

Recent Models of Massive Neutrinos

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Based on: JHEP 1307 (2013) 020 [arXiv: 1303.4573]

and work in progress



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Outline

- Some reminders
- Some new stuff
- Some consequences

Neutrino Mass and Mixing

- Data consistent with non-zero neutrino mass

Mass splittings

$$\Delta m_{21}^2 \simeq 7.6 \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{31}^2 \simeq 2.4 \times 10^{-3} \text{ eV}^2$$

Mixing angles

$$\sin^2 \theta_{12} \simeq 0.31$$

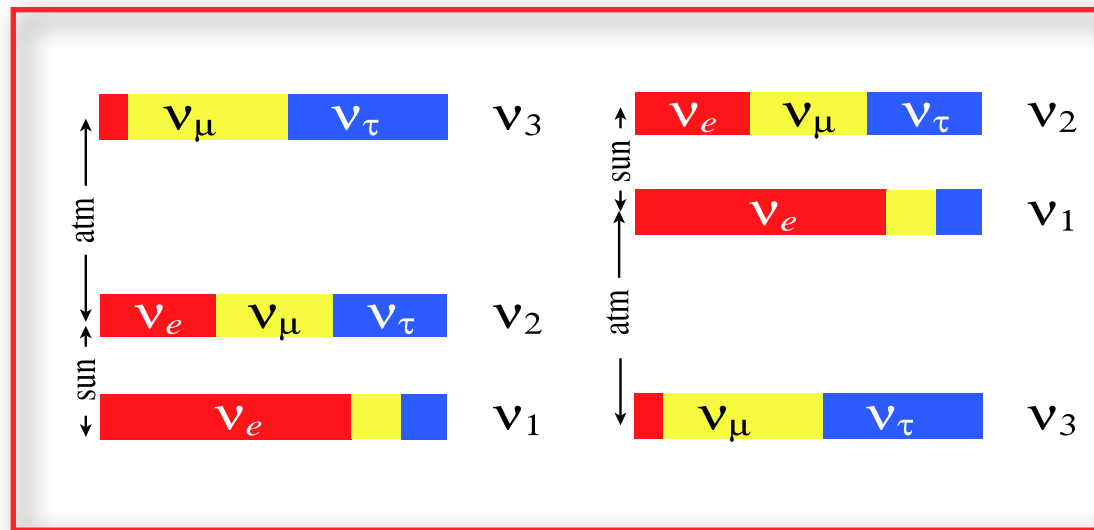
$$\sin^2 \theta_{23} \simeq 0.52$$

$$\sin^2 \theta_{13} \simeq 0.013$$

(Schwetz et al, 2012)

- Evidence for BSM physics
- Number of “known unknowns” remain

Neutrino Mass and Mixing



- Still unknown: mass hierarchy (normal or inverted)
overall mass scale

Still Unknown: Dirac or Majorana

- Dirac: add N_R

Boost (LT)

$$\nu_L = (\nu_L, \bar{\nu}_R) \qquad (\bar{\nu}_L, \nu_R) = N_R$$

CP CP

- Distinct anti-particle
- Lepton number symmetry conserved

Still Unknown: Dirac or Majorana

- Majorana: use “what we got”

$$\nu_L = (\nu_L, \bar{\nu}_R)$$

Boost (LT)

CP

- No distinct anti-particle: $\nu^c \propto \nu$
- Lepton number symmetry is broken

→ $0\nu\beta\beta$ decay

Still Unknown: Origin of Mass

- Masses are very small $\longrightarrow m_\nu \lesssim 0.1 \text{ eV} \sim 10^{-6} \times m_e$

Question: Why are they so tiny?

- Two main approaches  seesaw models
radiative models

- Point of this talk:

A new class

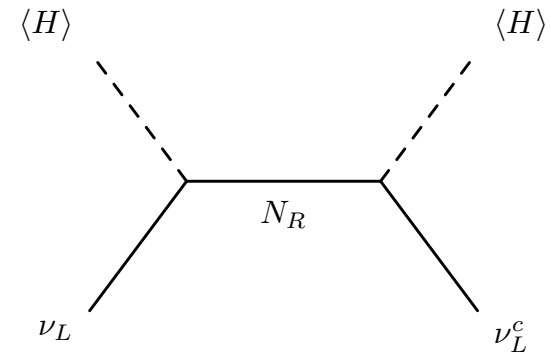


seesaw/radiative models

The Seesaw Mechanism

- Add heavy sterile neutrino

$$\mathcal{L} \supset \lambda_\nu L H N_R + M_R N_R^c N_R \longrightarrow$$



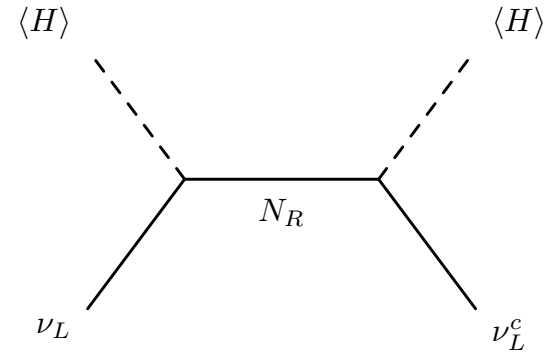
- Small mass due to heavy partner

$$m_\nu \simeq \lambda_\nu^2 \frac{\langle H \rangle^2}{M_R} \longrightarrow M_R \lesssim 10^{14} \text{ GeV}$$

