

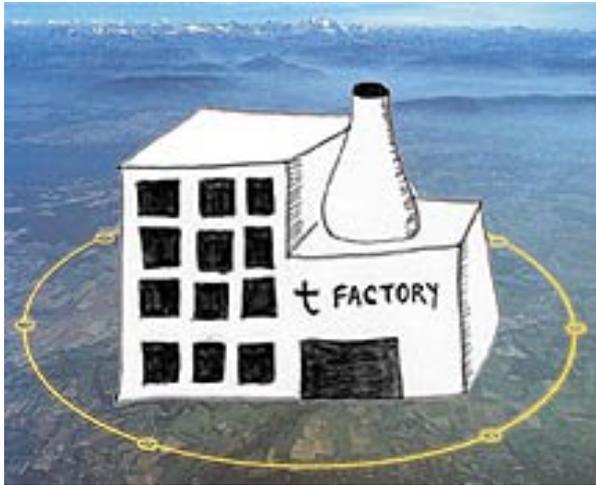


# Single top production cross section using the ATLAS detector at the LHC

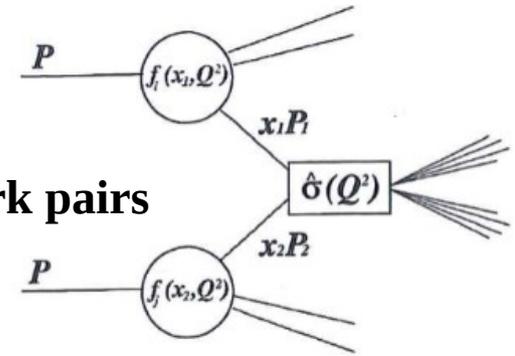
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on the behalf of the ATLAS collaboration  
LPSC, Université Grenoble-Alpes CNRS/IN2P3

DIS 2014  
Warsaw 29/4/2014

# Top production @ the LHC

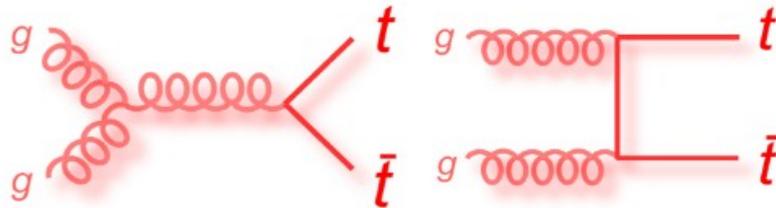


Proton-proton collisions  
at  $\sqrt{s} = 7$  TeV (2011) and 8 TeV (2012)  
generated **more than 5 millions of top quark pairs**  
via strong interaction

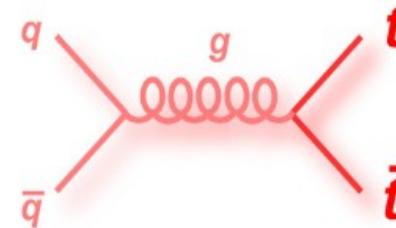


## TOP PAIR PRODUCTION

gluon-gluon fusion (90%)



quark-quark annihilation (10%)

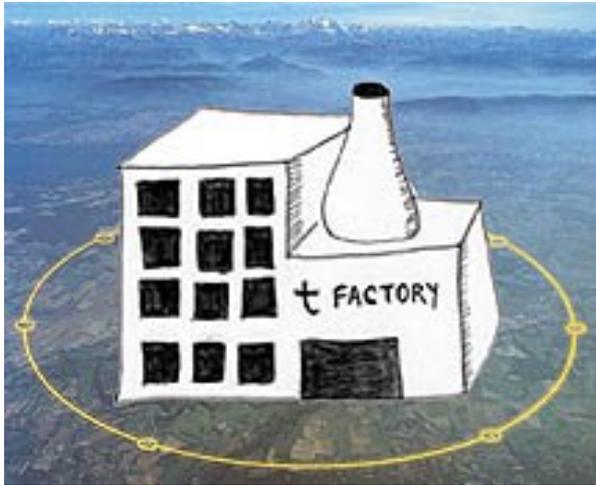


Predicted cross section at NNLO in QCD  
including resummation of NNLL soft gluon terms with top++2.0  
 $m(t)c^2 = 172.5$  GeV, uncertainty accounting for PDF and scale variations

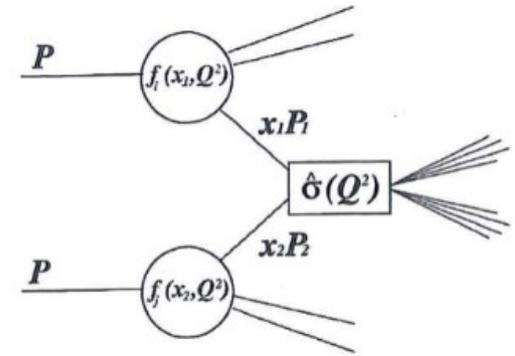
	LHC @ 7 TeV	LHC @ 8 TeV
$\sigma_{t\bar{t}}$	$177^{+10}_{-11}$ pb	$253^{+13}_{-15}$ pb

(Phys Lett B 710, 2012)

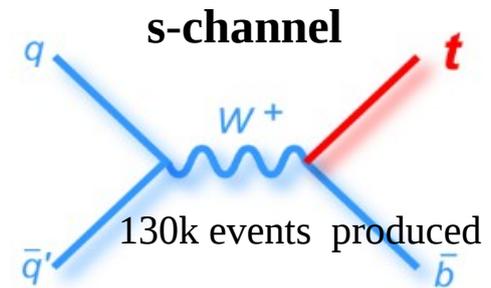
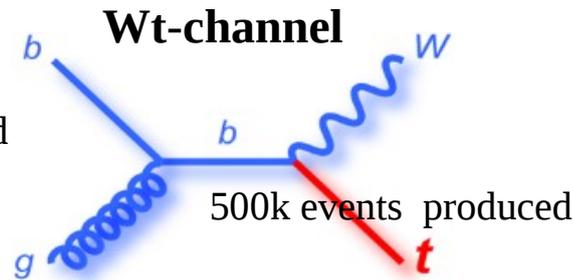
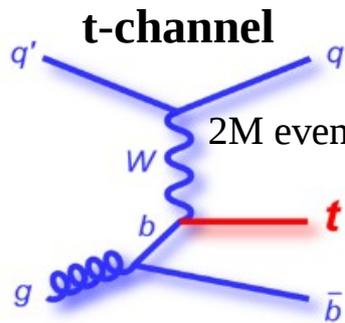
# Top production @ the LHC



Proton-proton collisions at  $\sqrt{s}=7$  TeV (2011) and 8 TeV (2012) generated **more than 2.5 millions of top quarks via electroweak interaction**



## SINGLE TOP PRODUCTION



Predicted cross sections at NLO in QCD + NNLO soft-gluon corrections  $m(t)c^2=172.5$  GeV, uncertainty accounting for PDF and scale variations

	LHC @ 7 TeV	LHC @ 8 TeV	
(Phys. Rev. D 83, 091503(R), 2011)	$\sigma_t$	$64.57^{+2.63}_{-1.74}$ pb	$87.76^{+3.44}_{-1.91}$ pb
	$\sigma_{Wt}$	$15.74^{+1.17}_{-1.21}$ pb	$22.37^{+1.52}_{-1.52}$ pb
(Phys. Rev. D 81, 054028, 2010)	$\sigma_s$	$4.63^{+0.20}_{-0.18}$ pb	$5.61^{+0.21}_{-0.21}$ pb

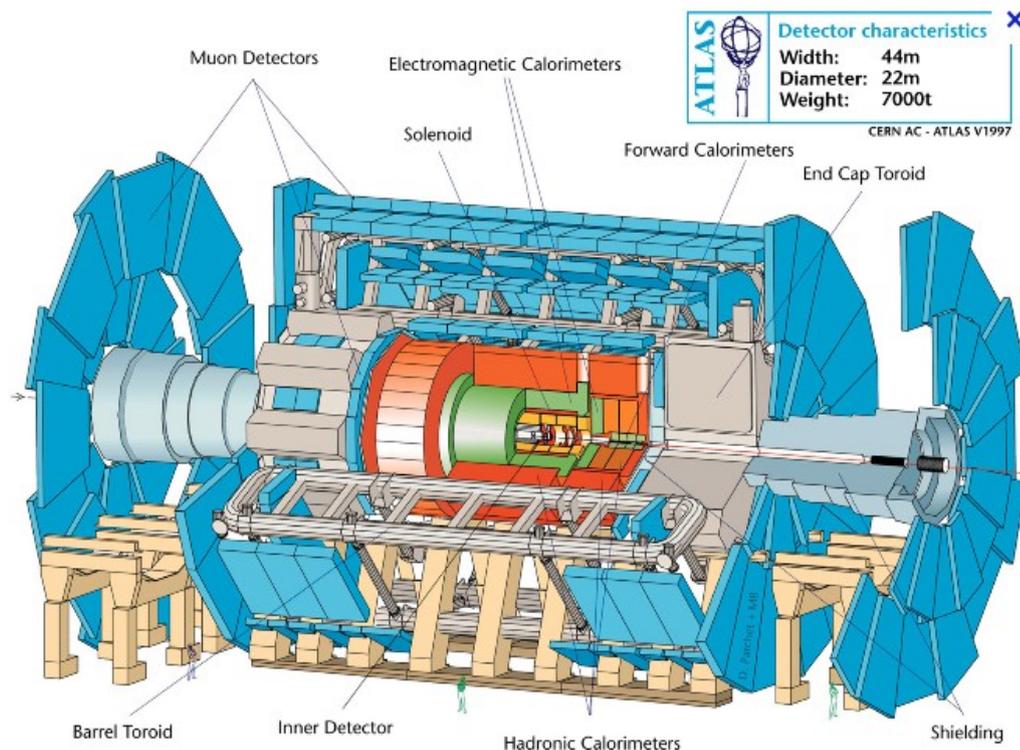
(Phys. Rev. D 82, 054018, 2010)

# Single top motivation

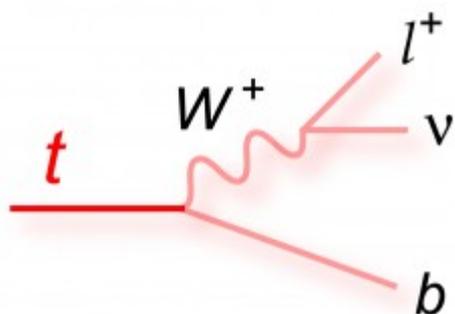
## Why looking for single top with the ATLAS detector?

