

Fits to $\tau \rightarrow K \pi \nu_\tau$ decays

Viabie information can be obtained from the decay spectra for exclusive τ -decay channels.

A first step in this direction is a reliable description of the

$\tau \rightarrow \nu_\tau K \pi$ decay spectrum:

(MJ, Pich, Portolés 2006/08)

(Boito, Escribano, MJ 2008)

(Passemar et al. 2006-11)

$$\frac{d\Gamma_{K\pi}}{d\sqrt{s}} = \frac{G_F^2 |V_{us}|^2 M_\tau^3}{32\pi^3 s} \left(1 - \frac{s}{M_\tau^2}\right)^2 \times$$

$$\left[\left(1 + 2 \frac{s}{M_\tau^2}\right) q_{K\pi}^3 |F_+^{K\pi}(s)|^2 + \frac{3\Delta_{K\pi}^2}{4s} q_{K\pi} |F_0^{K\pi}(s)|^2 \right].$$

To this end the $K\pi$ vector and scalar form factors $F_+^{K\pi}(s)$ and $F_0^{K\pi}(s)$ are required as an input.

A description of the $K\pi$ vector form factor can be obtained within chiral perturbation theory with resonances ($R\chi$ PT):

$$F_{+}^{K\pi}(s) = \frac{m_{K^*}^2}{m_{K^*}^2 - s - \kappa \operatorname{Re} \widetilde{H}_{K\pi}(s) - im_{K^*} \gamma_{K^*}(s)}.$$

The parameters of this model, namely m_{K^*} and γ_{K^*} , can be fitted from experimental data for p -wave $K\pi$ scattering, or from the τ data.

The physical parameters M_{K^*} and Γ_{K^*} can be inferred from the pole of $F_{+}^{K\pi}(s)$ in the complex s -plane.

Also a second resonance contribution can easily be included.

In the **elastic, single-channel** case a **subtracted** dispersive representation is **available** which is related to the **Omnès** solution:

$$\frac{F_+^{K\pi}(s)}{F_+^{K\pi}(0)} = \exp \left(\frac{\alpha_1 s}{M_\pi^2} + \frac{\alpha_2 s^2}{2M_\pi^4} + \frac{s^3}{\pi} \int_{s_{K\pi}}^{s_{\text{cut}}} \frac{\delta_{K\pi}(s')}{(s')^3 (s' - s - i0)} ds' \right)$$

where $\delta_{K\pi}(s)$ is the **elastic P-wave $K\pi$ phase shift**.

The **two** subtraction constants α_1 and α_2 are related to **slope** and **curvature** of the **vector** form factor:

$$\lambda'_+ = \alpha_1, \quad \lambda''_+ = \alpha_2 + \alpha_1^2.$$

The **scalar** form factor $F_0^{K\pi}(s)$ can be obtained from a dispersion relation analysis of **S-wave** $K\pi$ scattering data.

(MJ, Oller, Pich 2000/02)

The **scalar** form factors are defined by:

$$i \langle \Omega | \partial^\mu (\bar{s} \gamma_\mu u) | \Gamma \rangle = \frac{\Delta_{K\pi}}{\sqrt{2}} C_\Gamma F_\Gamma(s)$$

where the C_Γ are **Clebsch-Gordan** coefficients and

$$\Delta_{K\pi} = M_K^2 - M_\pi^2.$$

As yet, the form factors $F_i(s)$ have not been **measured** directly. Thus an **indirect** determination is required.

From **unitarity** we have the following relation:

$$\text{Im}F_k(s) = \sum_i \sigma_i(s) F_i(s) t_0^{ik}(s)^*$$

with $t_0^{ik}(s)$: **S-wave** $I=1/2$ scattering amplitudes.

The $F_i(s)$ also satisfy **dispersion** relations. In the **2-channel** case with $F_1 \equiv F_{K\pi}$ and $F_3 \equiv F_{K\eta'}$:

$$F_1(s) = \frac{1}{\pi} \int_{s_1}^{\infty} \frac{\sigma_1 F_1 t_0^{11}(s')^*}{(s' - s - i0)} ds' + \frac{1}{\pi} \int_{s_3}^{\infty} \frac{\sigma_3 F_3 t_0^{13}(s')^*}{(s' - s - i0)} ds'$$

$$F_3(s) = \frac{1}{\pi} \int_{s_1}^{\infty} \frac{\sigma_1 F_1 t_0^{13}(s')^*}{(s' - s - i0)} ds' + \frac{1}{\pi} \int_{s_3}^{\infty} \frac{\sigma_3 F_3 t_0^{33}(s')^*}{(s' - s - i0)} ds'$$

