

# **The effects of a neutron star translational and rotational motion in observable timing and evolution of radiopulsars**

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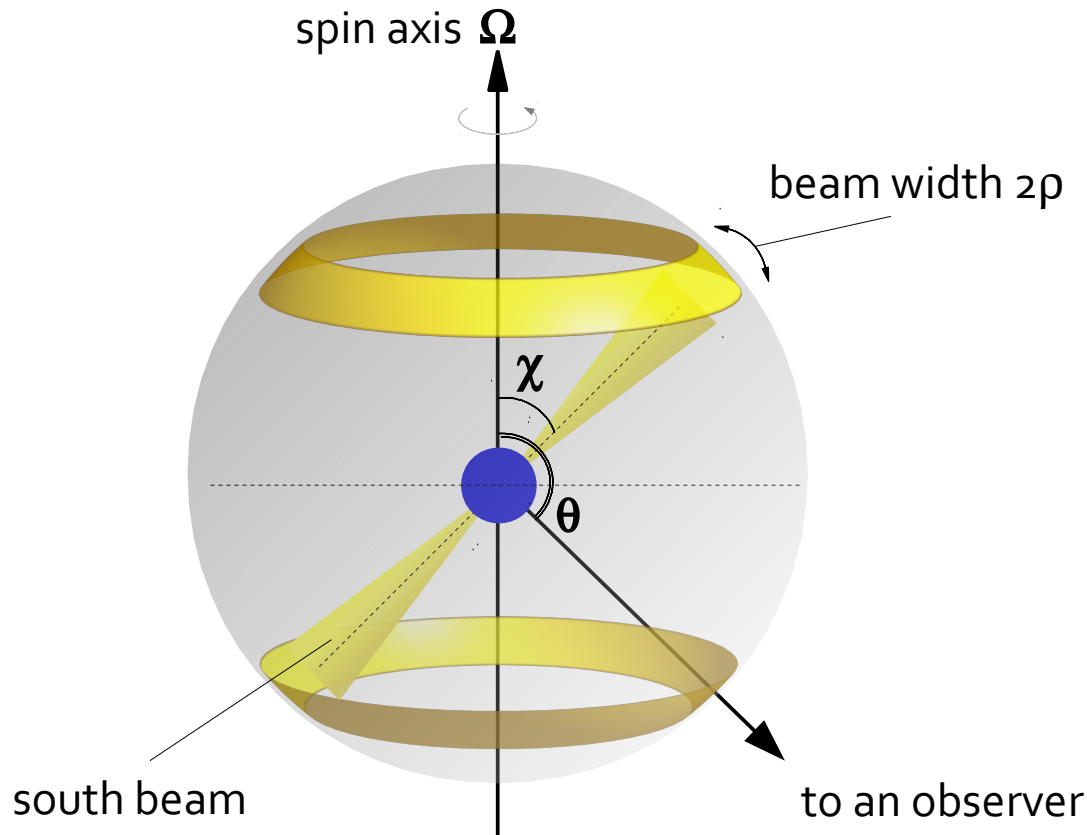
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# Outline

We've studied the influence of neutron stars complex rotation (possible free & force precession) and **space motion** on the observed properties pulsars ensemble: timing behaviour, distribution density on the P-Pdot diagram etc.

## Contents of this talk:

- Pulsars beaming fraction and observed fraction
- Anisotropy in spin – observer angles distribution
- Modeling a galaxy of pulsars
- Results
- Conclusions