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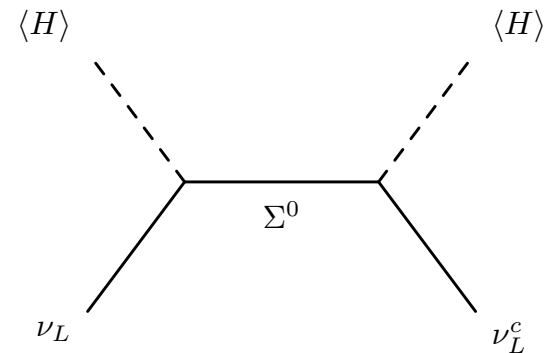
$$m_\nu \simeq \lambda_\nu^2 \frac{\langle H \rangle^2}{M_R} \longrightarrow M_R \lesssim 10^{14} \text{ GeV}$$

- Type-III: use fermion triplet

Foot, Lew, He, Joshi (1989)

$$\Sigma \sim (1, 3, 0) = \begin{pmatrix} \Sigma^0 & \Sigma^+/\sqrt{2} \\ \Sigma^-/\sqrt{2} & -\Sigma^0 \end{pmatrix}$$

$$\mathcal{L} \supset \lambda_\nu L \Sigma H + M_\Sigma \Sigma^c \Sigma \longrightarrow$$



The Type-II Seesaw Mechanism

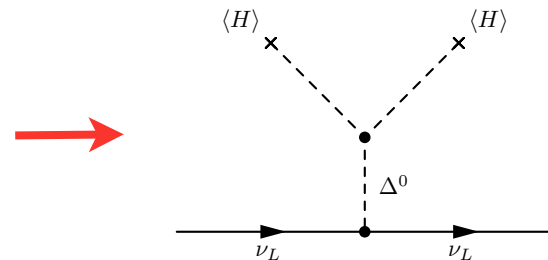
- Use scalar triplet $\Delta \sim (1, 3, 2) = \begin{pmatrix} \Delta^+ & \Delta^{++} \\ \Delta^0 & -\Delta^+ \end{pmatrix}$

$$\mathcal{L} \supset \lambda_\nu L^c \Delta L - M_\Delta^2 \Delta^2 + \mu H \Delta H \quad \longrightarrow \quad \begin{array}{c} \langle \Delta \rangle \\ \times \\ \vdots \\ \nu_L \quad \nu_L^c \end{array}$$

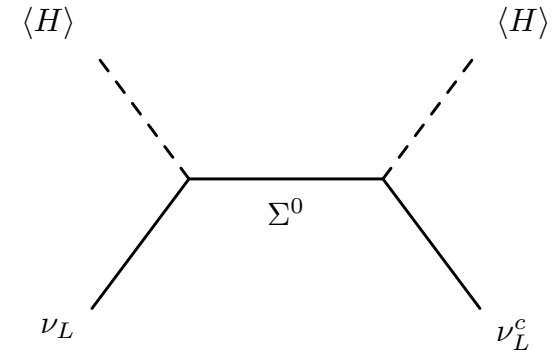
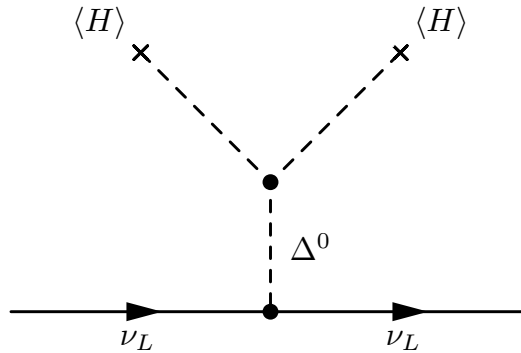
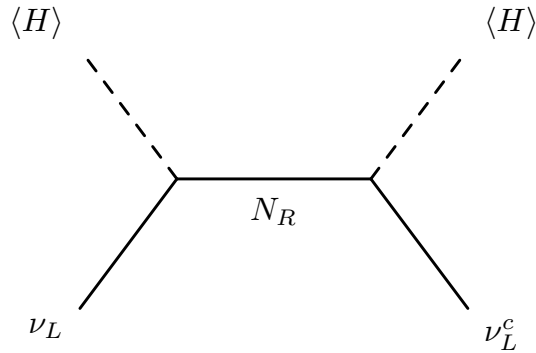
$$m_\nu = \lambda_\nu \langle \Delta \rangle \simeq \lambda_\nu \times \frac{\mu \langle H \rangle^2}{M_\Delta^2} \quad \longrightarrow \quad \text{VEV seesaw} \quad \langle H \rangle = \sqrt{\frac{-\mu_H^2}{\lambda}}$$

- Satisfies constraint $\rho = \frac{M_W^2}{\cos^2 \theta_W M_Z^2} \neq 1 \quad \longrightarrow \quad \langle \Delta \rangle \lesssim \mathcal{O}(\text{GeV})$

Heavy scalar can also explain small mass



Only Three d=5 Seesaws



- Non-renormalizable operator: $\mathcal{O}_\nu = \frac{1}{\Lambda} (LH)^2$
- $d = [\mathcal{O}_\nu] = 5$

$$\frac{1}{\Lambda} \sim \frac{1}{M_R}$$

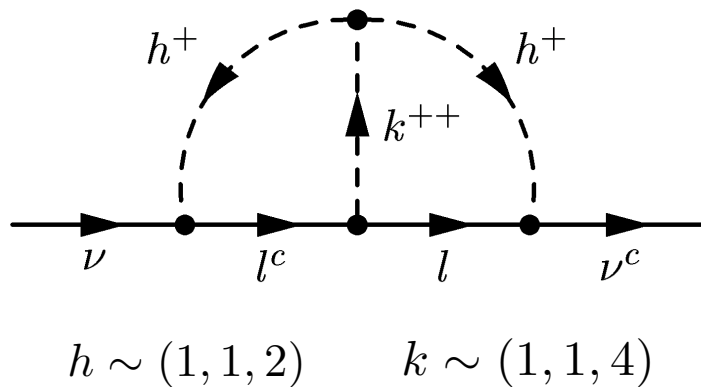
$$\frac{1}{\Lambda} \sim \frac{\mu}{M_\Delta^2}$$

$$\frac{1}{\Lambda} \sim \frac{1}{M_\Sigma}$$

Other main mechanism

Radiative Neutrino Mass

- Add extra scalars (Zee, Babu)
- Mass at one- (Zee) or two-loop level (Zee, Babu)



$$m_\nu \sim \frac{f^2 \lambda}{(16\pi^2)^2} \frac{\mu}{m_k} \frac{m_\ell^2}{m_k} \log(m_k^2/m_h^2)$$

