



# Missing Transverse Momentum at LHC

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Cairns  
Colonial Club Resort

# Outline

- Missing transverse momentum ( $E_T^{\text{miss}}$ ) at LHC:
  - Motivation for measuring  $E_T^{\text{miss}}$
  - How it is defined and reconstructed
- $E_T^{\text{miss}}$  challenge with LHC data:
  - Events cleaning
  - Pile-up
- Strategy for  $E_T^{\text{miss}}$  reconstruction and calibration:
  - Suppressing the pile-up effects
- How to study of  $E_T^{\text{miss}}$  performance:
  - Resolution
  - Scale
  - Tails
- $E_T^{\text{miss}}$  systematic uncertainties
- Conclusions

# Missing transverse momentum definition

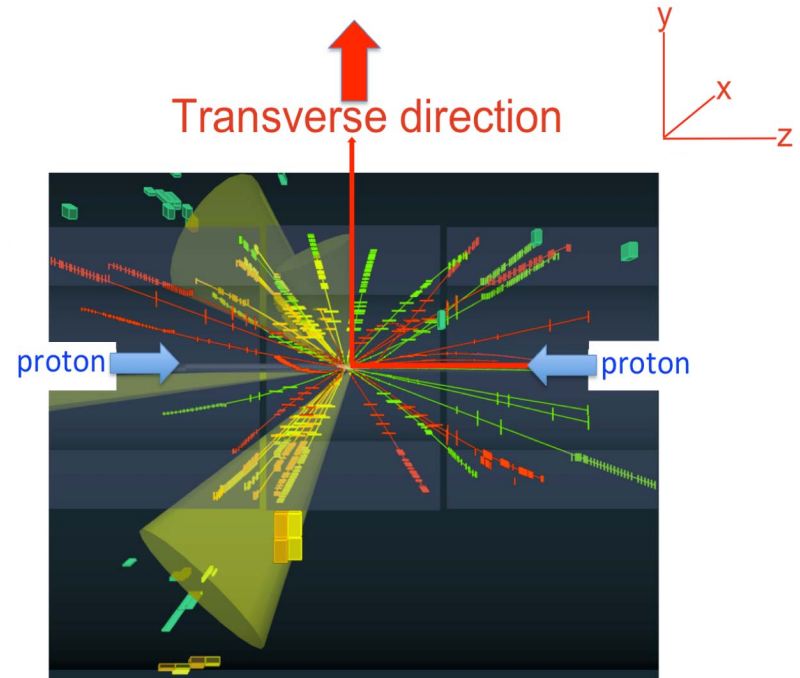
In a hadron collider event the **missing transverse momentum** ( $E_T^{\text{miss}}$ ) is defined as the event momentum imbalance in the plane transverse to the beam axis, where momentum conservation is expected. Such an imbalance may signal the presence of non detected particles, such as neutrinos and or new weakly-interacting particles

The two  $E_T^{\text{miss}}$  components,  $E_x^{\text{miss}}$  and  $E_y^{\text{miss}}$ , are calculated as the opposite sum of the momenta of all particles seen in the detector

$$\mathbf{E}_{x,y}^{\text{miss}} = - \sum_{\text{particles}} \mathbf{p}_{x,y}$$

$$E_T^{\text{miss}} = \sqrt{(E_x^{\text{miss}})^2 + (E_y^{\text{miss}})^2}$$

$$\sum \mathbf{E}_T = \sum_{\text{particles}} \mathbf{p}_T$$



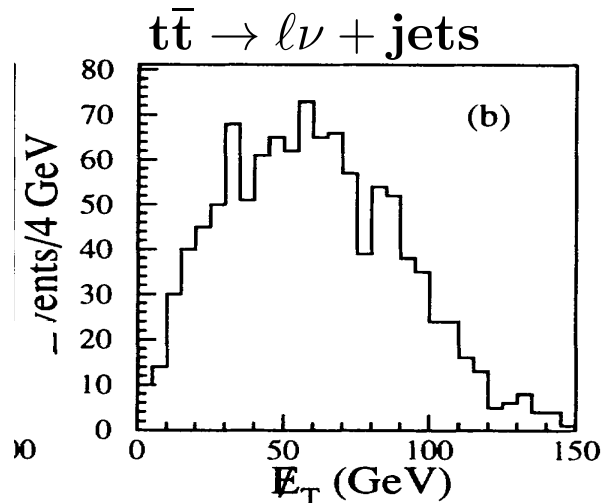
# Motivation for measuring $E_T^{\text{miss}}$

$E_T^{\text{miss}}$  is due to non-interacting particles  
in detector ( $\nu$ , LSP)  $\Rightarrow$  **True  $E_T^{\text{miss}}$**

$\rightarrow$  A very good measurement of  $E_T^{\text{miss}}$ , i.e. of  $p_T \nu$ , is a crucial requirement for the **study of many physics measurements**

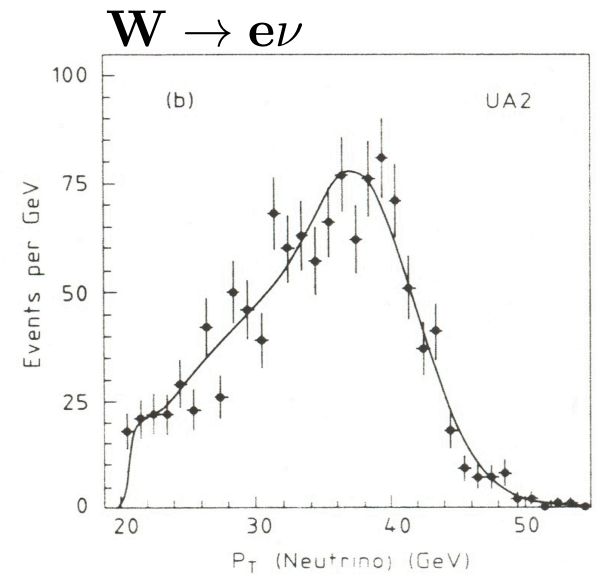
- $W \rightarrow l \nu$ , semi-leptonic top decays,  $Z \rightarrow \tau \tau$

Tevatron  $p\bar{p}$  collider 1.8 TeV  
CDF Collaboration  
Top quark discovery



Phys. Rev. D 50 (1994)

CERN  $p\bar{p}$  collider 630 GeV  
UA2 Collaboration  
Measurement of W mass



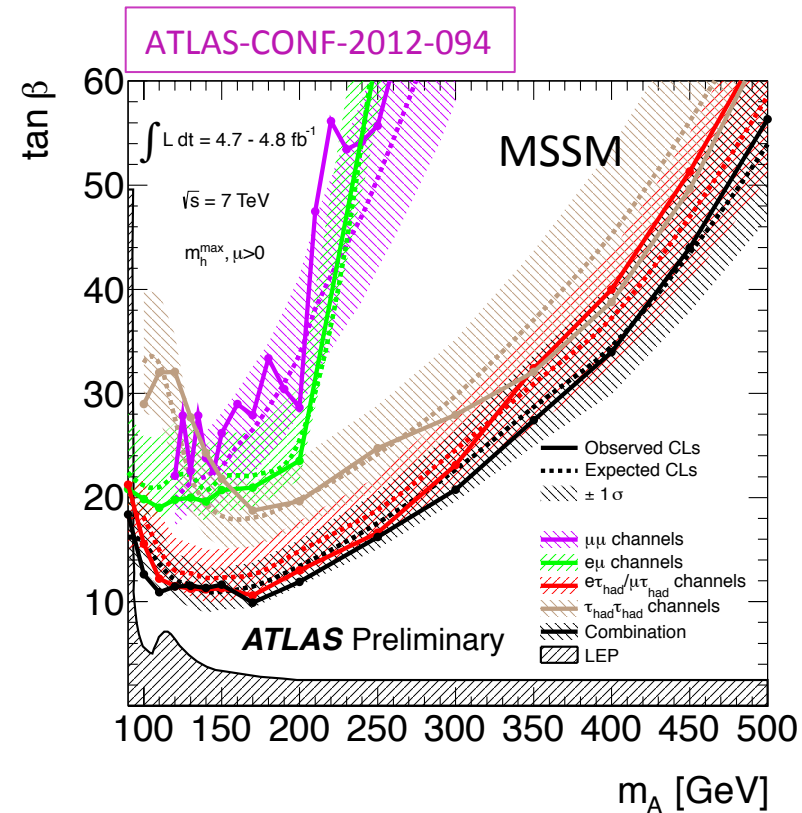
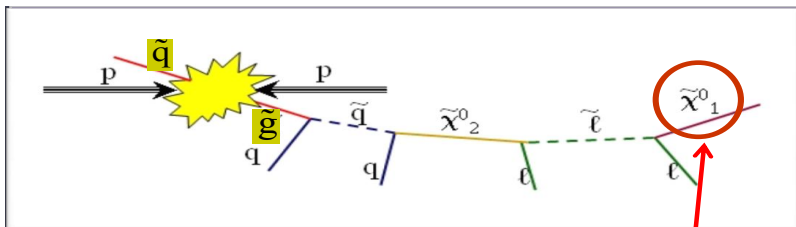
Phys. Lett. B 241 (1990)

# Motivation for measuring $E_T^{\text{miss}}$

$E_T^{\text{miss}}$  is due to non-interacting particles in detector ( $\nu$ , LSP)  $\Rightarrow$  **True  $E_T^{\text{miss}}$**

- $\rightarrow E_T^{\text{miss}}$  plays a major role for the physics at LHC
- $\rightarrow$  A very good  $E_T^{\text{miss}}$  measurement, i.e. of  $p_T \nu$  or of  $p_T(\text{lsp})$  is a crucial requirement for the study of many physics measurements **and for discovery physics**

- $W \rightarrow l\nu$ , top decays,  $Z \rightarrow \tau\tau$
- SM Higgs ( $H \rightarrow \tau\tau$ ,  $H \rightarrow WW \rightarrow l\nu l\nu / l\nu + \text{jets}$ )
- MSSM Higgs ( $A/H \rightarrow \tau\tau$ ,  $H^\pm \rightarrow \tau\nu$ )
- Higgs to invisible decays
- SUSY  $\rightarrow$  Large  $E_T^{\text{miss}}$  signature from lsp



In  $H \rightarrow \tau\tau$  events can reconstruct the invariant  $\tau\tau$  mass from the two  $E_T^{\text{miss}}$  components

# $E_T^{\text{miss}}$ reconstruction and calibration

$E_T^{\text{miss}}$  is a complex event quantity:

- It is calculated adding all significant signals from all detectors
  - Calorimeter input signals (from charged and neutral particles)
    - used to reconstruct high  $p_T$  physics objects (e,  $\gamma$ ,  $\tau$ , jets)
    - not used in high  $p_T$  physics objects
  - Muons
  - Reconstructed tracks (from charged particles)
- Avoid double counting
- Coherent Calibration

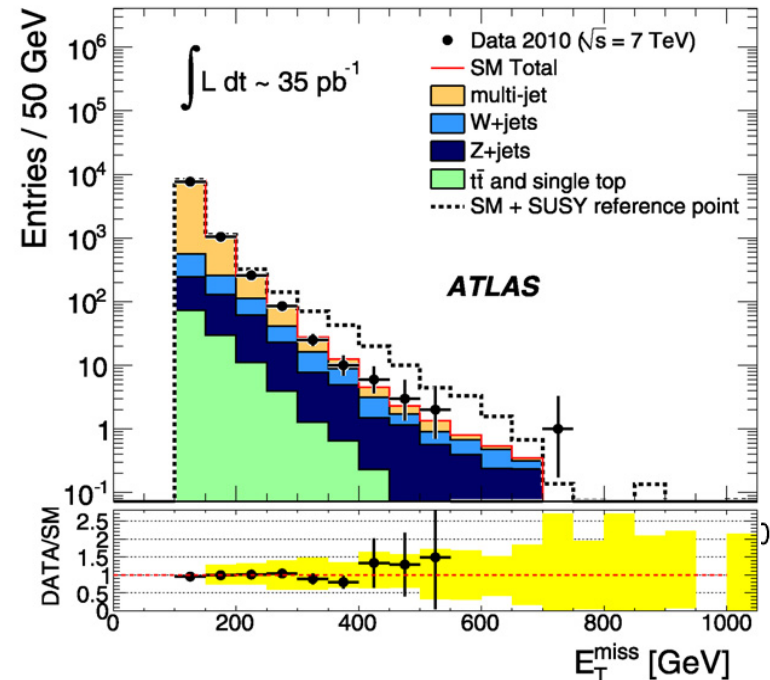
# Fake $E_T^{\text{miss}}$

$E_T^{\text{miss}}$  is due to non-interacting particles, BUT  $E_T^{\text{miss}}$  can also due to:

- Problems in detector:
    - dead, hot, noisy channels
  - Backgrounds:
    - cosmic rays, beam-halo, beam-gas
  - Cracks/gaps in the detector, azimuthal response variations
  - Energy lost in dead materials (cracks, cryostats..)
  - Noise, pile-up noise
  - Mis-measurements of muons, jets
- ⇒ “Fake”  $E_T^{\text{miss}}$

- ⇒ First require detailed understanding of instrumental  $E_T^{\text{miss}}$  sources
  - Event Cleaning
- ⇒ Then understand other source of “fake”  $E_T^{\text{miss}}$
- ⇒ Suppress pile-up at LHC !