### observed fraction

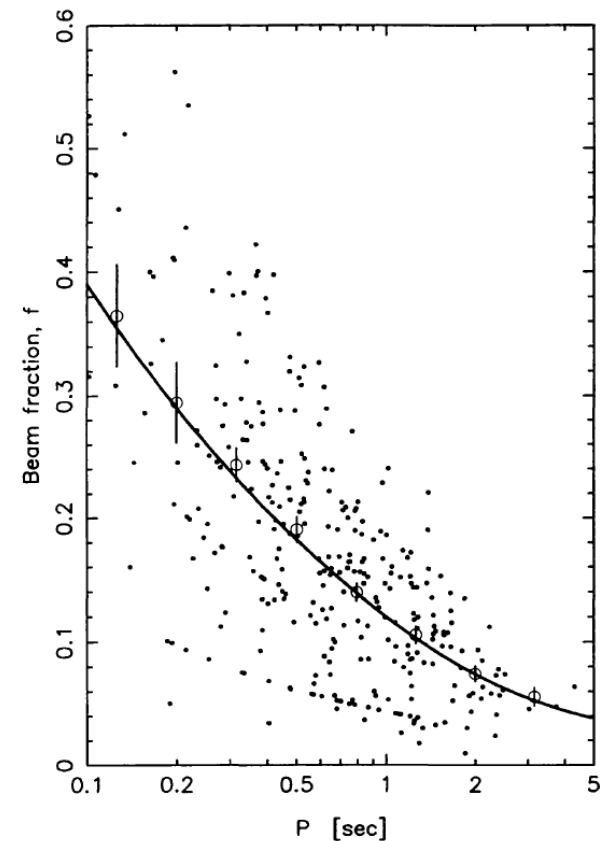
$$OF \equiv \frac{\text{number of pulsars directed to an observer}}{\text{total number of active pulsars}}$$

**BF of a pulsars group is an equivalent to their OF only in case of isotropically distributed  $\theta$**

### beaming fraction

$$BF \equiv \frac{\text{illuminated solid angle}}{4\pi} = f(\chi, \rho)$$

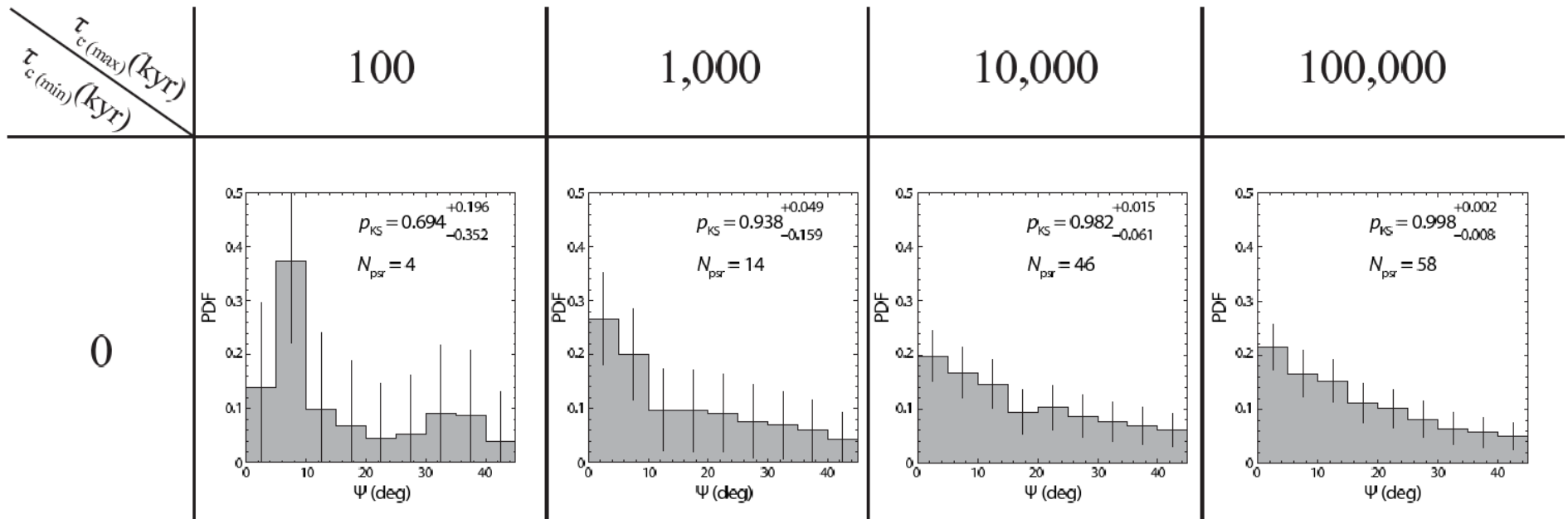
$$\langle BF \rangle = 0.09 \cdot \log^2 \left( \frac{P}{10 \text{ sec}} \right) + 0.03$$



(Tauris & Manchester 1998)

# Pulsars spin-velocity alignment

Several authors (Johnston et al. 2005, Rankin 2007, Noutsos et al. 2013) argued in favour of observational evidence of the **alignment** of the rotational and velocity vectors for isolated pulsars.



