

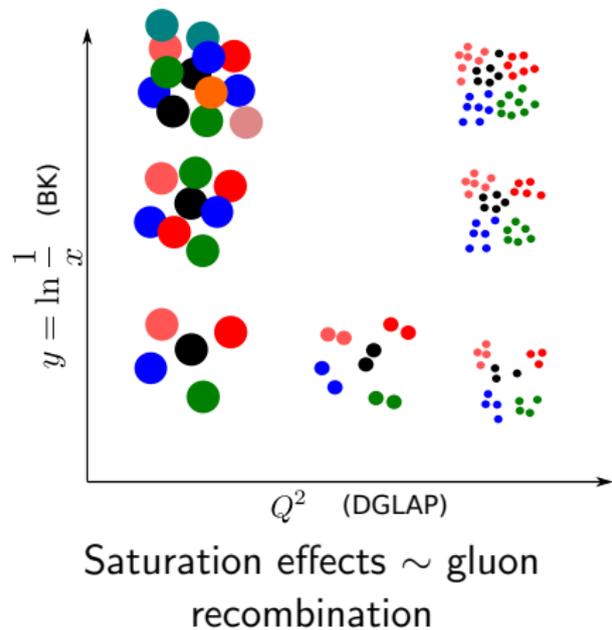
Diffractive vector meson production in ultraperipheral heavy ion collisions from the Color Glass Condensate

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Introduction

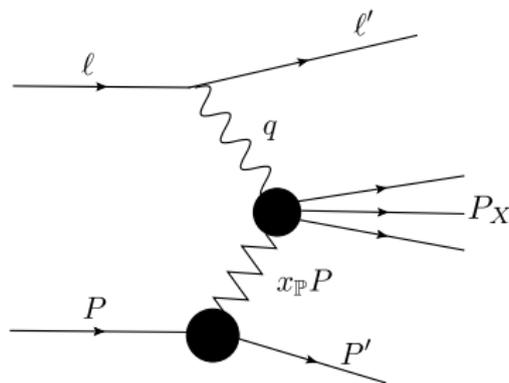


- The Color Glass Condensate framework describes many small- x processes accurately
- Necessary input: gluon density at $x = x_0$ (from DIS)
- There is very little small- x nuclear DIS data

eA collisions would be ideal, but before eRHIC/LHeC have to use something else
 \Rightarrow ultraperipheral AA

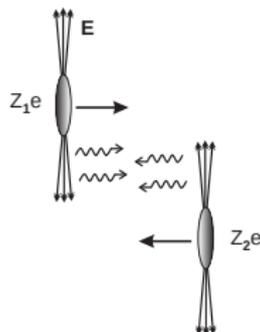
(Lepton-proton) Diffraction

Diffractive deep inelastic scattering (DDIS) = DIS with no exchange of quantum numbers (color).



- $e + p \rightarrow e + p + X$, proton interacts via "pomeron exchange"
- $x_{\mathbb{P}}$: fraction of proton momentum carried by the pomeron.
- $Q^2 = -q^2$: virtuality of the photon.

Ultrapерipheral AA collision



nucl-ex/0502005

$b \gtrsim 2R_A$: strong interactions suppressed, nucleus creates photon flux $n(\omega)$
 $\sigma \sim n(\omega)\sigma^{\gamma A}(\omega)$

Probes gluons with $x = M_V e^y / \sqrt{s}$

- Forward LHC: $x \sim 0.02$ and $x \sim 10^{-5}$.
- Midrapidity LHC: $x \sim 10^{-3}$

Dipole model is valid only at $x \lesssim 10^{-2} \Rightarrow$ at LHC limit $y \lesssim 2 \dots 3$.