Phys. Lett. B 701 (2011) 186



- QCD with “fake”  $E_T^{\text{miss}}$  are background for inclusive no-lepton SUSY events
- Can fake “new physics”
- understanding this background is crucial for SUSY searches !

# Data-quality requirements and Event cleaning

## Data-quality (detector level)

- Stable proton beams, nominal magnetic field conditions.
- NO detector problems: use only data with a fully functioning calorimeter, inner detector and muon spectrometer

## Cleaning (event level)

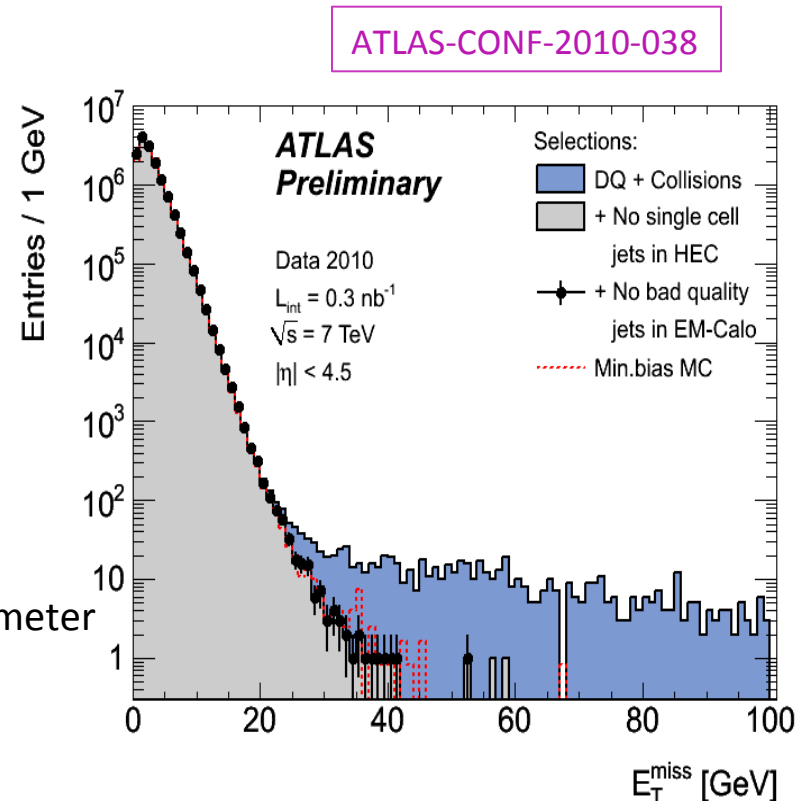
### Discard events with bad jets:

#### → Jets due to non-collision background

- Beam-gas events
- Beam-halo events
- Cosmic ray muons overlapping in-time with collision events

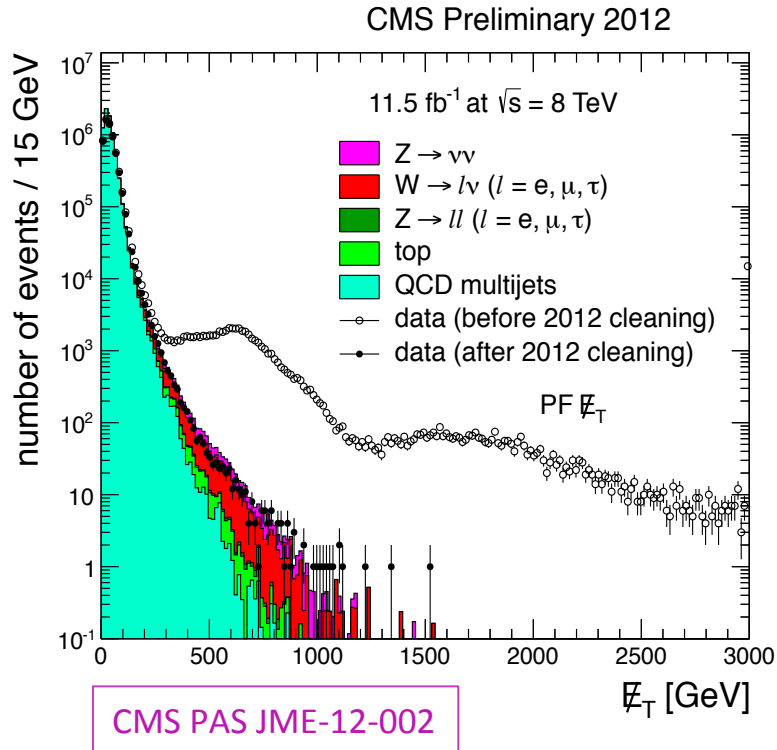
#### → Jets from calorimeter noise

- **Sporadic noise bursts** in the hadronic endcap calorimeter
  - few noisy calorimeter cells contribute to almost all of the jet energy.
- **Coherent noise** in the electromagnetic calorimeter.





# Event cleaning

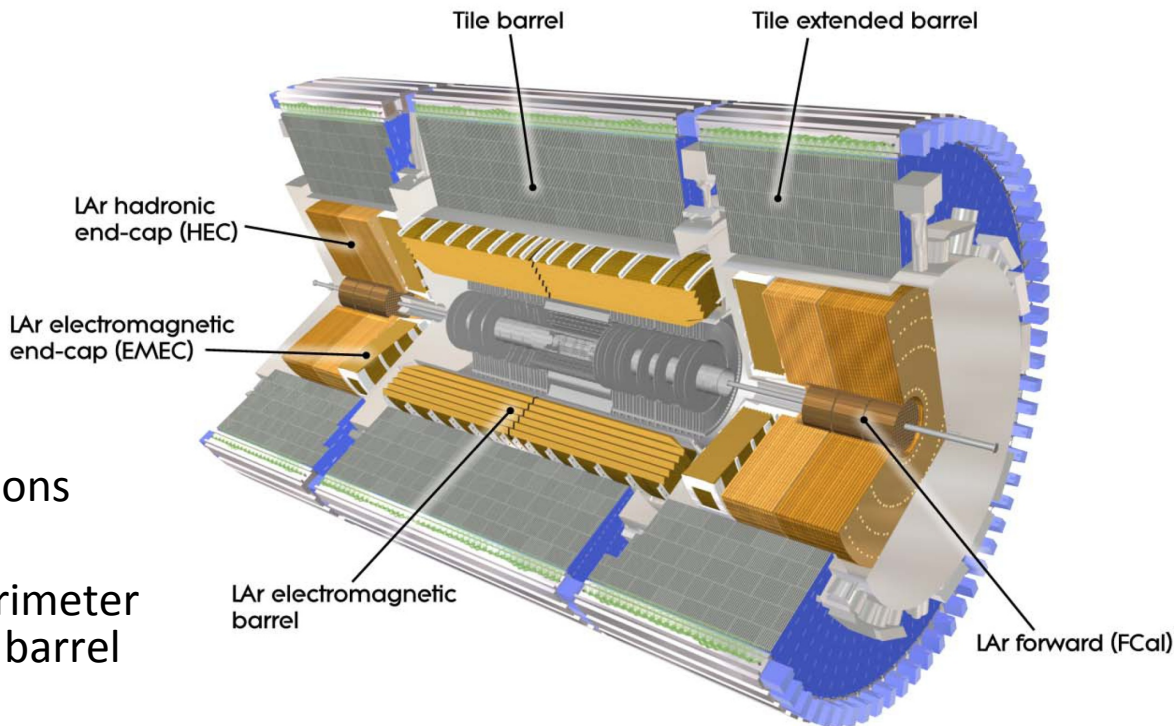


- Anomalous high  $E_T^{\text{miss}}$  events in data before 2012 cleaning mainly come from:
  - Misfires of the HCAL laser calibration system
  - Electronic noise in HCAL
  - Fake  $E_T^{\text{miss}}$  from track reconstruction
- Few remaining anomalous events are removed by applying jet identification cut
  - neutral hadron energy fraction of the jet  $< 0.90$
  - photon energy fraction  $< 0.95$

# ATLAS calorimeters

Main features for  $E_T^{\text{Miss}}$  reconstruction and calibration:

- Noise suppression
- Non-compensating ( $e/h > 1$ ):
  - Response to hadrons is lower than that to electrons and photons
  - Developed specific calibrations
- Dead material:
  - Energy loss before EM calorimeter and between EM and HAD barrel calorimeters:
    - dead material corrections
- Different technologies and many transition regions:
  - “Crack” regions:  $\eta \approx 1.4, 3.2$
- Magnetic field bending



## ATLAS Fiducial Regions

- Hadronic Calorimeter:
  - Barrel:  $|\eta| < 1.7$
  - Endcap:  $1.5 < |\eta| < 3.2$
- Electromagnetic Calorimeters
  - Barrel:  $|\eta| < 1.4$
  - Endcap:  $1.375 < |\eta| < 3.2$
- Forward:  $3.2 < |\eta| < 4.9$

$$\eta = -\log(\tan(\theta/2))$$

# Calorimeter input signals to $E_T^{\text{miss}}$

- **Hard signal in calorimeters**

Fully reconstructed & calibrated particles and jets

→ detector inefficiencies already corrected for physics objects

→ have to avoid mis-measured objects

- **Soft signals in calorimeters**

Signals not used in reconstructed physics objects

→ need to be included in  $E_T^{\text{miss}}$  to reduce scale biases and improve resolution

→ need to suppress noise (electronic and pile-up noise)

- Applying symmetric or asymmetric noise cuts to cell signals can introduce a bias

- Topological clustering applies more reasonable noise cut

→ need to be calibrated

- low-pT particles can easily be lost due to magnetic field or because their energy does not survive noise cuts → use tracks to correct for calorimeter inefficiency

- **Need to avoid double counting**

Same signal can only be used in one physics object

→ Veto  $E_T^{\text{miss}}$  contribution from already used signals

**UA2:** The neutrino transverse momentum was estimated from the transverse component of the momentum balance of the electron and of the calorimeter calibrated cells.  
**To avoid double counting the cells in the electron core are not used.**

# Fake $E_T^{\text{miss}}$ from jets mis-measurement

Fake  $E_T^{\text{miss}}$  can be created by mis-measurement of any objects: electrons, photons, taus, muons and hadronic jets.

In particular jet mis-measurement can be a dangerous source of fake  $E_T^{\text{miss}}$ ,  
→ suppression strategies are needed

- **Mis-measured jets in cracks** → event topology analysis
- **Jet leakage from the calorimeters or fluctuations** in large jet energy deposits in non-instrumented regions → check energy sharing between calorimeters
- **Jets mis-calibration** → compare with track jets

→ Generates  $E_T^{\text{miss}}$  pointing to this jet:

- study angular correlation between  $E_T^{\text{miss}}$  and jets

→ Careful analysis of full event topology

# Calorimeter noise-suppressed input signals to $E_T^{\text{miss}}$

**Topoclusters:** group of calorimeter cells topologically connected optimized for electronic noise and pile-up suppression

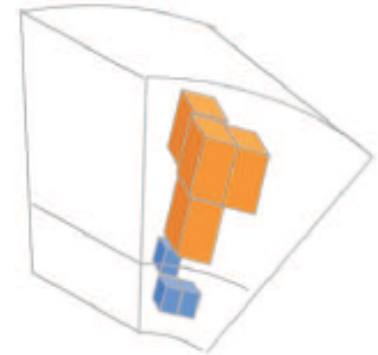
• Cluster cells in 3D via noise-driven thresholds:

- Seed:  $|E_{\text{cell}}| > 4 \sigma_{\text{noise}}$
- Neighbours:  $|E_{\text{cell}}| > 2 \sigma_{\text{noise}}$
- Perimeter cells  $|E_{\text{cell}}| > 0$
- $\sigma_{\text{noise}} = \sqrt{(\sigma_{\text{noise}}^{\text{electronic}})^2 + (\sigma_{\text{noise}}^{\text{pile-up}})^2}$

$\Phi$		0	0	0			
	0	0	2	0	0		
	0	2	2	2	0	0	
	0	2	4	2	2	0	
	0	2	2	2	0	0	
	0	0	0	0	0		
							$\eta$

**Topoclusters calibration (Local Hadron Weighting- LCW)**

- Classifications as “em-like” or “hadron-like” clusters based on cluster shape variables: energy density and depth.
  - Hadronic weights, derived from pion MC simulation, applied to “hadron-like” clusters.
  - Corrections for dead material and out of cluster
- No bias. Cells with very small signals can survive based on the signals in neighboring cells
- Improve correspondence between clusters and stable particles
- Intrinsically noise and pile-up suppressed, but contribution from pile-up fluctuations can survive, more pile-up suppression techniques needed



# Input Muon signals to $E_T^{\text{miss}}$

- **Hard signal in muon spectrometer**

Fully reconstructed & calibrated muons

→ Any muons which are not reconstructed, badly measured or fake can be a source of fake  $E_T^{\text{miss}}$ !