- ▶ Precision test of the Standard Model
  - t-channel ,  $Wt$  associated production, s-channel cross section measurements
  - extraction of the CKM matrix element  $|V_{tb}|$
- ▶ PDF constraints
  - $t\bar{t}$  cross section ratio
- ▶ Direct sensitivity to new physics
  - Flavor Changing Neutral Currents (t-channel final state)
  - Heavy resonances i.e  $W'$  boson (s-channel final state)

only analyses reported here are cited



# Object modelling



**Electrons:** isolated electromagnetic calorimeter clusters matched to inner detector tracks

**Muons:** muon spectrometer  $p_T$  associated to ID track

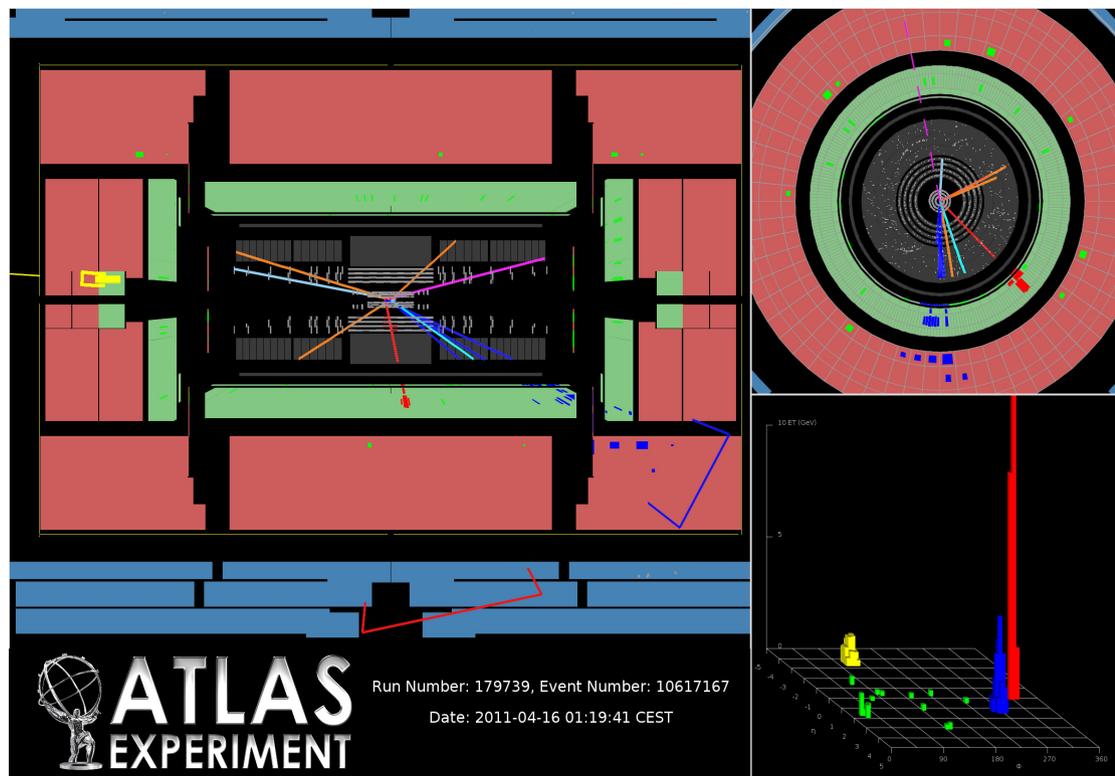
**Associated experimental systematic uncertainties:**  
Lepton energy/momentum scale & resolution  
Lepton trigger & ID efficiency

$E_T^{\text{miss}}$ : extracted from total energy for particles travelling perpendicular to the beam axis

Missing momentum scale & resolution

**b-tagged jets:** hadronic calorimeter clusters + multivariate combination of topological variables and b-tagging algorithms

Jet energy scale, several component (JES)  
Jet energy resolution, reconstruction efficiency  
Jet vertex fraction, b-tagging and mis-tagging efficiency



# Precision measurements

# t-channel cross section

8 TeV analysis,  $\int L dt = 20.3 \text{ fb}$

(ATLAS-CONF-2014-007)

## Event selection

- 1 isolated  $e/\mu$ ,  $p_T > 25 \text{ GeV}$
- 2 jets, 1 b-tagged,  $p_T > 30 \text{ GeV}$
- $E_T^{\text{miss}} > 30 \text{ GeV}$ ,  $m_T(W) > 50 \text{ GeV}$

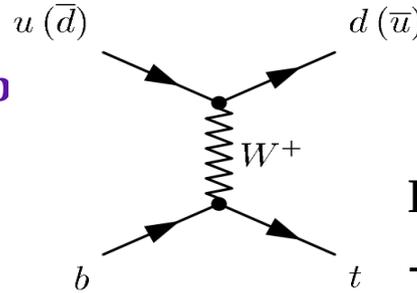
## Multivariate analysis

t-channel discrimination realized by combining 14 kinematic variables into a Neural Network discriminant

## Systematic uncertainty

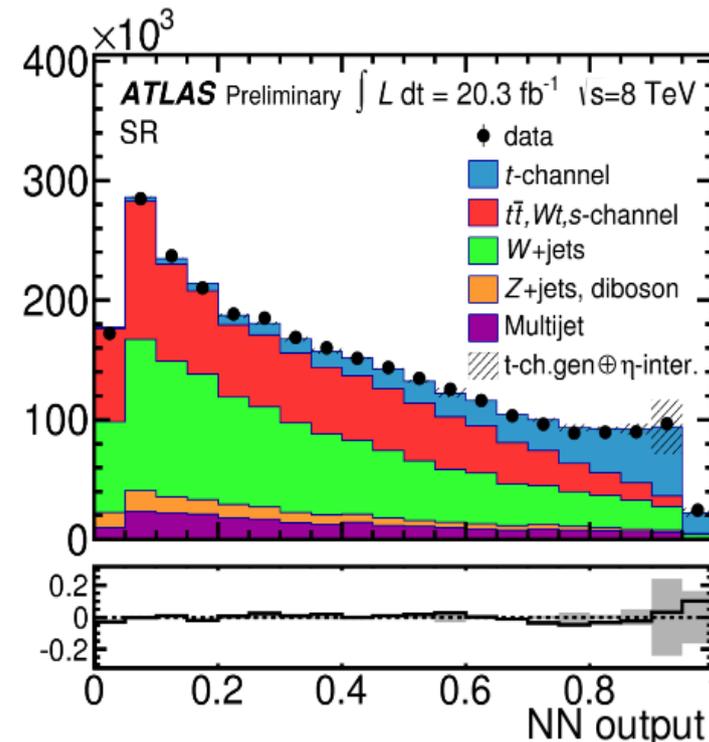
Pseudo experiments estimate the impact of background normalization, luminosity, object modelling (JES  $\eta$  intercalibration being the greatest), PDF & generator.

↓  
New procedure to minimize them



## Backgrounds

- Dominant contributions from  $t\bar{t}$  and  $W$ +jets events, examined in specific control regions
- multijet data driven normalization
  - $E_T^{\text{miss}}$  fit for  $e$ +jets final state,
  - matrix method for  $\mu$ +jets final state



# t-channel cross section

## Fiducial cross section

$$\sigma_{\text{fid}} = \frac{\epsilon_{\text{corr,sel}}}{\epsilon_{\text{corr,fid}}} \cdot \frac{\hat{\nu}}{\mathcal{L}} = R_f \cdot \frac{\hat{\nu}}{\mathcal{L}}$$

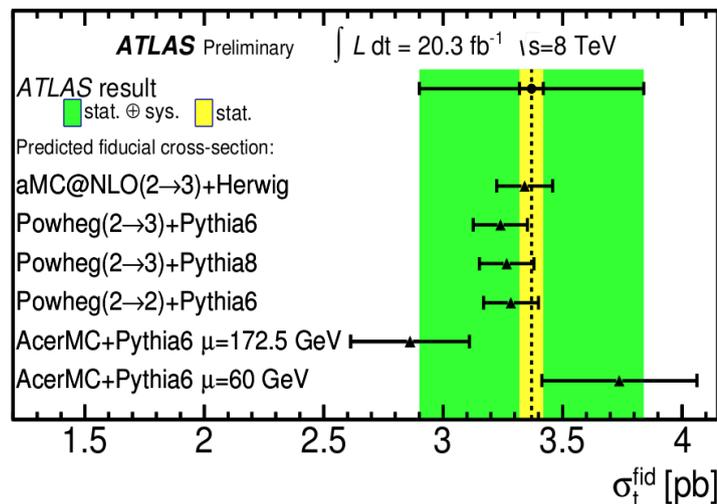
Fraction of events selected by offline selection which are in the fiducial volume (vice-versa)

Fit to the maximum likelihood to the NN output distribution

Fiducial volume within the detector acceptance: stable particles with lifetime > 30 ps

Object	Cut
Electrons	$p_T > 25 \text{ GeV}$ and $ \eta  < 2.5$
Muons	$p_T > 25 \text{ GeV}$ and $ \eta  < 2.5$
Jets	$p_T > 30 \text{ GeV}$ and $ \eta  < 4.5$
Lepton ( $\ell$ ), Jets ( $j_i$ )	$p_T > 35 \text{ GeV}$ , if $2.75 <  \eta  < 3.5$
	$\Delta R(\ell, j_i) > 0.4$
$E_T^{\text{miss}}$	$E_T^{\text{miss}} > 30 \text{ GeV}$
Transverse W-boson mass	$m_T(W) > 50 \text{ GeV}$
Lepton ( $\ell$ ), jet with the highest $p_T$ ( $j_1$ )	$p_T(\ell) > 40 \text{ GeV} \left(1 - \frac{\pi -  \Delta\phi(j_1, \ell) }{\pi - 1}\right)$