Coherent and incoherent diffraction

Diffraction off the nucleus:

- Coherent diffraction: nucleus remains intact

$$\frac{d\sigma^{\gamma^*A \rightarrow VA}}{dt} \sim \langle |\mathcal{A}(x, Q^2, t)|^2 \rangle$$

- Quasielastic = coherent + incoherent

$$\frac{d\sigma^{\gamma^*A \rightarrow V(A^*+A)}}{dt} \sim |\langle \mathcal{A}(x, Q^2, t) \rangle|^2$$

- Incoherent, nucleus is allowed to break up

$$\frac{d\sigma^{\gamma^*A \rightarrow VA^*}}{dt} \sim |\langle \mathcal{A}(x, Q^2, t) \rangle|^2 - \langle |\mathcal{A}(x, Q^2, t)|^2 \rangle$$

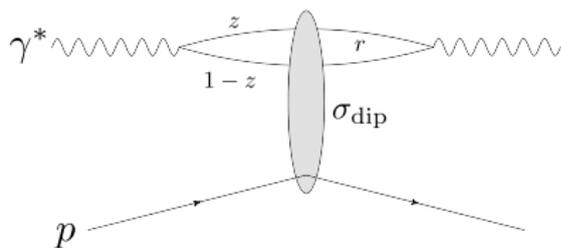
$\langle \rangle$ = Average over nucleon positions.

Dipole cross section

CGC: Dipole-proton cross section

$$\sigma_{\text{dip}}(x, r, \Delta) = 2 \int d^2b e^{ib \cdot \Delta} N(r, x, b)$$

Universal dipole amplitude N



- Total $\gamma^* p$:

$$\int d^2r dz |\Psi^\gamma(Q^2, r, z)|^2 \sigma_{\text{dip}}(x, r, \Delta = 0)$$

- Total diffraction:

$$\frac{1}{16\pi} \int d^2r dz |\Psi^\gamma(Q^2, r, z) \sigma_{\text{dip}}(x, r, \Delta)|^2$$

- Exclusive diffraction:

$$\frac{1}{16\pi} \left| \int d^2r dz \Psi^* \Psi^V(Q^2, r, z) \sigma_{\text{dip}}(x, r, \Delta) \right|^2$$

- Inclusive particle production (pp, pA):

$$\sim xg(x, Q^2) \int d^2r e^{ir \cdot p_T} [1 - N(r, x)]$$

- + Correlations, ...

Impact parameter dependent BK evolution is problematic (work in progress), use IPsat model (Kowalski, Teaney 2003; Rezaeian et al, 2013):

$$N(r, x, b) = 1 - \exp \left[-\frac{\pi^2}{2N_c} \alpha_s x g(x, \mu^2) T_p(b) r^2 \right]$$

- Fit to HERA data: initial condition for DGLAP evolution of $xg(x, \mu^2)$
- Proton profile T_p gaussian

Generalize for nuclei:

- $T_p(b) \rightarrow \sum_{i=1}^A T_p(b - b_i)$
- Average over different nucleon configurations from Woods-Saxon
⇒ coherent/incoherent diffraction

Quasielastic cross section can be computed using a factorized approximation (T. Lappi, H.M, 1011.1988)

