



# PIDRIX: PARTICLE IDENTIFICATION MATRIX FACTORIZATION

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# TRADITIONAL PARTICLE IDENTIFICATION

Measure things that carry information about particle identity ( $dE/dx$ ,  $1/\beta$ , invariant mass, Čerenkov radiation, calorimeter energy, etc)

Painstakingly model these things for each particle type (crystal ball, normal distribution, Student's t-distribution, etc)

Maximize the likelihood to determine our best estimates of shapes and yields

This step is really hard and error prone.  
**Can we just skip it?**

# PIDRIX

Measure things that carry information about particle identity ( $dE/dx$ ,  $1/\beta$ , invariant mass, Čerenkov radiation, calorimeter energy, etc)



Maximize the likelihood to determine our best estimates of shapes and yields

Yes!

# THE ASSUMPTIONS

- ❖ We have a number of different particles that we measure
  - ❖  $\pi$ ,  $K$ ,  $p$ ,  $e$ , etc but could also be “background”
- ❖ For a given particle type each dimension of measurement is uncorrelated
  - ❖ e.g. ToF measurement error is not correlated with  $dE/dx$  measurement error
  - ❖ Bin on momentum and pseudorapidity to remove physics correlations
- ❖ Our particle yields follow Poisson statistics
  - ❖ No yields less than zero

# TOWARDS SOMETHING MORE MATHEMATICAL

This means that we expect the observed density to be of the form

$$A_{total}(x, y) = \sum_i^r v_i(x) * u_i(y) \quad \text{where } v_i(x) \text{ and } u_i(y) \text{ are positive semidefinite and real}$$

## Get your bearings:

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$x, y$  are the measurement variables (dE/dx and  $1/\beta$  for instance)

$A_{total}(x, y)$  is a fit function

$v_i(x) * u_i(y)$  is the contribution to the fit function for the  $i^{\text{th}}$  particle

$r$  is the number of distinct particle types

# FURTHER...

We use histograms so let's discretize things

$$A_{total}(x, y) = \sum_i^r v_i(x) * u_i(y) \quad \text{becomes} \quad A_{m_x n} = U_{m_x r} V_{r_x n}$$

Now say our observation histogram is  $T$ ,  
then maximizing the log-likelihood is equivalent to  
minimizing the generalized Kullback-Leibler divergence:

$$D_{GKL}(T||A) = \sum_{i,j} T_{ij} \ln \frac{T_{ij}}{A_{ij}} - T_{ij} + A_{ij}$$

# NOW WHAT?

## Traditional Fitting

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Restrict the space of possible  $U$  and  $V$  matrices according to our models before minimizing.

e.g. “The columns of  $U$  must be Gaussians and the rows of  $V$  must be Student’s  $t$ -distributions.”

## Pidrix

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Don’t model anything. Just minimize the KL divergence.

i.e. Skip the hard part.

# MINIMIZING THE KULLBACK-LEIBLER DIVERGENCE

These update rules can quickly be seen to be stable when  $T=A$ .

Also, if  $U$  and  $V$  start out positive definite and real then they will also remain so (our non-negative yield constraint).

$$U_{ip} \leftarrow U_{ip} \frac{\sum_{\alpha} V_{p\alpha} T_{i\alpha} / A_{i\alpha}}{\sum_{\alpha} V_{p\alpha}}$$

$$V_{pj} \leftarrow V_{pj} \frac{\sum_{\alpha} U_{\alpha p} T_{\alpha j} / A_{\alpha j}}{\sum_{\alpha} U_{\alpha p}}$$



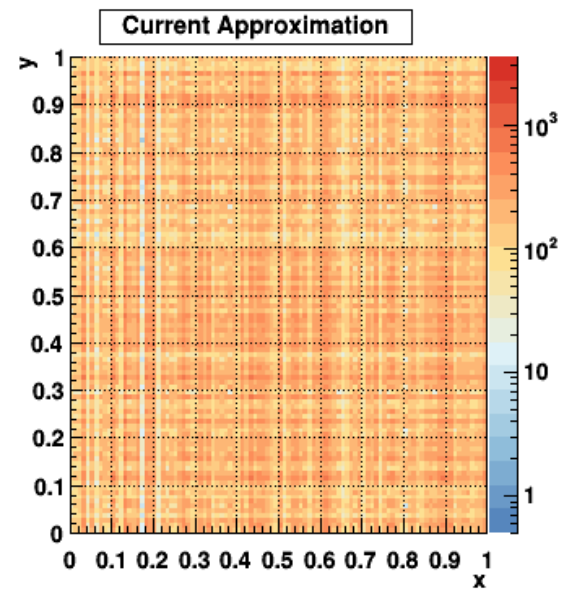
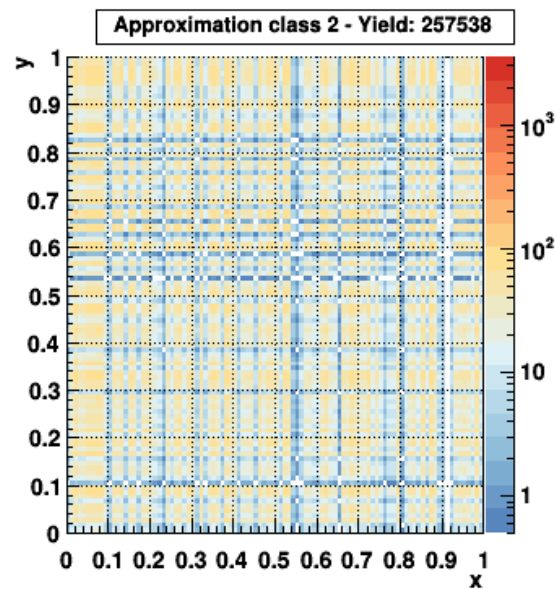
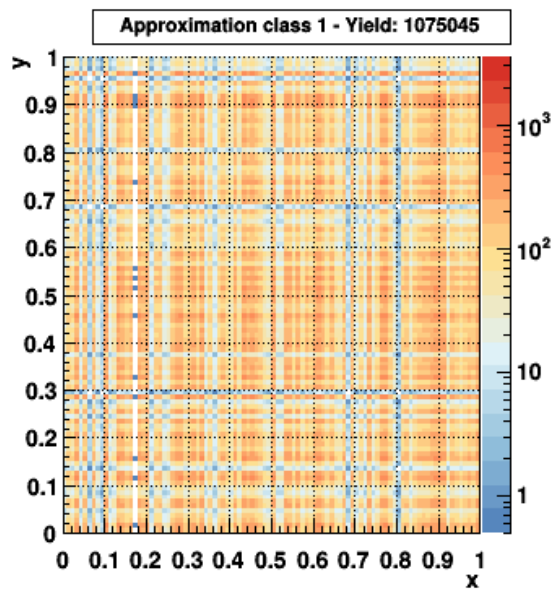
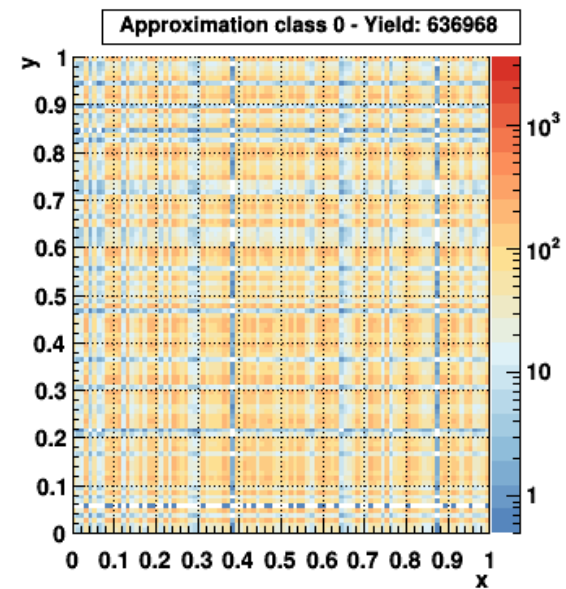
## HOW DO THEY WORK?

$$U_{ip} \leftarrow U_{ip} \frac{\sum_{\alpha} V_{p\alpha} T_{i\alpha} / A_{i\alpha}}{\sum_{\alpha} V_{p\alpha}}$$

- ❖ First: think of T/A (element-wise) like a residual
  - ❖ If an entry is less than one then we are underestimating that bin with our approximation, greater than one and we're overestimating
- ❖ For a given particle p (proton) we look at a specific i (dE/dx)
- ❖ Then we take our current normalized ToF distribution (the p<sup>th</sup> column of V) corresponding to protons and dot that with the row of T/A that corresponds to our current i value (dE/dx value)
- ❖ If this particle's ToF distribution has lots of yield where we are underestimating (overestimating) T then we'll increase (decrease) the dE/dx distribution for this particle at our current i

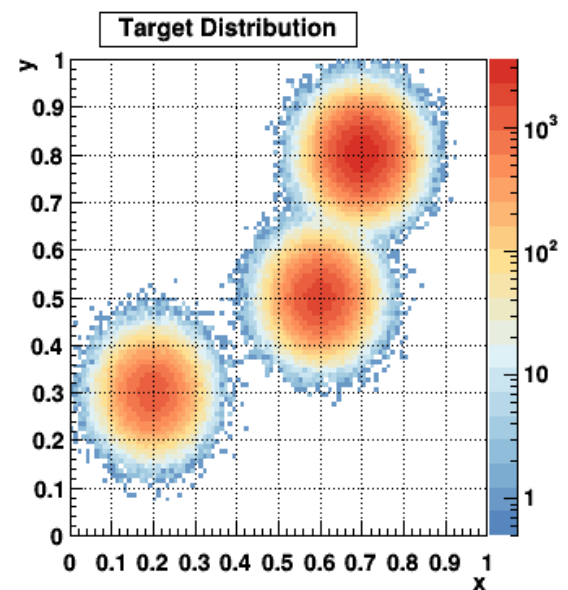
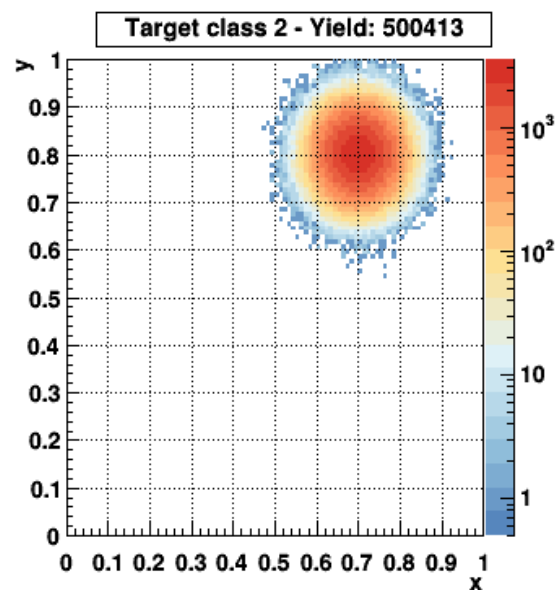
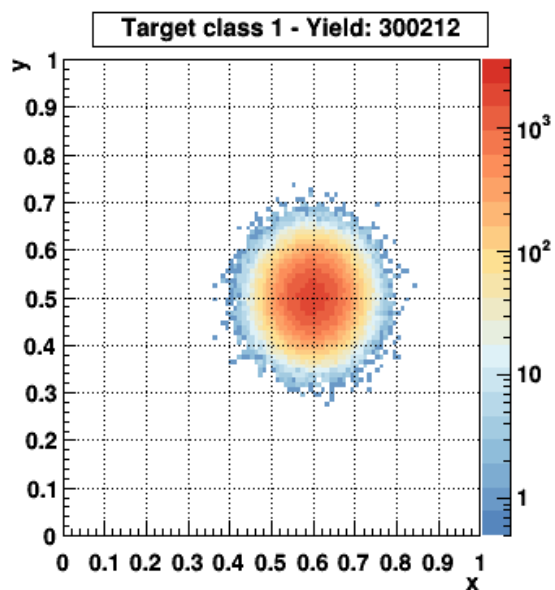
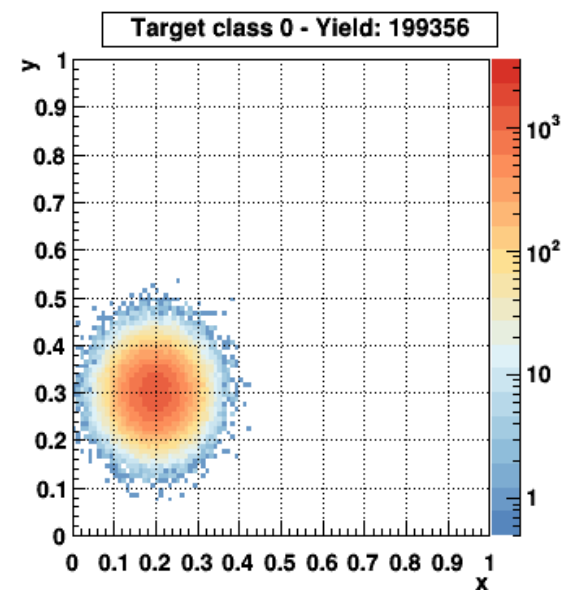
# NOW LET'S SEE IT IN ACTION