Measured fiducial cross-section vs predictions for different MC event generators



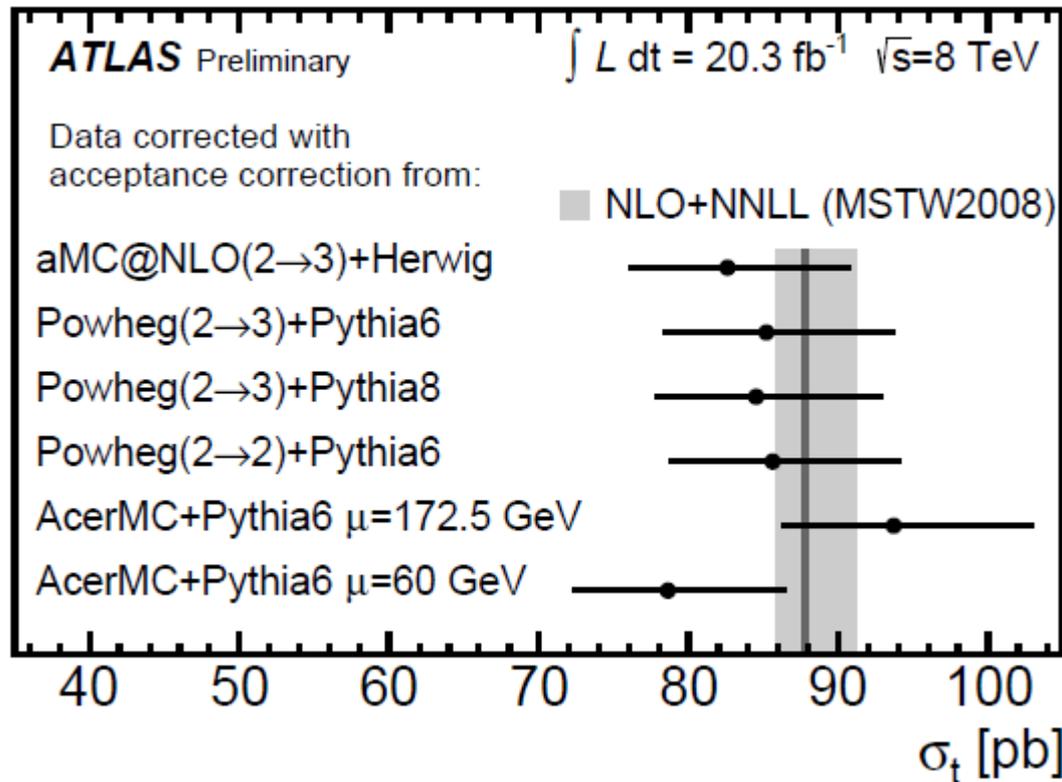
$$\sigma_{\text{fid}} = 3.37 \pm 0.05 \text{ (stat.)} \pm 0.47 \text{ (syst.)} \pm 0.09 \text{ (lumi.) pb}$$

Relative uncertainty: 14% vs 17% using the acceptances referred to the full phase space  
Dominant contributions: JES  $\eta$  intercalibration, t-channel generator

# t-channel cross section

## Extrapolation of the inclusive cross section and $Vtb$

The fiducial cross section is extrapolated to the full phase space where  $\epsilon_{\text{fid}}$  is the selection efficiency of the particle-level selection

$$\sigma = \frac{1}{\epsilon_{\text{fid}}} \cdot \sigma_{\text{fid}}$$


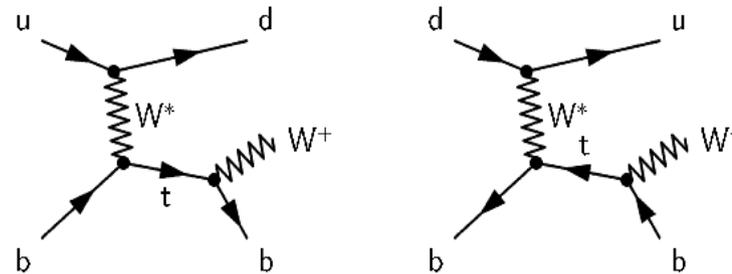
- Vertical line indicates the approximate NNLO cross-section, including the uncertainty
- Extrapolations using the acceptance of several MC generators and renormalization scales  
Uncertainty on the fiducial cross section +PDF

$$\sigma_t = 82.6 \pm 1.2 \text{ (stat.)} \pm 11.4 \text{ (syst.)} \pm 3.1 \text{ (PDF)} \pm 2.3 \text{ (lumi.) pb}$$

# t-channel cross section ratio

7 TeV analysis,  $\int L dt = 4.7 \text{ fb}^{-1}$

(ATLAS-CONF-2012-056)



## Motivations

- $R_t = \sigma_t(t)/\sigma_t(\bar{t})$  sensitive to the ratio of u-quark to d-quark PDFs in the range  $0.02 < x < 0.5$ 
  - further constraint on PDFs (complementary to W asymmetry)
- Smaller uncertainties than the total cross section because of partial cancellations

## Almost the same strategy than 8TeV t-channel, except:

- Inclusive selection of 2 or 3 jets events, looser cut  $m_T(W) > 30 \text{ GeV}$
- Data driven background estimates: multijet & W+jets
- Neural Network training done individually for 4 channels: 2jets+ $l^{+/-}$ , 3jets+ $l^{+/-}$
- 4 NN discriminants fitted simultaneously

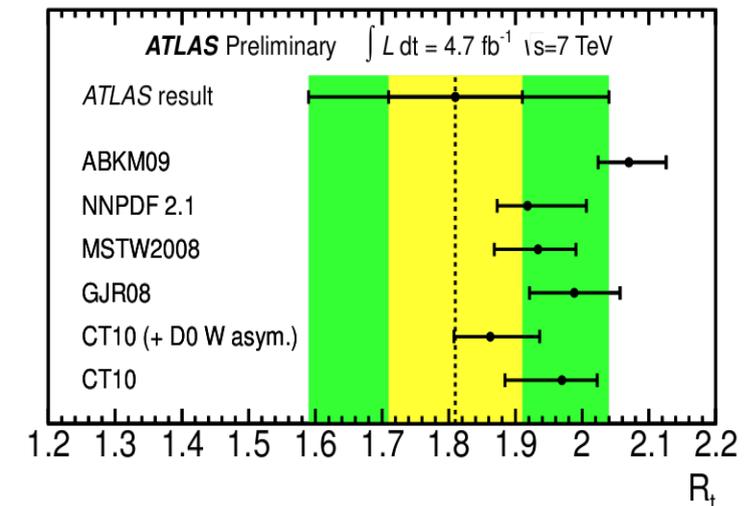
## Results

$$\sigma_t(t) = 53.2 \pm 1.7 \text{ (stat.)} \pm 10.6 \text{ (syst.) pb} = 53.2 \pm 10.8 \text{ pb}$$

$$\sigma_t(\bar{t}) = 29.5 \pm 1.5 \text{ (stat.)} \pm 7.3 \text{ (syst.) pb} = 29.5^{+7.4}_{-7.5} \text{ pb.}$$

$$R_t = 1.81 \pm 0.10 \text{ (stat.)}^{+0.21}_{-0.20} \text{ (syst.)} = 1.81^{+0.23}_{-0.22}$$

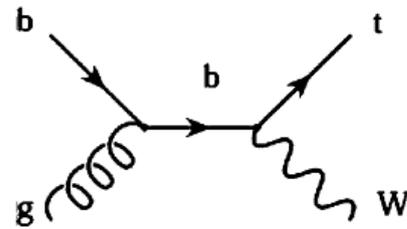
Dominant systematic uncertainties: JES, btag, generator



# Wt-channel cross section

8 TeV analysis,  $\int L dt = 20.3 \text{ fb}^{-1}$

(ATLAS-CONF-2013-100)



## Event selection

- 2 isolated opposite sign leptons  $p_T > 25 \text{ GeV}$   
 $e\mu$  final state only (best signal purity)
- 1 or 2 jets, 1 b-tagged,  $p_T > 30 \text{ GeV}$

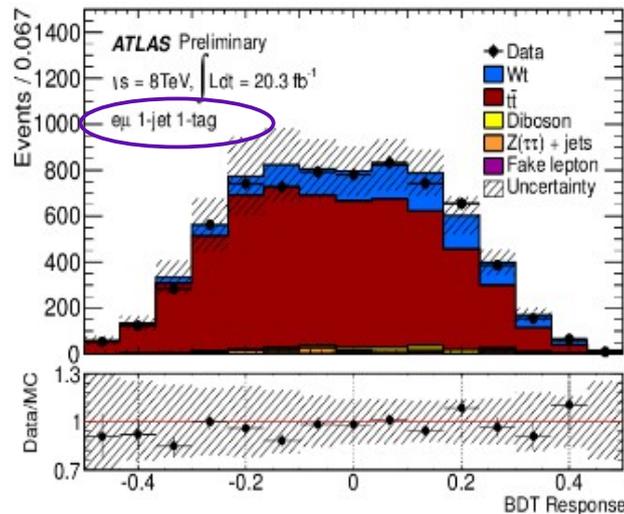
## Backgrounds

- Big contribution from top pair production,  $Z(\tau\tau)$ +jets & diboson events contribute also
- MC modeling for every source except multijet, estimated via the matrix method

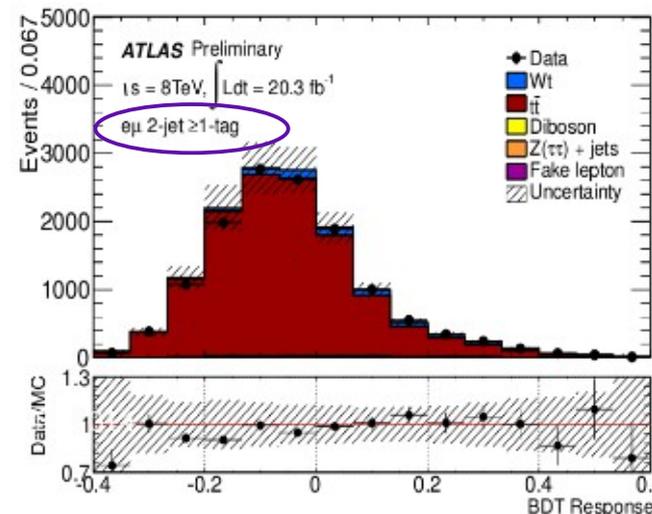
## Multivariate analysis

Boosted Decision Trees optimized separately for 1 jet and 2 jets events to account for differences in topology and kinematics

19 variables



20 variables



# Wt-channel cross section

## Results

Simultaneous maximum likelihood fit to the two BDT output distributions

$$\sigma_{wt} = 27.2 \pm 2.8 \text{ (stat)} \pm 5.4 \text{ (syst)} \text{ pb}$$

$$\text{Significance} = 4.2 \sigma$$

Main systematic uncertainties:

JES detector modelling component, btag,  $E_t^{\text{miss}}$  scale (all profiled), generator and parton shower

7 TeV analysis,  $\int L dt = 4.7 \text{ fb}^{-1}$

(Phys.Lett. B 716 2012)