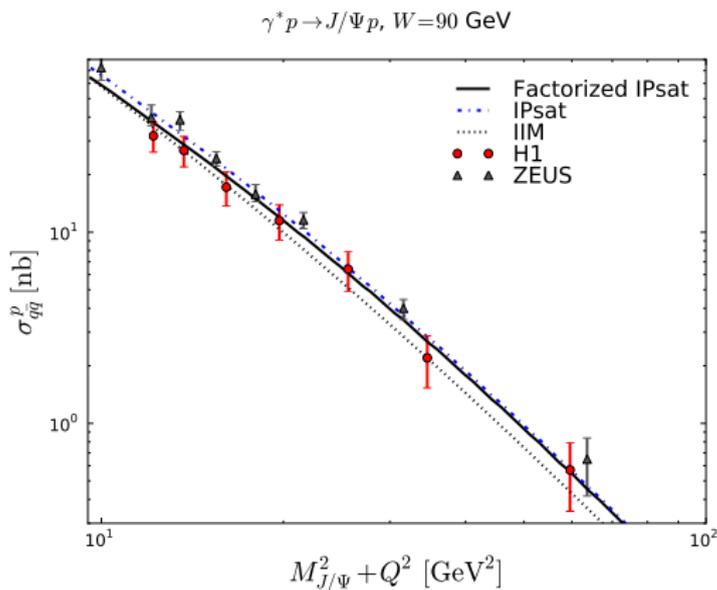
$$N(r, x, b) \approx T(b)N(r, x)$$

$$N(r, x) = 1 - \exp \left[-\frac{1}{2\pi B_p} \frac{\pi^2}{2N_c} \alpha_s x g(x, \mu^2) r^2 \right]$$

- Also large corrections from real part and skewness
- Compare "fIPsat" to IIM (Iancu, Itacura, Munier) model: study dependence on dipole model

Comparison with the HERA data

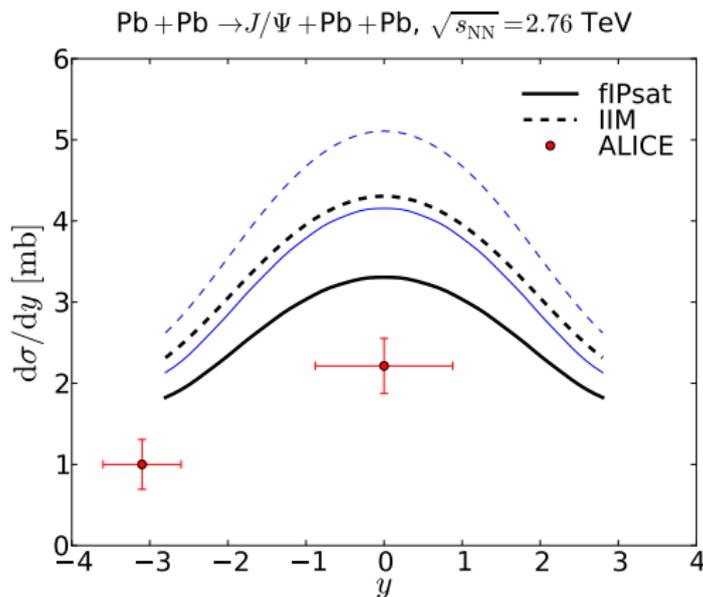
Compare with HERA $\gamma^* p \rightarrow J/\Psi p$ data:



T. Lappi, H. Mäntysaari, 1011.1988

Also a good description of the F_2 data.

Comparison with the ALICE data: coherent diffraction

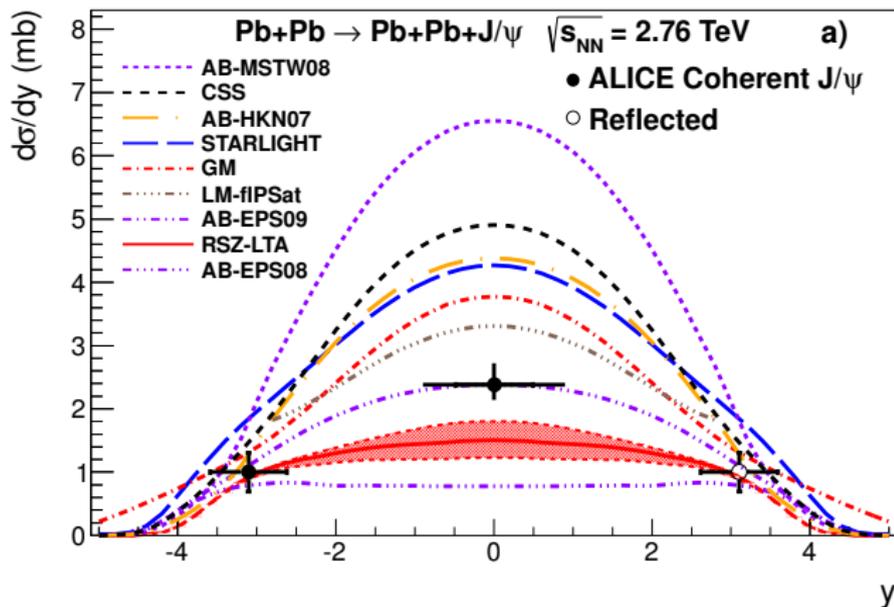


T. Lappi and H. Mäntysaari, 1301.4095; Data: ALICE 1310.7732

Different dipole and vector meson wavef models: change overall normalization, but shape is very similar.