- New physics *can* be light (ie observable)
- No more complex than Type-I (or III) seesaw

Two new scalars

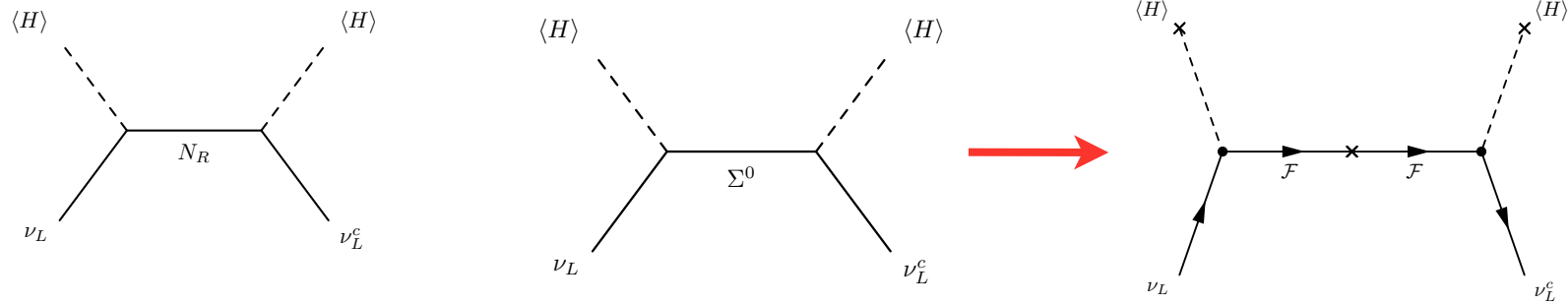
Two new fermions

Radiative mass or seesaw?

Only experiment can tell.

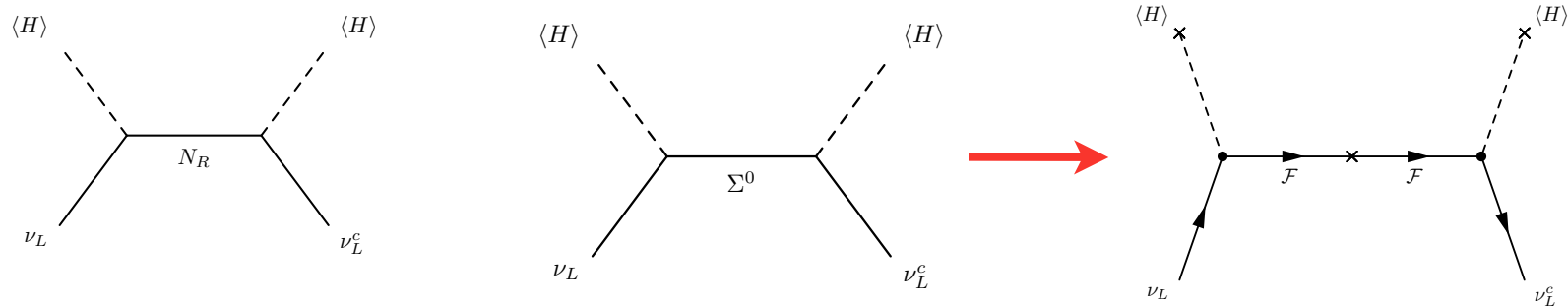
Generalized Seesaws

- Type-I and -III are similar:

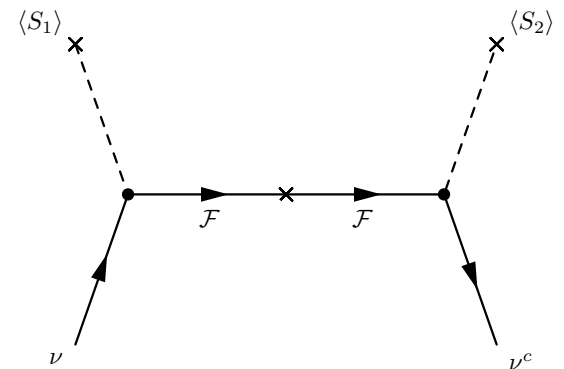


Generalized Seesaws

- Type-I and -III are similar:



- Generalize to generic diagram

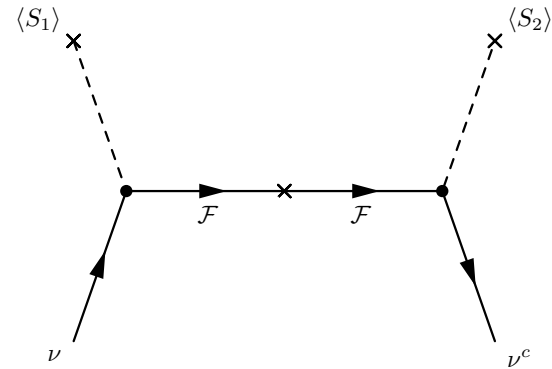


- Simple minded question: Are there more seesaws?