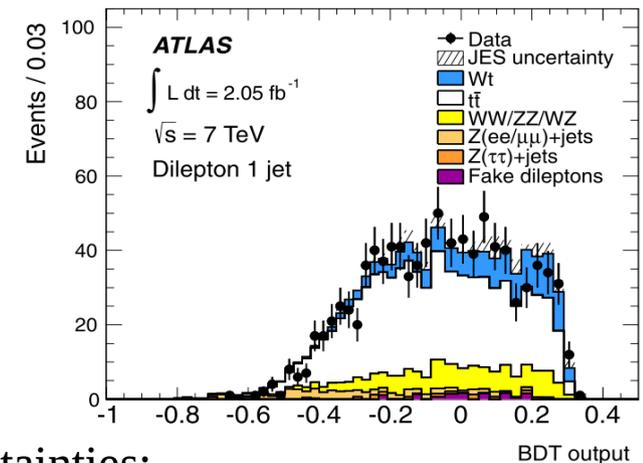
**Similar strategy than 8TeV Wt-channel, different selection & background modeling:**

- $e\mu/ee/\mu\mu$  opposite sign dilepton categories, cuts to reduce Z decays contamination
- 1, 2,  $\geq 3$  jets events to control the main backgrounds, no b-tagging requirement
- $E_T^{\text{miss}} > 50 \text{ GeV}$
- data-driven estimates for Drell-Yan processes

Simultaneous maximum likelihood fit of 3 BDTs distributions (optimized per jet-bin) → **FIRST EVIDENCE!**

$$\sigma_{wt} = 16.8 \pm 2.9 \text{ (stat)} \pm 4.9 \text{ (syst)} \text{ pb}$$

$$\text{Significance} = 3.4 \sigma$$

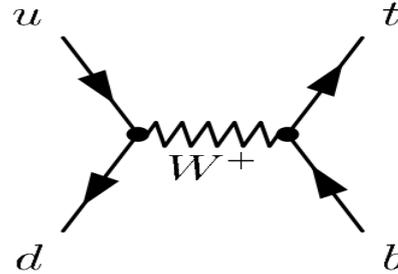


Main systematic uncertainties:

JES, btag, pile-up, generator and parton shower

# s-channel cross section

s-channel is the most challenging production mode at the LHC since it proceeds via  $q\bar{q}$  annihilation



7 TeV analysis,  $\int L dt = 0.7 \text{ fb}^{-1}$  (ATLAS-CONF-2011-118)

## Event preselection

- 1 isolated  $e/\mu$ ,  $p_T > 25 \text{ GeV}$
- 2 central jets, both b-tagged,  $p_T > 25 \text{ GeV}$
- $E_T^{\text{miss}} > 30 \text{ GeV}$ ,  $m_T(W) > 60 \text{ GeV}$  -  $E_T^{\text{miss}}$

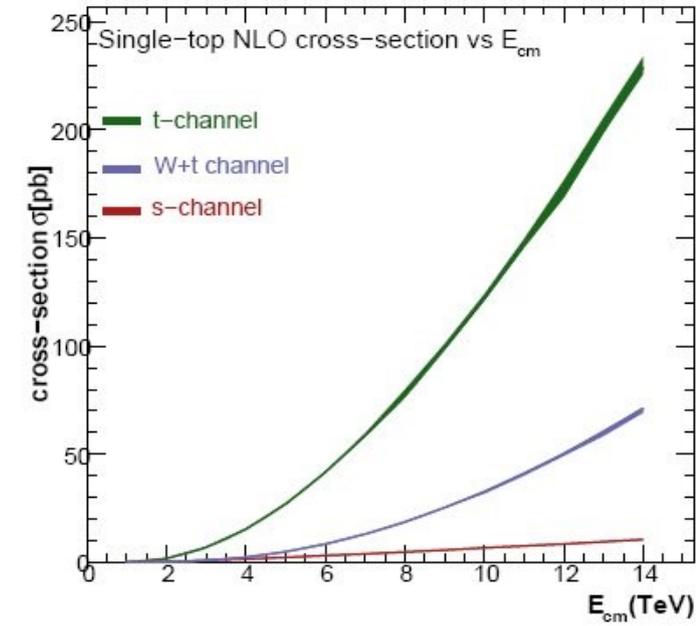
## Cut-based analysis

Most discriminating variables [  $m_T(W)$ ,  $p_T(\text{jet1}, \text{jet2})$ ,  $m(\text{top}, \text{jet1}/2)$ ,  $\Delta R(\text{jet1}, \text{jet2})$ ,  $\Delta R(\text{jet1}, l)$  ]  
employed to increase  $S/\sqrt{B}$  from 0.26 to 0.98

## Results

Counting experiment modelled by a likelihood function fitted to the data

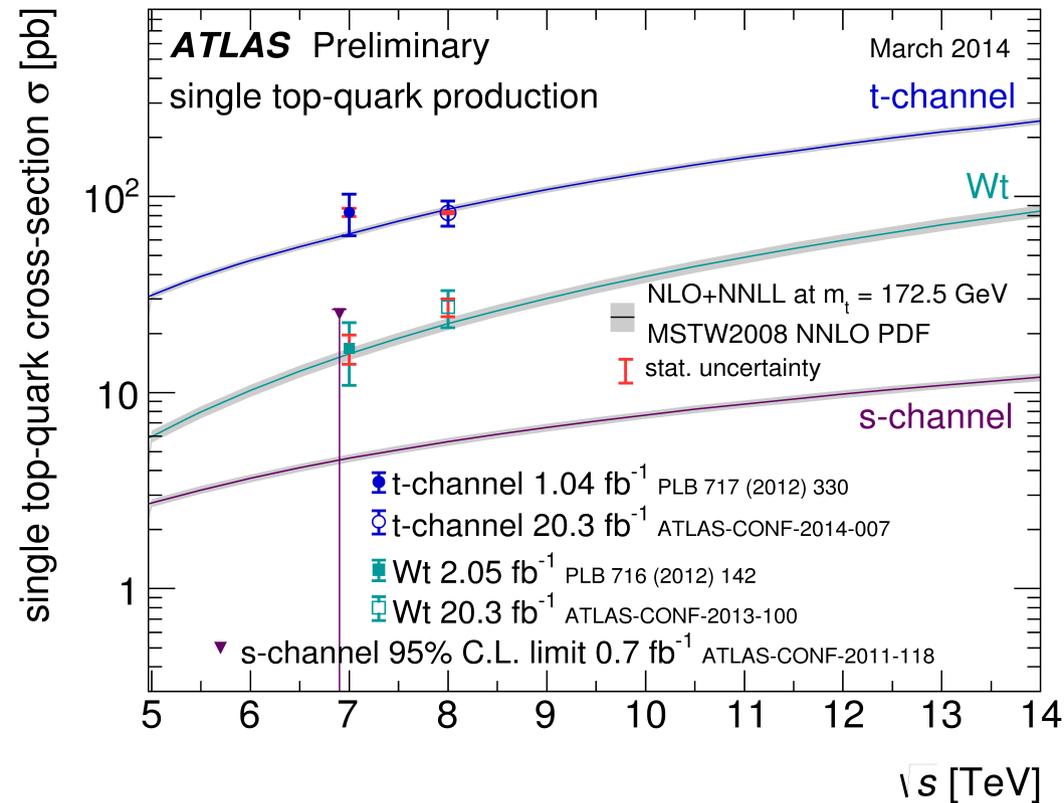
→ upper limit @ 95% CL:  $\sigma_s < 26.5 \text{ pb}$



## Backgrounds

- huge contamination from top pair production and  $W$ +jets
- data-driven estimates for multijet and  $W$ +jets events

# single top measurements



Cross section measurements for the 3 channels:

- no deviation wrt SM expectations
- constraints on BSM physics, since the production mechanisms are separately susceptible to new physics scenarios
- constraints on W-t-b vertex, which appears in the production and decay



Direct  $|V_{tb}|$  measurement, assuming it is  $\gg |V_{ts}|, |V_{td}|$

independently on the number of quark generations & the CKM matrix unitarity:

$$|V_{tb} f|^2 = \sigma_{\text{measured}} / \sigma_{\text{predicted}}$$

with  $f=1$  in the SM



$ V_{tb} $	t-channel	Wt-channel
$\sqrt{s}=7$ TeV	$1.13^{+0.14}_{-0.13}$ (11.9 %)	$1.03^{+0.16}_{-0.19}$ (17.0 %)
$\sqrt{s}=8$ TeV	$0.97^{+0.09}_{-0.10}$ (9.8 %)	$1.10 \pm 0.12$ (11.2 %)

# BSM Physics

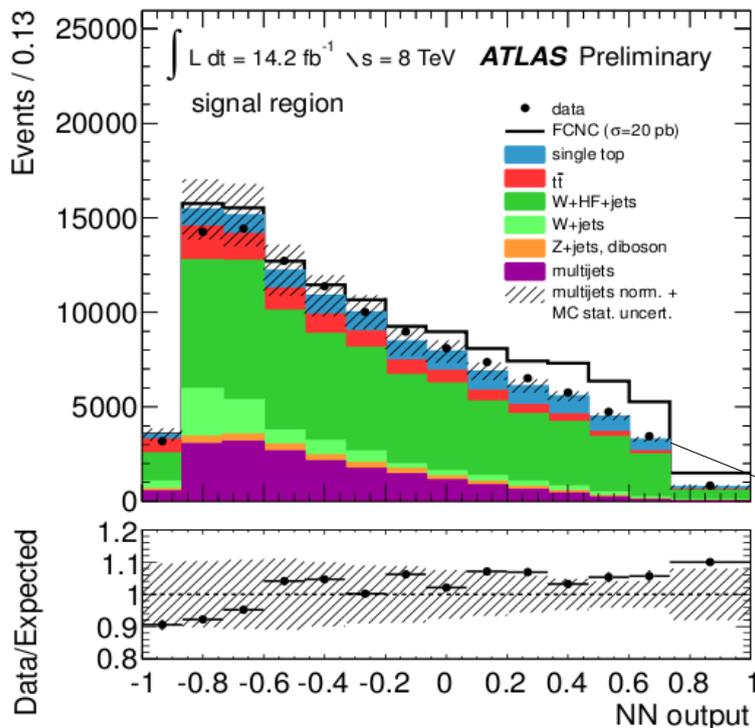
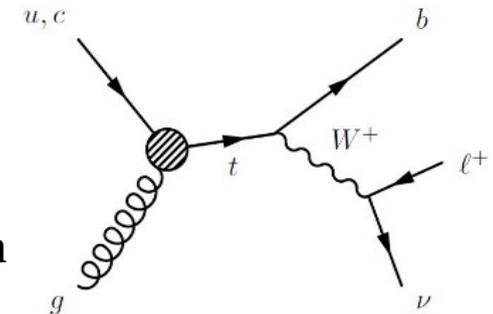
# FCNCs in single top production