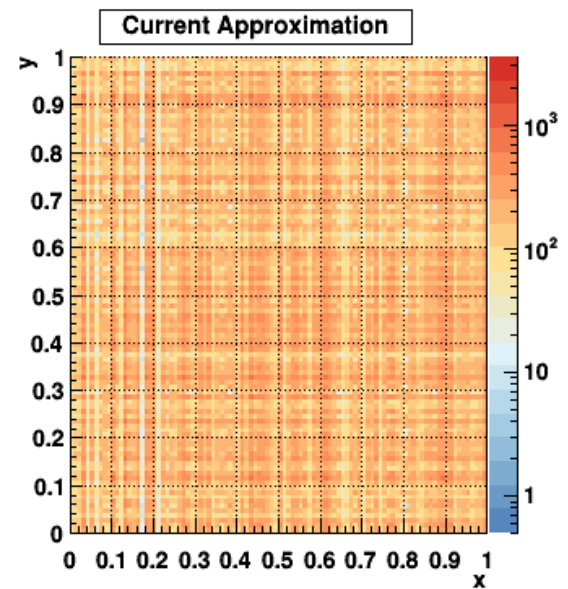
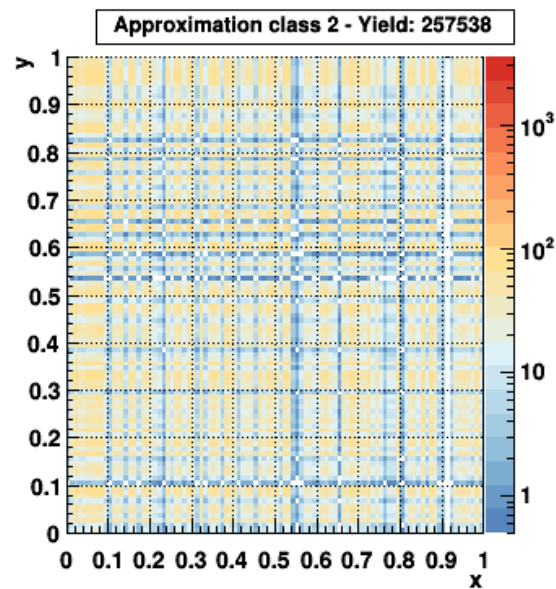
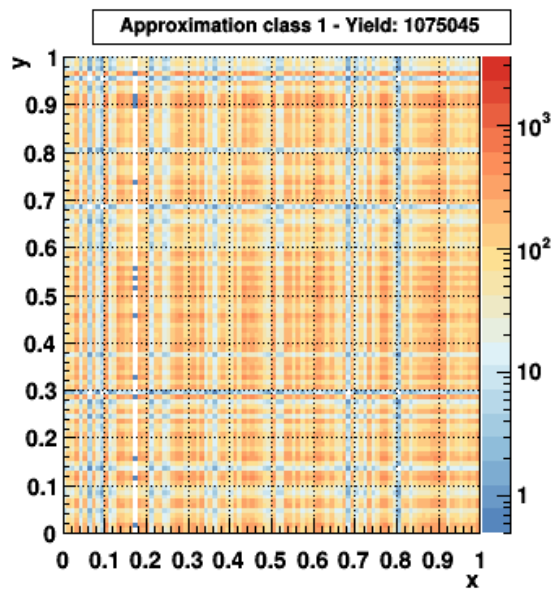
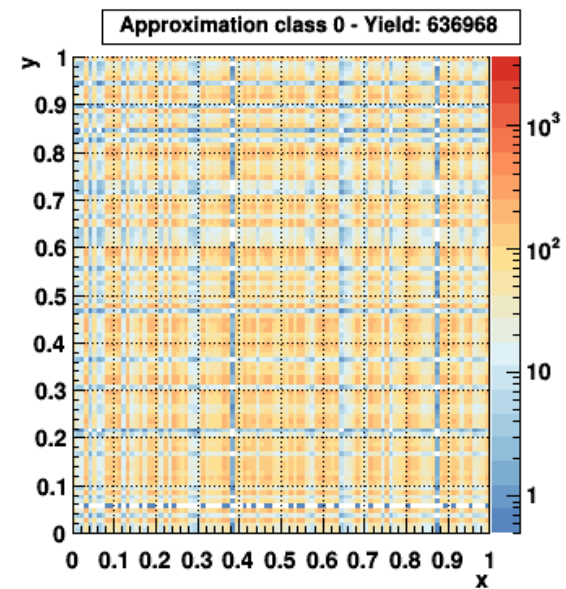
We'll start with random distributions and see how well they converge...



iteration 000000

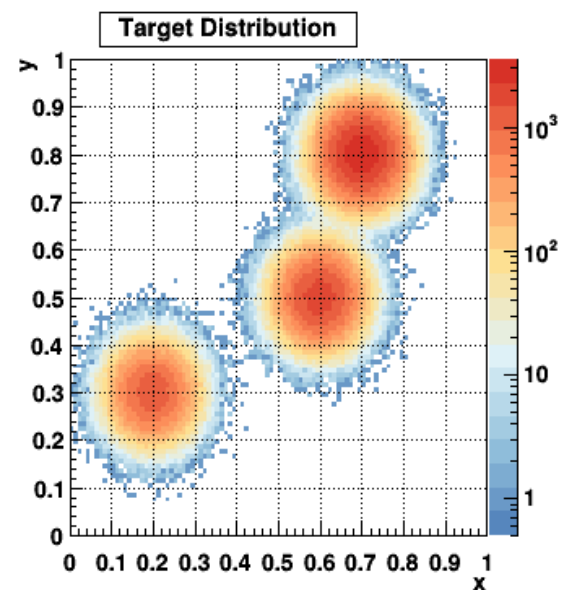
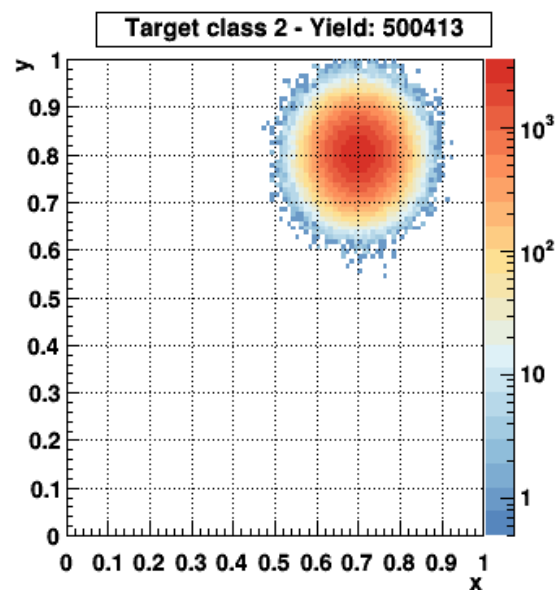
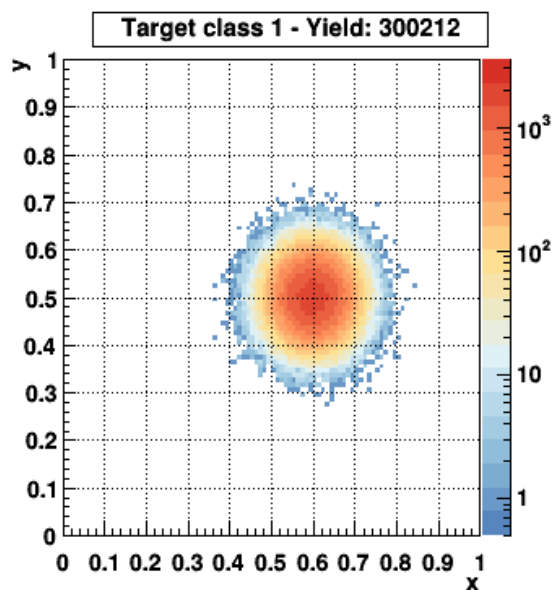
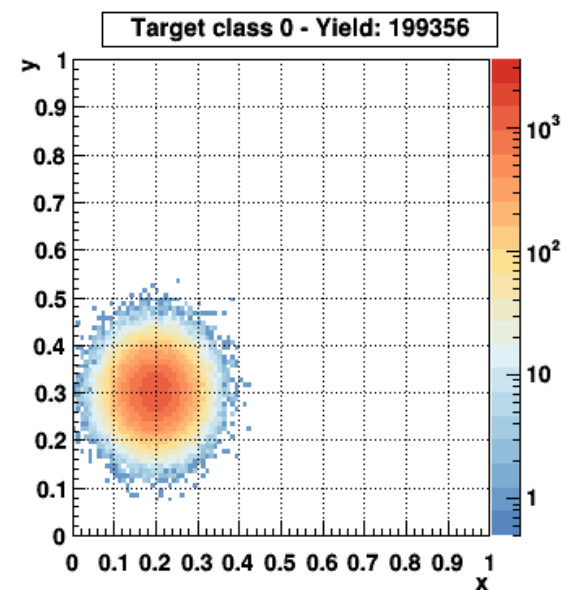
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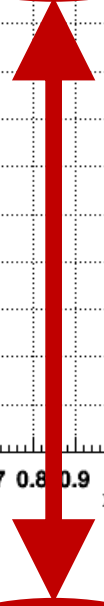
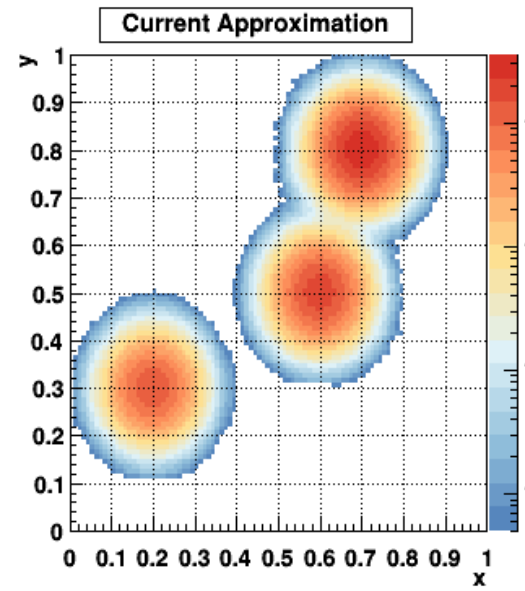
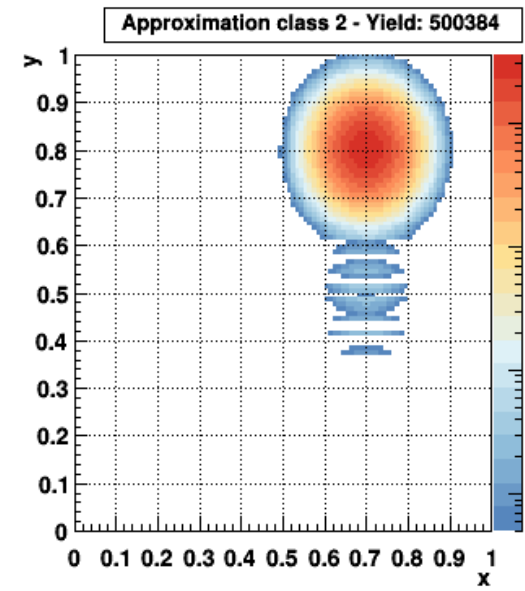
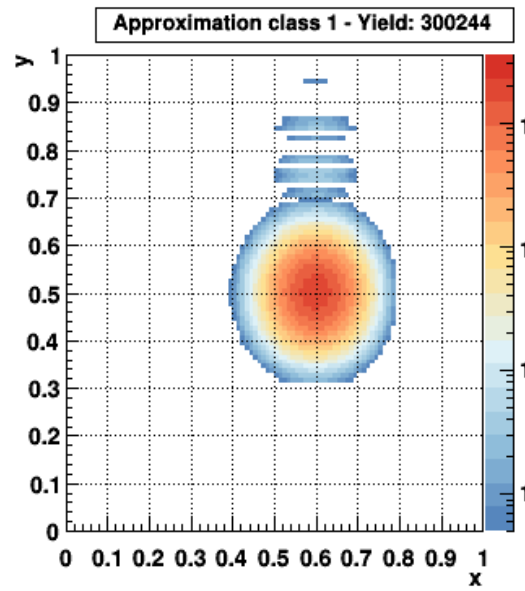
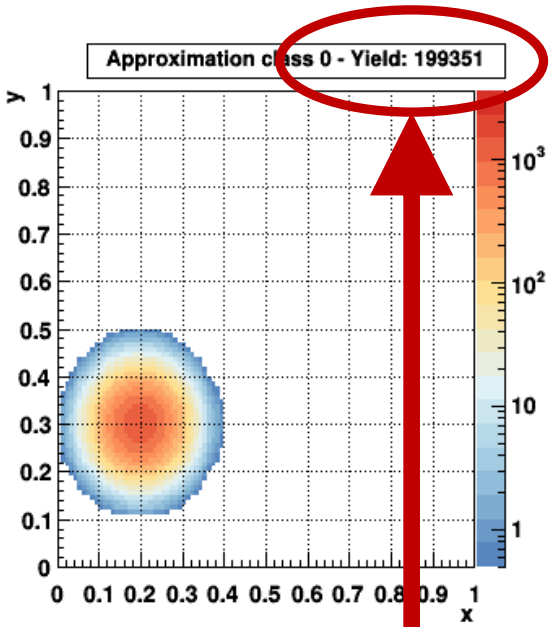




