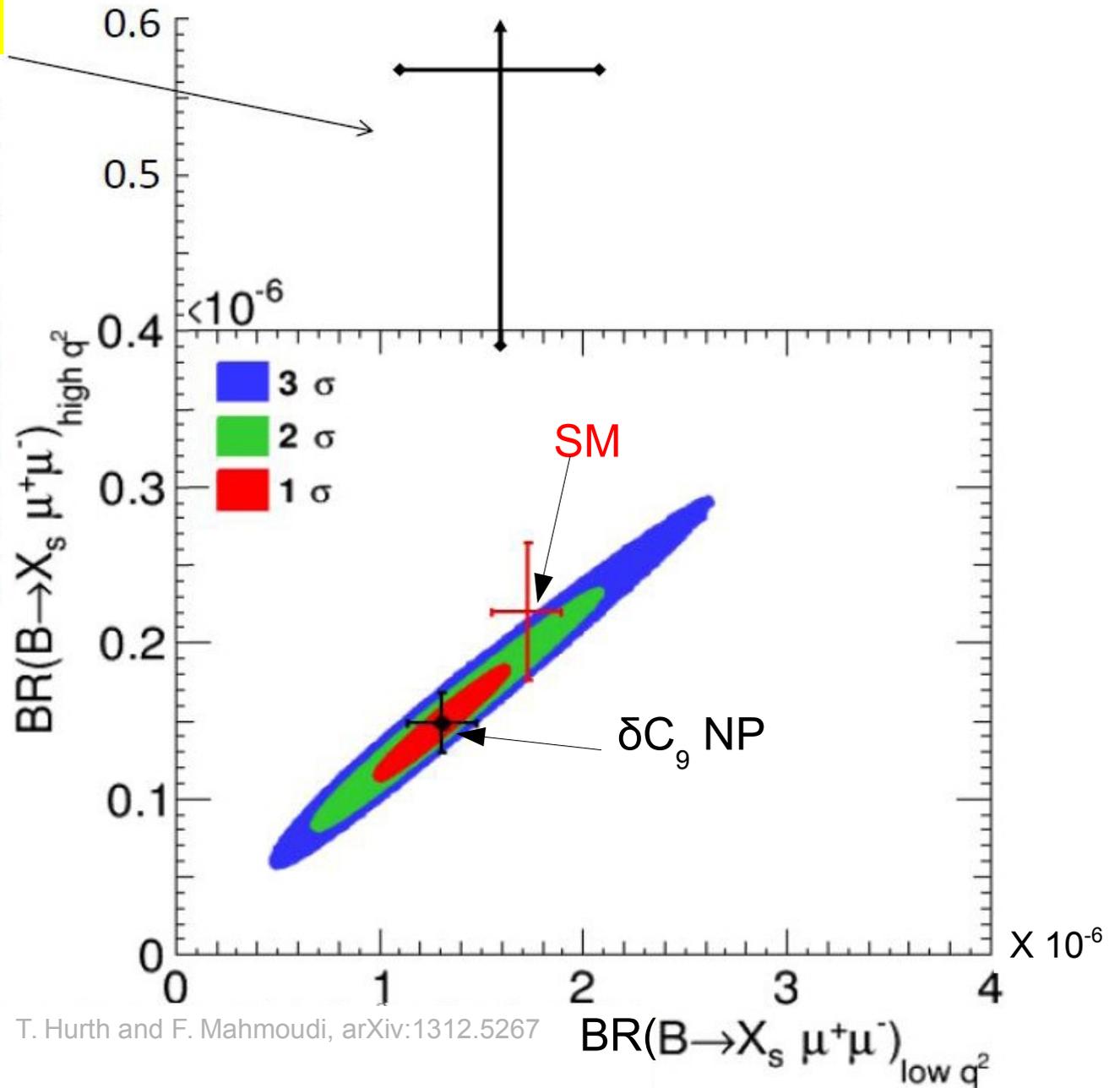
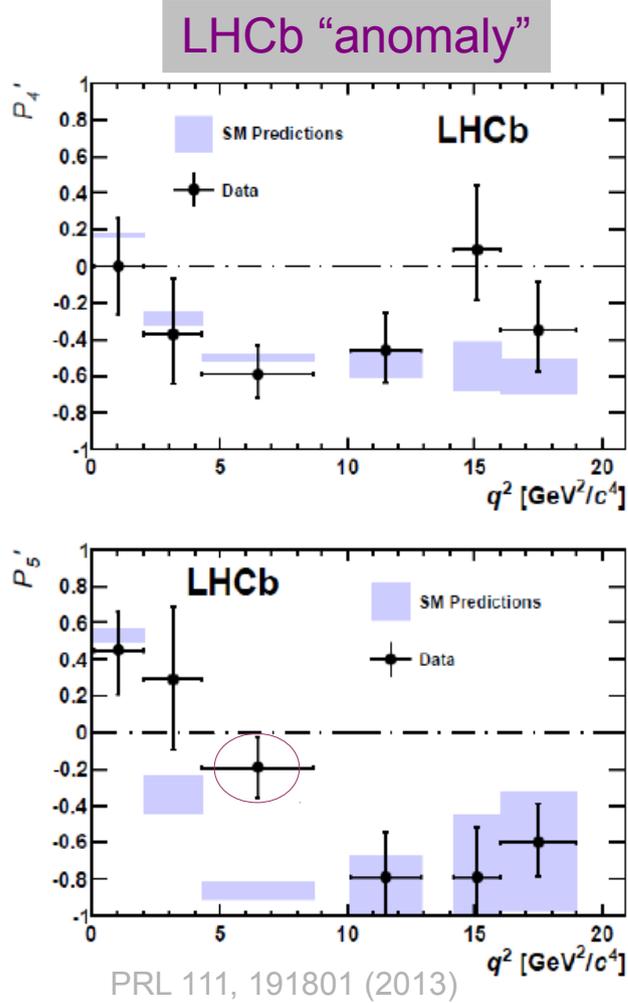