Apparently not answered

Generalized Seesaws

- Generic diagram



- Features:
 - Two external scalars
 - Heavy intermediate fermion
 - Mass insertion on internal line

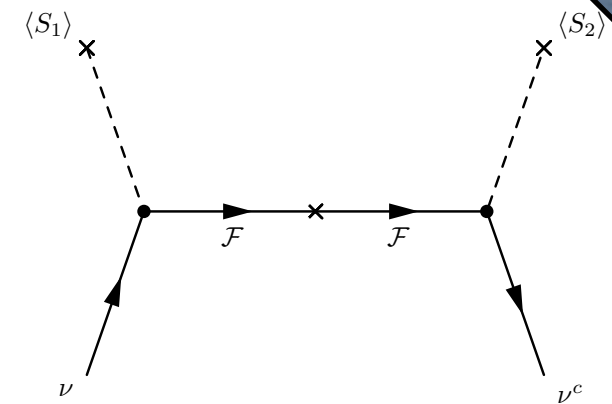


Majorana
Dirac

- Neutrino mass $\longrightarrow m_\nu \simeq \lambda_1 \lambda_2 \frac{\langle S_1 \rangle \langle S_2 \rangle}{M_{\mathcal{F}}}$

Rules of The Game

- Minimal particle extension
- Don't rely on new symmetries
- Small VEV's are naturally suppressed
- Determine all possibilities



VEV seesaw
(cf Type-II)

Results

Model	S_1	\mathcal{F}	S_2	Mass Insertion	$[\mathcal{O}_\nu]$	Ref.
(a)	(1, 2, 1)	(1, 1, 0)	–	Majorana	$d = 5$	Type I
(b)	(1, 2, 1)	(1, 3, 0)	–	Majorana	$d = 5$	Type III
(c)	(1, 4, –3)	(1, 3, 2)	(1, 2, 1)	Dirac	$d = 7$	Babu et al (2009)
(d)	(1, 4, 1)	(1, 5, 0)	–	Majorana	$d = 9$	Kumericki et al (2012)
(e)	(1, 3, 0)	(1, 4, –1)	(1, 5, –2)	Dirac	$d = 9$	KM (2013)
(f)	(1, 4, –3)	(1, 5, 2)	(1, 4, 1)	Dirac	$d = 9$	Picek et al (2011)

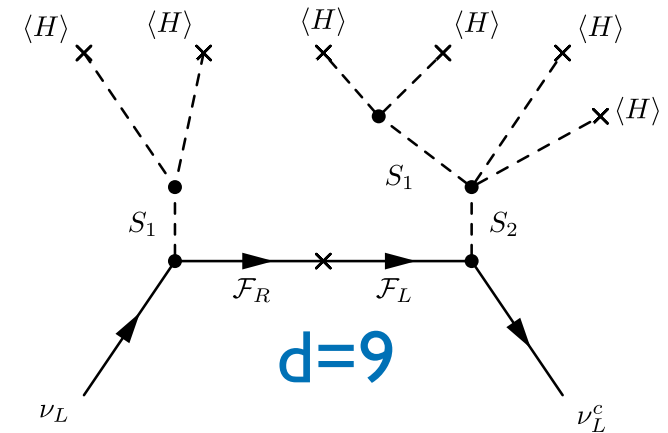
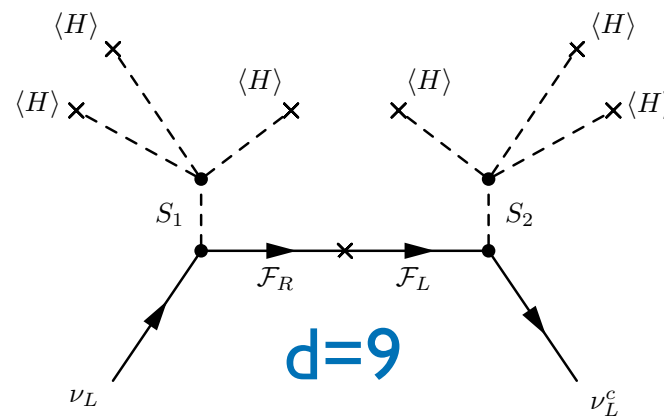
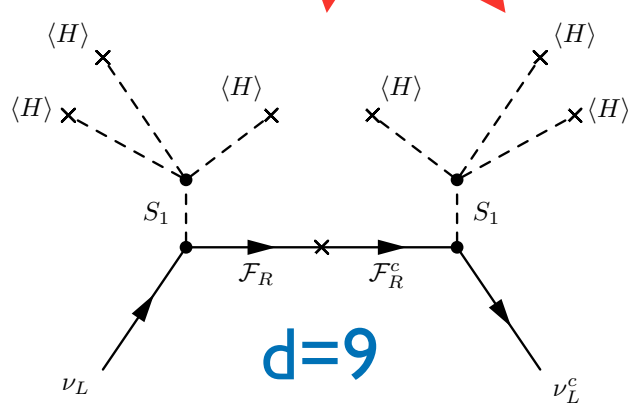
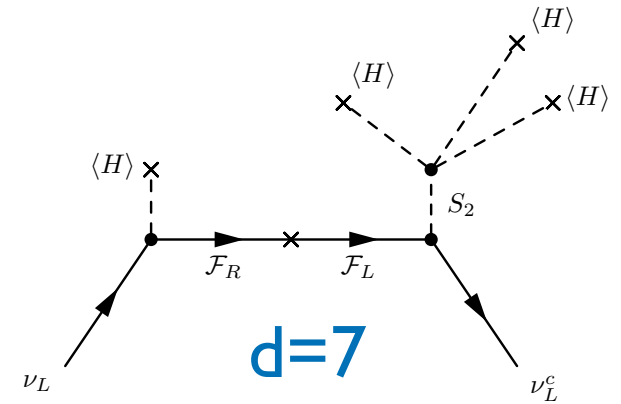
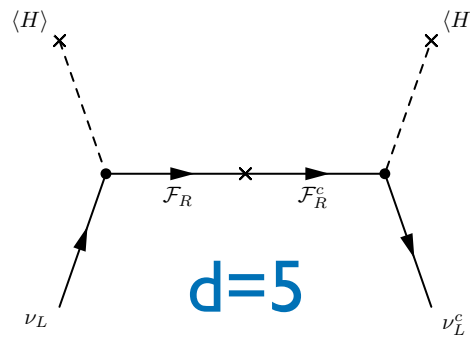
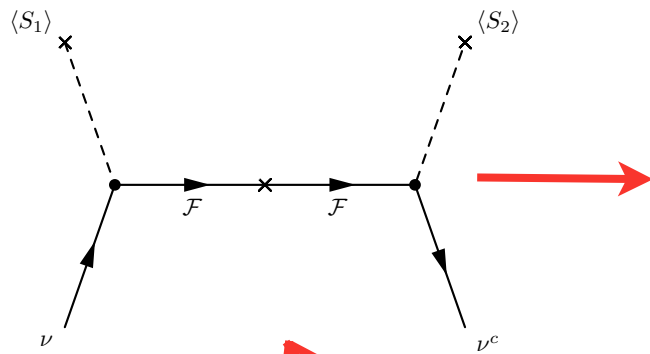
- Six with $d \leq 9$