- includes all muons reconstructed in muon spectrometer
  - use tracks in the region of inefficiency of muon spectrometer
- Choose best measurement
- Apply quality criteria to avoid bad-measured muons
- Have to avoid fake muons:
  - Fake muons from jet punch-through

→ Muons may generate isolated or embedded soft calorimeter signals

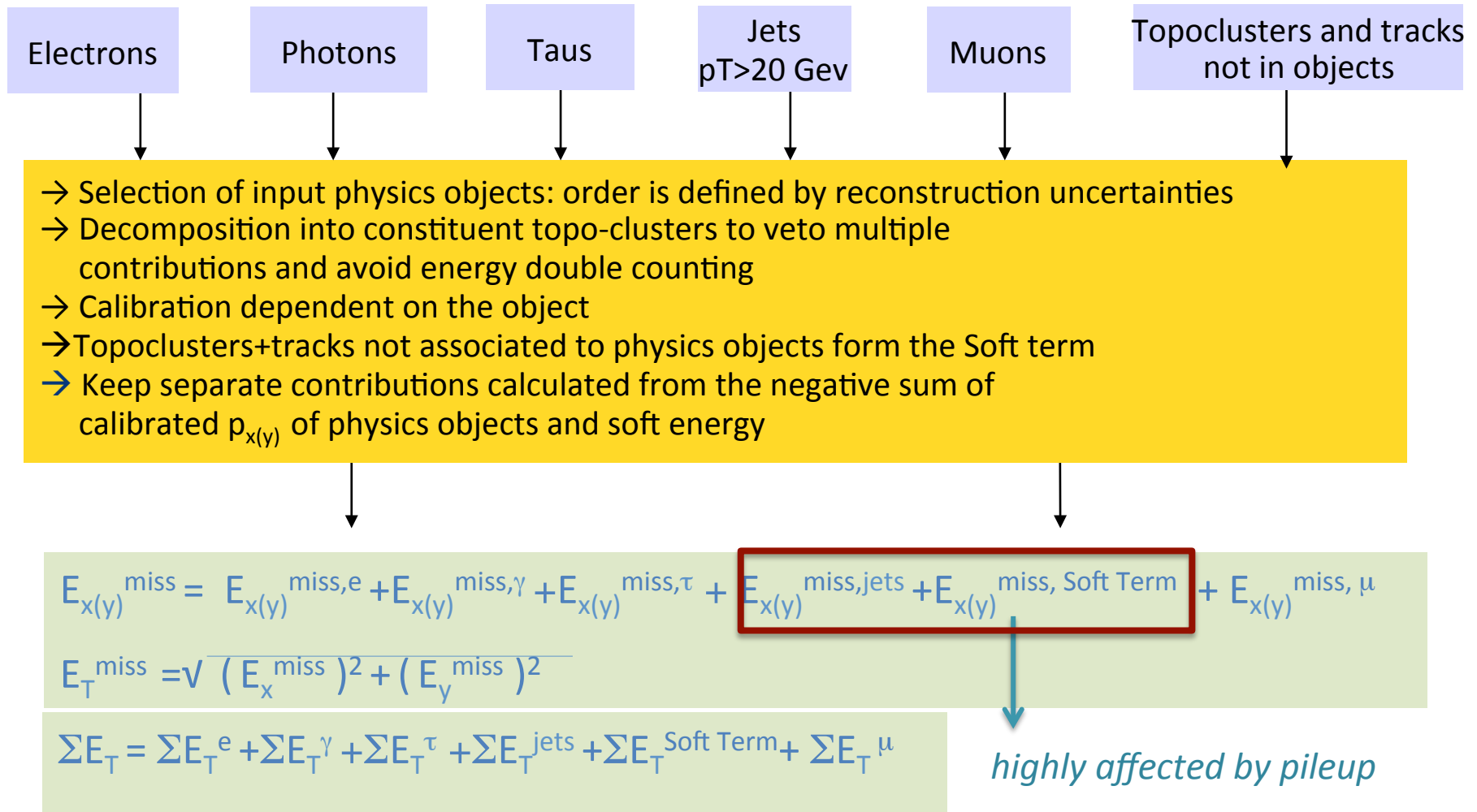
- Care needed to avoid double counting

subtract muon energy deposited in calorimeters when the combined muon momentum (from muon spectrometer and inner detector) is used)

**CDF:**  $E_T^{\text{miss}}$  is defined to be the negative of the vector sum of transverse energy in calorimeter towers.

For events with muon candidates, the vector sum of the calorimeter transverse energy is corrected by vectorially subtracting the energy deposited by the muon and then adding the PT of the muon candidate as measured in the CTC.

# $E_T^{\text{miss}}$ reconstruction and calibration in ATLAS



Very *flexible algorithm*: different definitions and calibrations for physics objects are allowed  
 Coherence with Physics analysis

# The ATLAS Soft Term algorithm

## (1) Track selection

All reconstructed tracks

Apply quality criteria

Veto on tracks associated to high physics objects

Veto on tracks associated to TopoClusters already used

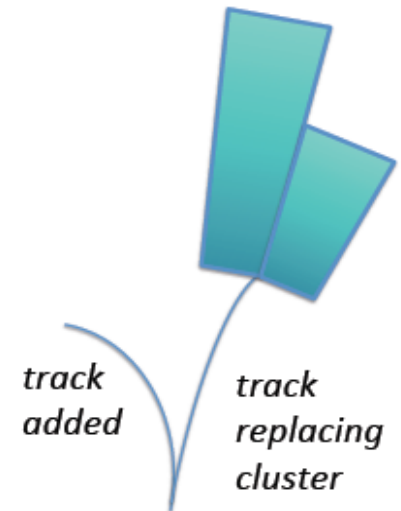
Add good tracks to  $E_T^{\text{miss}}$  calculation

## (2) Cluster removal

All TopoClusters not associated to physics objects

Veto on TopoClusters associated to good tracks

Add remaining TopoClusters to  $E_T^{\text{miss}}$  calculation



- Improve calculation of the low contribution to Soft Term
- Tracks are added to recover the contribution from low- $p_T$  particles which do not reach the calorimeter or do not seed a TopoCluster.
- No association with PV => no pile-up suppression at this level

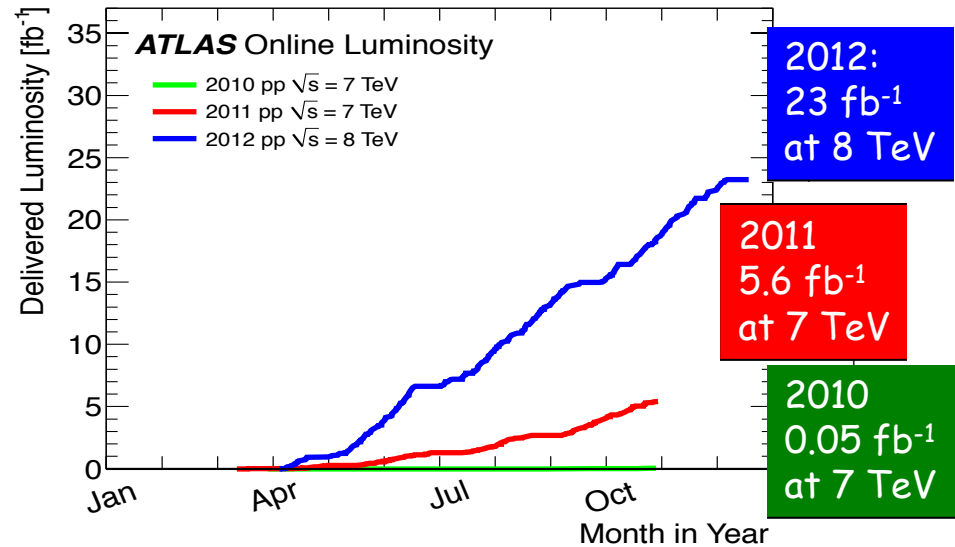


# The pile-up

→ The LHC luminosity increased from 2010 to 2012

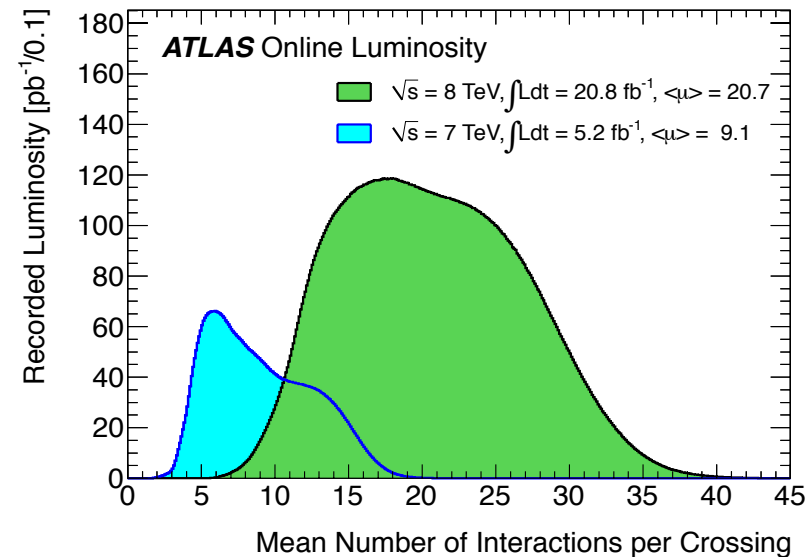
**LHC**  
Up to  
 $L_{\text{peak}} 7.7 \cdot 10^{33} \text{ cm}^{-2} \cdot \text{s}^{-1}$   
in 2012 at 8 TeV

**ATLAS**  
~90% of delivered  
luminosity used in  
physics analyses



→ The Pile-up, i.e. the contribution of additional pp collisions superimposed to the hard physics process, is the price to pay for this!

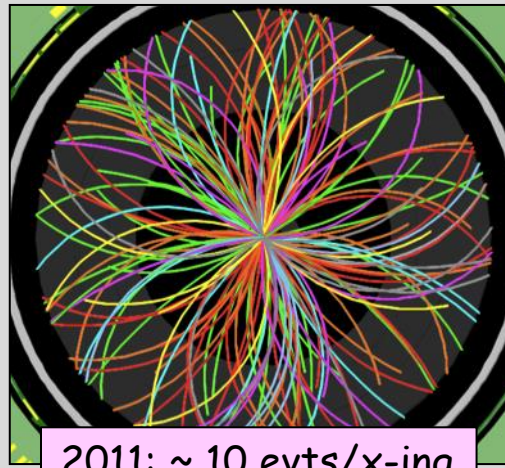
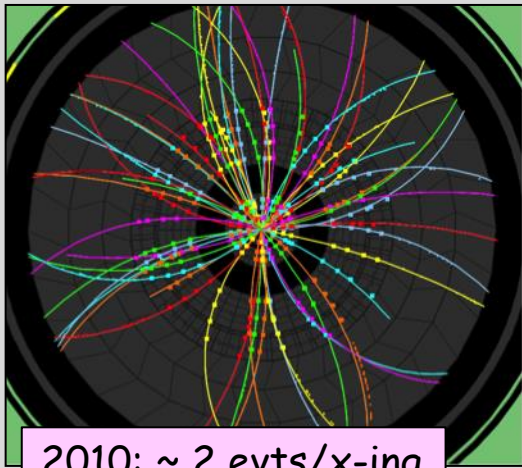
$$\langle \mu \rangle = L \times \sigma_{\text{inel}} / N_{\text{bunch}} \times f_{\text{LHC}}$$



# The pile-up

Pile-up is one of the main challenges for  $E_T^{\text{miss}}$  at LHC

- $E_T^{\text{miss}}$  has the largest acceptance (coverage area) of any given reconstructed quantity.
- Considerable contribution to  $E_T^{\text{miss}}$  (and jets) fluctuations from pile-up



# Pile-up suppression in jets in ATLAS

Jets are corrected for pile-up using Jet Area

$$p_T^{\text{corr}} = p_T - \rho A$$

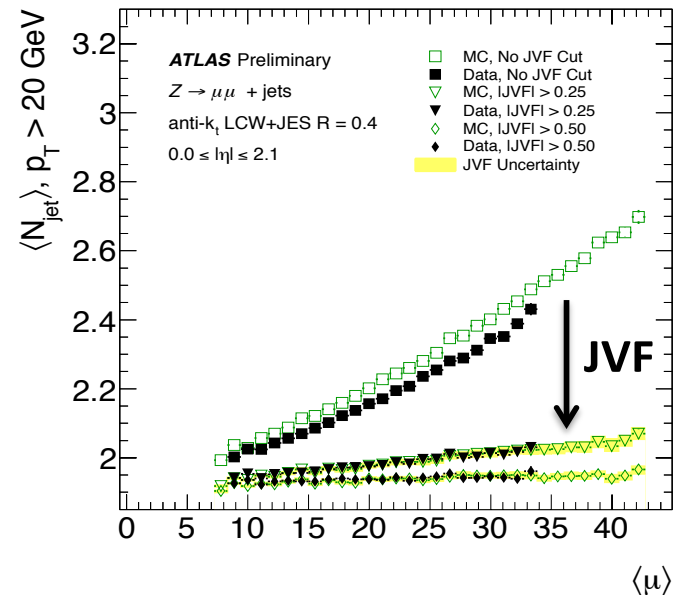
**Jet Area** based on the idea that noise (pile-up) has a lower  $p_T$  density ( $\rho$ ) than signal:

- **event  $p_T$  density  $\rho$**  is calculated from all jets ( $k_t$  jets) as median ( $p_T^{\text{jet}}/A^{\text{jet}}$ )
- Each jet is then corrected subtracting  $\rho A$  where **A is the jet area**
- Jet-by-jet subtraction => improves jet resolution
- Captures event-by-event fluctuations
- Data driven method: no dependence on pile-up modelling

Some pile-up jets remain after subtraction  
→ Further pile-up suppression using tracks associated with the primary vertex (**Jet Vertex Fraction JVF**)

$$\text{JVF} = \frac{\sum_{\text{tracks}_{\text{jet}, \text{PV}}} p_T}{\sum_{\text{tracks}_{\text{jet}}} p_T}$$

If  $p_T^{\text{jet}} < 50$  GeV and  $|\eta| < 2.4$ , keep jet only if  $|\text{JVF}| > 0$



# Pile-up suppression in soft term in ATLAS

Soft Term is very similar to pile-up, so any correction should be based on PV association or on exploiting the small difference between signal and pile-up

**STVF** is a correction based on

*fraction of tracks from PV:*

$$\text{STVF} = \frac{\sum_{\text{tracks}_{\text{SoftTerm}, \text{PV}}} p_T}{\sum_{\text{tracks}_{\text{SoftTerm}}} p_T}$$

use *tracks not matched to high- $p_T$  physics objects*

PV is the the first primary vertex

(vertex with  $\max \sum p_{T, \text{trk}}^2$ )

SoftTerm scaled by the “soft term vertex fraction” **STVF**

- **Limitations:** calculated in limited coverage (ATLAS ID  $|\eta| < 2.5$ ) and does not take into account neutral contributions

- **Jet Area** is based on the idea that noise (pileup) has a lower density ( $\rho$ ) than signal
- Similar to pile-up subtraction in jets
- Here “jet” means: jet  $0 < p_T < 20 \text{ GeV}$

For each event compute the  $p_T$  density  $\rho_{\text{ref}}$  (estimate of event-by-event pile-up activity)  
→ **reclusterize jets** from topoclusters and tracks from soft term with  $k_t$  algorithm

$$E_{T, \text{jet}}^{\text{corr}} = \begin{cases} 0 & E_{T, \text{jet}} \leq N \rho_{\text{ref}} A_{\text{jet}} \\ E_{T, \text{jet}} - \rho_{\text{ref}} A_{\text{jet}} & E_{T, \text{jet}} > N \rho_{\text{ref}} A_{\text{jet}} \end{cases}$$

+ Filter jets asking for  $|\text{JVF}| > 0.25$

- captures event-by-event fluctuations
- jet-by-jet correction

# $p_T^{\text{miss}}$ reconstructed from tracks in ATLAS

Sum of  $p_T$  of all **tracks from the PV** that pass the following standard criteria:

- $p_T > 500$  MeV
  - at least 1 pixel hit and 6 SCT hits
  - $|\eta| < 2.5$
  - $|d_0| < 1.5$  mm
  - $|z_0 \sin\theta| < 1.5$  mm
- PV association*

+

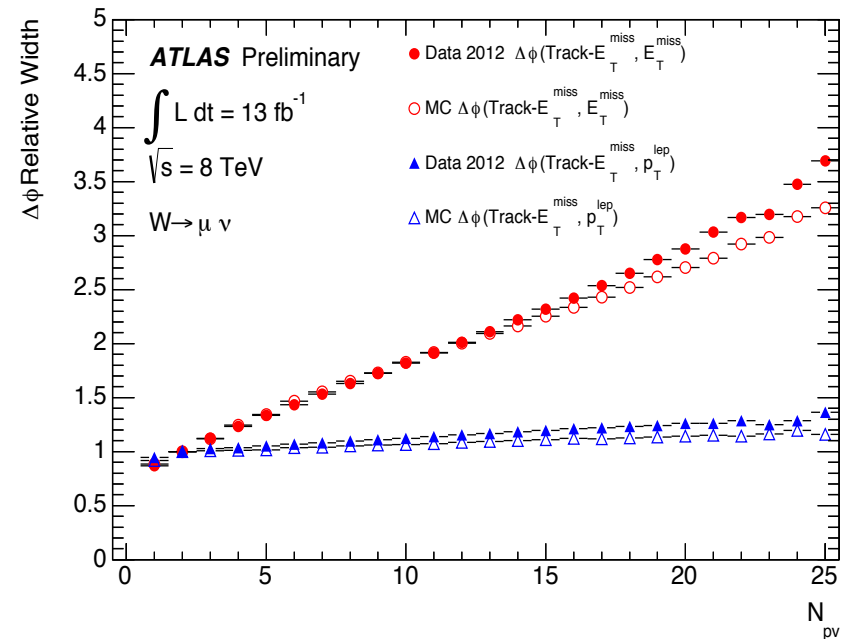
**Tracks of any muons** passing:

- Staco Combined
- $p_T > 6$  GeV
- $|\eta| < 2.5$

**Tracks of any electrons** passing:

- Medium++
- $p_T > 10$  GeV
- $|\eta| < 2.47$

- Good **stability vs pileup**
- **Limited acceptance** ( $|\eta| < 2.5$ , no neutrals)  $\rightarrow$  performance degradation in events with jets



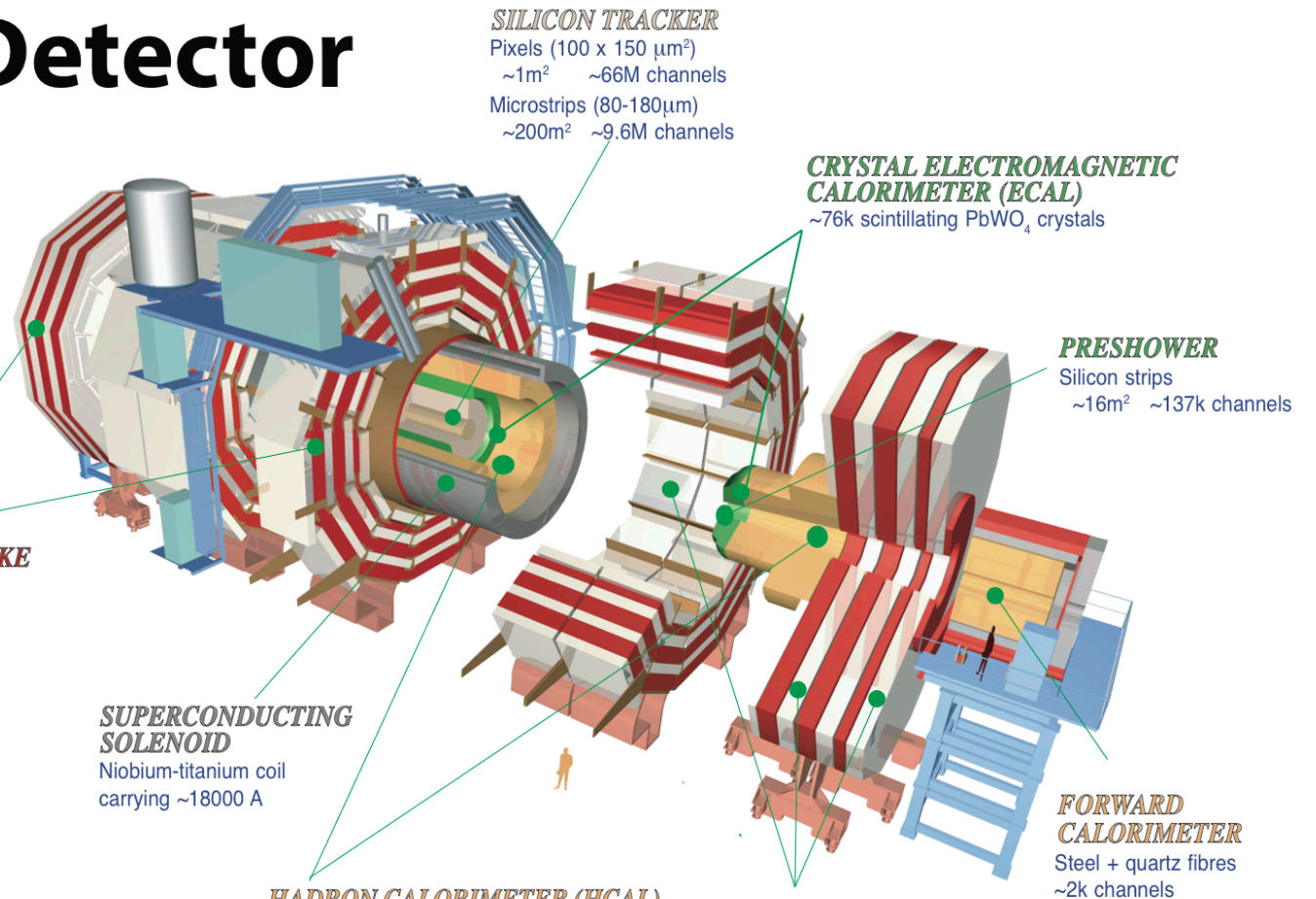
Width of  $\Delta\phi(p_T^{\text{miss}}, p_T^\mu)$  in  $W \rightarrow \mu\nu$

$$p_{T_{\text{miss}}} = \sqrt{p_{x_{\text{miss}}}^2 + p_{y_{\text{miss}}}^2}$$

# Compact Muon Solenoid (CMS)

## CMS Detector

Pixels  
 Tracker  
 ECAL  
 HCAL  
 Solenoid  
 Steel Yoke  
 Muons



**Total weight** : 14000 tonnes  
**Overall diameter** : 15.0 m  
**Overall length** : 28.7 m  
**Magnetic field** : 3.8 T

# MET reconstruction algorithms

## Particle-Flow (PF) MET

$$\vec{E}_T = - \sum_{\text{pf-candidates}} \vec{p}_T$$

- ★ negative of the vector sum over all transverse momentum of PF-candidates
- ★ used in most current CMS analyses

### No-PU PF MET

New

### MVA PF MET

- ★ divide PF particles into: particles from hard scattering and particles from pile-up
- ★ contribution from “pile-up” particles is scaled down
- ★ re-calculate MET from two particles categories above

- ★ multivariate regression (BDT) that produces a correction for the hadronic recoil
- ★ 5 MET variables calculated from PF particles
- ★ Trainings have been done to optimize the MET resolution

# The No-PU PF MET algorithm



★ **Principle:** divide PF particles into two categories

- **PF particles from hard scatter interaction (HS particles):** leptons/photons, PF particles within jets of  $p_T > 30$  GeV and pass the MVA PU-jet ID, charged hadrons not clustered within jets of  $p_T > 30$  GeV and associated to the HS vertex
- **PF particles from pile-up (PU particles):** charged hadrons that are neither within jets of  $p_T > 30$  GeV nor associated to the HS vertex, neutral PF particles within jets of  $p_T > 30$  GeV, PF particles within jets of  $p_T > 30$  GeV and fail the MVA PU-jet ID

★ PF particles from pile-up are scaled down by a factor :

$$S_F = \frac{\sum_{\text{HS-charged}} p_T}{\sum_{\text{HS-charged}} p_T + \sum_{\text{PU-charged}} p_T}.$$

★ No-PU PF MET is computed from :

$$\vec{E}_T = - \left[ \sum_{\text{leptons}} \vec{p}_T + \sum_{\text{HS-jets}} \vec{p}_T + \sum_{\text{HS-charged}} \vec{p}_T + S_F \cdot \left( \alpha \cdot \sum_{\text{PU-charged}} \vec{p}_T + \beta \cdot \sum_{\text{neutrals}} \vec{p}_T + \gamma \cdot \sum_{\text{PU-jets}} \vec{p}_T + \delta \cdot \vec{\Delta}_{\text{PU}} \right) \right].$$

$\alpha, \beta, \gamma, \delta$  optimized on  $Z \rightarrow \mu\mu$  to get the best MET resolution



# $E_T^{\text{miss}}$ Performance evaluation

Performance has to be studied in terms of:

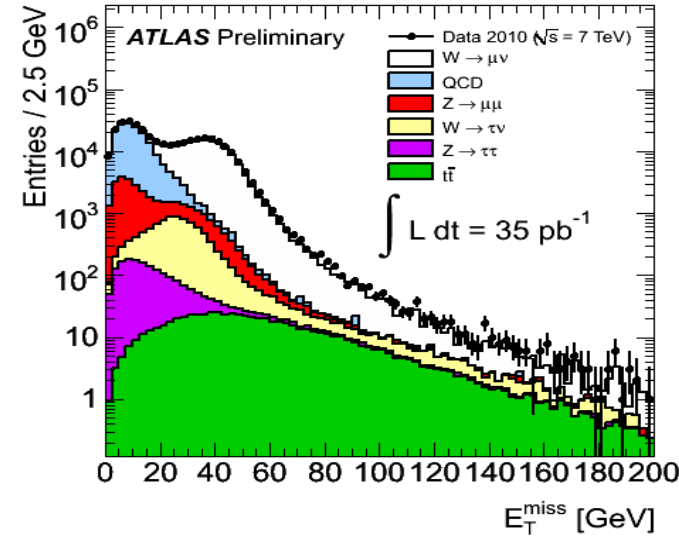
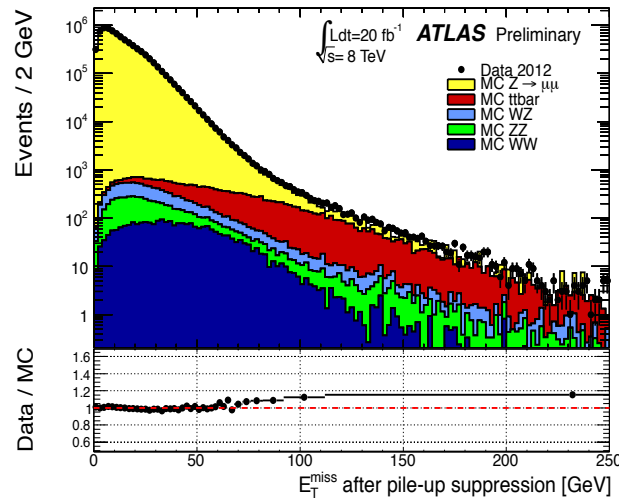
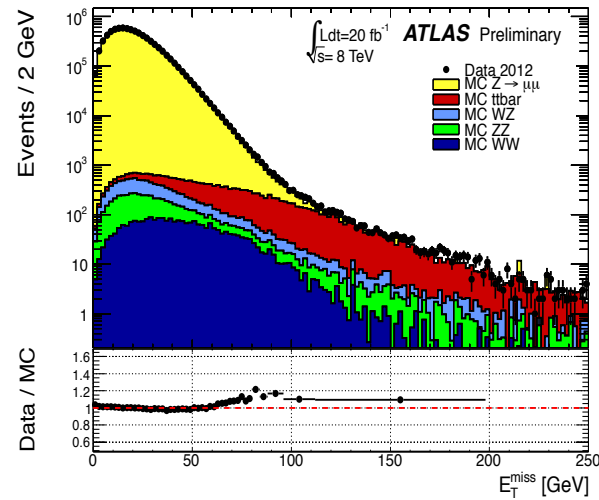
- Resolution (important for mass reconstruction in decays to non detected particles)
- Scale (important for mass reconstruction and when applying threshold cuts on  $E_T^{\text{miss}}$ )
- Tails (NO fake  $E_T^{\text{miss}}$ )
  
- Agreement between data and MC simulation

# Data-MC comparison

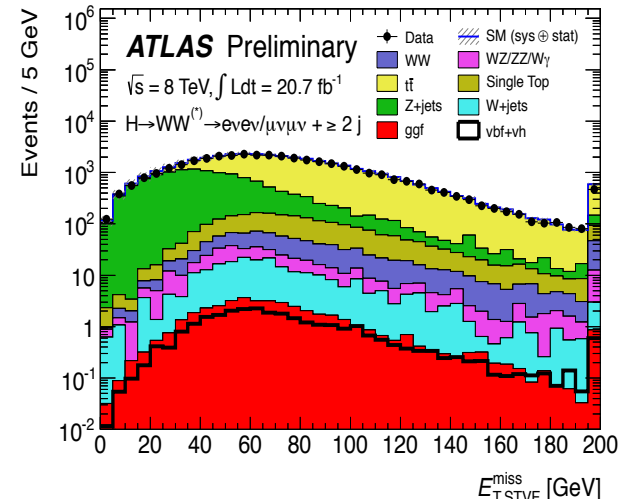
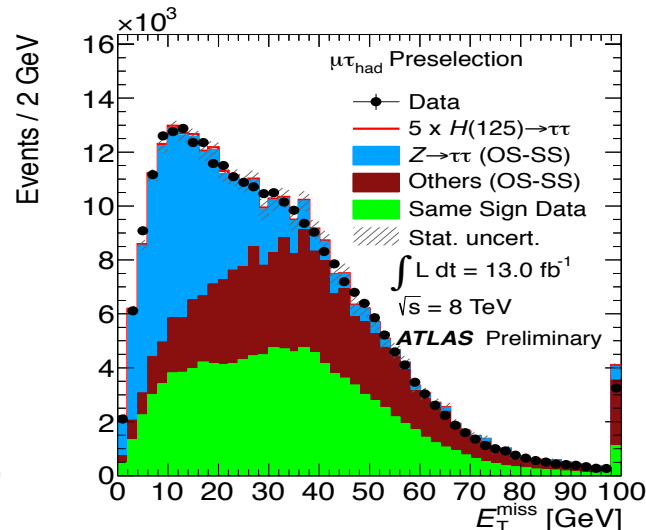
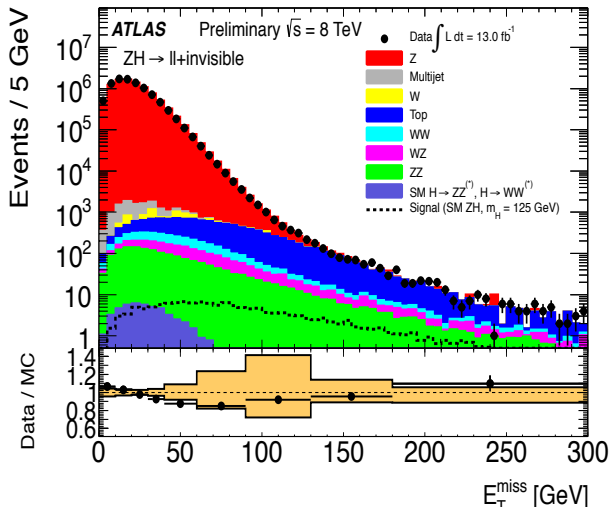
Important to understand  $E_T^{\text{miss}}$  both in data and simulation

- Check if the data are well described by MC simulation
- Check that there are no tails from fake  $E_t^{\text{miss}}$  in events where no true  $E_T^{\text{miss}}$  is expected

# Data-MC comparison ATLAS



Events in tails are compatible with physics signal candidates involving real  $E_T^{\text{miss}}$ . Very few additional tails after pile-up suppression



MC simulation describes data well.

# $E_T^{\text{miss}}$ Resolution

- The  $E_T^{\text{miss}}$  measurement is obtained from what is seen in the detector, so it depends on the total transverse energy measured in the detector, mainly on the total transverse energy in calorimeters (the muon momenta are better measured)

The resolution of the two  $E_T^{\text{miss}}$  components is estimated from the width of the distributions  $(E_x^{\text{miss}} - E_x^{\text{miss, True}}, E_y^{\text{miss}} - E_y^{\text{miss, True}})$

- It is studied as a function of  $\Sigma E_T$
- The stability of the resolution vs the pile-up can be studied looking at its dependence on  $\langle \mu \rangle$  or on the number of reconstructed vertices  $N_{\text{pv}}$ .

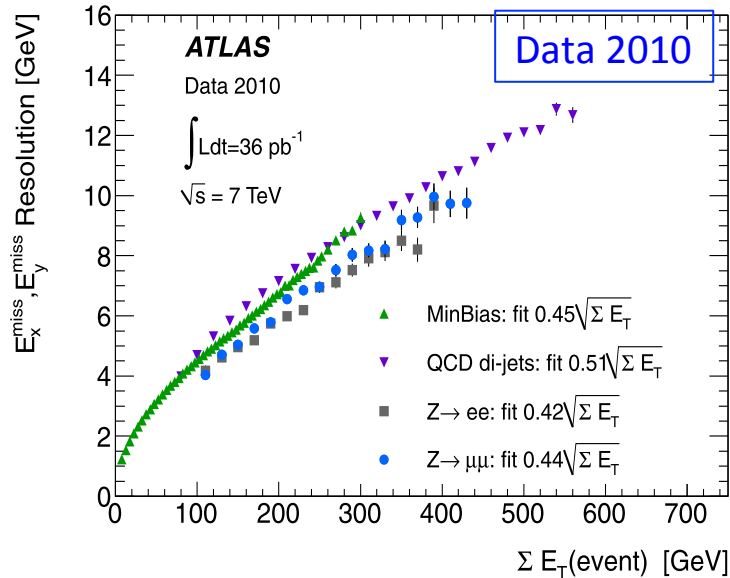
Can be studied in:

- data in events with NO true  $E_t^{\text{miss}}$
- in events with true  $E_T^{\text{miss}}$ , the resolution can be directly studied only in MC

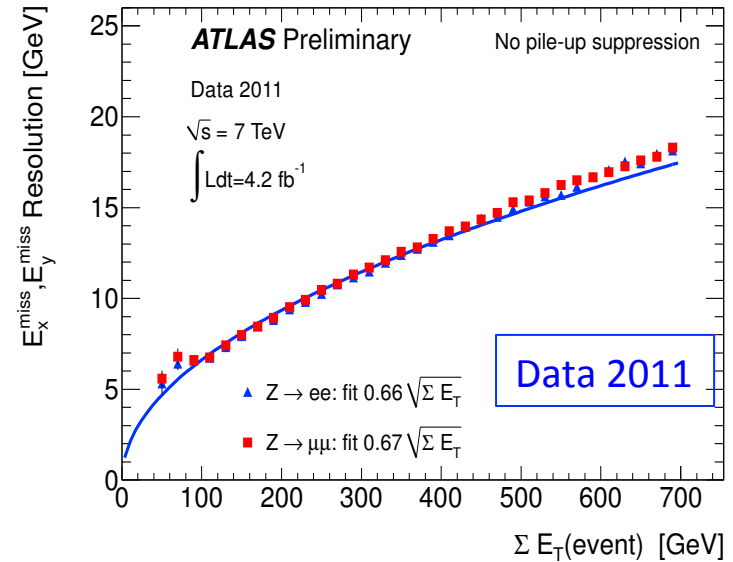
→ Importance of  $E_T^{\text{miss}}$  resolution in mass reconstruction in  $H \rightarrow \tau\tau$

# Resolution vs $\Sigma E_T$ in data at 7 TeV

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$E_T^{\text{miss}}$  resolution highly affected by pile-up

→ fitting with  $\sigma = k \cdot \sqrt{\Sigma E_T}$ .

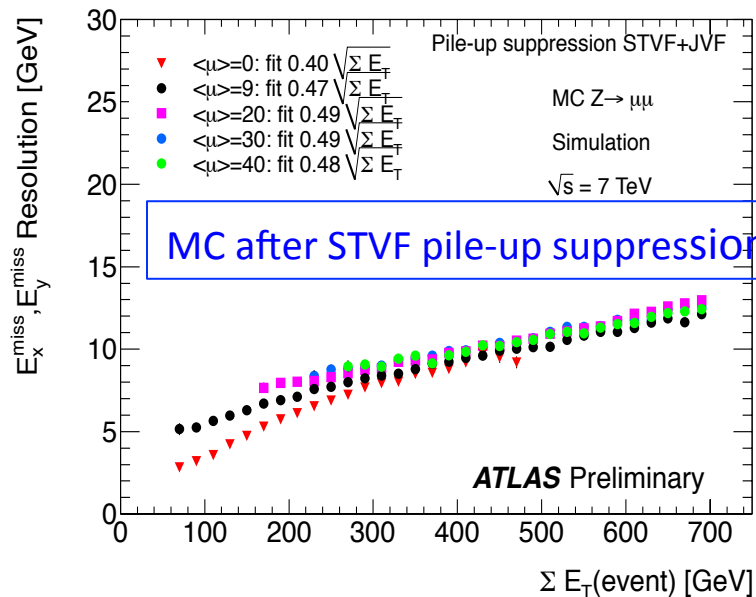
$k \sim 0.5 \text{ GeV}^{1/2}$  in 2010,  $\sim 0.7 \text{ GeV}^{1/2}$  in 2011

# Effect of pile-up on Resolution

Study of resolution with increasing pile-up conditions:

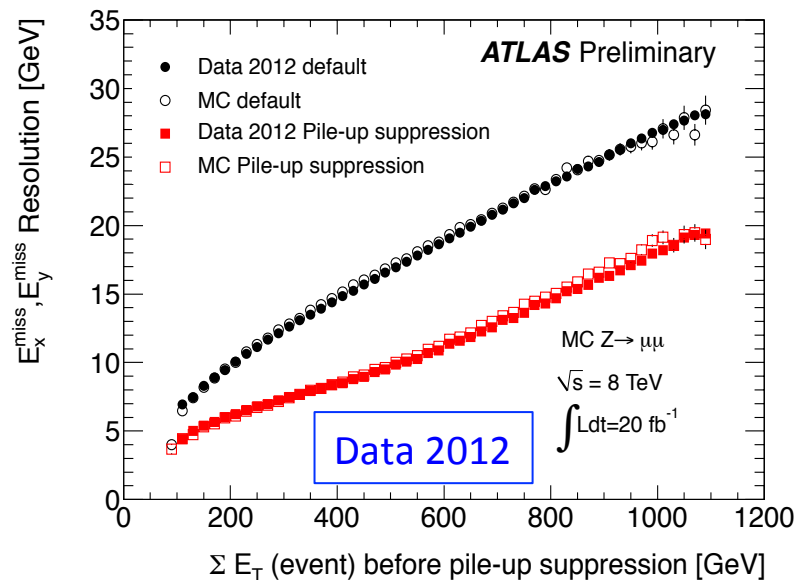
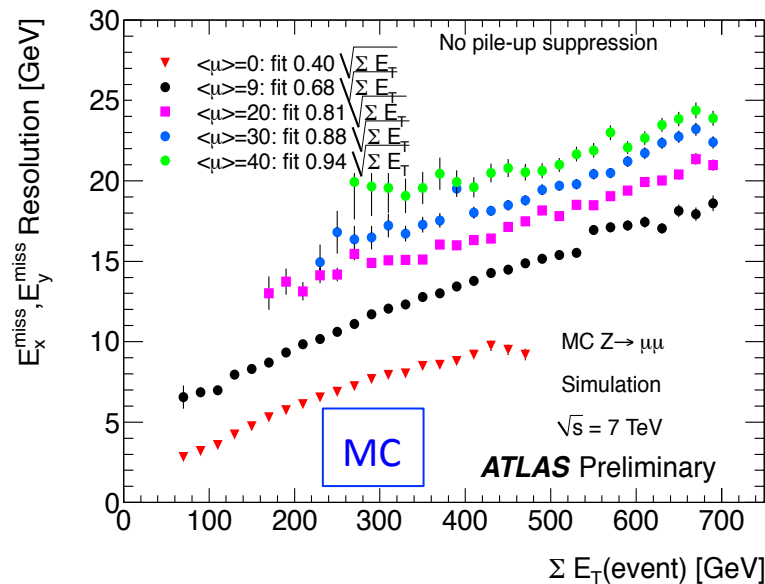
- $E_{x,y}^{\text{miss}}$  resolution doubles from  $\langle\mu\rangle = 0$  to 20 (2012 pile-up conditions)

$$\rightarrow \sigma(E_x^{\text{miss}}, E_y^{\text{miss}}) = (0.4 \text{ GeV}^{1/2} + 0.09 \text{ GeV}^{1/2} \cdot \sqrt{\mu}) \cdot \sqrt{\Sigma E_T}$$



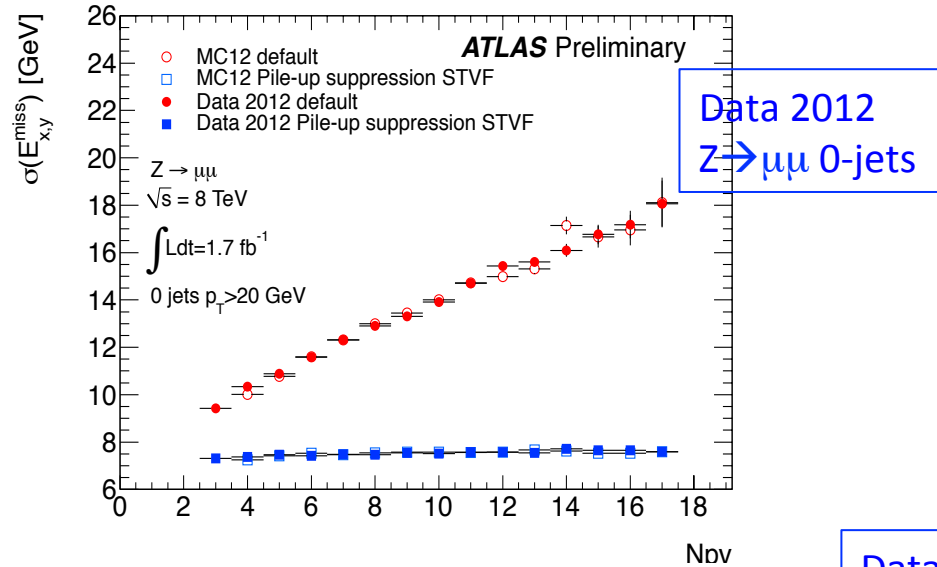
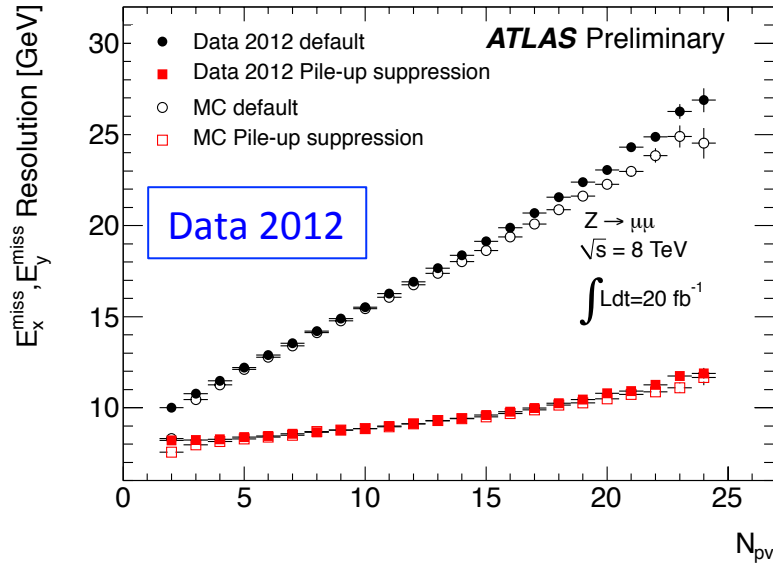
- $E_{x,y}^{\text{miss}}$  resolution closer to the resolution in after pile-up suppression with **STVF**

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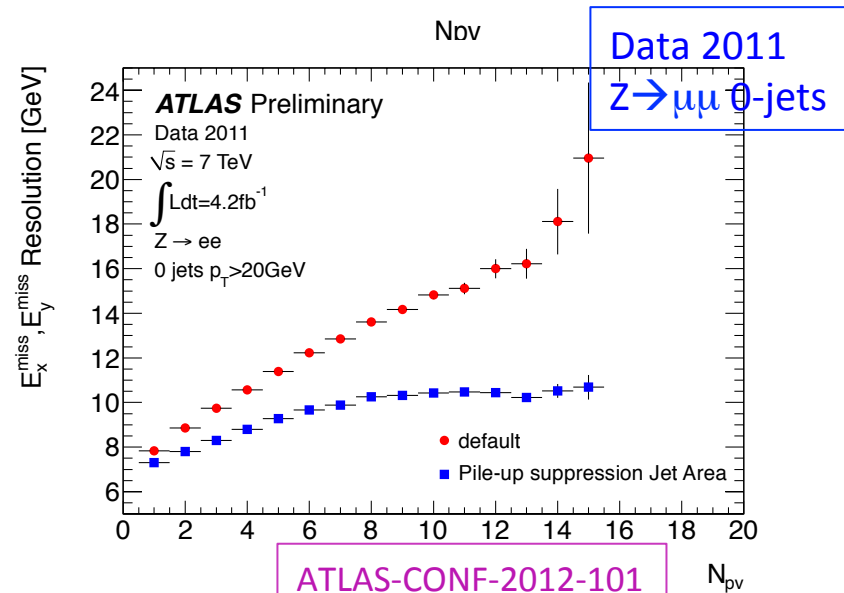
# Resolution after pile-up suppression in data at 8 TeV

**STVF method:** Resolution less dependent on  $N_{pv}$



**Jet area method:**

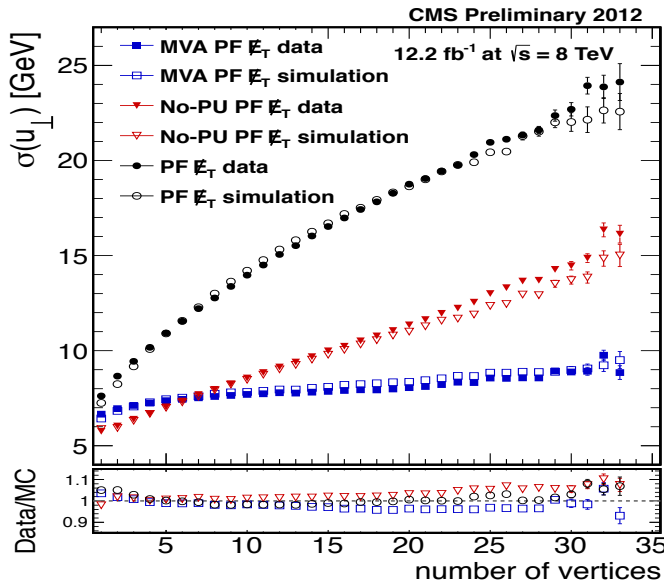
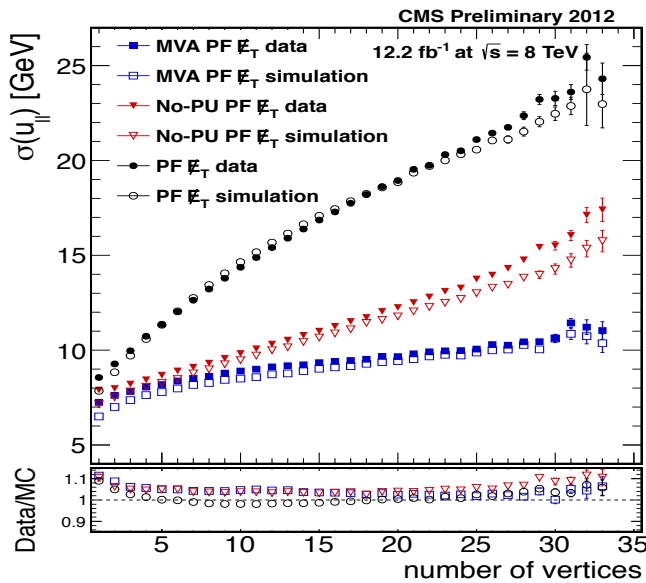
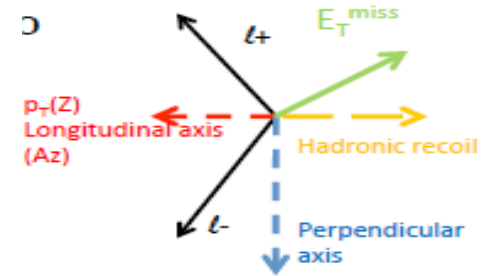
- Improves the resolution but some residual dependence on  $N_{pv}$  still present
- Some more improvement combining with a track based filter (JVF) on  $k_t$  jets



# Resolution after pile-up suppression CMS

CMS shows the resolution of the hadronic recoil

$Z \rightarrow \mu\mu$  events



CMS PAS JME-12-002

- Resolution less dependent on  $N_{pv}$  after pile-up suppression



# H → ττ mass reconstruction

The full reconstruction of the H → ττ mass

$$\mathbf{m}_{\tau\tau} = \sqrt{(2(\mathbf{E}_{\tau\text{had}} + \mathbf{E}_{\nu 1})(\mathbf{E}_{\text{lep}} + \mathbf{E}_{\nu 2})(1 - \cos\theta))}$$

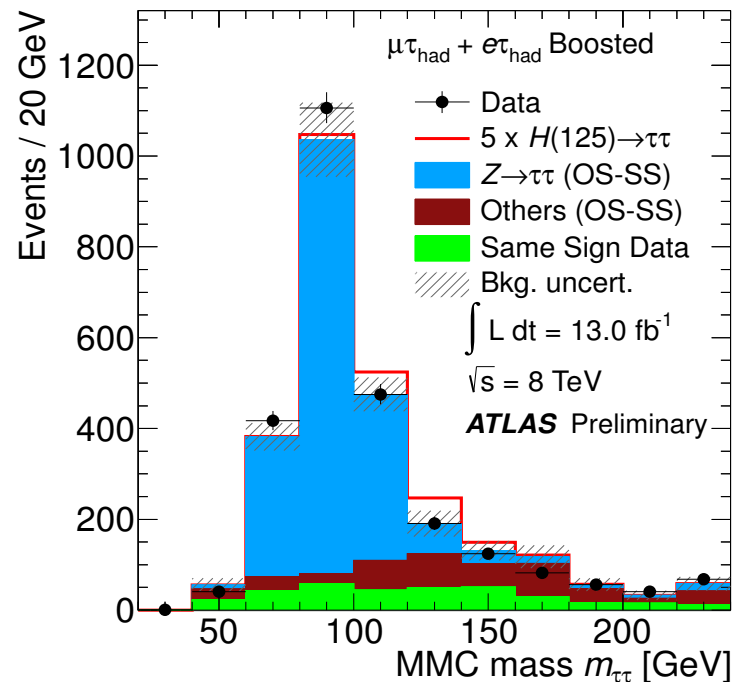
requires solving equations with more unknown than constraints.