8 TeV analysis,  $\int L dt = 14.2 \text{ fb}^{-1}$

(ATLAS-CONF-2013-063)

## Motivations

- FCNC highly suppressed by SM, but enhanced in many BSM scenarios
- $qg \rightarrow t \rightarrow l\nu b$  single top production has good sensitivity and distinct kinematical properties from the SM process: softer top-quark  $p_T$ , higher  $W$   $p_T$ , strongly charge asymmetric production



## Analysis strategy

- Same event selection & background modelling of the t-channel cross section analysis @ 8 TeV
- NLO generator (ME<sub>TOP</sub>) simulates strong FCNC processes
- Neural Network classifier built from 13 variables
- Binned likelihood fit to the NN output distribution to extract the FCNC contribution

No excess observed,  
here signal shape scaled to 20 pb

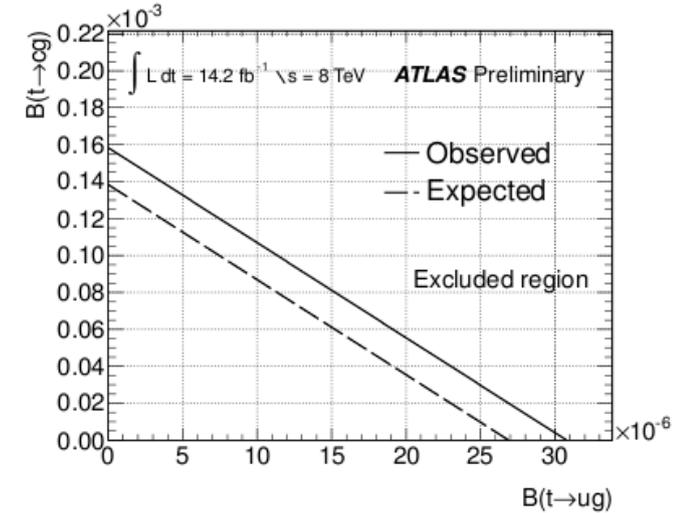
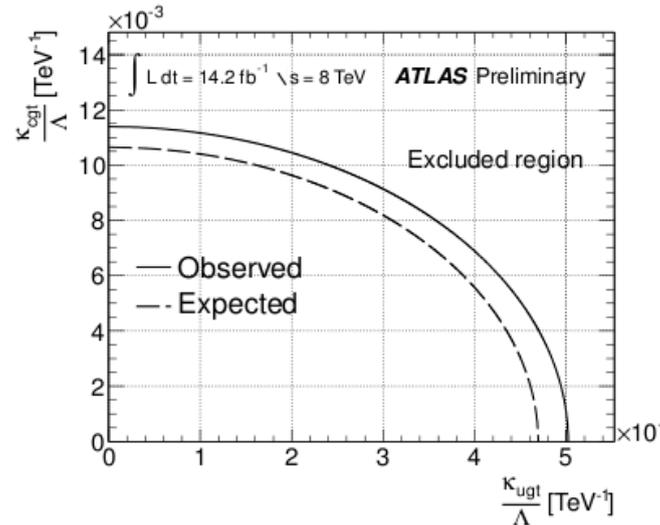
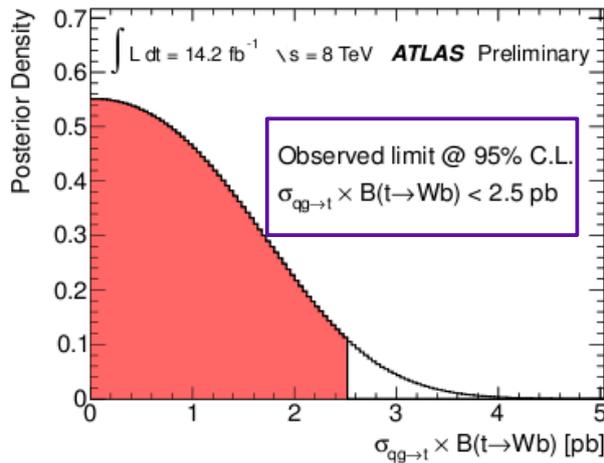
# FCNCs in single top production

Results

(A)

(B)

(C)



(A) Upper limit on cross section \* branching fraction

(B) With the hypotheses of pure left-handed coupling &  $B(t \rightarrow Wb)=1$   
 $\Rightarrow$  limits on the coupling strength  $\sigma_{qg \rightarrow t} = 100351.97 \left(\frac{\kappa_{ugt}}{\Lambda}\right)^2 + 19496.79 \left(\frac{\kappa_{cgt}}{\Lambda}\right)^2 \text{ pb}$

(C) For small coupling constants ( $k < 0.1$ )  
 $\Rightarrow$  limits on the branching fractions  $B(t \rightarrow qg) = \frac{\Gamma_{t \rightarrow qg}}{\Gamma_{t \rightarrow bW}} = 1.2205 \times \left(\frac{\kappa_{qgt}}{\Lambda}\right)^2$

7 TeV analysis,  $\int L dL = 4.7 \text{ fb}^{-1}$  (Phys.Lett. B 712 2012)

Analogous analysis, less stringent limits [ by a factor 3 on (A) and (C), by a factor 2 on (B) ]

# W' boson search

8 TeV analysis,  $\int L dt = 14.2 \text{ fb}^{-1}$

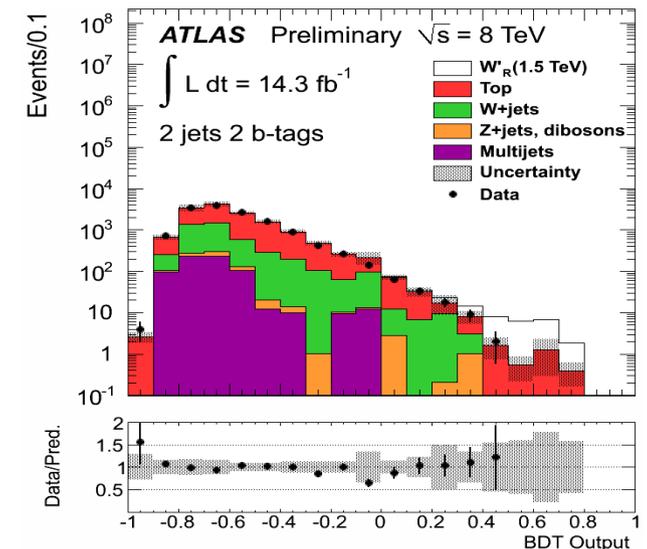
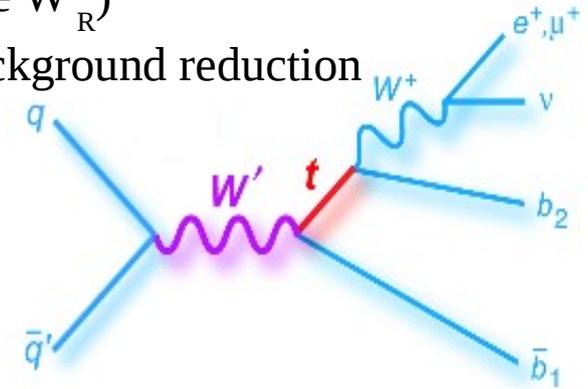
(ATLAS-CONF-2013-050)

## Motivations

- Predicted by many BSM theories
- $W' \rightarrow tb$  signature allows to assess models with leptophobic  $W'$  (i.e  $W'_R$ )
- Semileptonic final state  $\Rightarrow$  cleaner experimental signature and background reduction

## Analysis strategy

- Almost same event selection & background modelling than s-channel cross section analysis @ 7 TeV +  $m(tb) > 270 \text{ GeV}$
- LO generator (MADGRAPH) rescaled to NLO XS calculations (Sullivan Phys. Rev. D 86, 075018) provided for various  $W'$  chiralities and masses (0.5 - 3 TeV, steps of 250 GeV)
- Signal discrimination via two Boosted Decision Trees with 2-jets & 3-jets events. In the training:  $W'_R$  (1.75 TeV), which minimizes the expected limit
- Simultaneous fit of the BDT output distribution in 2 channels for each value of the hypothesized  $W'_{R/L}$  mass  
 $\Rightarrow$  CLs limit @ 95% CL

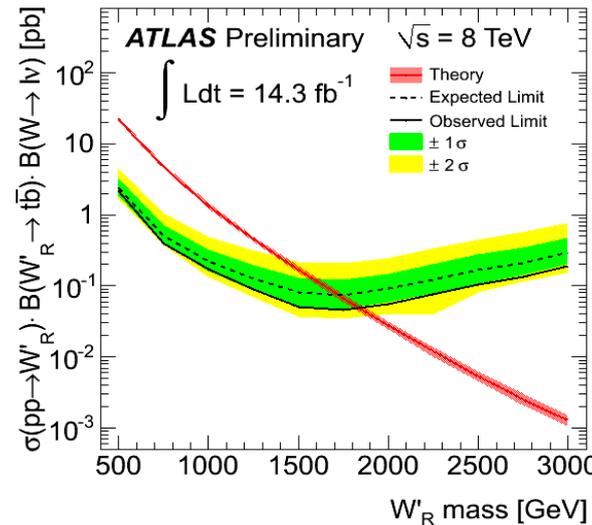
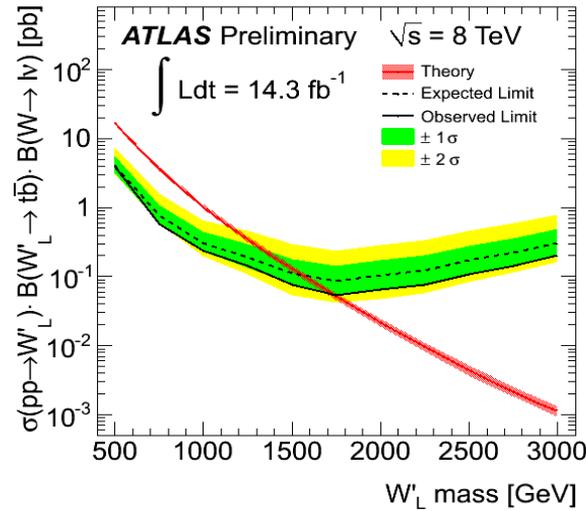


# W' boson search

## Results

Exclusion limits on W' cross section and mass:

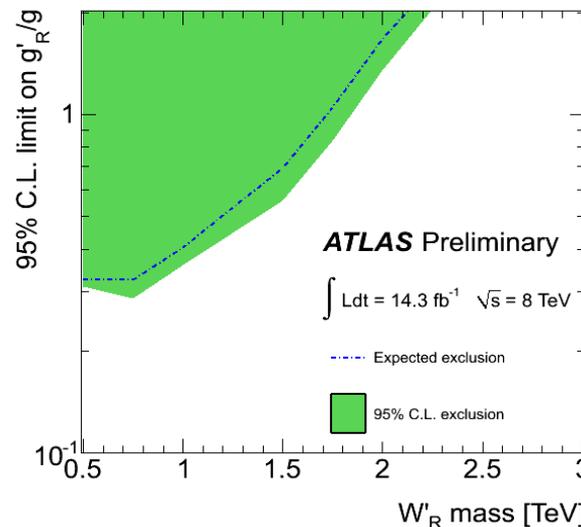
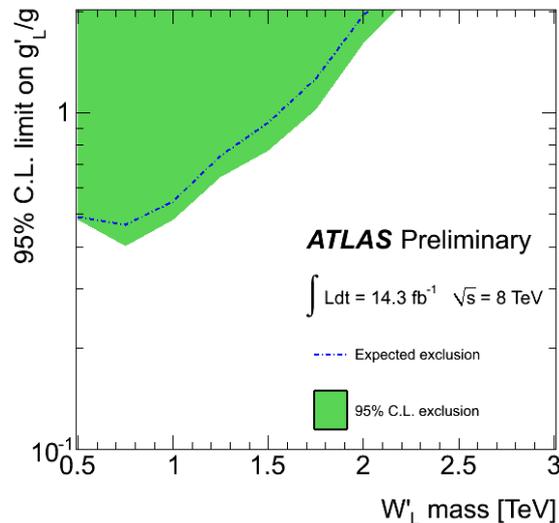
$$m(W'_R) > 1.84 \text{ TeV} \quad m(W'_L) > 1.74 \text{ TeV}$$



$$\sigma(W') \sim (g'/g_{SM})^2$$

The hypothesis  $g'$  for the mass  $m_{W'}$  is excluded if  $\sigma(g', m_{W'}) > \sigma(m_{W'})$  for a given W' chirality

Exclusion limits on W' couplings (model-independent):



Limits are shown up to a  $g'/g_{SM}$  value of 2, for which the approximation holds

# Conclusion

- ▶ Single top production in ATLAS:
  - observed with high precision in the t-channel
  - evidence in the  $Wt$  associate production
  - upper limit for the s-channel
- ▶ All the measured cross section are in overall agreement with Standard Model NLO QCD + NNLL predictions and, thus, help to constrain new physics
- ▶ BSM searches sharing single top final states show no evidence for  $W'$  bosons or FCNCs

...Next LHC phase will push the energy frontier upward,  
and the higher luminosity will enlarge the statistics  
⇒ improvement of the precision of the analyses involving t-channel and  $Wt$  final states !

# Backup slides

# Data driven bkg estimates

## Multijet: MATRIX METHOD

Modelled from tight and loose data samples merged with the MM weight (derived from the efficiencies and fake rates measured on data)

$$N_{fake}^{tight} = \frac{\epsilon_{fake}}{\epsilon_{real} - \epsilon_{fake}} (N^{loose} \epsilon_{real} - N^{tight})$$

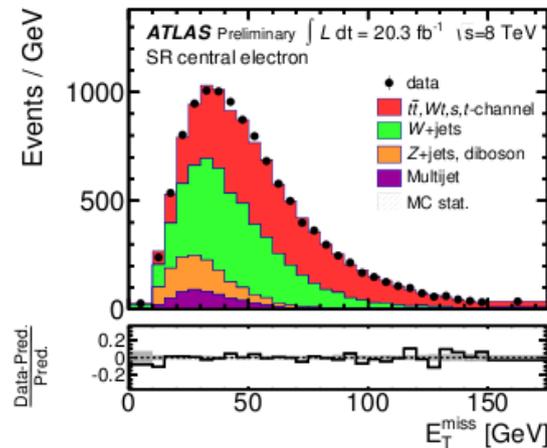
gives the normalization

$$w_{tight} = \epsilon_{fake} \frac{\epsilon_{real} - 1}{\epsilon_{real} - \epsilon_{fake}}$$

is applied for the shapes

## Multijet: $E_T^{miss}$ FIT

Normalization: binned maximum likelihood fit on the observed  $E_T^{miss}$  in a control region



- Simulated multijet events with jet-electron model for e+jets final states
- Data driven multijet events with anti-muon method in  $\mu$ +jets final states

## W+jets:

Monte Carlo template with data driven normalization:

- Fit to the discriminant output distribution ( $t\bar{t}$  ratio)

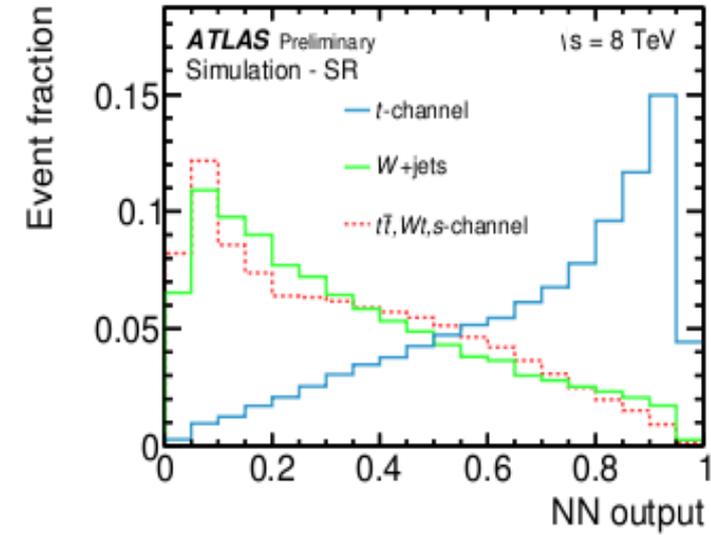
- Overall scale factor (W') obtained per jet-bin and lepton selection in a CR defined by events containing 2 b-tag jets &  $m(tb) < 270$  GeV

$$SF_i = \frac{N_{Data,i}^{CR} - N_{other,i}^{CR}}{N_{Wjets,i}^{CR}}$$

# t-channel cross section

Input variables used for the NN ordered by their importance:

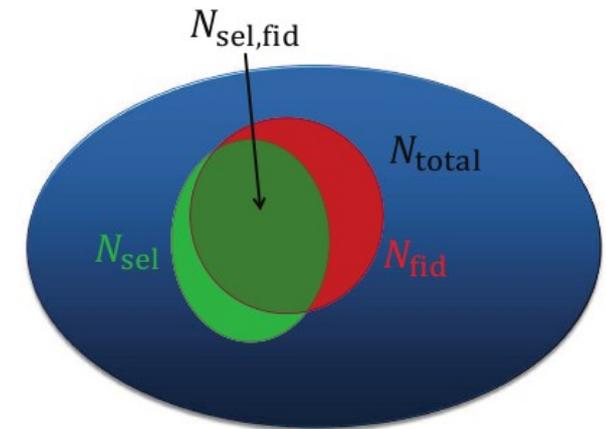
Variable	Definition
$ \eta(j) $	pseudorapidity of the light quark (untagged) jet ( $j$ )
$m(\ell\nu b)$	top-quark mass reconstructed from the charged lepton, neutrino and $b$ -quark jet
$m(jb)$	invariant mass of the tagged ( $b$ ) and light quark jet ( $j$ )
$m_T(W)$	transverse mass of the reconstructed $W$ boson
$m(\ell b)$	invariant mass of the lepton ( $\ell$ ) and the tagged jet ( $b$ )
$\eta(\ell\nu)$	pseudorapidity of the reconstructed $W$ -boson
$\cos \Theta(\ell, j)_{\ell\nu b \tau.f.}$	cosine of the angle $\theta$ between the charged lepton and the light quark (untagged) jet ( $j$ ) in the rest frame of the reconstructed top quark
$H_T(\ell, \text{jets}, E_T^{\text{miss}})$	scalar sum of the transverse momenta of the jets, the charged lepton and the missing transverse momentum
$E_T^{\text{miss}}$	transverse missing momentum
$\Delta R(\ell\nu b, \ell)$	$\Delta R$ of the reconstructed top quark and the charged lepton
$p_T(\ell\nu)$	transverse momentum of the reconstructed $W$ -boson
$\eta(\ell\nu b)$	pseudorapidity of the reconstructed top quark
$\eta(b)$	pseudorapidity of the $b$ -quark jet ( $b$ )
$p_T(\ell\nu b)$	transverse momentum of the reconstructed top quark



$$\sigma = \frac{\hat{\nu}}{\epsilon \mathcal{L}} = \frac{1}{\epsilon_{\text{fid}}} \cdot \frac{\epsilon_{\text{corr,sel}}}{\epsilon_{\text{corr,fid}}} \cdot \frac{\hat{\nu}}{\mathcal{L}} = \frac{1}{\epsilon_{\text{fid}}} R_f \sigma_{\text{fid}}$$