- Two more with $d = 11$

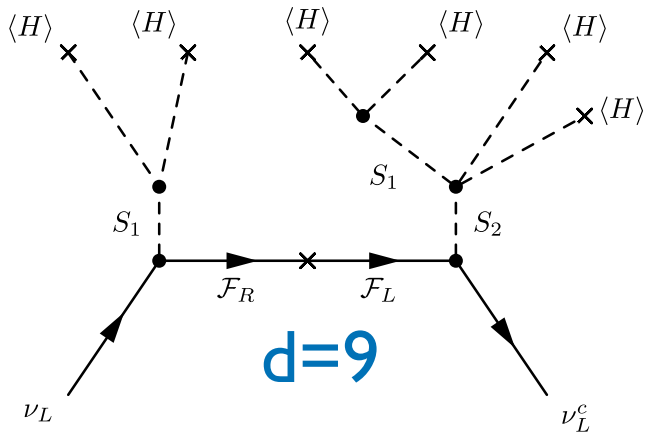


But that's all!

Tree-Level Seesaws with $d > 5$



Tree-Level Seesaws with $d > 5$



$$[\mathcal{O}_\nu] = d \quad \longrightarrow \quad \mathcal{O}_\nu = \frac{1}{M^{d-4}} L^2 H^{d-3}$$

$$\longrightarrow \quad m_\nu \sim \frac{\langle H \rangle^{d-3}}{M^{d-4}}$$

For example: $d = 5 \quad \longrightarrow \quad m_\nu \sim \frac{\langle H \rangle^2}{M} \quad \longrightarrow \quad M \lesssim 10^{14} \text{ GeV}$

or: $M \sim \text{TeV} \quad \longrightarrow \quad \lambda \sim 10^{-6}$

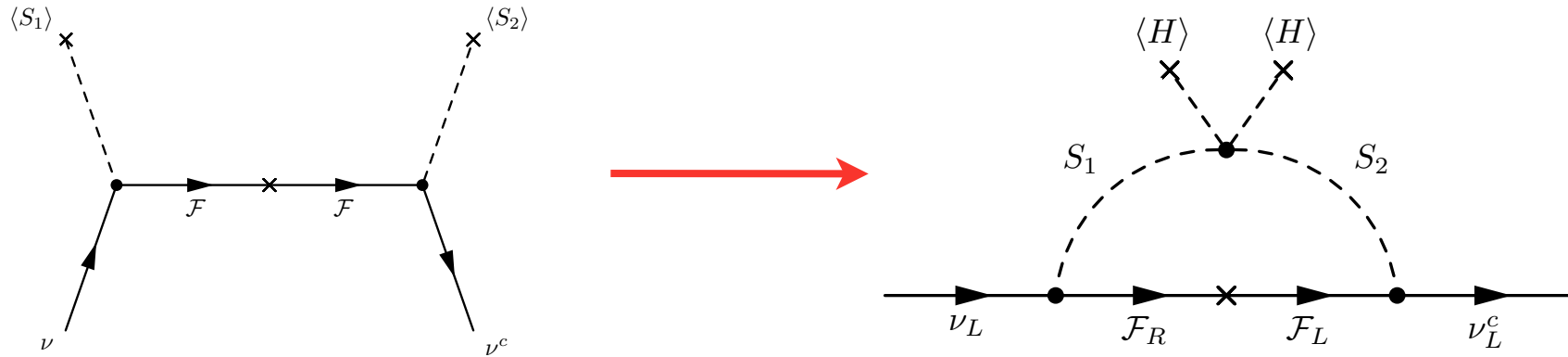
$d = 9 \quad \longrightarrow \quad m_\nu \sim \frac{\langle H \rangle^6}{M^5} \quad \longrightarrow \quad M \lesssim 10^7 \text{ GeV}$