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$\chi^2/\text{NDF} = 9488.33$



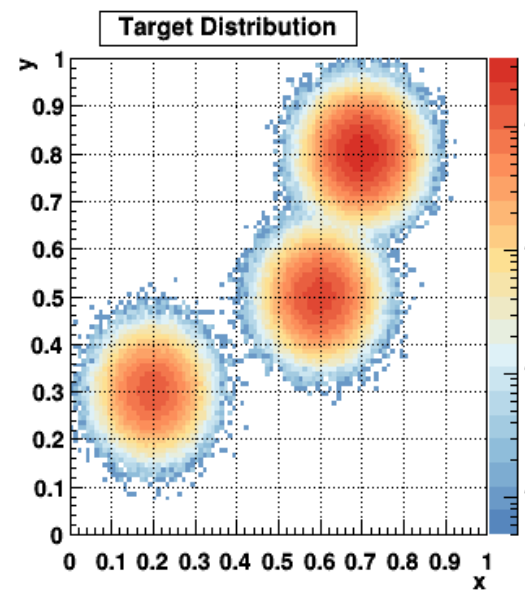
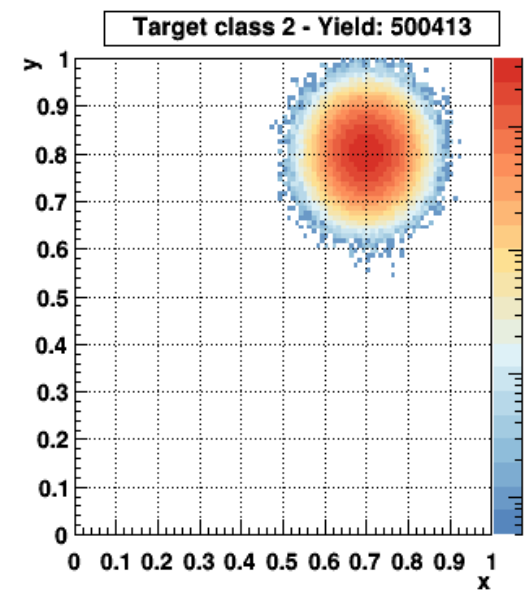
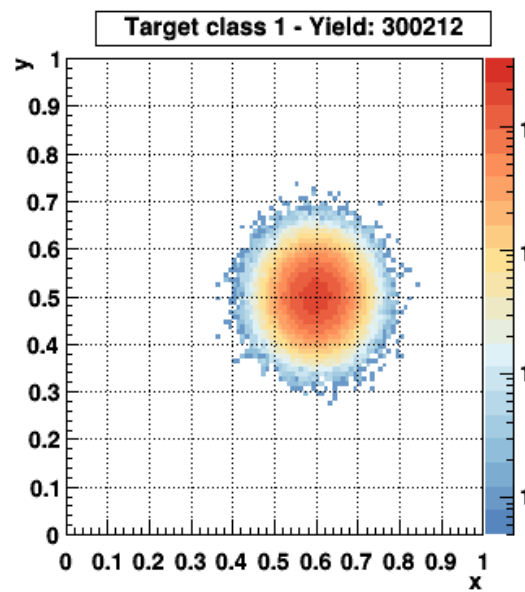
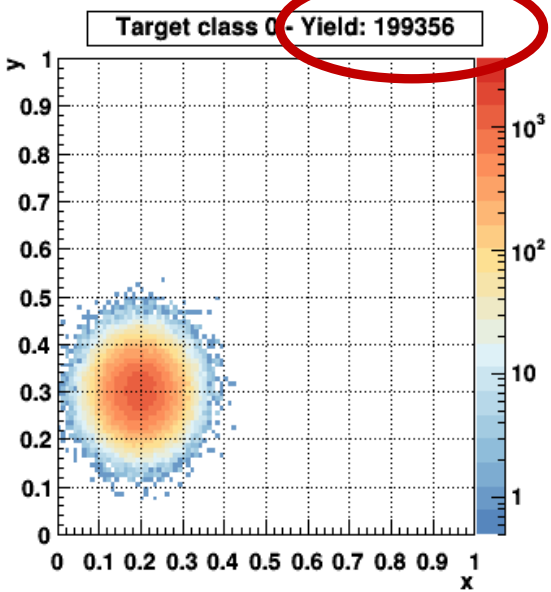


**Spot on!**

iteration 050000

$\chi^2/\text{NDF} = 0.95$

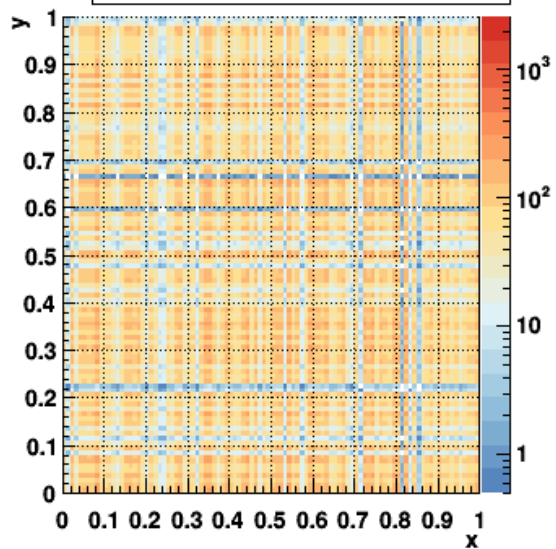
**Not too shabby!**



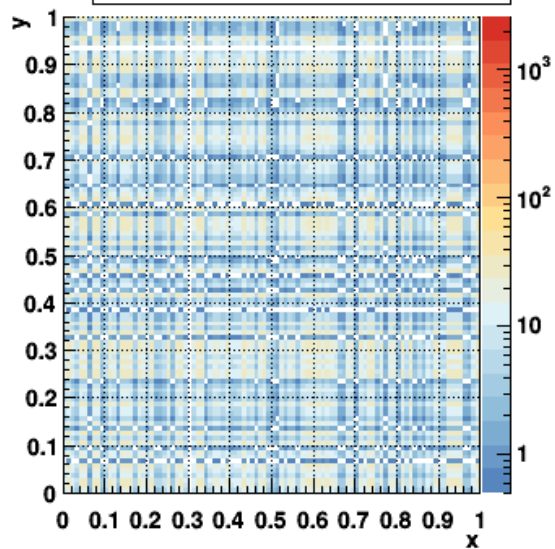
THAT ONE WAS TOO EASY!

Modeling separated Gaussians is trivial,  
what about more complicated distributions?