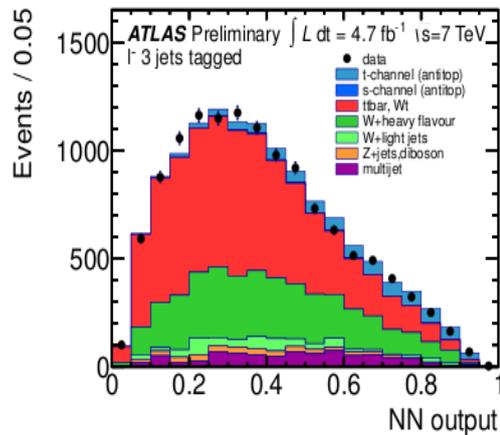
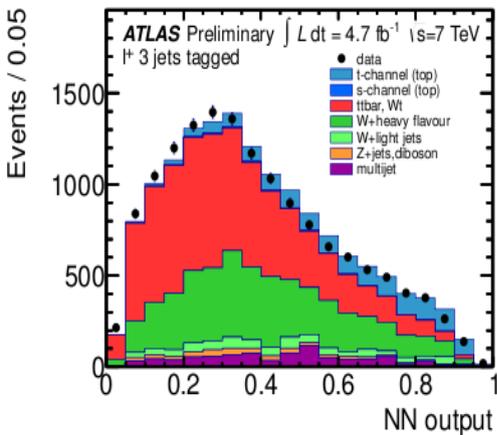
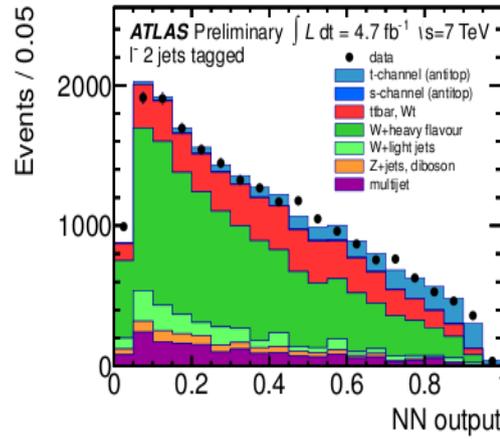
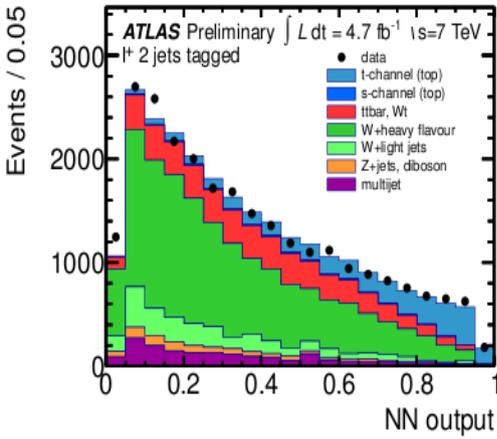
$$\Rightarrow \sigma_{\text{fid}} = \frac{\epsilon_{\text{corr,sel}}}{\epsilon_{\text{corr,fid}}} \cdot \frac{\hat{\nu}}{\mathcal{L}} = R_f \frac{\hat{\nu}}{\mathcal{L}}$$

$$\epsilon = \frac{N_{\text{sel}}}{N_{\text{total}}}, \quad \epsilon_{\text{fid}} = \frac{N_{\text{fid}}}{N_{\text{total}}}, \quad \epsilon_{\text{corr,sel}} = \frac{N_{\text{sel,fid}}}{N_{\text{sel}}}, \quad \epsilon_{\text{corr,fid}} = \frac{N_{\text{sel,fid}}}{N_{\text{fid}}},$$



# t-channel cross section ratio

NN distributions:

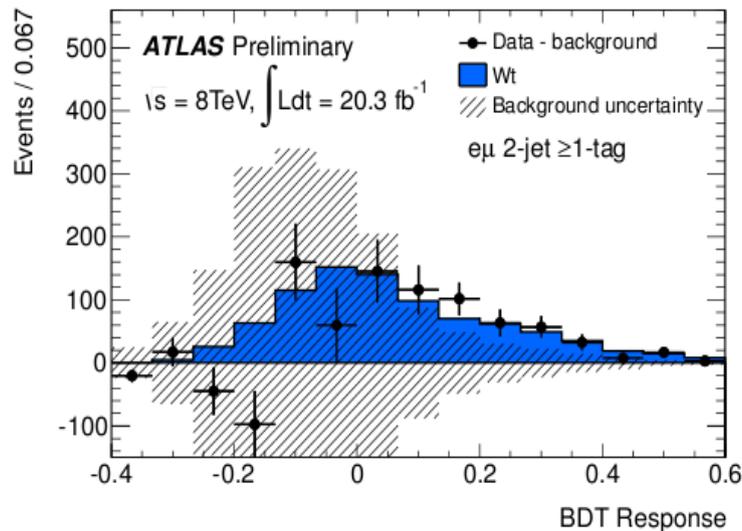
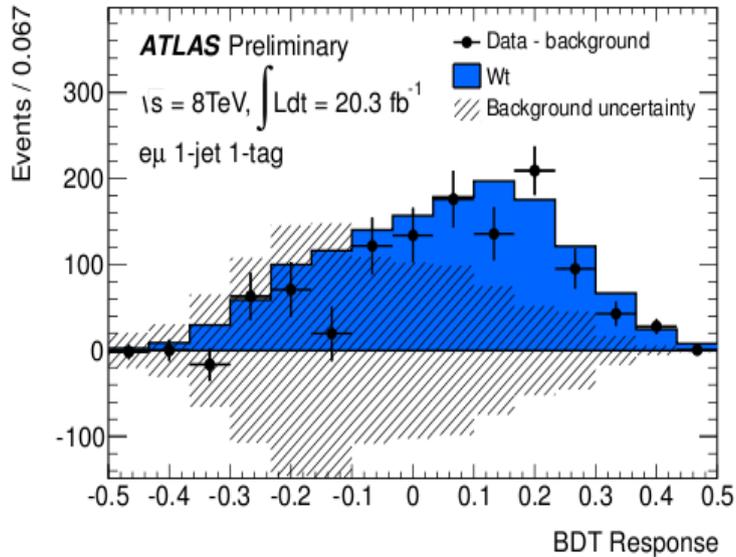


Systematic breakdown:

Source	$\Delta\sigma_t/\sigma_t(t)$ [%]	$\Delta\sigma_t/\sigma_t(\bar{t})$ [%]	$\Delta R_t/R_t$ [%]
Data statistics	$\pm 3.2$	$\pm 5.0$	$\pm 5.5$
MC statistics	$\pm 2.4$	$\pm 3.7$	$\pm 3.7$
Multijet normalisation	+1.1 / -2.0	+3.1 / -4.2	+3.9 / -3.7
Other background normalisation	$\pm 3.4$	$\pm 1.3$	$\pm 4.5$
Jet energy scale	$\pm 16.4$	$\pm 19.5$	+3.7 / -3.6
Jet energy resolution	$\pm 3.4$	$\pm 4.3$	$\pm 0.8$
Jet reconstruction efficiency	+0.7 / -0.4	$\pm 0.3$	+0.8 / -0.5
$b$ -tagging efficiency scale-factor	$\pm 5.9$	$\pm 8.5$	$\pm 2.5$
Mistag efficiency scale-factor	$\pm 0.8$	$\pm 2.0$	$\pm 2.7$
$b/\bar{b}$ acceptance	$\pm 1.0$	$\pm 1.0$	$\pm 0.4$
$E_T^{\text{miss}}$ modeling	+0.6 / -0.9	+0.9 / -1.3	+0.8 / -0.7
Lepton efficiencies	$\pm 2.9$	$\pm 2.9$	$\pm 0.3$
Lepton energy resolution	+0.5 / -0.8	+1.1 / -1.4	+1.0 / -1.1
Electron energy scale	+0.3 / -0.5	+0.7 / -0.8	$\pm 0.6$
PDF	$\pm 3.3$	$\pm 4.5$	+1.1 / -1.2
W+jets shape variation	+0.6 / -0.5	$\pm 0.5$	$\pm 0.7$
Top MC generator	$\pm 7.1$	$\pm 7.1$	$\pm 0.7$
ISR / FSR	$\pm 0.7$	$\pm 3.5$	$\pm 4.2$
Luminosity	$\pm 3.9$	$\pm 3.9$	$\pm 0.4$
Total Systematic	$\pm 20.0$	+24.7 / -24.9	+11.5 / -11.1
Total	$\pm 20.2$	+25.2 / -25.4	+12.8 / -12.4

# Wt-channel cross section

BDT distributions, bkg subtraction:



Systematic breakdown:

Source	$\Delta\sigma/\sigma$ [%]	
	observed	expected
Data statistics	7.1	8.6
MC statistics	2.8	3.5
Experimental uncertainties		
Lepton modeling	2.4	2.4
Jet identification	0.2	0.6
Jet energy scale	10	12
$b$ -jet energy scale	5.0	6.3
Jet energy resolution	0.7	0.2
$E_T^{\text{miss}}$ scale	4.1	5.0
$E_T^{\text{miss}}$ resolution	4.5	5.3
Flavor tagging	8.4	9.4
Theory uncertainties		
$Wt/t\bar{t}$ overlap modeling	1.4	1.6
PDF	2.5	3.2
Background normalization	3.6	4.4
ISR/FSR	5.9	6.0
$Wt$ generator and PS	11	11
$t\bar{t}$ generator and PS	7.5	9.2
Luminosity	3.7	3.9
Total (syst)	20	23
Total (syst+stat)	21	24

# s-channel cross section

Cut-based analysis:

Selection	Signal	Background	$S/\sqrt{B}$
Preselection Only	104	153802	0.26
Number of tagged jets=2	18	415	0.88
$30 < m_{top, jet2} < 247 \text{ GeV}/c^2$	17	349	0.91
$p_T(jet1, jet2) < 189 \text{ GeV}/c$	17	346	0.91
$m_T(W) < 111 \text{ GeV}/c$	17	318	0.95
$0.43 < \Delta R(b - jet1, lepton) < 3.6$	17	308	0.97
$123 < m_{top, jet1} < 788 \text{ GeV}/c^2$	17	302	0.98
$0.74 < \Delta R(b - jet1, b - jet2) < 4.68$	16	269	0.98

Systematic breakdown:

Source	$\Delta\sigma/\sigma$ [%] cut-based
Data statistics	$\pm 100$
MC statistics	$\pm 70$
<i>b</i> -tagging	-30/+20
Jet and lepton modeling	-20/+10
MC generator modeling	-60/+20
Multijets normalization	$\pm 40$
Others	-10/+30
Luminosity	$\pm 50$
All systematics	-110/+90
Total uncertainty	-160/+150

# W' boson search

## Theories predicting massive spin-1 charged bosons

- extra-dims : KK excitations of the W
- technicolor : EWSB without scalars but new bosons
- L-R symmetry : right-handed counter-part of the W
- little-Higgs : cancels mH divergence, several new particles including a W

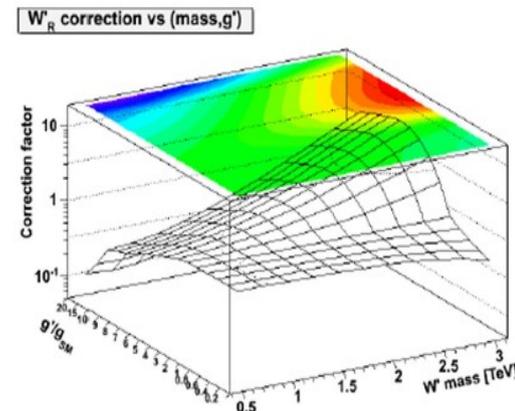
- Effective operator formalism: model independent description

$$\mathcal{L} = \frac{V'_{ij}}{2\sqrt{2}} \bar{f}_i \gamma_\mu \left( g'_{i,j}{}^{IR} (1 + \gamma^5) + g'_{i,j}{}^{IL} (1 - \gamma^5) \right) W'^\mu f_j + h.c.$$