or: $M \sim \text{TeV} \quad \longrightarrow \quad \lambda \sim 10^{-2}$

Larger d means: **Will be probed/ruled-out first**

Loop Masses

- All the new models have: $V(H, S_1, S_2) \supset \lambda S_1 S_2 H H$



- Neutrino mass: $M_\nu = M^{tree} + M^{loop}$

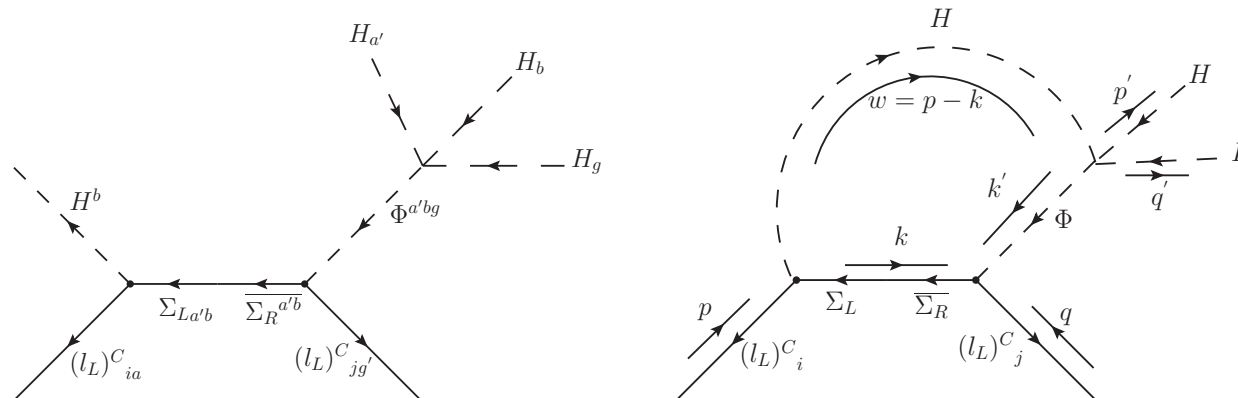
They are seesaw/radiative models

Seesaw/Radiative Models

- All have doubly charged fields
- Some have triply charged fields (!)
- Consider the d=7 model (Babu et al, 2009)

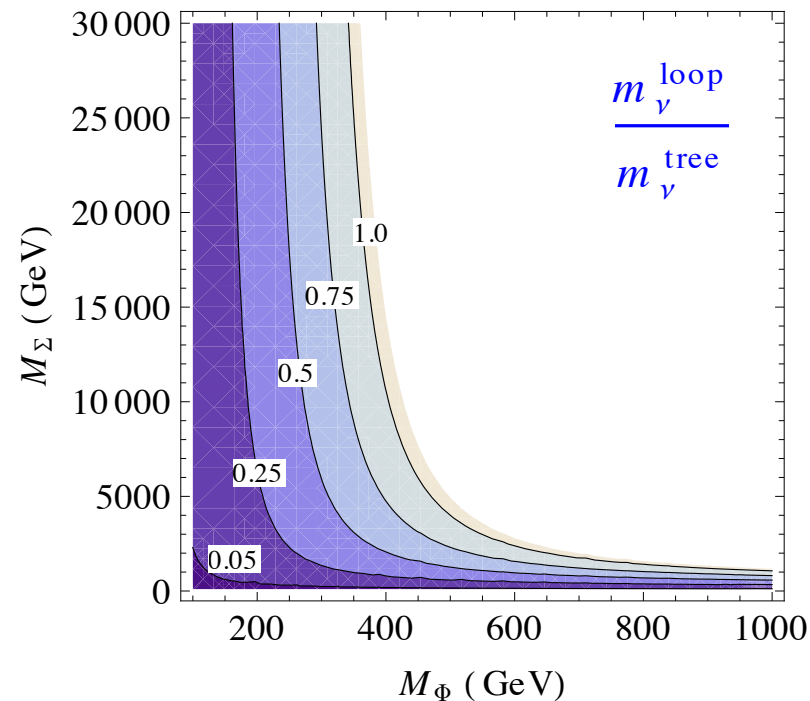
Beyond SM fields: $\Sigma \sim (1, 3, 2) = (\Sigma^{++}, \Sigma^+, \Sigma^0)$

$\Phi \sim (1, 4, 3) = (\Phi^{+++}, \Phi^{++}, \Phi^+, \Phi^0)$



Seesaw/Radiative Models

$$M_\nu = M^{tree} + M^{loop}$$

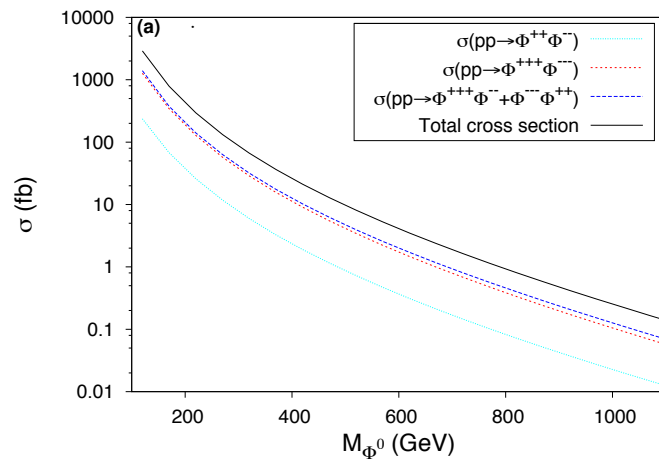


(Bambhaniya et al, 2013)