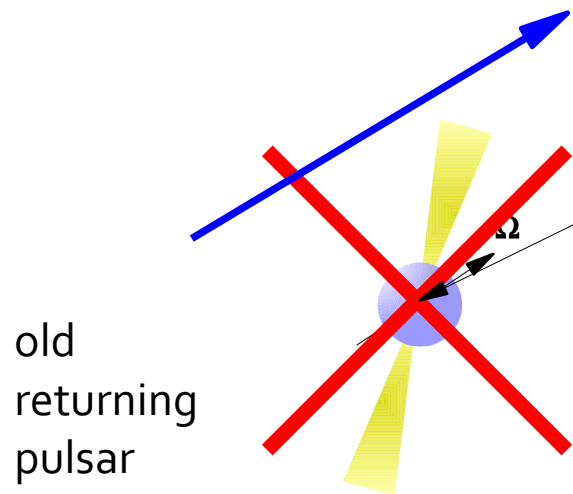
(Noutsos et al. 2013: distributions of spin-velocity angle  $\Psi$  for different age intervals)

In this case angle  $\theta$  between spin axis and direction to an observer is distributed anisotropically.

So, OF is not equivalent to BF any more!

$$\tau_{cross} \sim \frac{W}{2\mu_{tot}P} =$$

$$10^5 \times \left(\frac{W}{30ms}\right) \cdot \left(\frac{P}{1s}\right)^{-1} \cdot \left(\frac{\mu_{tot}}{30mas/yr}\right)^{-1} \text{ years}$$

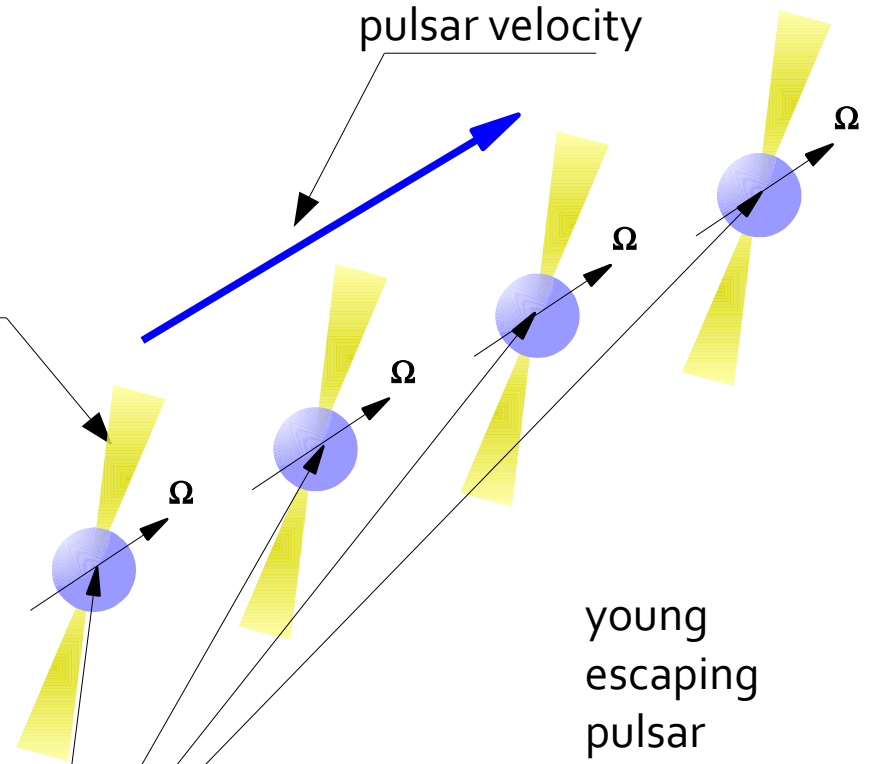


an observer



beam

pulsar velocity



young  
escaping  
pulsar

The line of sight in this situation  
will be mostly directed to south  
magnetic poles of pulsars

# What have we actually done?

beaming excess

$$BE = \frac{OF - BF}{BF}$$

How large can it be?

How is it distributed over P-Pdot diagram?