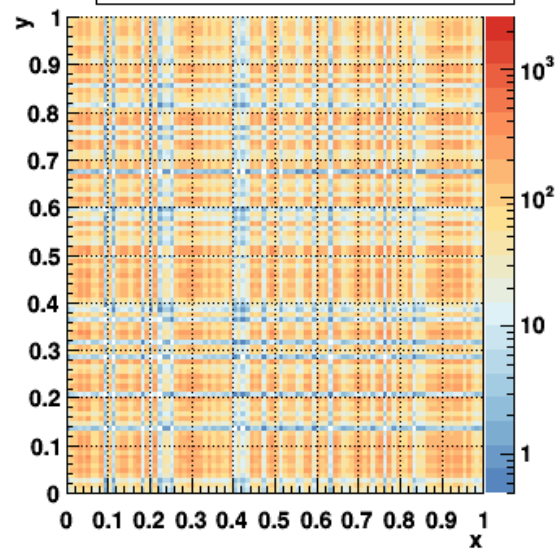
Approximation class 0 - Yield: 632716



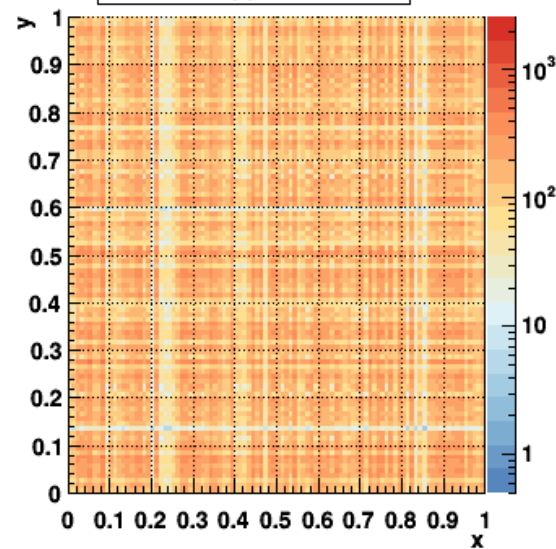
Approximation class 1 - Yield: 87523



Approximation class 2 - Yield: 810465



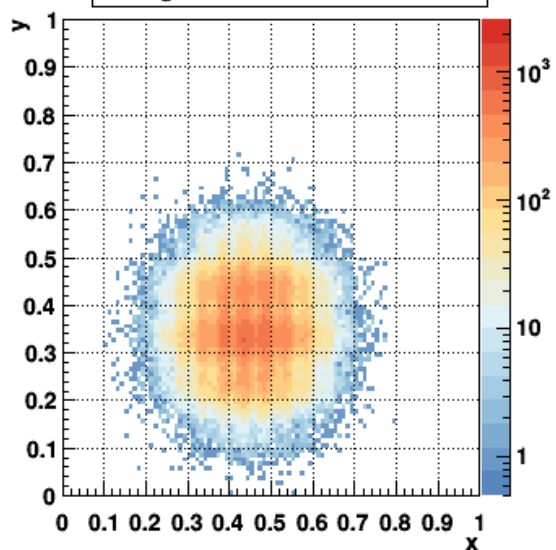
Current Approximation



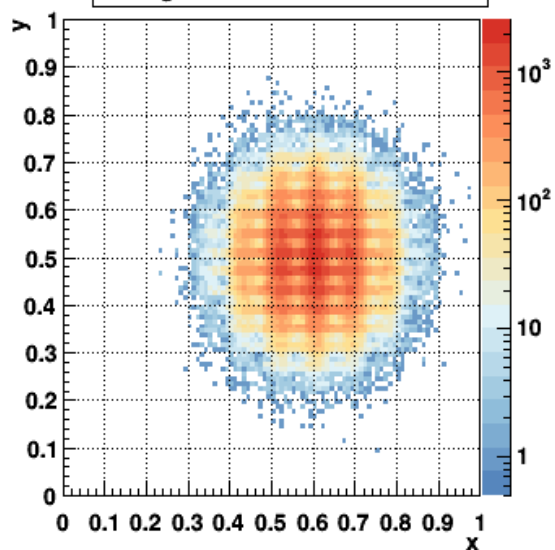
iteration 000000

$\chi^2/\text{NDF} = 5236.58$

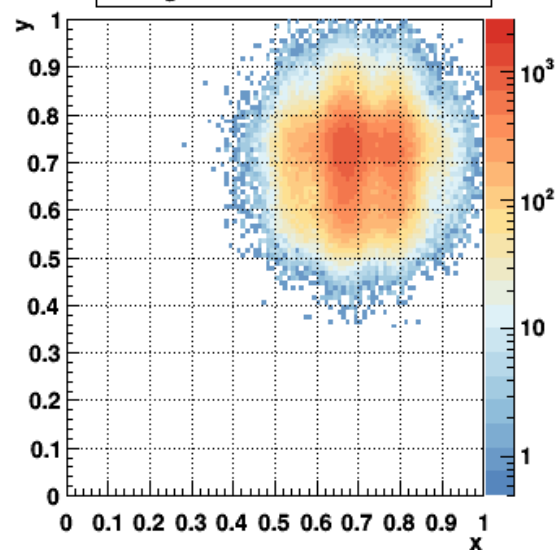
Target class 0 - Yield: 200333



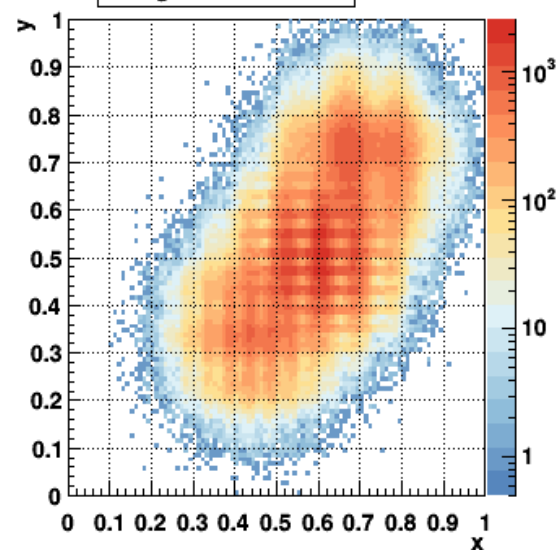
Target class 1 - Yield: 500323



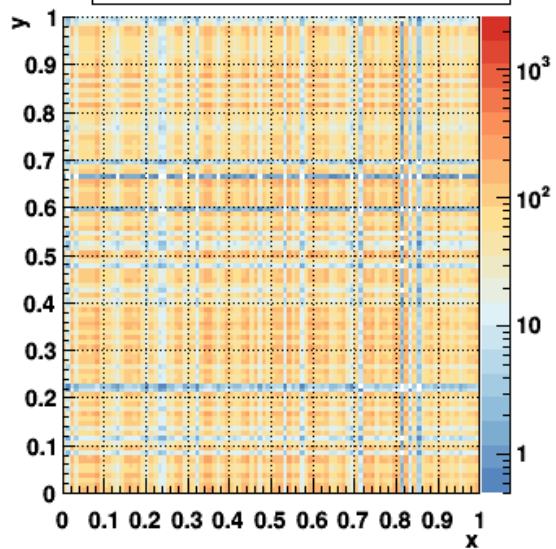
Target class 2 - Yield: 299299



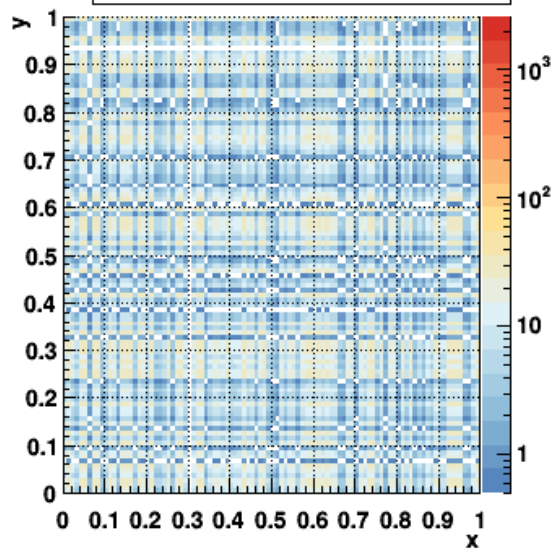
Target Distribution



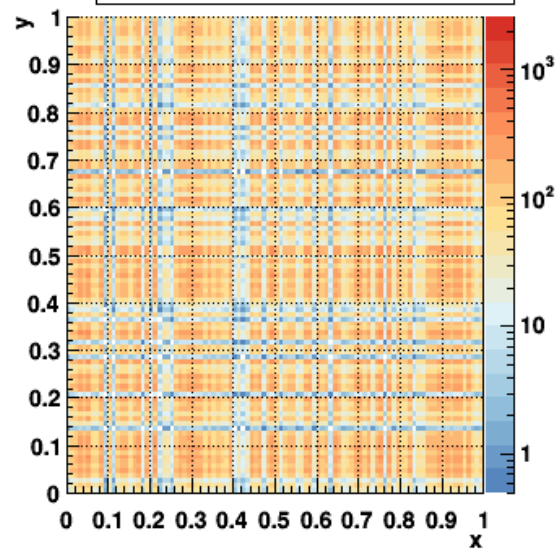
Approximation class 0 - Yield: 632716



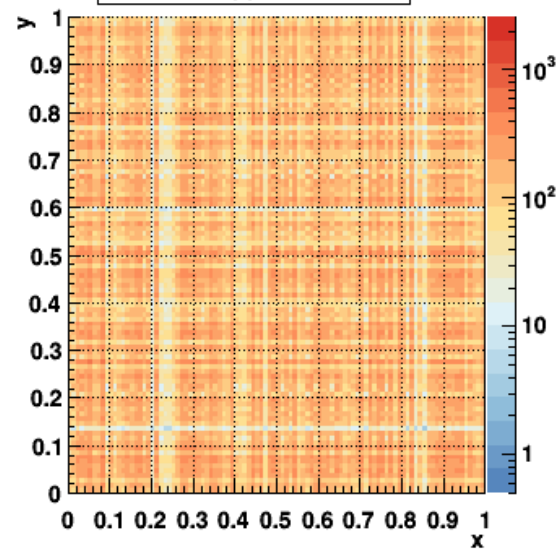
Approximation class 1 - Yield: 87523



Approximation class 2 - Yield: 810465



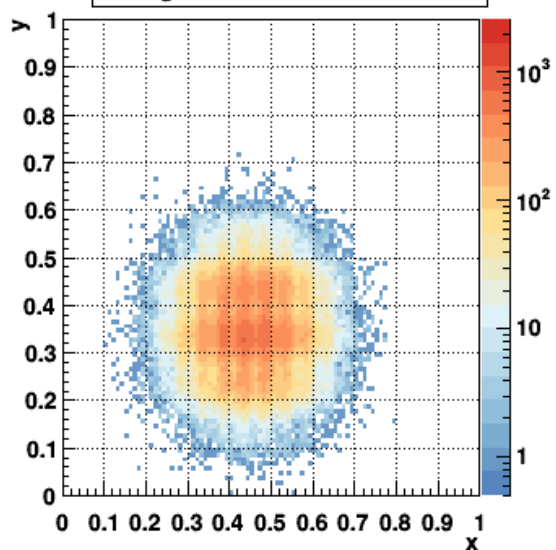
Current Approximation



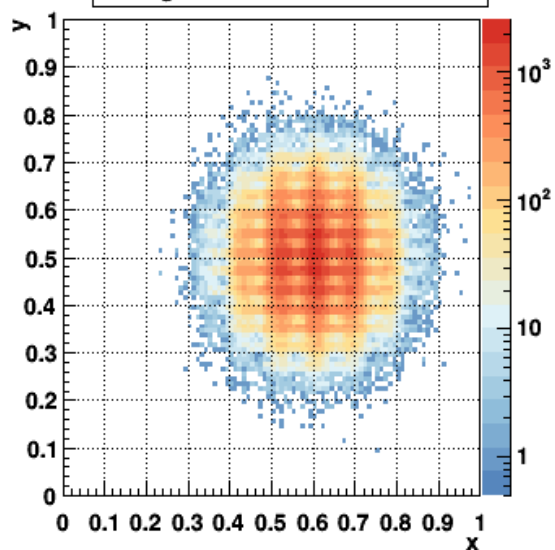
iteration 000000

$\chi^2/\text{NDF} = 5236.58$

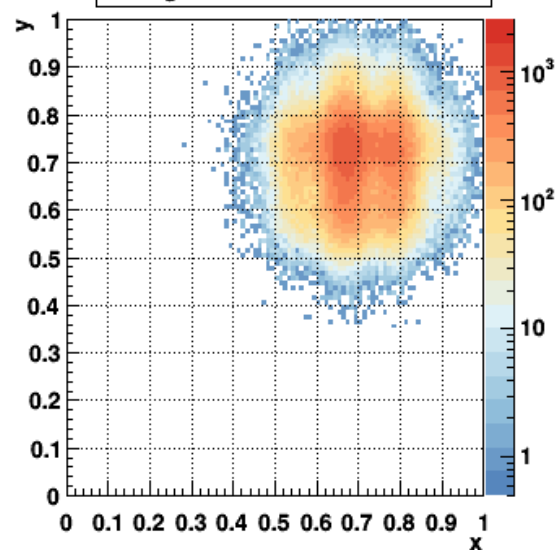
Target class 0 - Yield: 200333



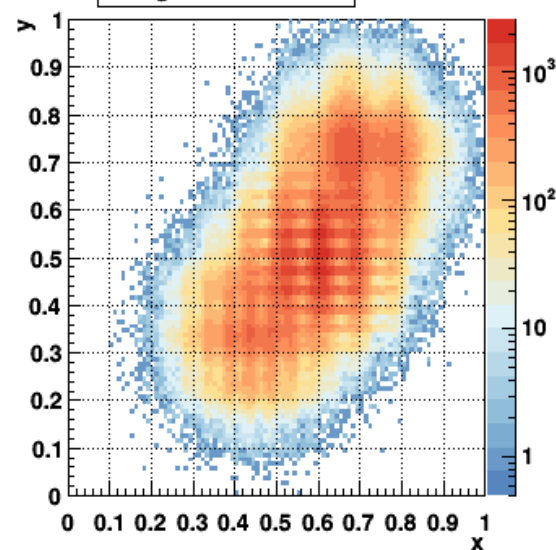
Target class 1 - Yield: 500323



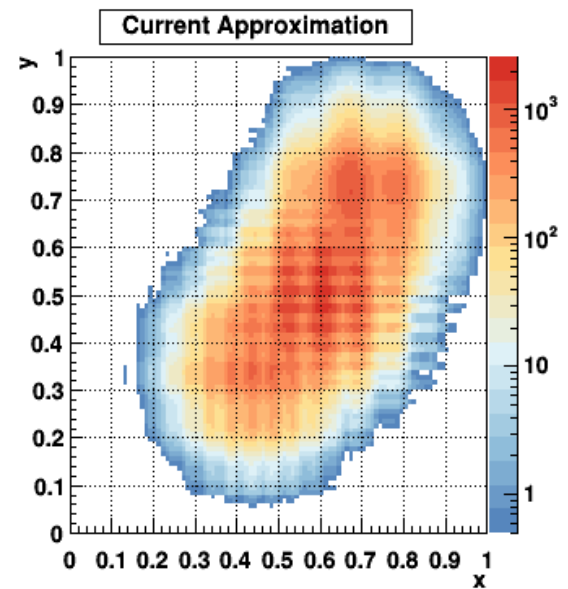
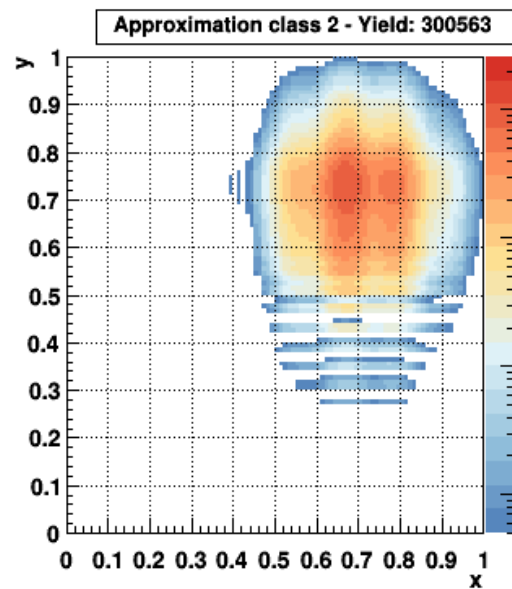
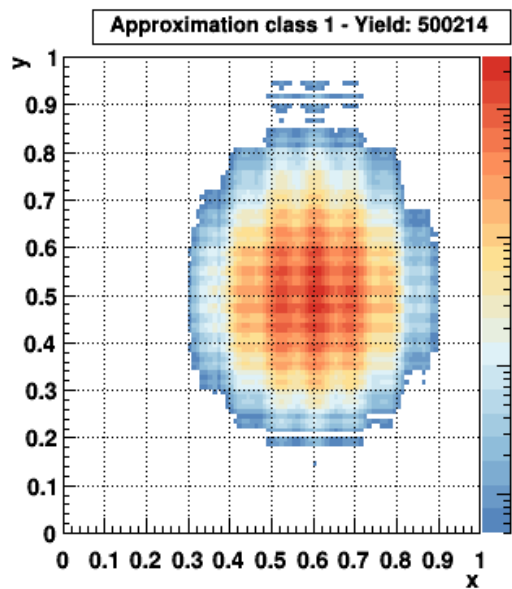
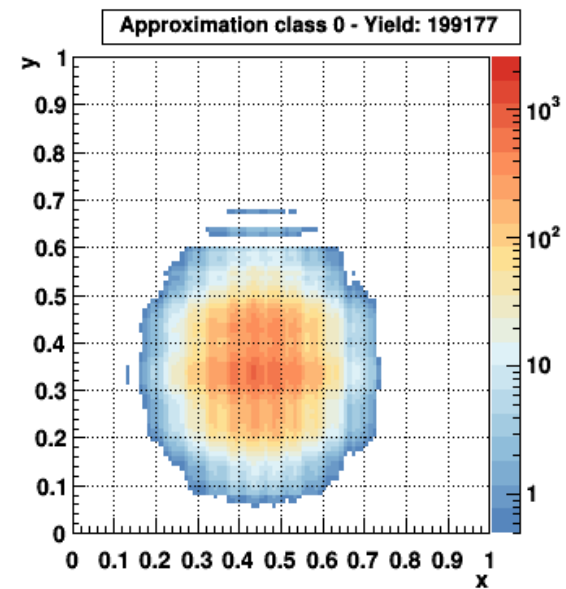
Target class 2 - Yield: 299299