- **Collinear approximation** assumes that the visible and undetectable τ decay products are collinear
- **Missing Mass Calculator (MMC)** scans over the neutrino directions and picks the most likely value of  $m_{\tau\tau}$ , according to the simulated probability functions from the τ decay.

A good  $E_{\tau}^{\text{miss}}$  resolution is crucial for the  $m_{\tau\tau}$  reconstruction.

- $E_{\tau}^{\text{miss}}$  after pile-up suppression gives a better invariant  $m_{\tau\tau}$  reconstruction with the collinear approximation (efficiency and resolution)
- Improved MMC results with  $E_{\text{tmiss}}$  after pile-up suppression with STVF.

$$\begin{aligned} E_{T_x} &= p_{\text{mis}_1} \sin \theta_{\text{mis}_1} \cos \phi_{\text{mis}_1} + p_{\text{mis}_2} \sin \theta_{\text{mis}_2} \cos \phi_{\text{mis}_2} \\ E_{T_y} &= p_{\text{mis}_1} \sin \theta_{\text{mis}_1} \sin \phi_{\text{mis}_1} + p_{\text{mis}_2} \sin \theta_{\text{mis}_2} \sin \phi_{\text{mis}_2} \\ M_{\tau_1}^2 &= m_{\text{mis}_1}^2 + m_{\text{vis}_1}^2 + 2\sqrt{p_{\text{vis}_1}^2 + m_{\text{vis}_1}^2} \sqrt{p_{\text{mis}_1}^2 + m_{\text{mis}_1}^2} \\ &\quad - 2p_{\text{vis}_1} p_{\text{mis}_1} \cos \Delta\theta_{vm_1} \\ M_{\tau_2}^2 &= m_{\text{mis}_2}^2 + m_{\text{vis}_2}^2 + 2\sqrt{p_{\text{vis}_2}^2 + m_{\text{vis}_2}^2} \sqrt{p_{\text{mis}_2}^2 + m_{\text{mis}_2}^2} \\ &\quad - 2p_{\text{vis}_2} p_{\text{mis}_2} \cos \Delta\theta_{vm_2} \end{aligned}$$



# $E_T^{\text{miss}}$ Scale

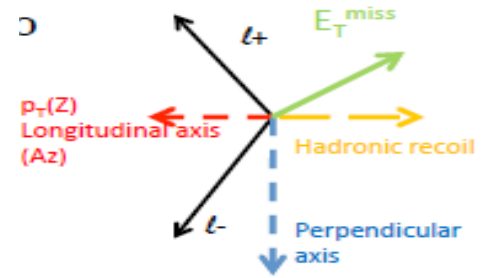
It is important to have a reconstructed  $E_T^{\text{miss}}$  as much as possible close to the the  $E_T^{\text{miss, True}}$ .

This can be checked:

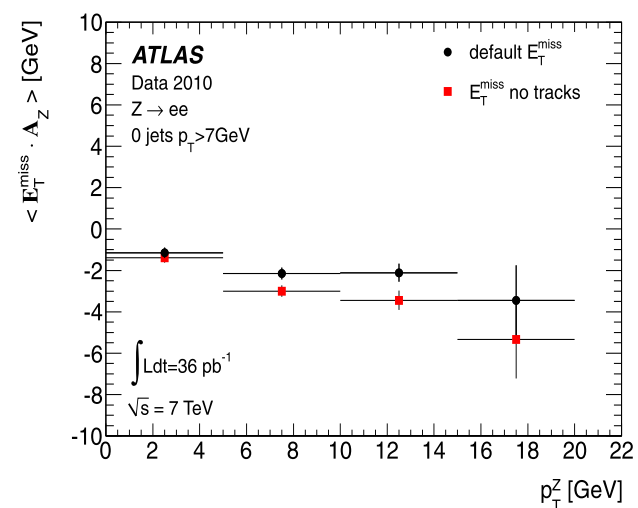
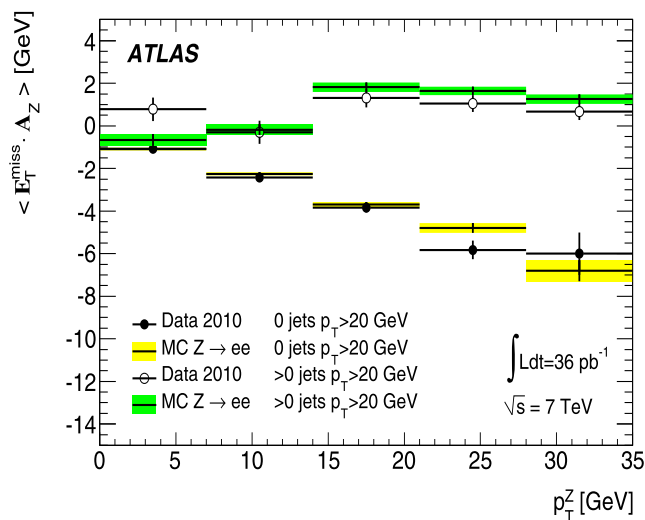
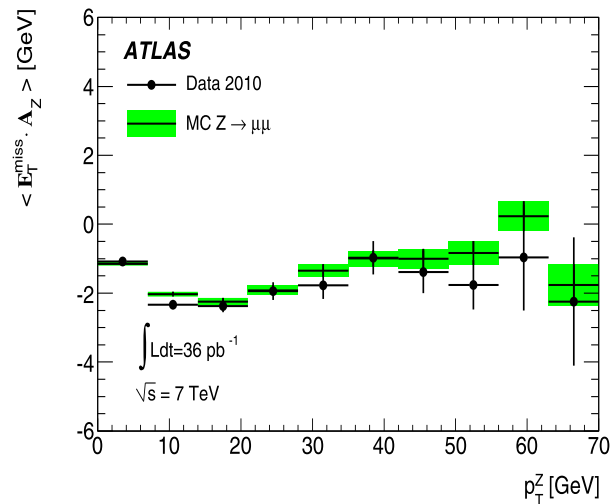
- in data  $Z \rightarrow ll$  events from the projection of the  $E_T^{\text{miss}}$  along the transverse direction of the Z boson
- in events with true  $E_T^{\text{miss}}$ , the linearity  $(E_T^{\text{miss}} - E_T^{\text{miss, True}}) / E_T^{\text{miss, True}}$ , can be studied in MC events
- reconstructing the mass in  $W \rightarrow lv$  and  $Z \rightarrow \tau\tau$  each of which contain true  $E_T^{\text{miss}}$  from unobserved neutrinos.

# $E_T^{\text{miss}}$ scale in $Z \rightarrow \ell\ell$ events

$E_T^{\text{miss}}$  projection onto the  $p_T^Z$



The longitudinal axis is defined by the vectorial sum of the 2 leptons momenta and it is sensitive to the balance between the muons and the hadronic recoi. If the leptons perfectly balance the hadronic recoil the projection of  $E_T^{\text{miss}}$  along the longitudinal axis ( $A_z$ ) should be zero  $\rightarrow E_T^{\text{miss}}$  Diagnostic plot



The hadronic recoil is under-estimated mainly in events with NO jets dominated by Soft Term

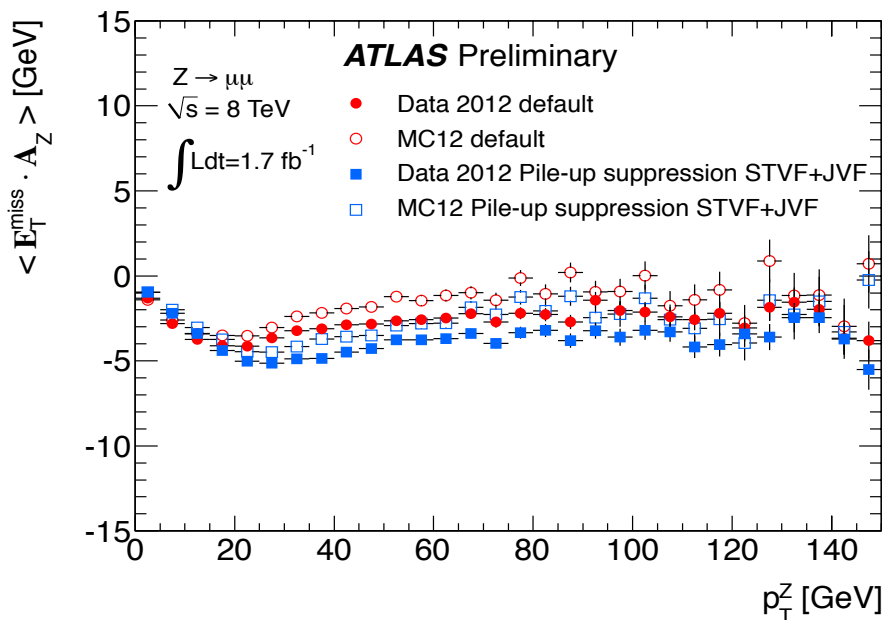
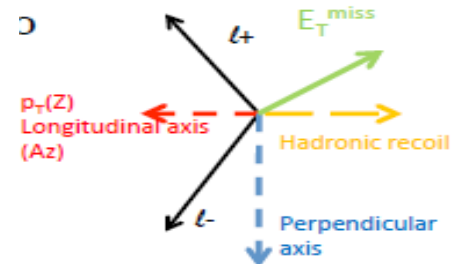
Adding tracks improves the hadronic recoil

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# $E_T^{\text{miss}}$ scale: effect of pile-up

The pile-up also affect the  $E_t^{\text{miss}}$  scale:

- Some more bias is observed in the  $E_T^{\text{miss}}$  projection onto the  $p_T^Z$  (Diagnostic plot)

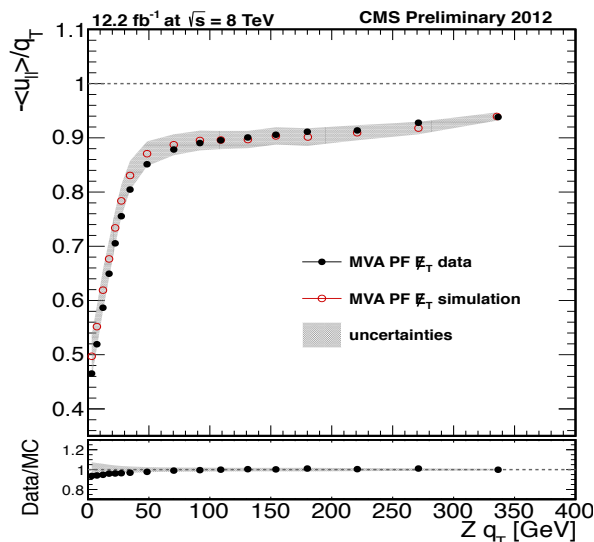
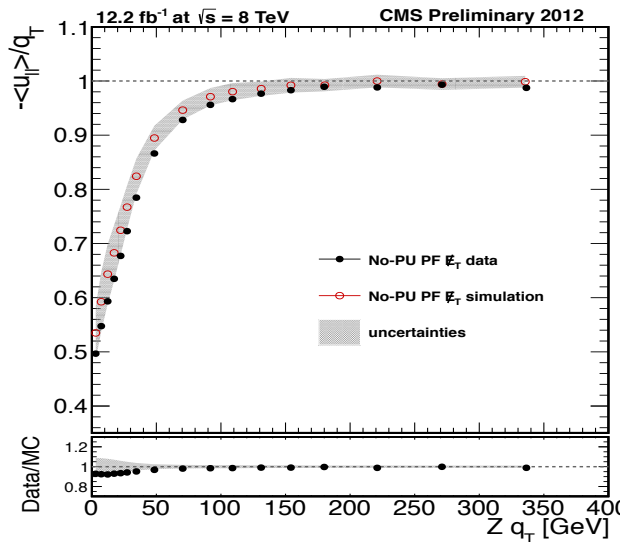


- After pile-up suppression with STVF method: the bias increases
- Smaller bias with Jet Area pile-up suppression

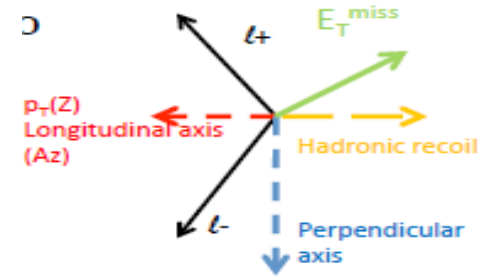
# Scale after pile-up suppression CMS

CMS shows the projection of the hadronic recoil onto  $p_T^Z$

$Z \rightarrow \mu\mu$  events



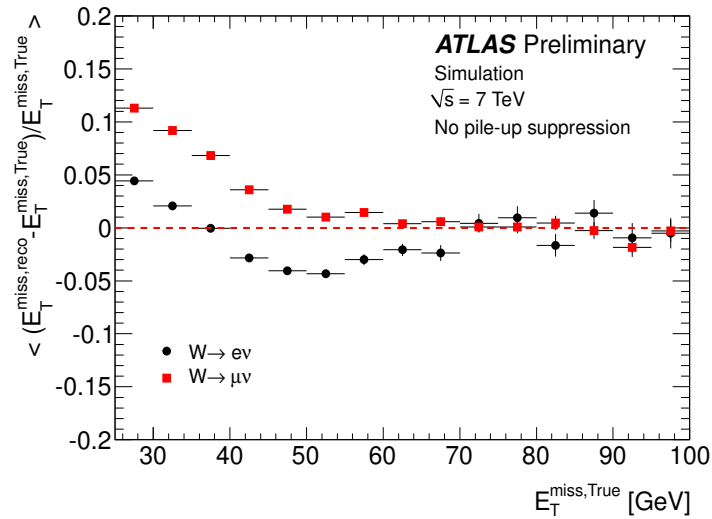
CMS PAS JME-12-002



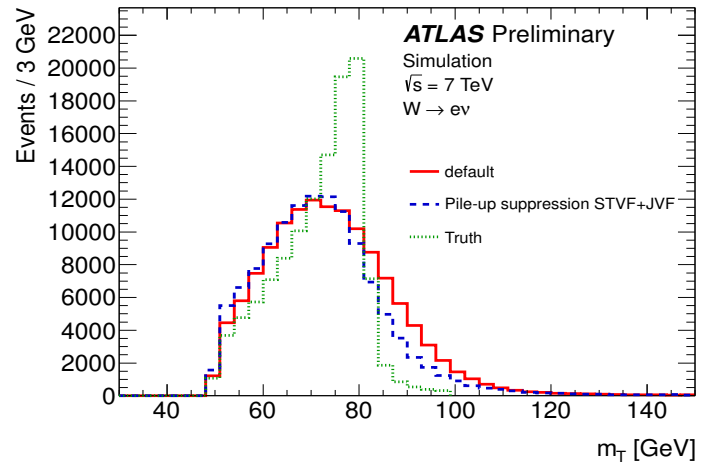
- The MVA response is around 0.9 because the BDT training data for the MVA used in this study is optimized for the improved resolution rather than for the unity response.