# Summary

- With the full data set of 471 M  $B\bar{B}$  pairs, BABAR continues to produce exciting results in radiative B decays
- $A_{CP}(B \rightarrow X_s \gamma)$  is the most precise result to date and agrees with the SM
- Mixing-induced CPV parameter for  $B^0 \rightarrow K_s \rho^0 \gamma$  is found to be compatible with the SM
- $\mathcal{B}(B \rightarrow X_s \ell^+ \ell^-)$  and  $A_{CP}(B \rightarrow X_s \ell^+ \ell^-)$  are generally consistent with the SM, however some tensions exist for the total  $\mathcal{B}$  and partial  $\mathcal{B}$  with  $q^2 > 14.2 \text{ GeV}^2$ 
  - The high  $q^2$   $\mathcal{B}$  disfavors the  $\delta C_9$  NP interpretation of LHCb “anomaly”
- To continue searching for NP, high-statistics data sample from Belle II and LHCb are needed for conclusive answers

# Backup Slides

# Our $\mathcal{B}(B \rightarrow X_s \ell^+ \ell^-)$ results



T. Hurth and F. Mahmoudi, arXiv:1312.5267

TABLE I:  $B \rightarrow X_s e^+ e^-$ ,  $B \rightarrow X_s \mu^+ \mu^-$  and  $B \rightarrow X_s \ell^+ \ell^-$  partial BFs (in units of  $10^{-6}$ ) and  $A_{CP}$  by  $q^2$  ( $\text{GeV}^2/c^4$ ) and  $m_{X_s}$  ( $\text{GeV}/c^2$ ) bin. The number in parentheses after each result is the multiplier which is applied to the measured semi-inclusive rate to account for unreconstructed and  $m_{X_s} > 1.8 \text{ GeV}/c^2$  final states. Estimated contributions from the vetoed charmonium  $q^2$  regions are included in both the total and  $m_{X_s}$  binned results, but not in the total  $A_{CP}$ . The first uncertainties are statistical, the second experimental systematics and the third model-dependent systematics associated with the multiplicative factor. There are no model-dependent  $A_{CP}$  systematics and  $A_{CP}$  is not measured as a function of  $m_{X_s}$ ; the multiplicative factors are not used in calculating the total  $A_{CP}$ .