Target Distribution

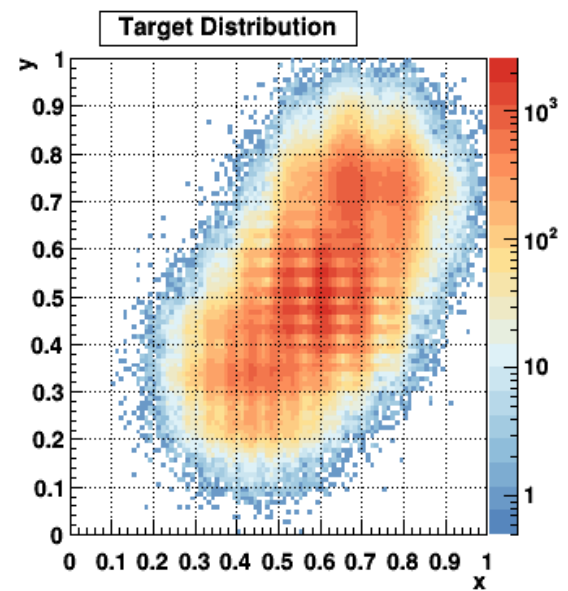
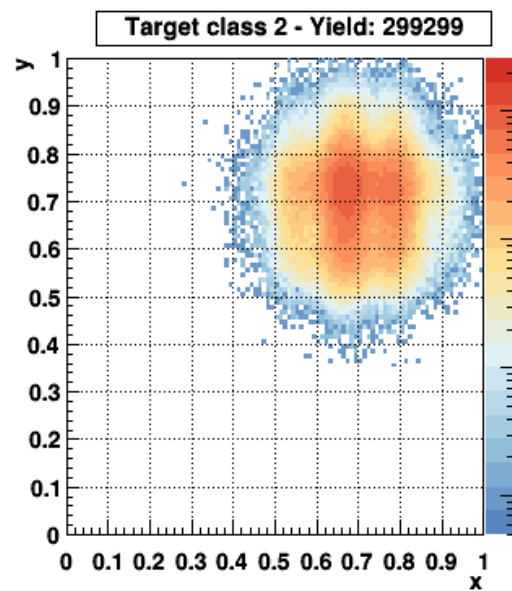
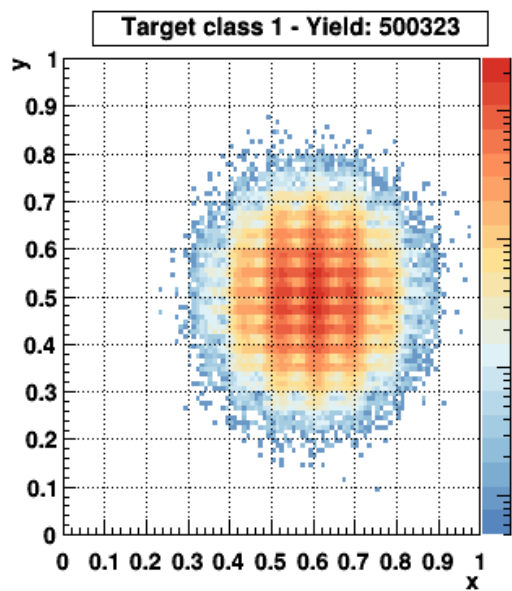
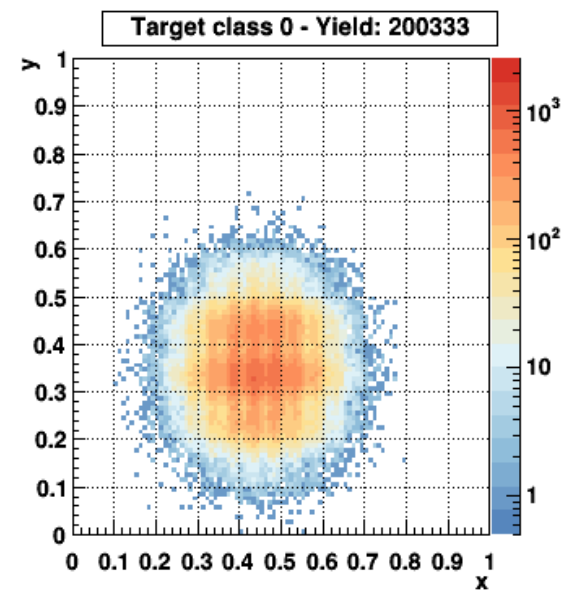






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$\chi^2/\text{NDF} = 2.44$



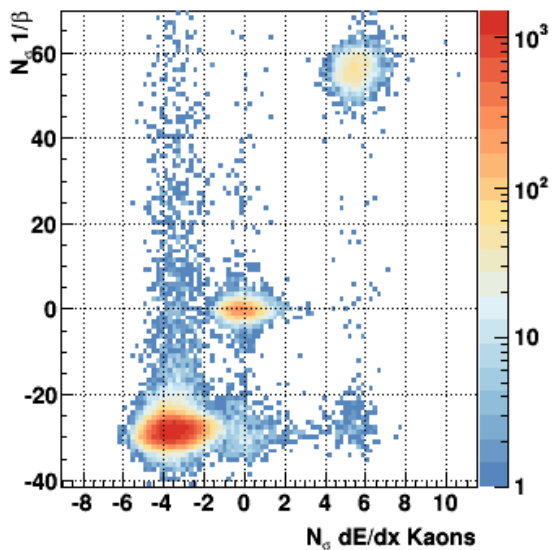
# BETTER... BUT STILL UNREALISTIC

How does it perform with less contrived examples?

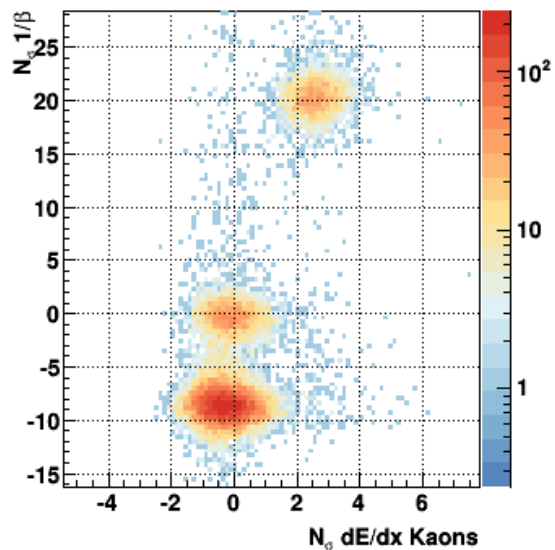
# SIMULATED DATASET: 1 BILLION TRACKS

- ❖  $\pi/K/p$  momentum spectra approximates those from Au+Au 62.4 GeV
- ❖ Time-of-Flight and  $dE/dx$  measurements are made for each particle
  - ❖ Non-Gaussian Shapes
  - ❖ Landau distribution for 15-45  $dE/dx$  hits with highest 30% of measurements rejected for each track
  - ❖ Possibility of similar momentum tracks merging and having higher  $dE/dx$  measurements
  - ❖ Possibility of ToF measurement mismatches between different particles
  - ❖ Student's t-distribution for ToF measurement resolution
- ❖ Momentum has limited resolution (2% on curvature)
- ❖ Finite momentum bin width size (50 MeV bins)
- ❖ **It's dirty, it's messy, and it's realistic**

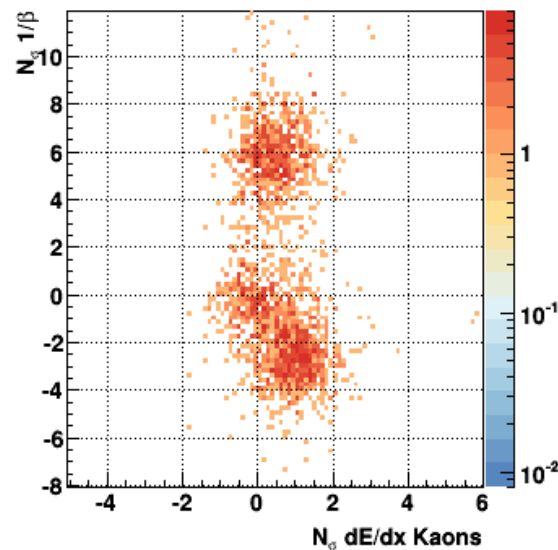
Simulated Data for  $p = 0.525$  GeV



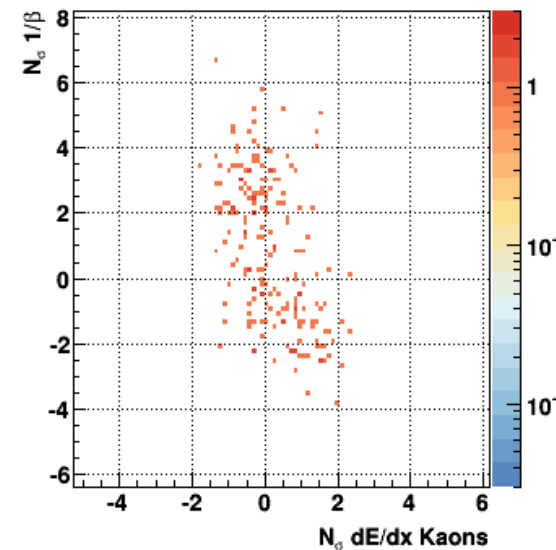
Simulated Data for  $p = 1.025$  GeV



Simulated Data for  $p = 2.025$  GeV

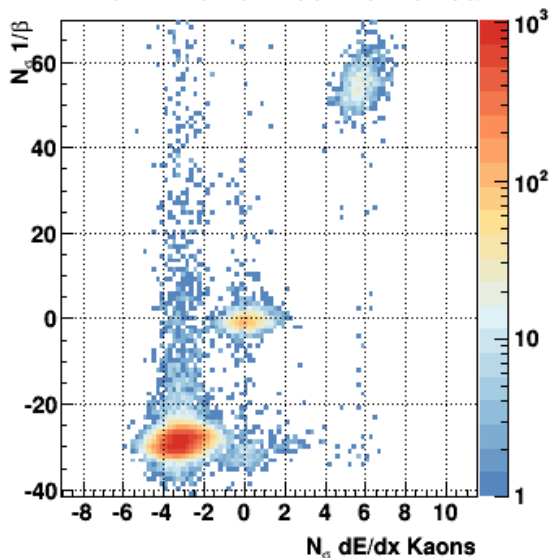


Simulated Data for  $p = 3.025$  GeV

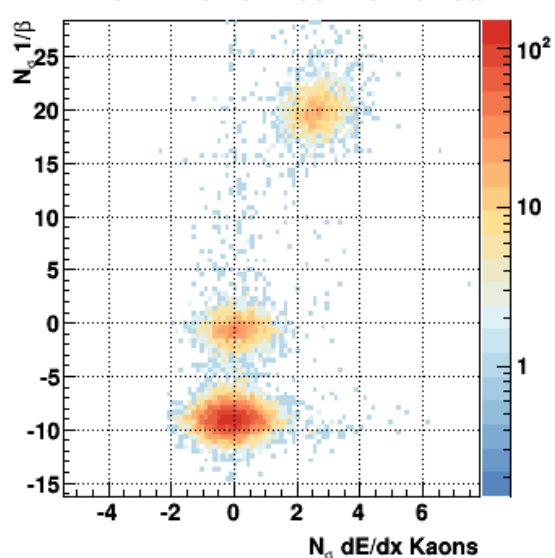


Pretty similar, I'm not trying to brush anything under the rug here.

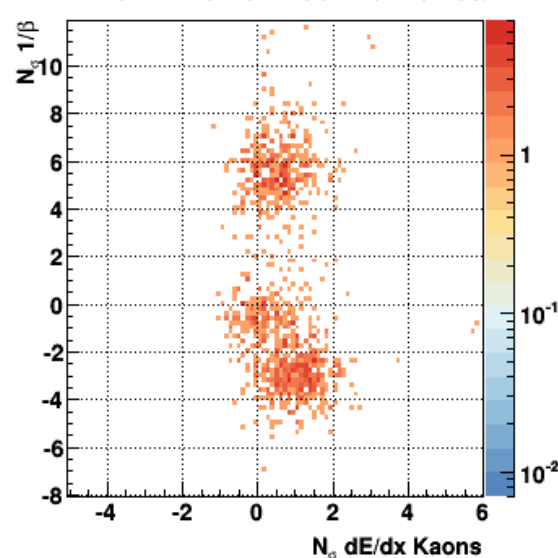
STAR Raw 62.4 GeV Au+Au Data



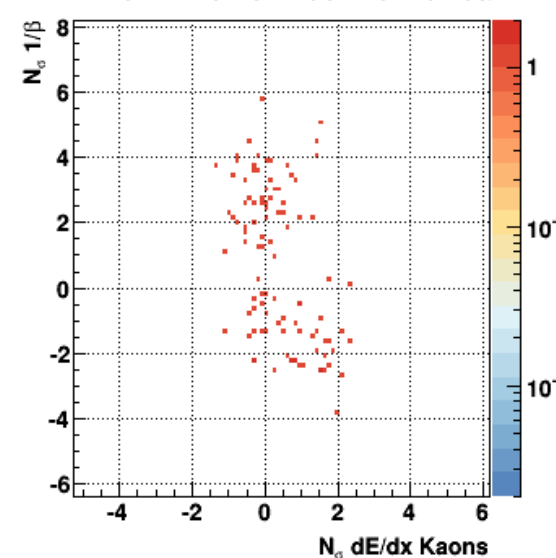
STAR Raw 62.4 GeV Au+Au Data



STAR Raw 62.4 GeV Au+Au Data



STAR Raw 62.4 GeV Au+Au Data



## NOW WE'LL COMPARE...

We're about to look at the ratio of Gaussian model fits and Pidrix fits over the true yields from simulations.