RHIC midrapidity: $d\sigma/dy = 109\mu\text{b}$, experimental $76 \pm 34\mu\text{b}$.

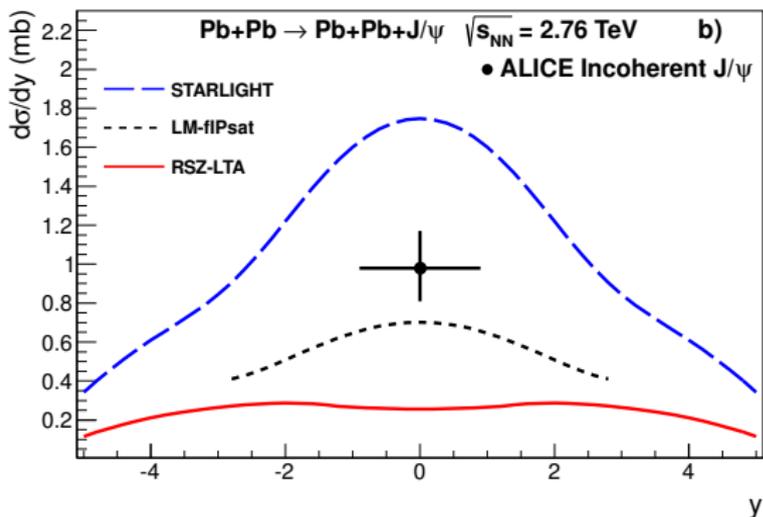
Coherent diffraction, model comparison



ALICE, 1310.7732

Unshadowed model (AB-MSTW08) clearly fails \Rightarrow saturation effects seen

Comparison of predictions (incoherent diffraction)



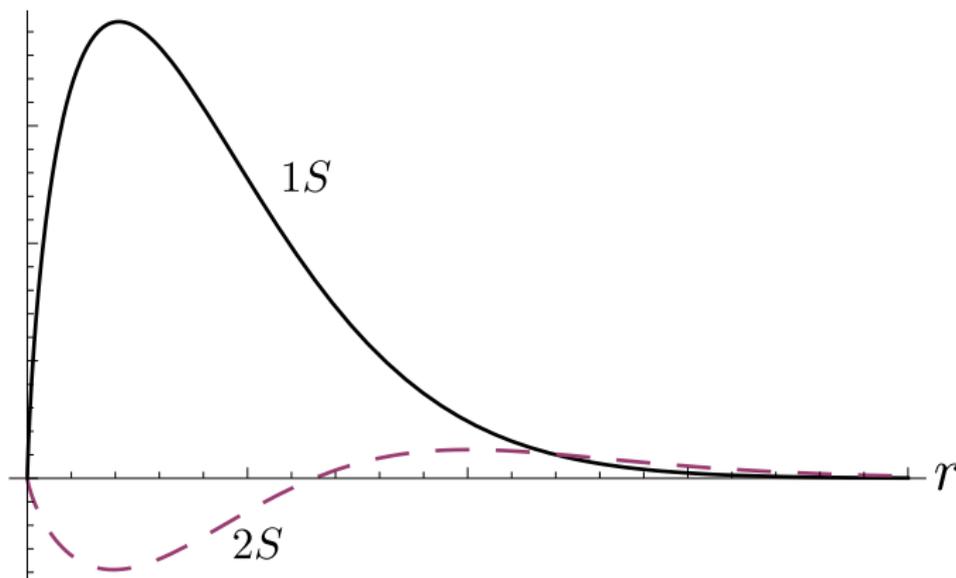
ALICE, 1305.1467

Normalization dependence on the J/ψ wave function model: $\sim 25\%$

$\Psi(2S)$ production

$\Psi(2S)$ wave function has a node \Rightarrow large suppression compared to J/Ψ