# $E_T^{\text{miss}}$ scale: Linearity and transverse mass reconstruction in $W \rightarrow \text{ln}$

$$\frac{(E_T^{\text{miss}} - E_T^{\text{miss, True}})}{E_T^{\text{miss, True}}},$$



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The bias at low  $E_T^{\text{miss, True}}$  is due to the finite resolution of  $E_T^{\text{miss}}$ .

→ The reconstructed  $E_T^{\text{miss}}$  is positive by definition, so the relative difference is positive when the  $E_T^{\text{miss, True}}$  is small.

The reconstructed  $m_T = \sqrt{2p_T^\ell E_T^{\text{miss}}(1 - \cos\phi)}$ , is closer to the Truth after pile-up suppression

# Systematic Uncertainties

The estimation of the systematic uncertainties on the  $E_T^{\text{miss}}$  measurement is needed for physics.

# Evaluation of $E_T^{\text{miss}}$ systematic uncertainty

$E_T^{\text{miss}}$  makes use of reconstructed objects, so its systematic uncertainty can be calculated from the uncertainty on each object and from the uncertainty on the Soft Term

The contribution of each term varies for different channels

- in Z and W events the contribution of Soft term is important.

Evaluation of systematic uncertainty on Soft Term (scale and resolution) with two methods:

- from data/MC ratio in  $Z \rightarrow \mu\mu$  events with **NO jets**
  - Scale uncertainty from  $E_T^{\text{miss}}$  projection onto the  $p_T^Z$
  - Resolution uncertainty from resolution
- from the balance between the SoftTerm and  $p_T^{\text{Hard}}$  in  $Z \rightarrow \mu\mu$  events
  - Scale and resolution from decomposition of  $E_t^{\text{miss,SoftTerm}}$  along  $p_T^{\text{Hard}}$  and its orthogonal direction

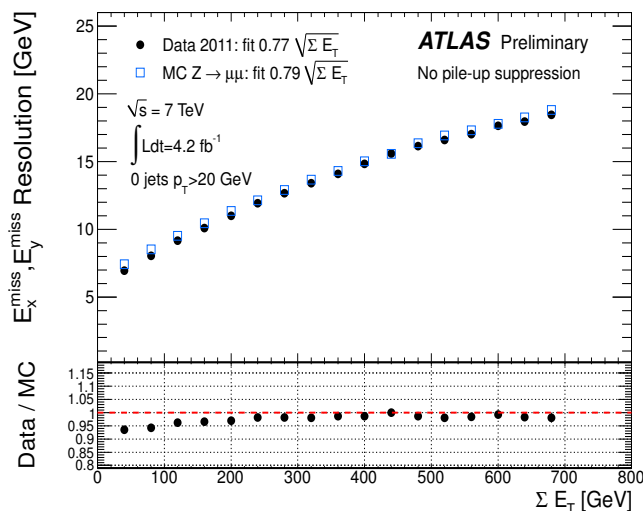
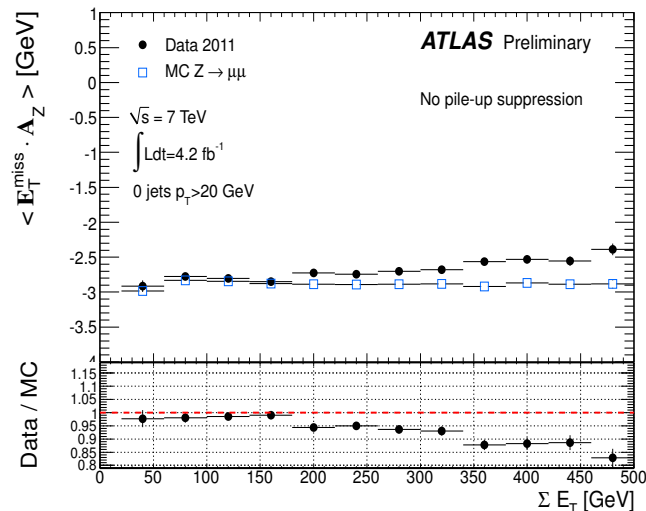
Scale uncertainty from the topocluster energy uncertainty (from e/p)

- difficult to determine clustering efficiency and scale in busy environment
- the cluster energy uncertainty in the forward region is conservatively estimated, since the uncertainty cannot be evaluated using tracks



# Evaluation of $E_T^{\text{miss}}$ systematic uncertainty

$E_T^{\text{miss,SoftTerm}}$  uncertainty from Data/MC ratio in  $Z \rightarrow \mu\mu$  events with NO jets  $> 20\text{GeV}$

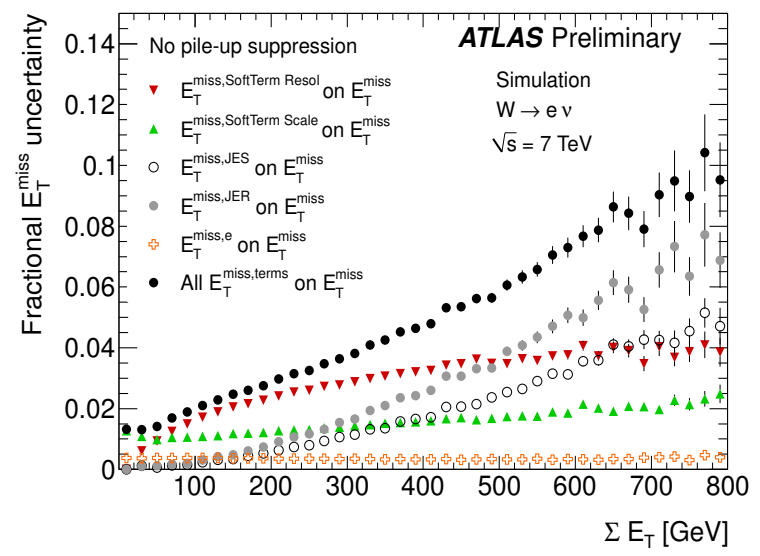


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- Uncertainty on scale from Data/MC ratio in Diagnostic plot (5%)
- Uncertainty on resolution from Data/MC ratio in Resolution plot (2%)

→ The overall systematic uncertainty in  $W \rightarrow e\nu$  events, calculated propagating the systematic uncertainties from all objects and SoftTerm and adding in quadrature is increasing with  $\Sigma E_T$  and around 3% in average (in 2011).

→ Systematic uncertainty is topology dependent



# Evaluation of $E_T^{\text{miss}}$ scale uncertainty in-situ

Use  $m_T$  distribution in data  $W \rightarrow l\nu$  to evaluate scale and resolution

→ Scale/smear the  $m_T$  distribution in MC with:

$$E_T^{\text{miss, smeared}} = \alpha E_T^{\text{miss, True}} \cdot \text{Gauss}(0, k \cdot \sqrt{\Sigma E_T})$$

and compare with data

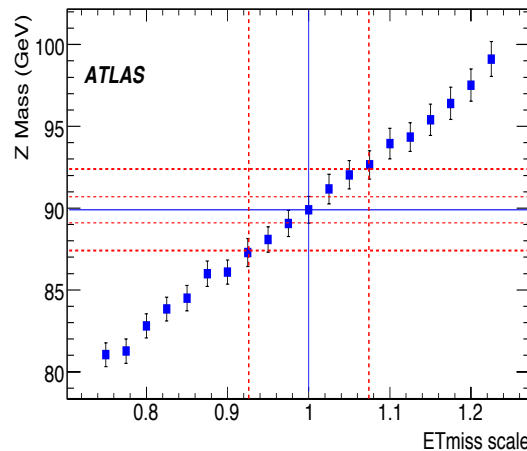
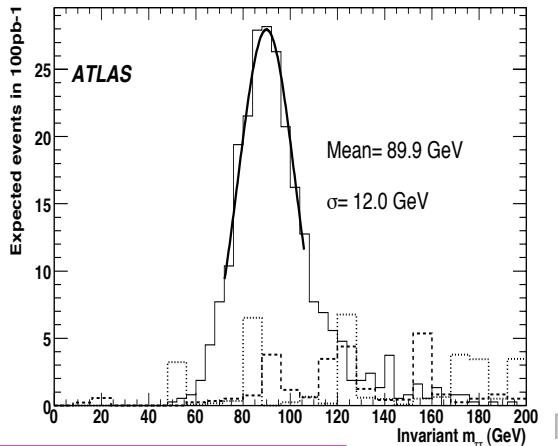
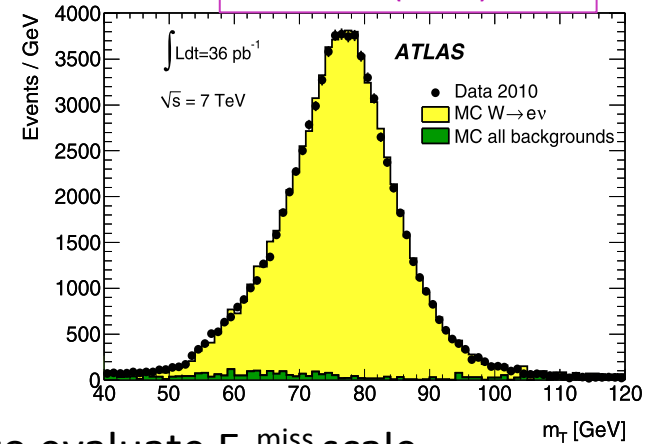
Can determine the  $E_T^{\text{miss}}$  absolute with a global uncertainty of about 2% (integrated luminosity of 36pb<sup>-1</sup>)

Use reconstructed invariant mass in  $Z \rightarrow \tau\tau \rightarrow \text{lepton-hadron}$  to evaluate  $E_T^{\text{miss}}$  scale

→ Use collinear approximation to get  $E_{\nu 1}$  and  $E_{\nu 2}$  to reconstruct invariant mass

$$m_{\tau\tau} = \sqrt{(2(\mathbf{E}_{\tau\text{had}} + \mathbf{E}_{\nu 1})(\mathbf{E}_{\text{lep}} + \mathbf{E}_{\nu 2})(1 - \cos\theta))}$$

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Check dependence of invariant mass on  $E_T^{\text{miss}}$  scale variations to get systematic uncertainty on the  $E_T^{\text{miss}}$  scale

# Conclusions

- $E_T^{\text{miss}}$  is an event complex quantity, calculated from signals in all detectors and is affected by detector acceptance, problems and non-collision background in the detector and by noise and pile-up
- ATLAS uses a flexible algorithm that allows to use physics objects with their proper calibration and soft energy contributions
- CMS uses a particle-flow algorithm both for high and low  $p_T$  contributions
- The pile-up conditions at increased LHC luminosity gives a deterioration in the  $E_T^{\text{miss}}$  performance
- Pile-up suppression methods are needed to mitigate pile-up mainly in jets and in the soft term to reduce the pile-up impact especially on the resolution
- A good  $E_T^{\text{miss}}$  performance in terms of resolution, scale and tails is crucial for many physics analyses
- The  $E_T^{\text{miss}}$  uncertainty is calculated from the uncertainties on the scale and the resolution of each physics object and of soft term and can also be calculated in-situ using  $W \rightarrow l\nu$  and  $Z \rightarrow \tau\tau$  events
- Dedicated optimisation of all these techniques needed to face the new challenge in 2015 data taking at very high luminosity.