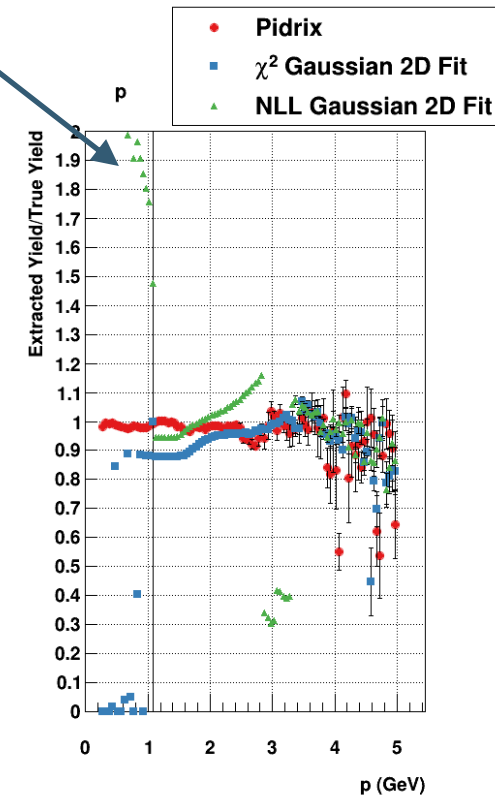
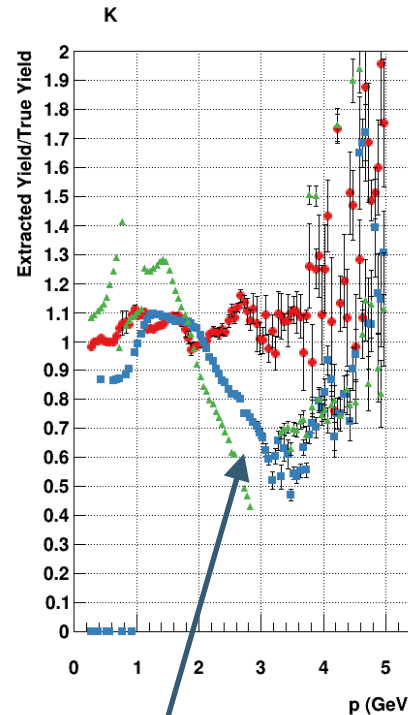
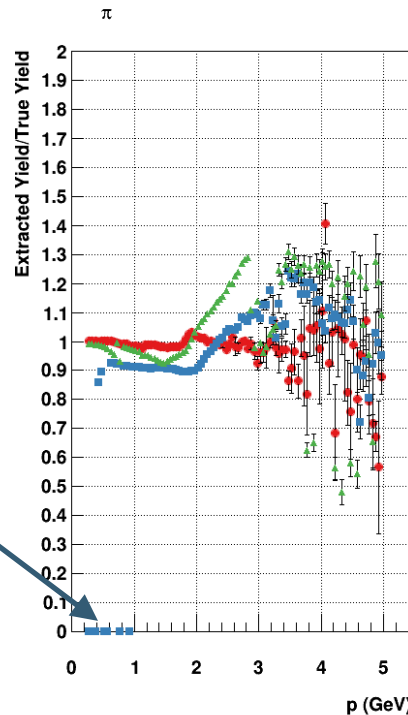
# Wow!

Yuck, failed fits. What do we have to fudge to get these to work?

***Pidrix outperforms the model fits, especially at high  $p$***

We don't realize when we make this mistake because the spectra still looks realistic.

Off the charts!



# WHAT ABOUT ERRORS?

- ❖ Perform Gibb's sampling on elements of  $U$  and  $V$ 
  - ❖ Interpolate between neighboring elements to set scale of fluctuations (to enforce symmetry of the transition)
    - ❖ Add one to interpolation to avoid zeros from propagating
  - ❖ Add Gaussian value with width of 10% of the fluctuation scale
  - ❖ Reject negative values
- ❖ Use k-means clustering between samples to associate particles
  - ❖ Various norms ('means and yields', symmetrized Kullback-Leibler, etc)
- ❖ This results in likelihood distributed  $U$  and  $V$  matrices
  - ❖ Directly compute standard deviations of yields, means, single particle distributions, or whatever you want

# HOW DO THEY COMPARE?

What factor would the yield error have to be multiplied by to be correct?

	Pidrix	Chi-squared	Log-likelihood
$\pi$	4.6	750000	77
K	7.2	230000	160
P	4.2	270000	360

Minuit errors don't mean much when the model isn't exactly right.

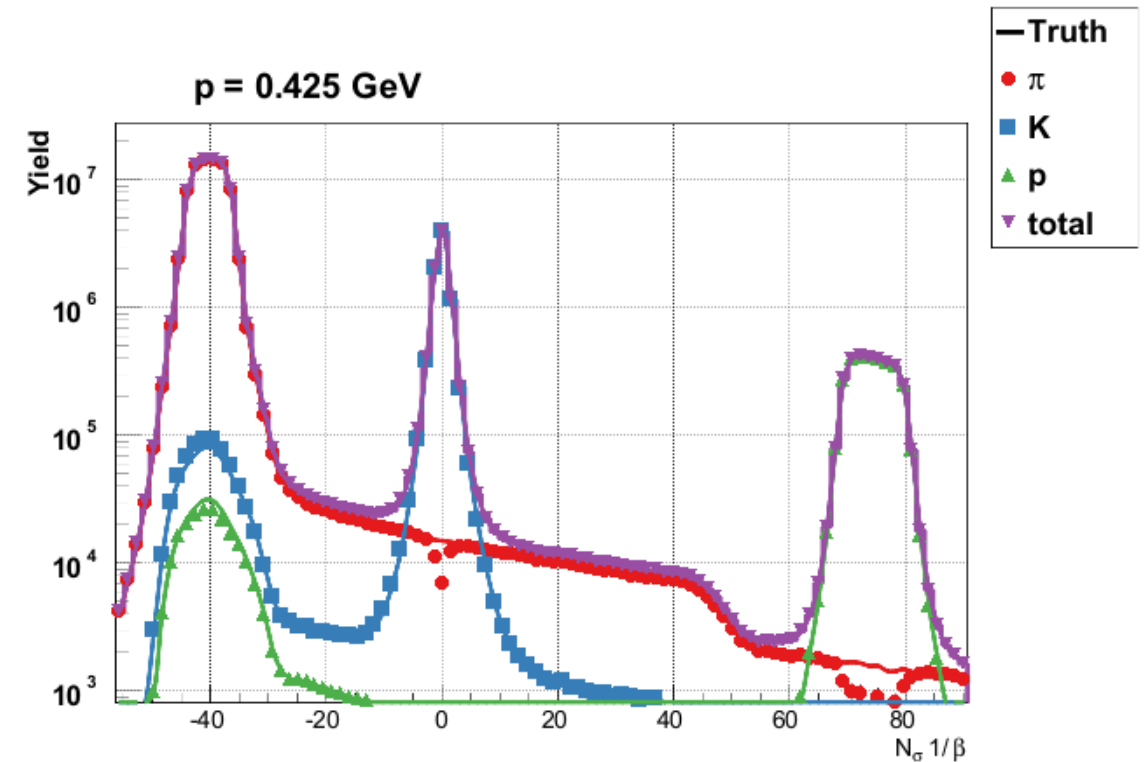
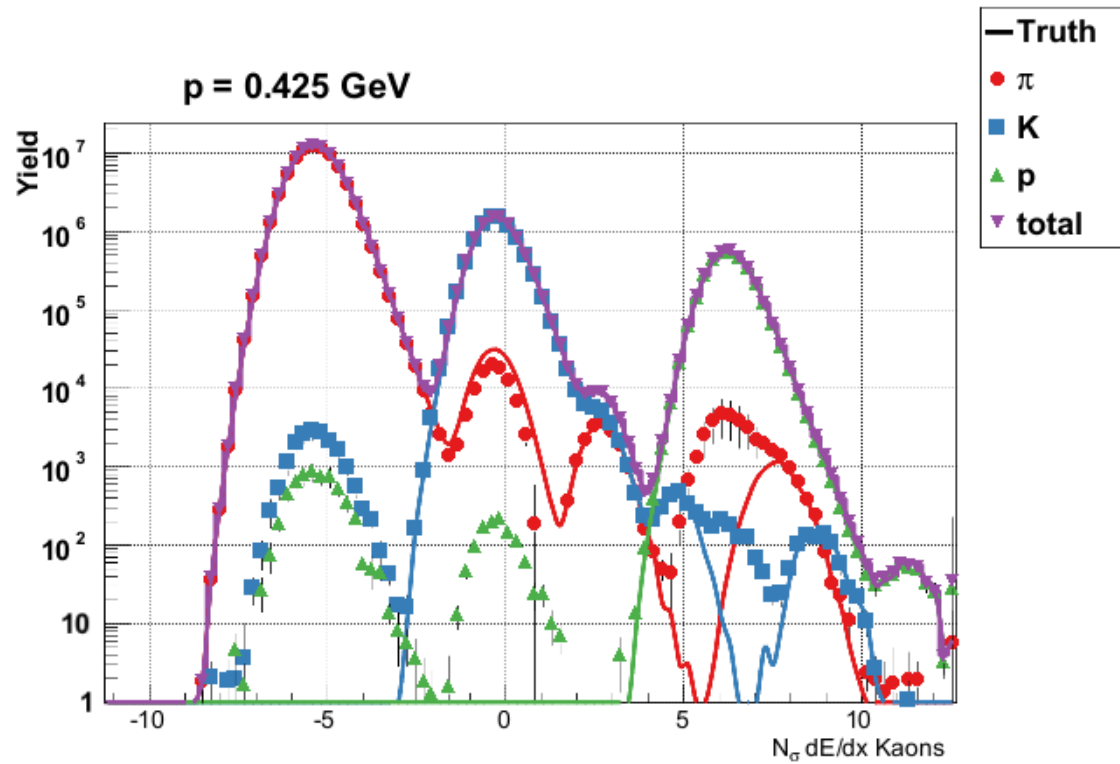
Pidrix gives far more accurate errors though still not great.

**Small errors aren't better if they're wrong!**



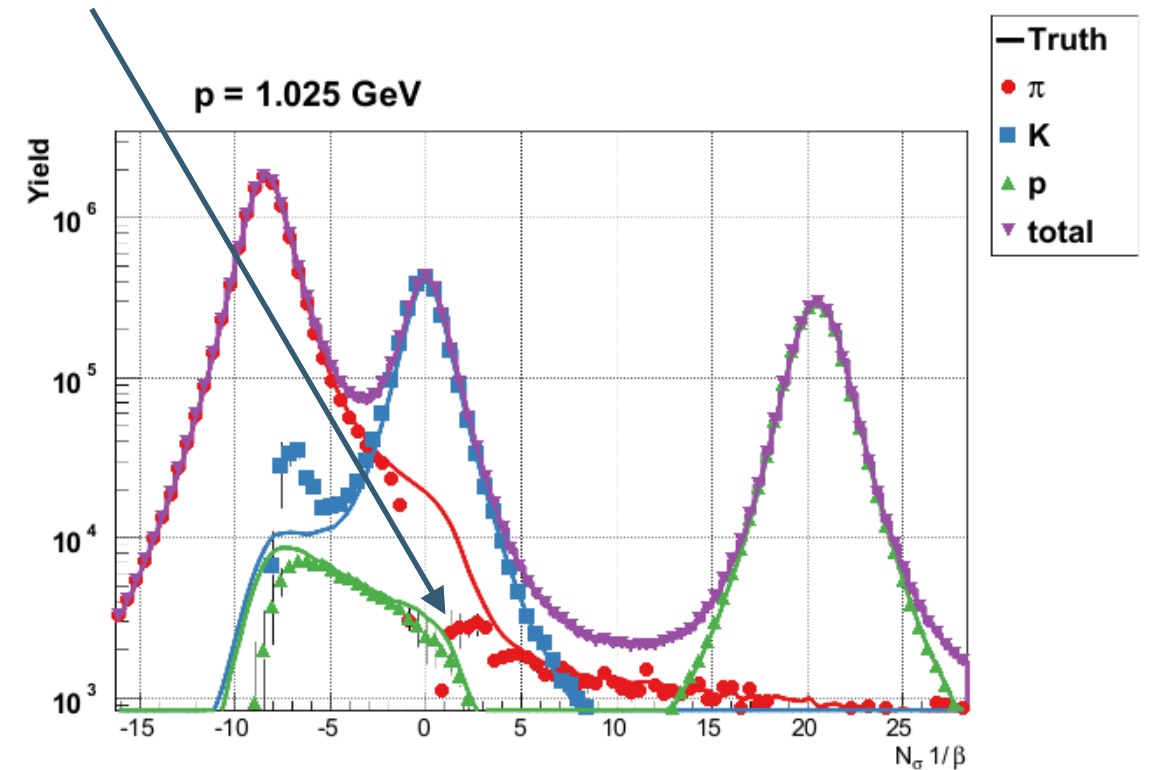
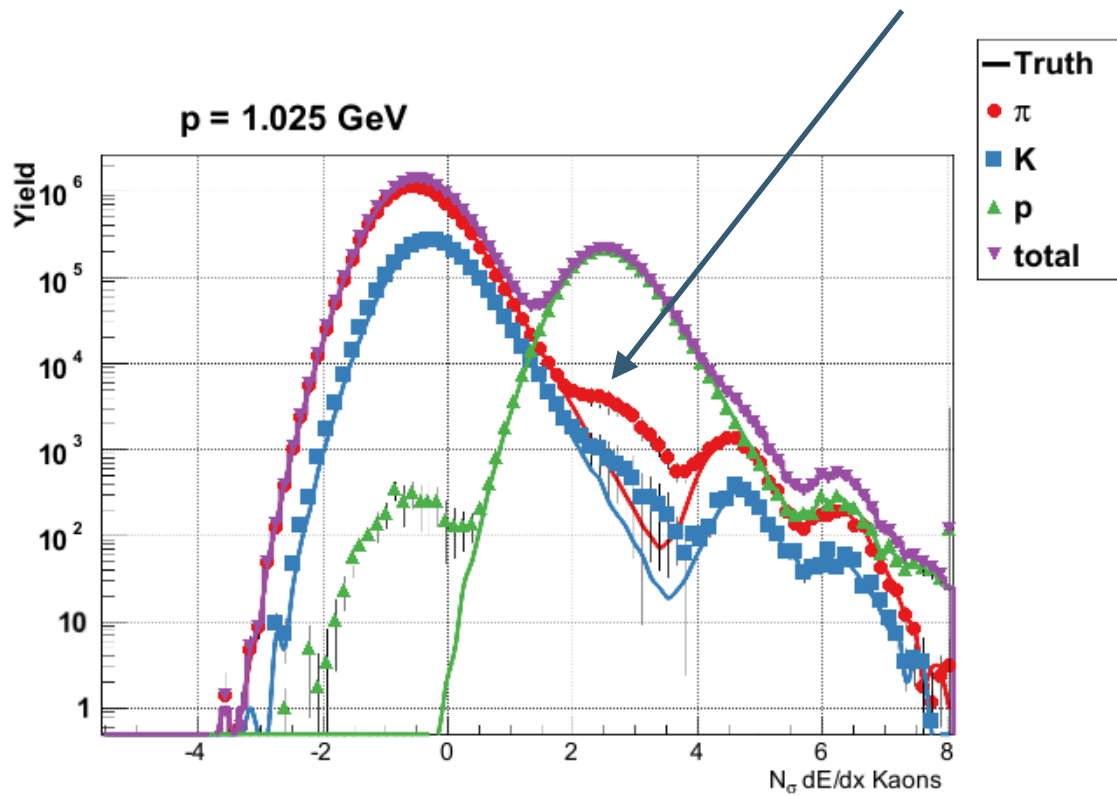
# BACK TO U AND V...