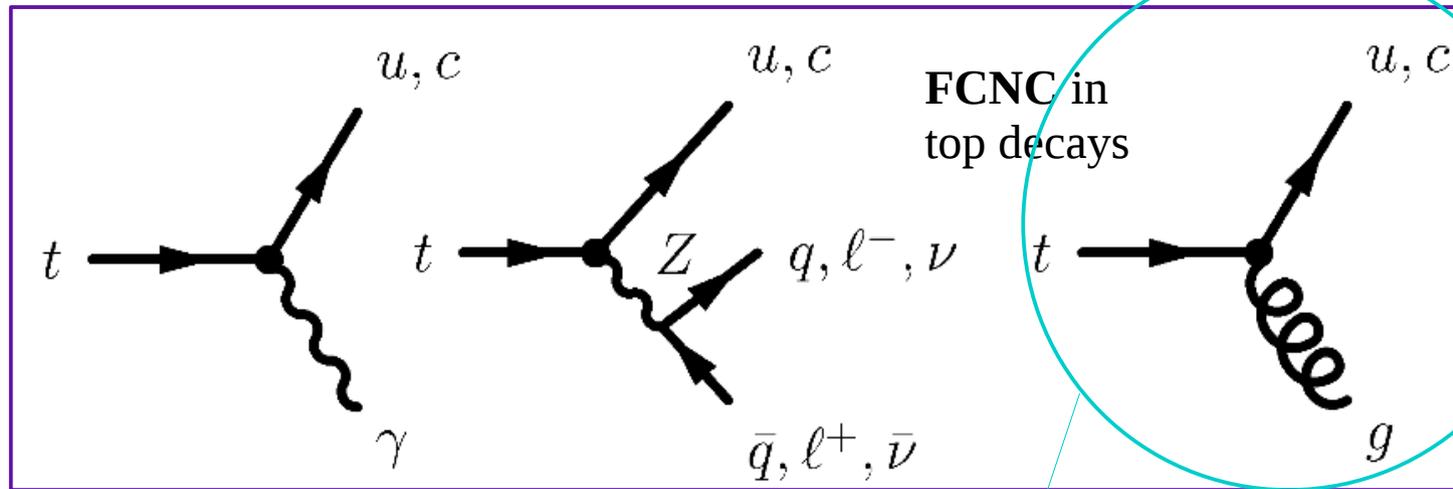
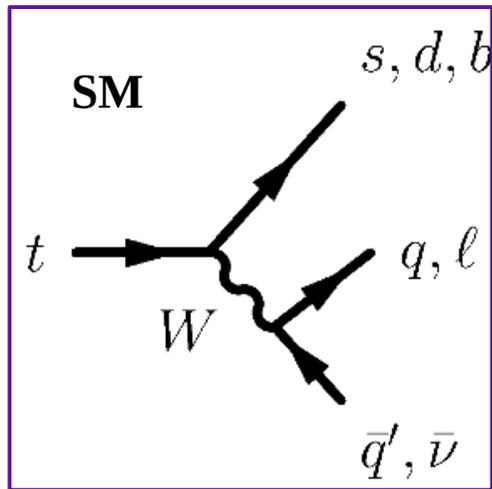
- In this effective theory  $m(W')$  independent of  $g$ , which can differ a lot from the SM  $g$  (model renormalisable up to  $g'=5g$ )

- $W' \rightarrow tb$  cross section roughly scales as  $\sim (g'/g)^2$  + correction factor due to  $W'$  width effects more important for high masses and high  $g'/g$

$\Rightarrow$  limit studied up to  $g'/g=2$ , safe since signal with  $= 4 * \text{standard width}$ , still lower than experimental resolution



# FCNCs in single top production

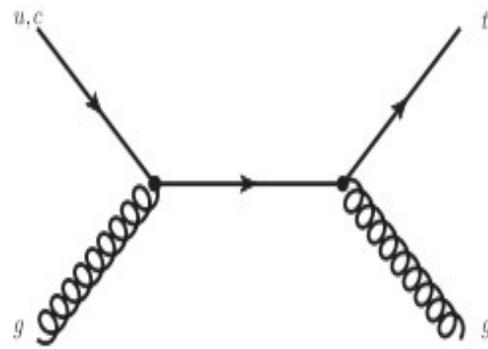
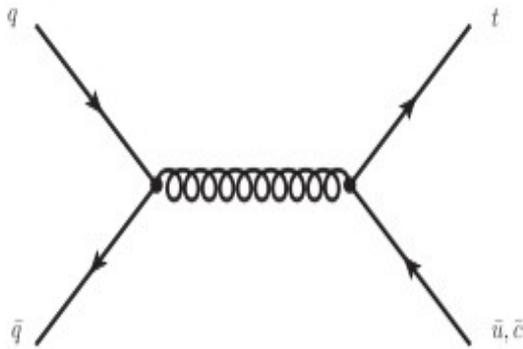
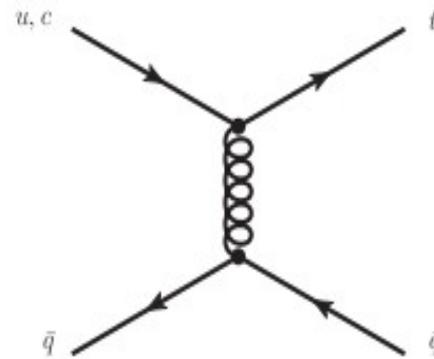
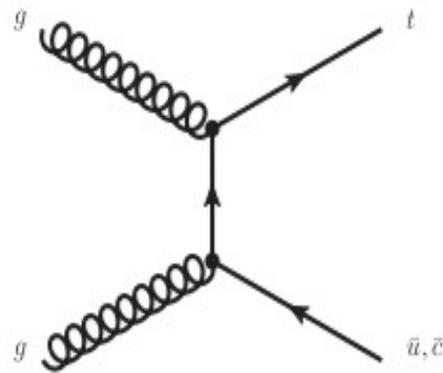
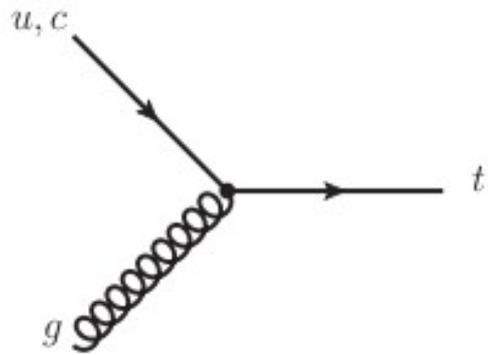


T-g-q coupling better studied in production mechanism,  
Since final state signatures dominated by multijet background

Table 1: Theoretical values for the branching fractions of FCNC top quark decays as predicted by the SM, the quark-singlet model (QS), two-Higgs doublet models with (2HDM II) or without flavour-conservation (2HDM III), minimal supersymmetric model (MSSM), SUSY with R-parity violation ( $\mathcal{R}$ SUSY), models with extra quarks (Extra q) and Topcolour-assisted Technicolour model (TC2) [29,30].

Process	SM	QS	2HDM II	2HDM III	MSSM	$\mathcal{R}$ SUSY	Extra q	TC2
$t \rightarrow uy$	$3.7 \times 10^{-16}$	$\sim 10^{-8}$	—	—	$2 \times 10^{-6}$	$1 \times 10^{-6}$	—	—
$t \rightarrow uZ$	$8 \times 10^{-17}$	$\sim 10^{-4}$	—	—	$2 \times 10^{-6}$	$3 \times 10^{-5}$	—	—
$t \rightarrow ug$	$3.7 \times 10^{-14}$	$\sim 10^{-7}$	—	—	$8 \times 10^{-5}$	$2 \times 10^{-4}$	—	—
$t \rightarrow cy$	$4.6 \times 10^{-14}$	$\sim 10^{-8}$	$\sim 10^{-7}$	$\sim 10^{-7}$	$2 \times 10^{-6}$	$1 \times 10^{-6}$	$\sim 10^{-8}$	$\sim 10^{-7}$
$t \rightarrow cZ$	$\sim 1 \times 10^{-14}$	$\sim 10^{-4}$	$\sim 10^{-8}$	$\sim 10^{-6}$	$2 \times 10^{-6}$	$3 \times 10^{-5}$	$\sim 10^{-4}$	$\sim 10^{-5}$
$t \rightarrow cg$	$4.6 \times 10^{-12}$	$\sim 10^{-7}$	$\sim 10^{-5}$	$\sim 10^{-4}$	$8 \times 10^{-5}$	$2 \times 10^{-4}$	$\sim 10^{-7}$	$\sim 10^{-5}$

# FCNCs in single top production



LO & NLO Feynman diagrams  
( $2 \rightarrow 1$ ) ( $2 \rightarrow 2$ )

# FCNCs in single top production

- Effective operator formalism: model independent description

$$\mathcal{L} = \mathcal{L}^{SM} + \frac{1}{\Lambda} \cancel{\mathcal{L}^{(5)}} + \frac{1}{\Lambda^2} \mathcal{L}^{(6)} + O\left(\frac{1}{\Lambda^3}\right) \quad L^{(\text{eff})} \rightarrow L^{(\text{SM})} \text{ in the low energy limit}$$

- For calculating branching fractions, FCNCs considered as additional channels for top decay:

$$\mathcal{B}(t \rightarrow qg) = \frac{\Gamma_{t \rightarrow qg}}{\Gamma_{t \rightarrow bW} + \Gamma_{t \rightarrow qg}} \quad q = u, c$$

for small coupling constants ( $k < 0.1$ ):  $\mathcal{B}(t \rightarrow qg) = \frac{\Gamma_{t \rightarrow qg}}{\Gamma_{t \rightarrow bW}} \longrightarrow$  result (C)

- Production cross section:

$$\sigma_{qg \rightarrow t} = \sum_{q=u,c} \left( \frac{K_{gqt}}{\Lambda} \right)^2 (b_{qL} |f_q^L|^2 + b_{qR} |f_q^R|^2)$$

assuming only a pure left-handed coupling ( $f_q^L=1, f_q^R=0$ ) and considering the  $l^+$  together:

$$\sigma_{qg \rightarrow t} = 100351.97 \left( \frac{K_{ugt}}{\Lambda} \right)^2 + 19496.79 \left( \frac{K_{cgt}}{\Lambda} \right)^2 \text{ pb} \longrightarrow \text{result (B)}$$