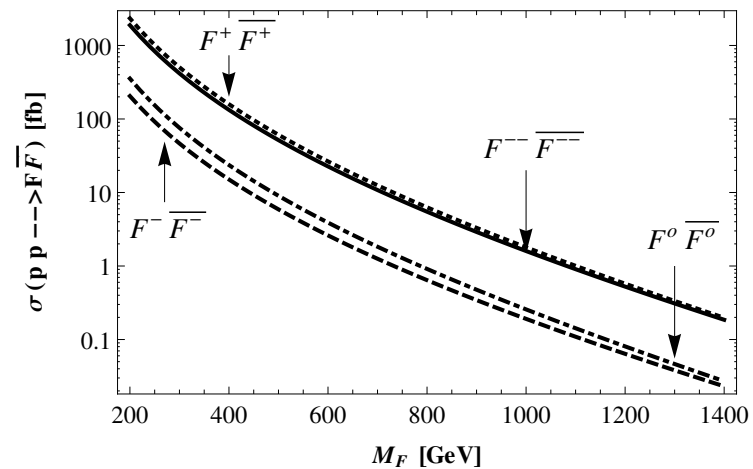
- Seesaw region “more accessible” at the LHC
- Similar results for all the new models

Seesaw/Radiative Models

- LHC production: controlled by E/W interactions



(Bambhaniya et al, 2013)



Fermions in new d=9 model

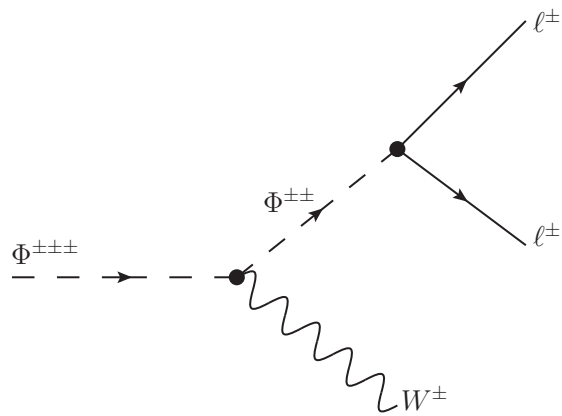
For $\sqrt{s} = 14$ TeV

- Similar results for all the models

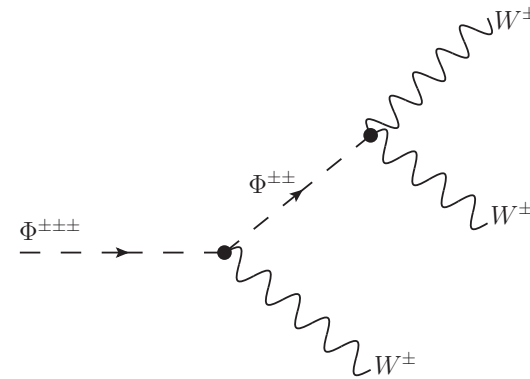
Triply Charged Scalar

- Decays

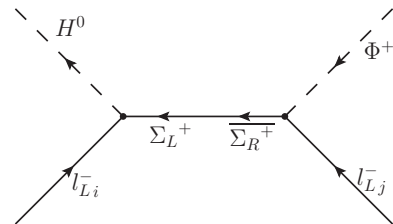
$$\Phi^{+++} \rightarrow W^+ \ell^+ \ell^+$$



$$\Phi^{+++} \rightarrow W^+ W^+ W^+$$

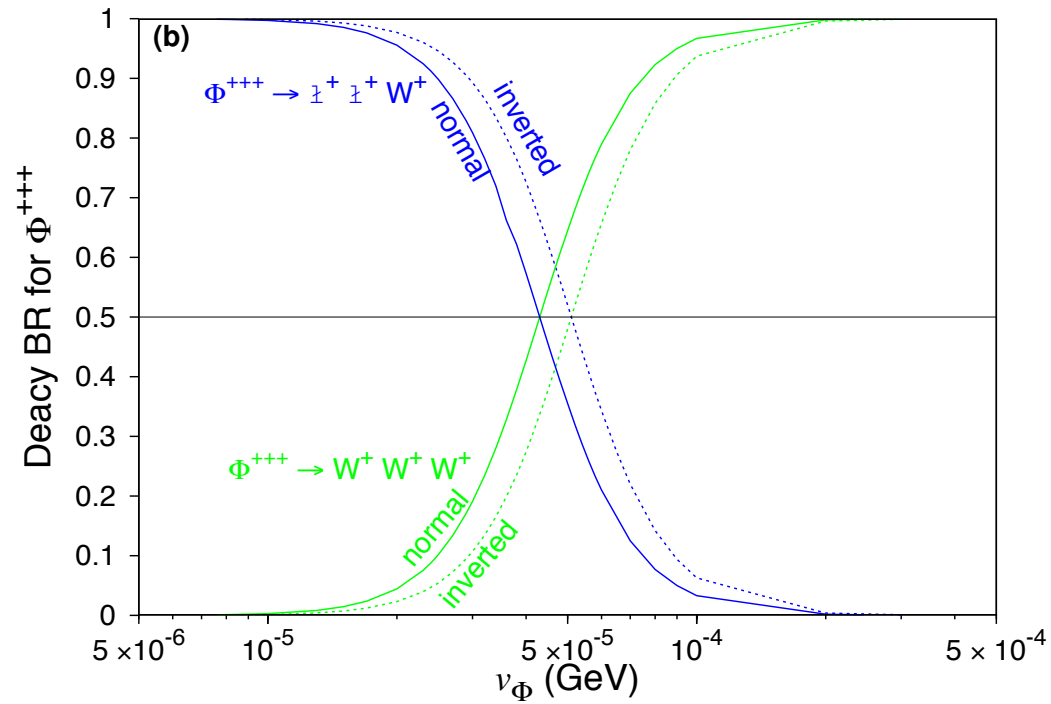


- Sensitive to neutrino mass



This is only for the d=7 model

Triply Charged Scalar



(Bambhaniya et al, 2013)

Detection

- Detection: multi-lepton final states (3,4,5,6)
- Including: same sign tri-lepton events
- Some model dependencies: e.g. Branching fractions etc
- General trend: “Some” signal has little SM background