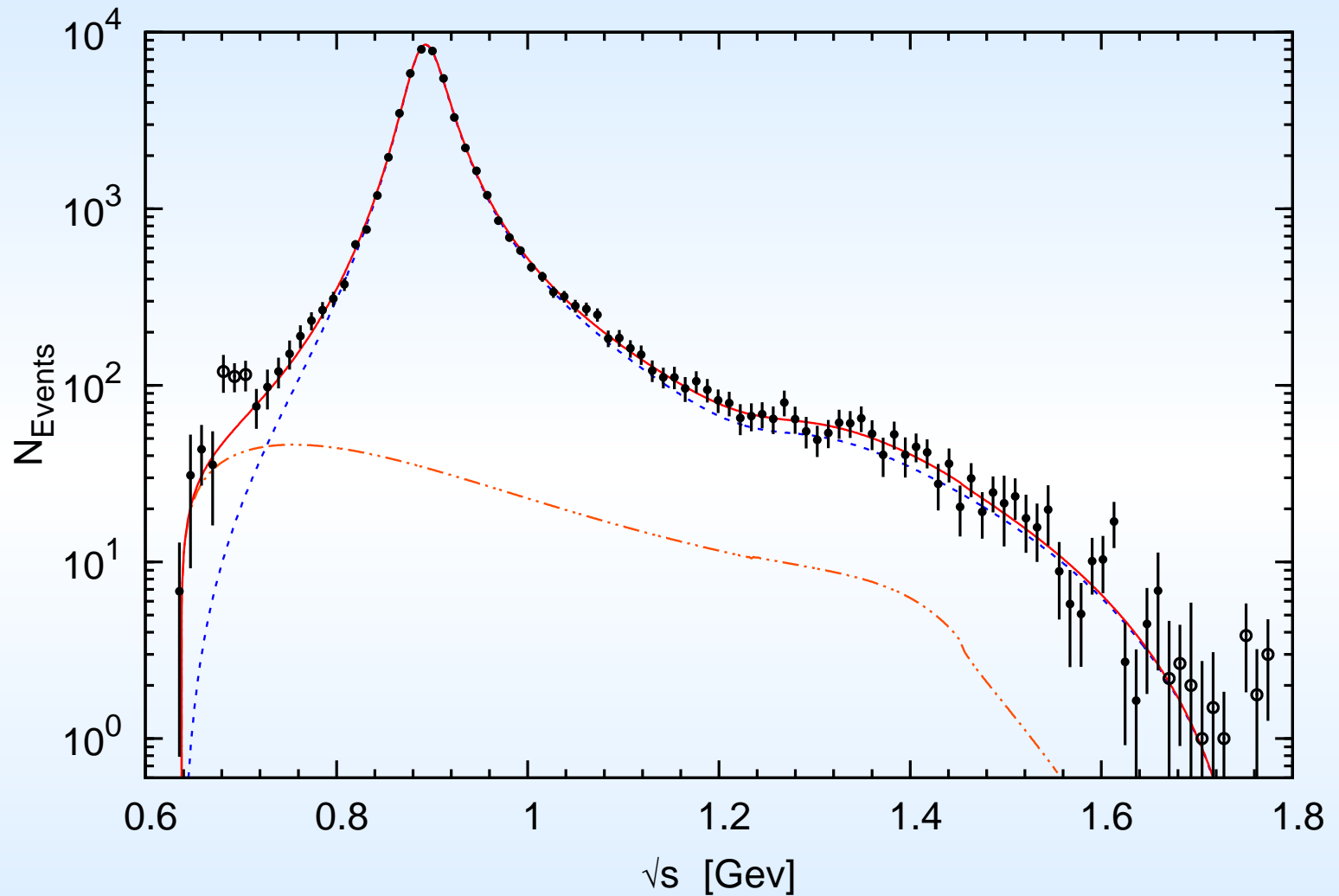
Several **ingredients** are required for solving the set of **coupled** integral equations:

☞ An **input** for the scattering amplitudes is obtained by fitting an Ansatz from **resonance** ChPT to **experimental** data for **S -wave $K\pi$** scattering. (M.J., Oller, Pich 2000/02)

☞ Two **integration** constants are also required. These can be chosen to be: $F_{K\pi}(0) = 0.972 \pm 0.012$ and

$$\frac{F_{K\pi}(\Delta_{K\pi})}{F_{K\pi}(0)} = \frac{F_K}{F_\pi F_{K\pi}(0)} + \frac{\Delta_{\text{CT}}}{F_{K\pi}(0)} = 1.2346(53)$$

☞ The last relation follows from the ratio of **leptonic K** and π decays, as well as $|V_{us}|F_{K\pi}(0)$ from K_{l3} decays.



$$M_{K^*} = 892.0 \pm 0.5 \text{ MeV}, \quad \Gamma_{K^*} = 46.5 \pm 1.1 \text{ MeV}$$

As a prediction of the model, we obtain the slope and the curvature of the vector form factor $F_+^{K\pi}(s)$:

$$\lambda'_+ = (25.49 \pm 0.31) \cdot 10^{-3}, \quad \lambda''_+ = (12.22 \pm 0.14) \cdot 10^{-4}.$$

The can be compared to earlier experimental determinations:

Collaboration	$\lambda'_+ [10^{-3}]$	$\lambda''_+ [10^{-3}]$
ISTRA 04	24.9 ± 1.6	0.84 ± 0.41
KTEV 04	20.64 ± 1.75	3.20 ± 0.69
NA48 04	28.0 ± 2.4	0.2 ± 0.5
KLOE 06	25.5 ± 1.8	1.4 ± 0.8

- Employing a **dispersive** representation of the $K\pi$ form factors, a **satisfactory** description of the $\tau \rightarrow \nu_\tau K\pi$ **decay** spectrum can be obtained.
- While a **coupled-channel** analysis is available for the **scalar** $K\pi$ form factors, our model for the **vector** form factors is **purely** elastic.
- To my **mind**, fits to **experimental** data should be done in a **two-way** approach: on the **one** hand, **experimentalists** can **try** to fit **theoretical** models provided by **theorists**.
- On the **other** hand, it would be **very** useful if **unfolded** distributions with **correlations** would be made available.

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Thank You for Your attention !