- Meson-photon wave function overlap, $z = 0.5$:

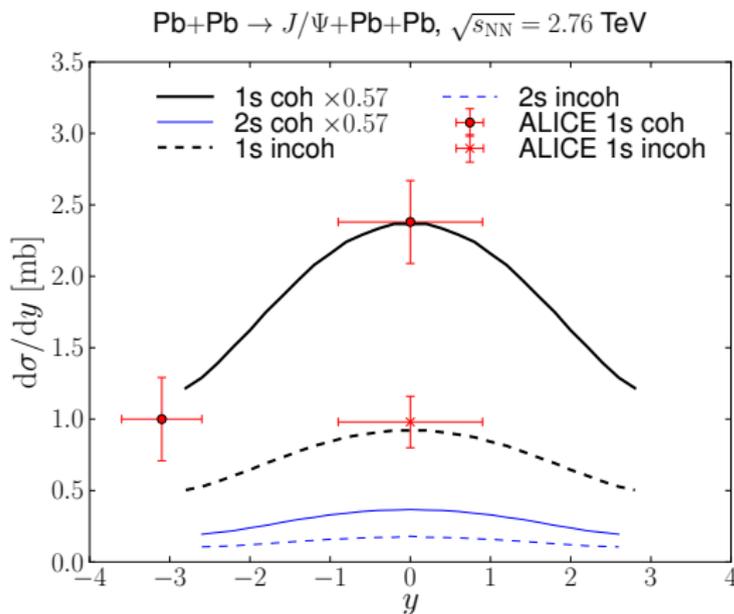


$\Psi(2S)$ production

$\Psi(2S)$ wave function has a node \Rightarrow large suppression compared to J/Ψ

- 2S/1S Ratio depends on event type:

$\gamma p < \text{coherent } \gamma A < \text{incoherent } \gamma A$

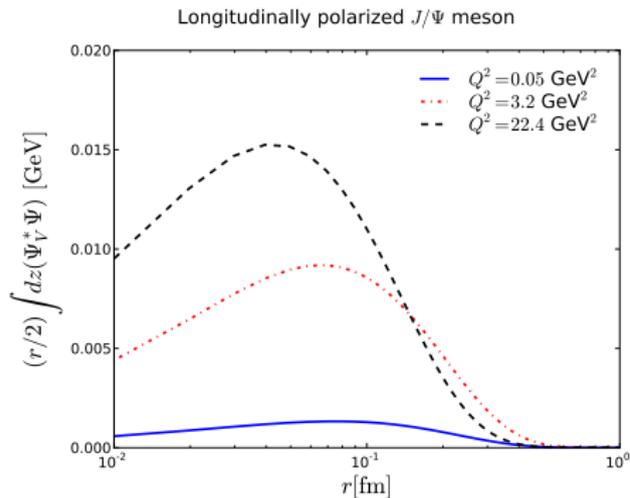
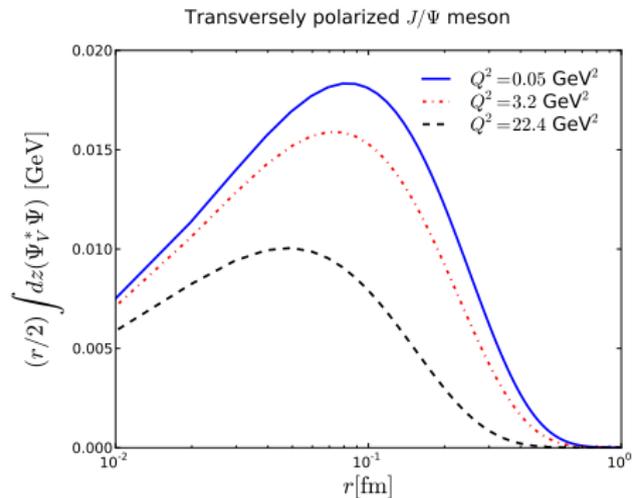