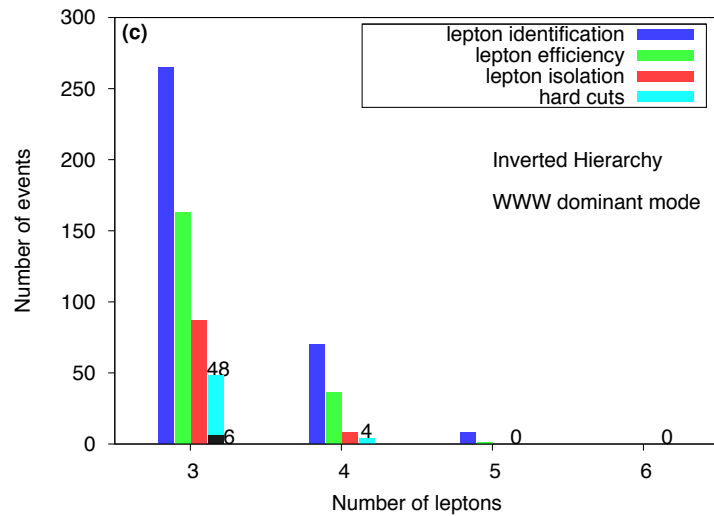
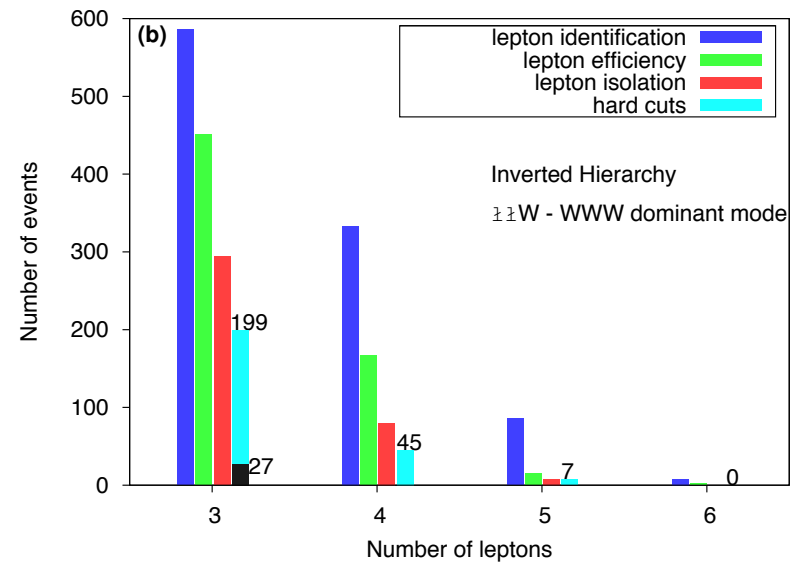
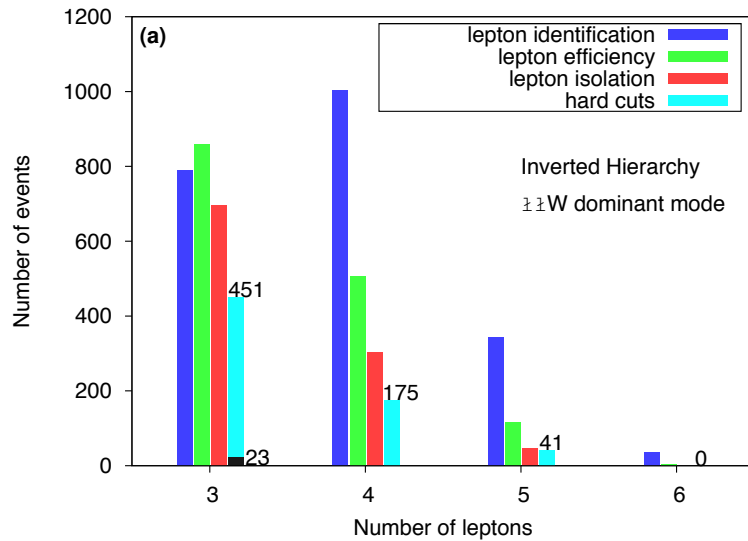
Background for the d=7 Model

- After cuts

\Downarrow processes \ multi-lepton channel \Rightarrow	3ℓ (fb)	$SS3\ell$ (fb)	4ℓ (fb)
$t\bar{t}$	18.245	–	–
$t\bar{t}(Z/\gamma^*)$	1.121	7.066×10^{-4}	0.069
$t\bar{t}W^\pm$	0.656	3.836×10^{-3}	–
$t\bar{t}t\bar{t}$	–	1.327×10^{-4}	–
$t\bar{t}b\bar{b}$	–	$< 10^{-4}$	–
$W^\pm(Z/\gamma^*)$	10.590	–	–
$(Z/\gamma^*)(Z/\gamma^*)$	1.287	–	0.047
TOTAL	31.899	4.675×10^{-3}	0.116

(Bambhaniya et al, 2013)

Signal Events




- Integrated luminosity: 100 fb^{-1}
- Note: SS3L --- tiny background
- Significance > 10
- Similar for normal hierarchy

Mass Hierarchy Dependence

4ℓ	$eeee$	$ee\mu\mu$	$e\mu\mu\mu$	$e\mu\mu\mu$	$\mu\mu\mu\mu$	Total events
IH	14	47	69	29	16	175
NH	1	1	23	40	61	126

- But not all the models have this feature

Conclusion

- Small mass: Seesaw or radiative
- New models: seesaw/radiative
- Seesaw: $d > 5$  lighter new physics
- Probe at LHC: Multi-lepton signals