Bin	Range	$B \rightarrow X_s e^+ e^-$	$B \rightarrow X_s \mu^+ \mu^-$	$B \rightarrow X_s \ell^+ \ell^-$	$A_{CP B \rightarrow X_s \ell^+ \ell^-}$
$q_0^2$	$1.0 < q^2 < 6.0$	$1.93^{+0.47+0.21}_{-0.45-0.16} \pm 0.18$ (1.71)	$0.66^{+0.82+0.30}_{-0.76-0.24} \pm 0.07$ (1.78)	$1.60^{+0.41+0.17}_{-0.39-0.13} \pm 0.18$	$-0.06 \pm 0.22 \pm 0.01$
$q_1^2$	$0.1 < q^2 < 2.0$	$3.05^{+0.52+0.29}_{-0.49-0.21} \pm 0.35$ (1.96)	$1.83^{+0.90+0.30}_{-0.80-0.24} \pm 0.20$ (2.02)	$2.70^{+0.45+0.21}_{-0.42-0.16} \pm 0.35$	$-0.13 \pm 0.18 \pm 0.01$
$q_2^2$	$2.0 < q^2 < 4.3$	$0.69^{+0.31+0.11}_{-0.28-0.07} \pm 0.07$ (1.73)	$-0.15^{+0.50+0.26}_{-0.43-0.14} \pm 0.01$ (1.80)	$0.46^{+0.26+0.10}_{-0.23-0.06} \pm 0.07$	$0.42^{+0.50}_{-0.42} \pm 0.01$
$q_3^2$	$4.3 < q^2 < 6.8$	$0.69^{+0.31+0.13}_{-0.29-0.10} \pm 0.05$ (1.53)	$0.34^{+0.54+0.19}_{-0.50-0.15} \pm 0.03$ (1.59)	$0.60^{+0.27+0.10}_{-0.25-0.08} \pm 0.05$	$-0.45^{+0.44}_{-0.57} \pm 0.01$
$q_4^2$	$10.1 < q^2 < 12.9$	$1.14^{+0.42+0.22}_{-0.40-0.10} \pm 0.04$ (1.16)	$0.87^{+0.51+0.11}_{-0.47-0.08} \pm 0.03$ (1.18)	$1.02^{+0.32+0.10}_{-0.30-0.07} \pm 0.04$	
$q_5^2$	$14.2 < q^2$	$0.56^{+0.19+0.03}_{-0.18-0.03} \pm 0.00$ (1.02)	$0.60^{+0.31+0.05}_{-0.29-0.04} \pm 0.00$ (1.02)	$0.57^{+0.16+0.03}_{-0.15-0.02} \pm 0.00$	
$q_{45}^2$	$q_4^2 \cup q_5^2$	—	—	—	$0.19^{+0.18}_{-0.17} \pm 0.01$
$m_{X_s,1}$	$0.4 < m_{X_s} < 0.6$	$0.69^{+0.18+0.04}_{-0.17-0.03} \pm 0.00$ (1.00)	$0.74^{+0.25+0.04}_{-0.23-0.04} \pm 0.00$ (1.00)	$0.71^{+0.15+0.03}_{-0.14-0.03} \pm 0.00$	
$m_{X_s,2}$	$0.6 < m_{X_s} < 1.0$	$1.20^{+0.34+0.10}_{-0.33-0.07} \pm 0.00$ (1.00)	$0.76^{+0.44+0.08}_{-0.40-0.07} \pm 0.00$ (1.00)	$1.02^{+0.27+0.06}_{-0.25-0.05} \pm 0.00$	
$m_{X_s,3}$	$1.0 < m_{X_s} < 1.4$	$1.60^{+0.72+0.27}_{-0.69-0.19} \pm 0.05$ (1.18)	$0.65^{+1.16+0.27}_{-1.08-0.25} \pm 0.02$ (1.18)	$1.32^{+0.61+0.19}_{-0.58-0.15} \pm 0.05$	
$m_{X_s,4}$	$1.4 < m_{X_s} < 1.8$	$1.88^{+0.76+0.71}_{-0.73-0.47} \pm 0.12$ (1.91)	$0.19^{+1.35+0.70}_{-1.25-0.50} \pm 0.10$ (1.91)	$1.36^{+0.67+0.50}_{-0.63-0.34} \pm 0.12$	
Total	$0.1 < q^2$	$7.69^{+0.82+0.50}_{-0.77-0.33} \pm 0.50$	$4.41^{+1.31+0.57}_{-1.17-0.42} \pm 0.27$	$6.73^{+0.70+0.34}_{-0.64-0.25} \pm 0.50$	$0.04 \pm 0.11 \pm 0.01$