Data: ALICE, 1310.7732, 1305.1467

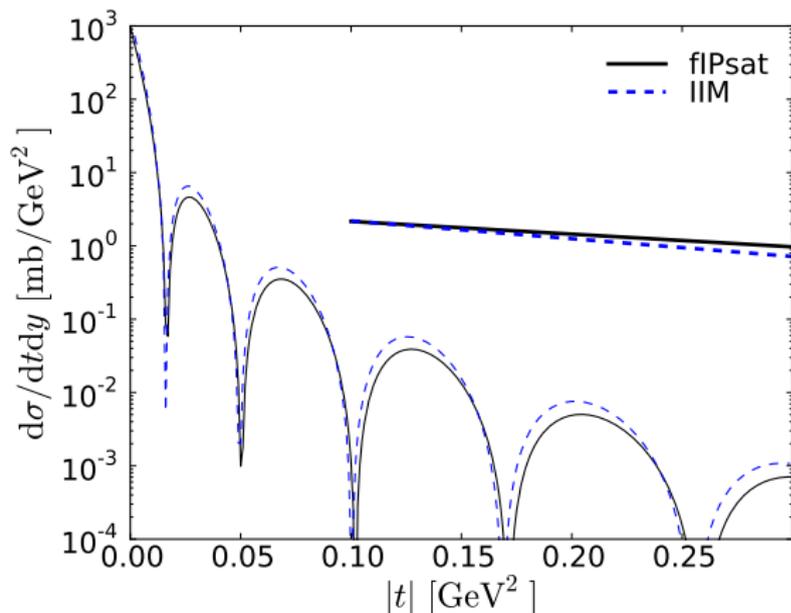
- Ultraperipheral heavy ion collisions make it possible to study γA diffraction at high energy
- Coherent and incoherent photoproduction measurements provide independent model constraints
- Dipole model description of incoherent and coherent diffraction in $\gamma^* A$
 - Here used IPsat parametrization fit to HERA
 - Absolute normalization has largish model dependence
 - Rapidity evolution of $d\sigma/dy$ is more precise prediction
- Prediction for $\Psi(2S)$ production and $\Psi(2S)/J/\Psi$ ratio
- Work in progress: use BK-evolved dipole amplitude consistently with the HERA F_2 and F_{2c} data

BACKUPS

Wave function overlap in J/Ψ production:

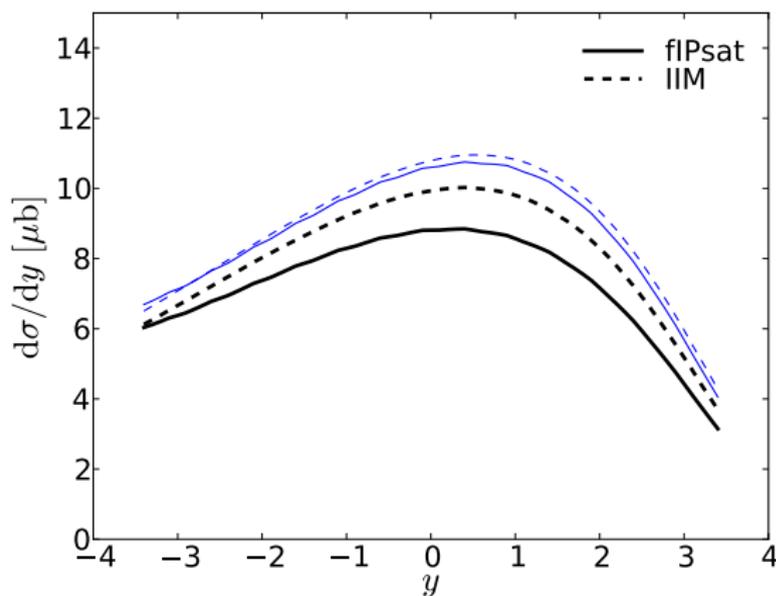


Differential cross section



T. Lappi, H. Mäntysaari, 1301.4095

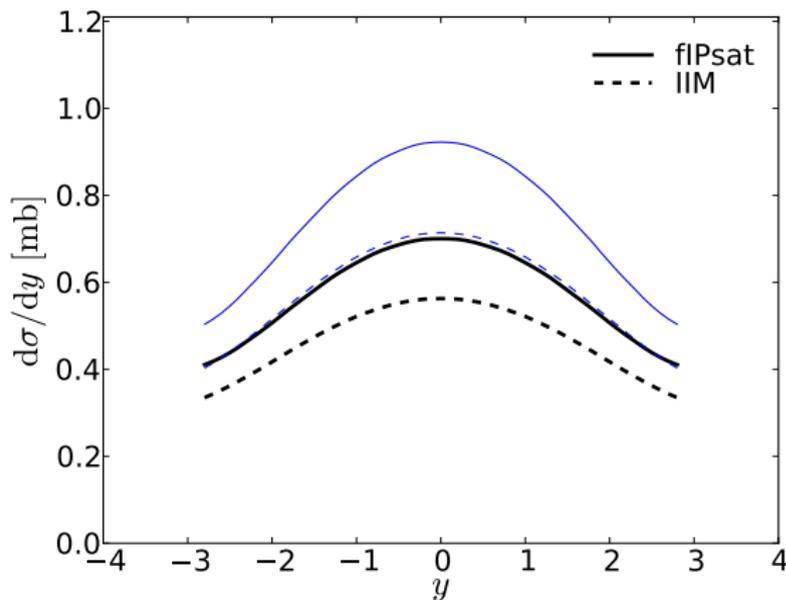
Assuming proton profile function $T_p(b) \sim e^{b^2/(2B_p)} \Rightarrow$ incoherent cross section $\sim e^{-B_p t}$: probes spatial distribution of gluons in proton!



T. Lappi, H. Mäntysaari, 1301.4095 CMS frame

As the photon flux $\sim Z^2$, dominant process is the one where the nucleus emits the photon \Rightarrow probes mostly proton structure.

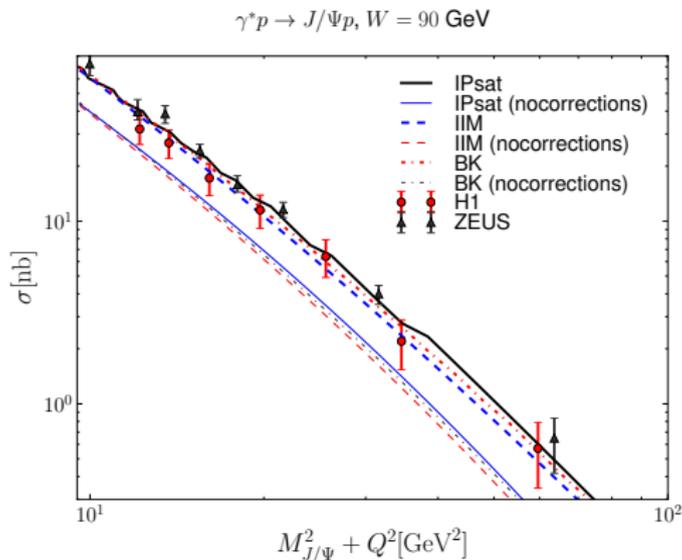
Prediction for incoherent diffraction



T. Lappi, H. Mäntysaari, 1301.4095

Again overall normalization uncertainty, but $\sigma(y=0)/\sigma(y=2) = 1.4$
more precise prediction.

Fit HERA σ_r : get automatically good description of σ_r^{charm} . Assume factorized impact parameter profile and $\sigma = \frac{1}{B_p} d\sigma/dt|_{t=0}$



Problem: large $\sigma_0 \sim 50 \text{ mb}$