Their columns and rows have physical meaning.



# SOMETIMES MINOR MIXING ISSUES UNDER PEAKS

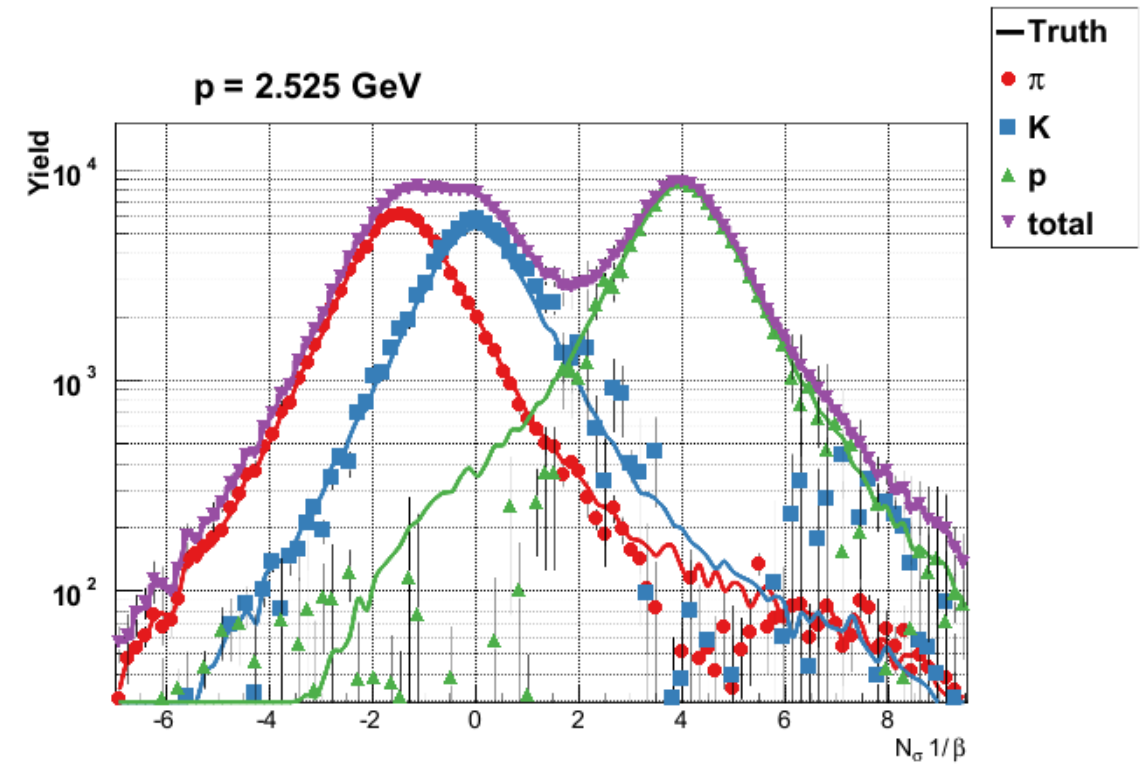
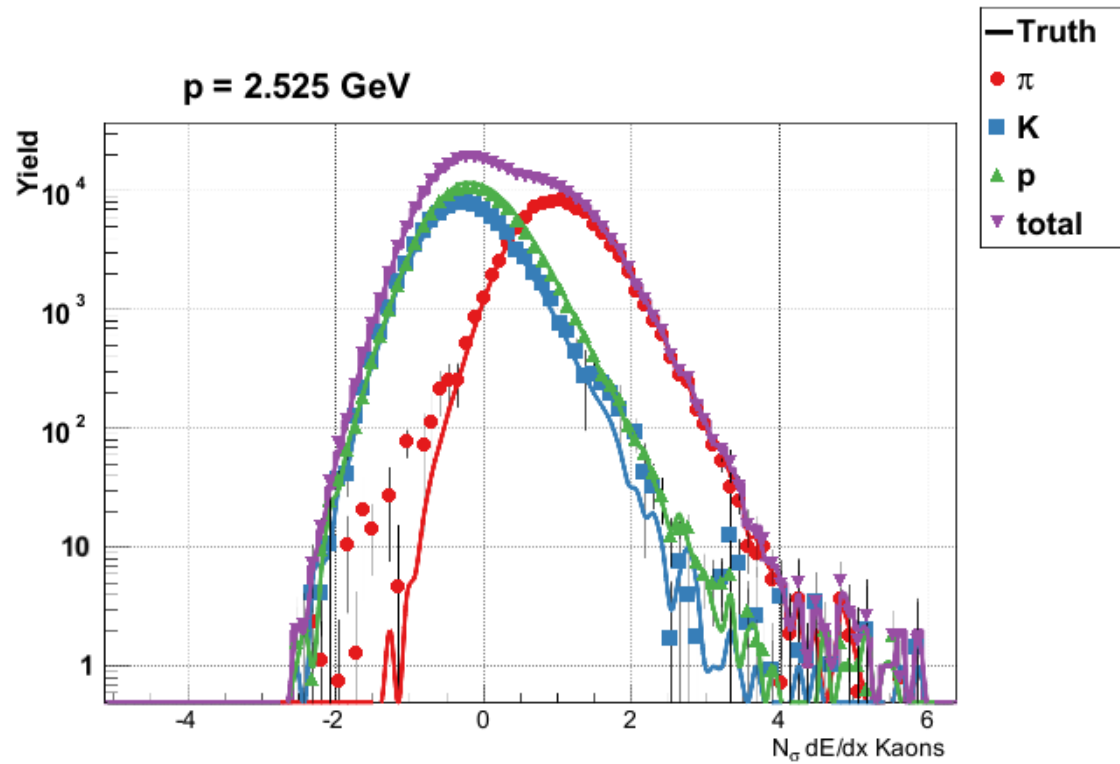
Is sampling not fully exploring the space?



Or does momentum smearing drive this?

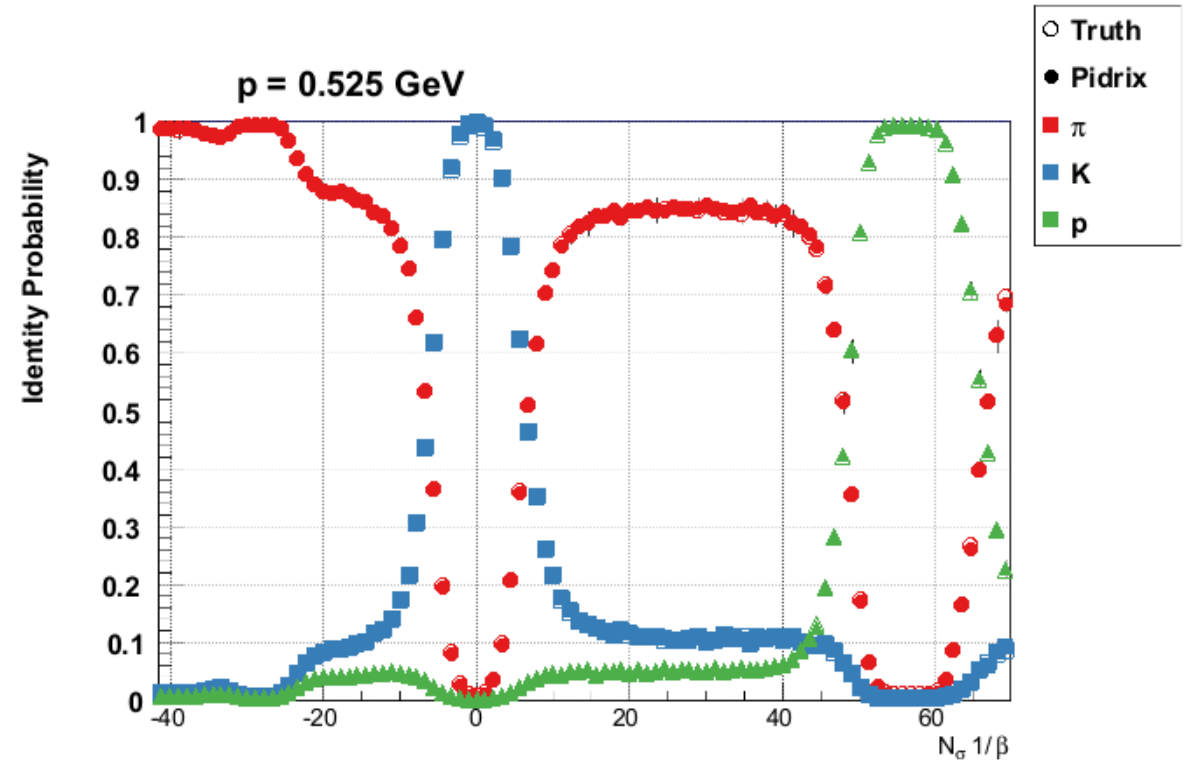
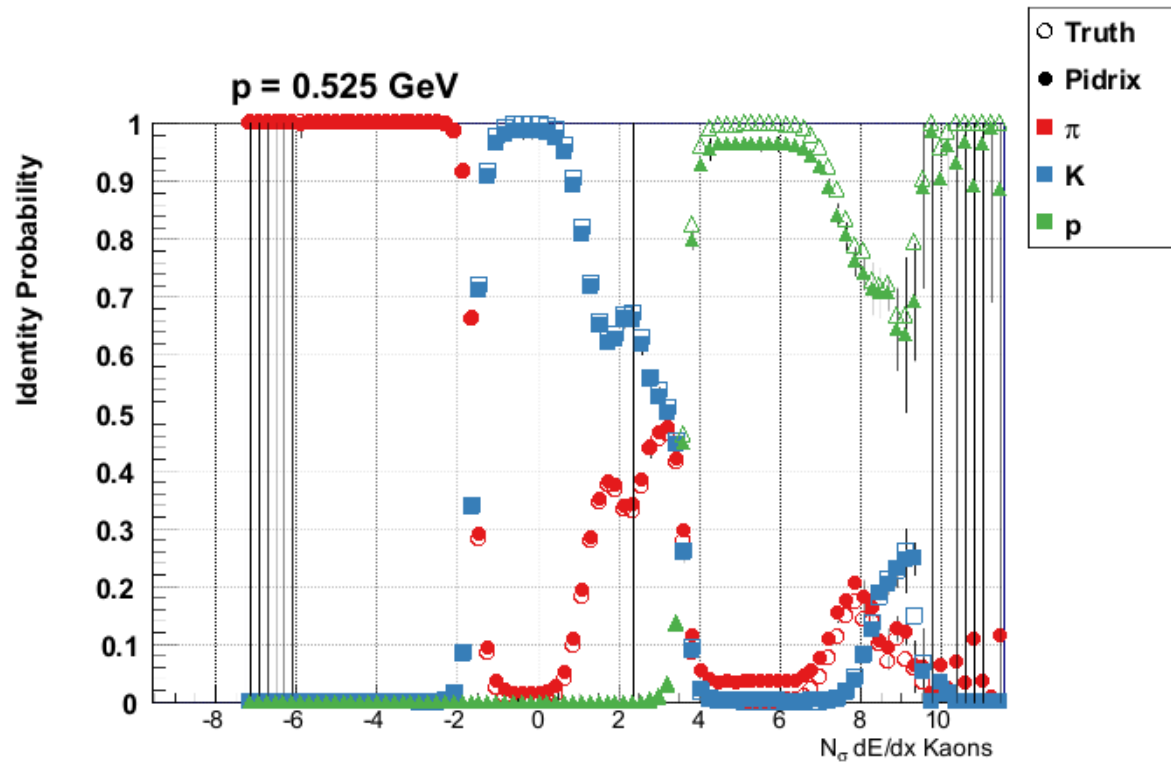
# PIDRIX PARTICULARLY EXCELS AT HIGH P

This region is extremely difficult to model.



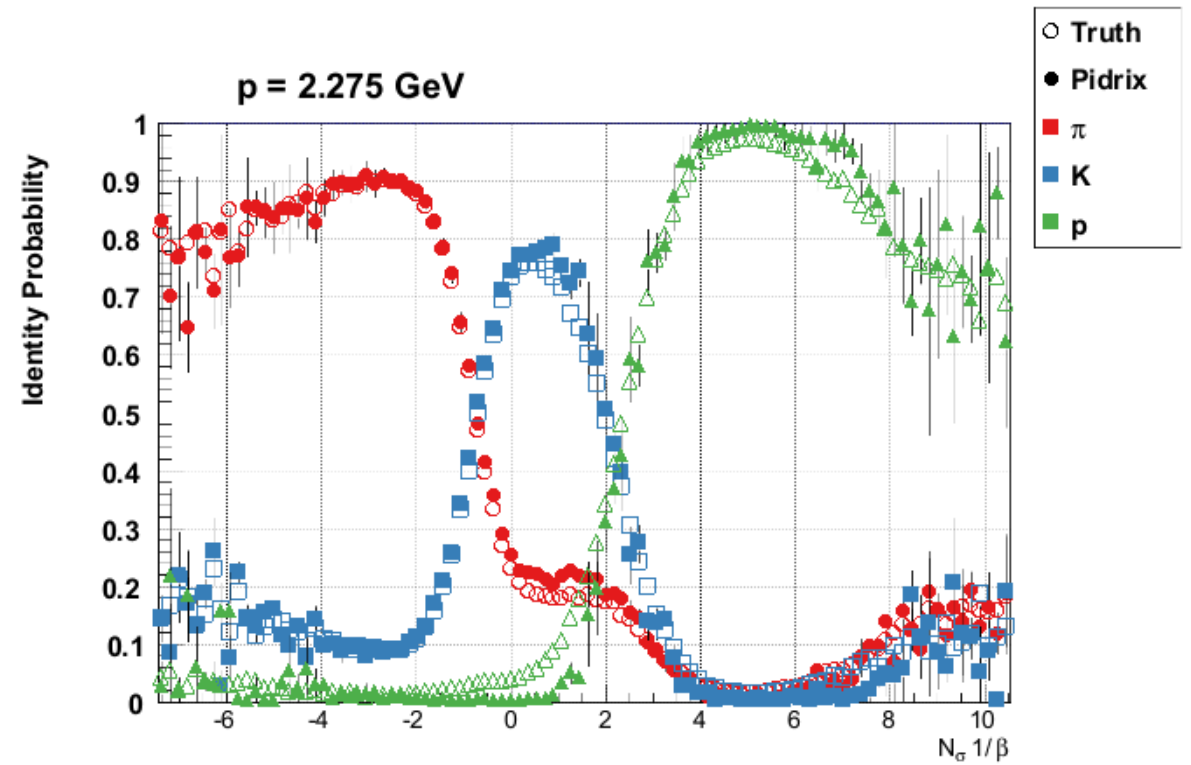
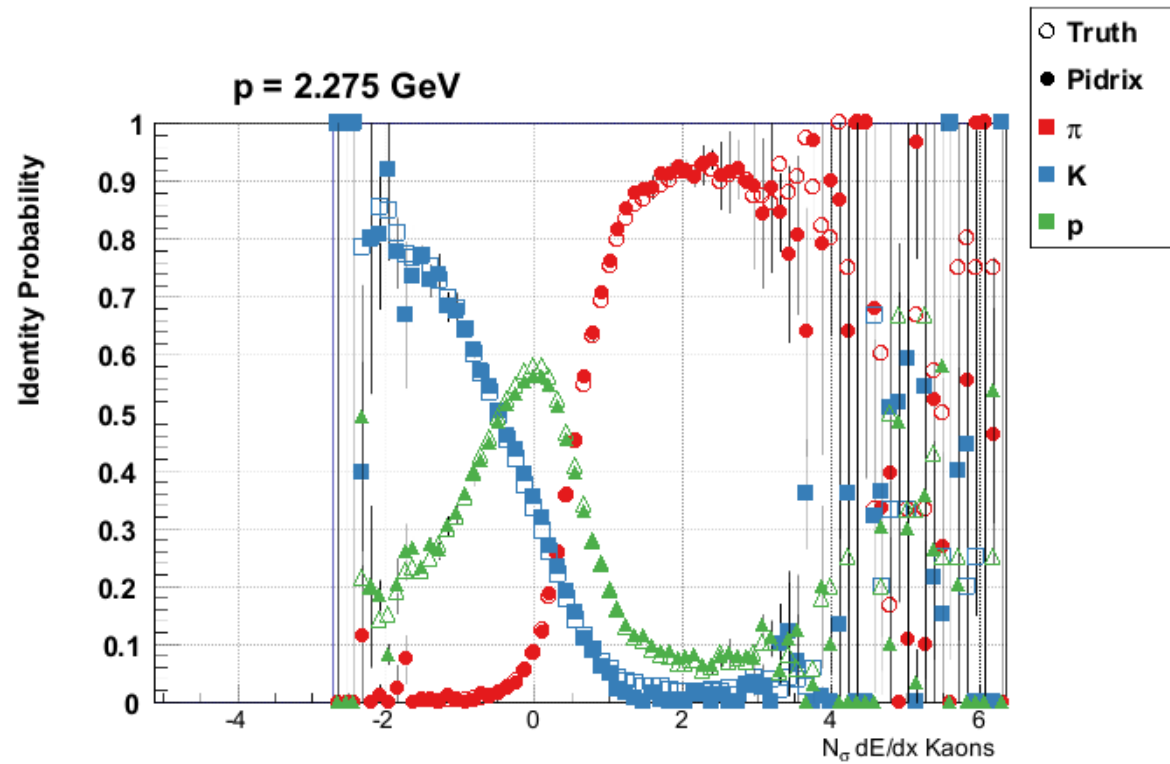
# PURITY AND EFFICIENCY

Want some help picking your cuts for high purity?



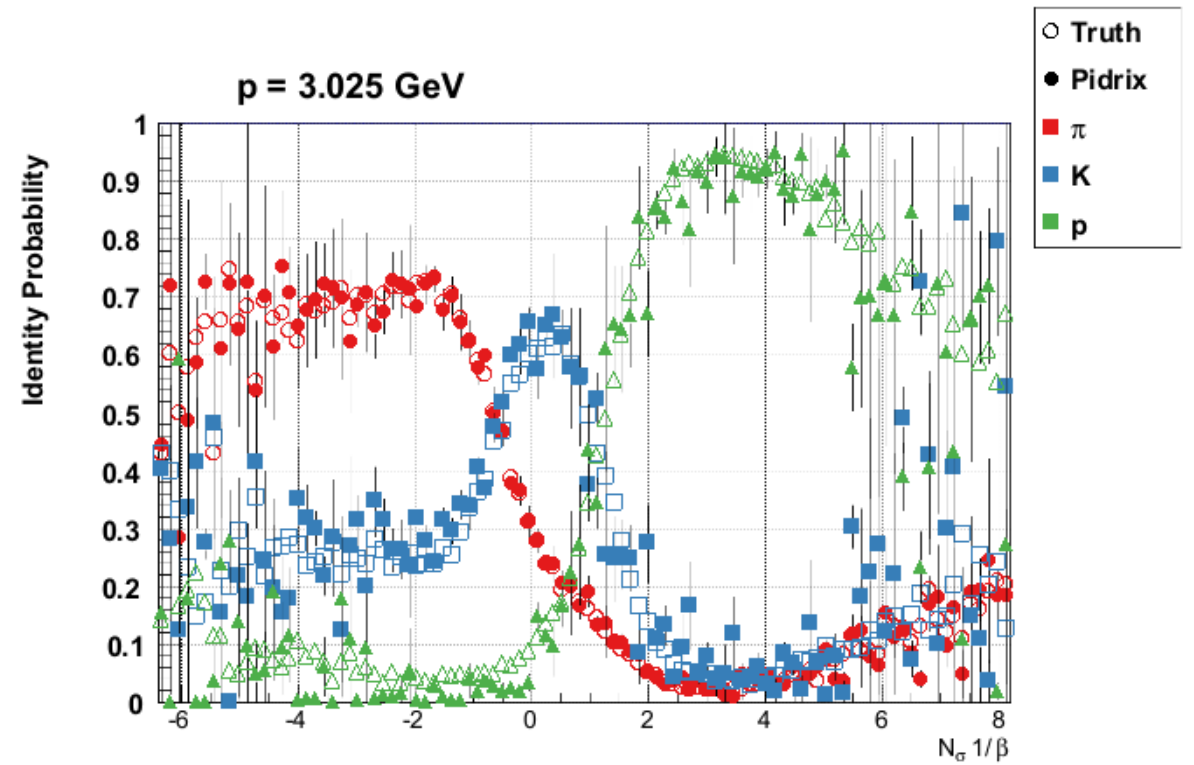
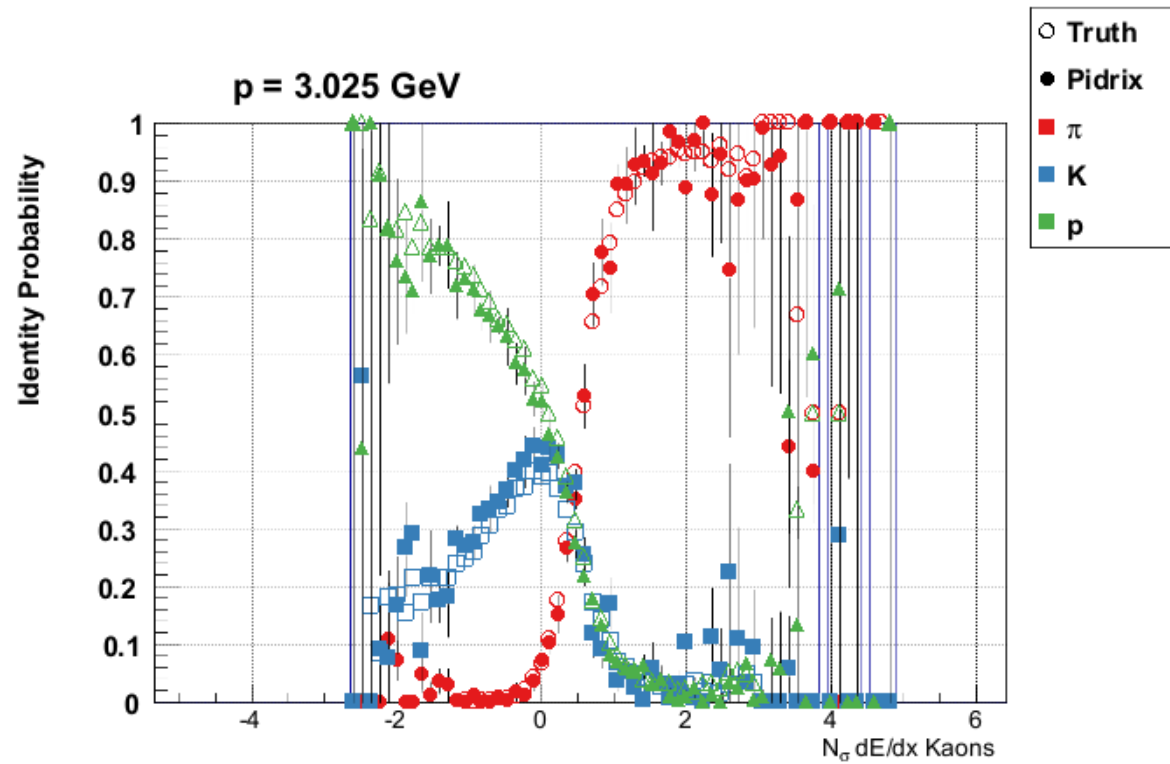
# PURITY AND EFFICIENCY

Or calculating what your cut efficiency is?



# PURITY AND EFFICIENCY

These can be extracted directly from U and V.



# REITERATION

- ❖ We're doing a negative log-likelihood minimization
  - ❖ Just not using a model and being clever about how we do this
- ❖ The worst assumption is that the measurements are uncorrelated
  - ❖ Even when this is broken by momentum resolution and finite bin width it still works
  - ❖ Progress is being made to accommodate breaking this assumption
- ❖ The same code applies to different cuts, binnings, and even experiments
  - ❖ No model means no model tuning
- ❖ It seems to work very well!

# LESS IS MORE?

- ❖ A new method for particle identification has been described
  - ❖ Faster, easier, and less assumption oriented than model based fitting
  - ❖ Extremely promising accuracy observed so far
- ❖ FOSS library is available:
  - ❖ <https://github.com/sangaline/pidrix>
  - ❖ Documentation is lacking... examples coming soon
- ❖ Interested in using this method?
  - ❖ Please email [esangaline@gmail.com](mailto:esangaline@gmail.com)
  - ❖ Especially if your experiment has  $>2$  PID measurements for each track!