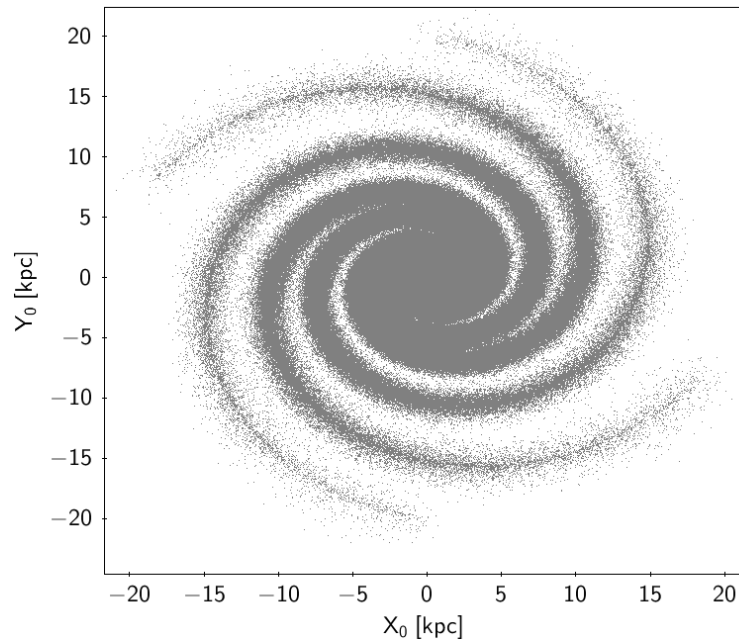
How does it depend on the constituents of pulsars evolution model?

- We have investigated the properties of  $BE$  for a set of models of pulsars evolution using population synthesis approach.
- We have focused on the evolution of spin-velocity and spin-observer angles. Also we have modeled the evolution of magnetic inclination angle  $\chi$  and pulsar beam width  $\rho$ .
- Moreover, we have made sure that pulsar spin axis conserves its direction in space during pulsar life in spite of braking torque affecting a neutron star.

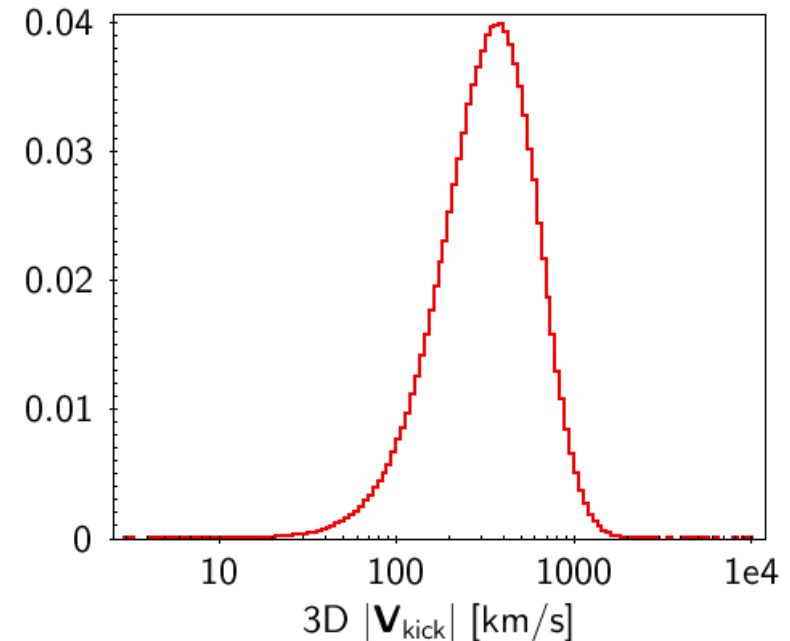
# Model of pulsars ensemble evolution

Our population code (GALP) basically reproduces the best model of Faucher-Giguere & Kaspi (2006):

Galactic structure & gravitational potential



Double-sided exponent  
with  $\langle V_{3D} \rangle \sim 380$  km/s



- Vector of kick velocity is distributed isotropically over the sky
- Initial periods  $P_o$  [sec]  $\sim normal(0.3, 0.15)$
- Initial magnetic fields  $\log B$  [Gs]  $\sim normal(12.55, 0.55)$
- Ages  $t$  [Gyr]  $\sim uniform(0, 1)$
- Approximation of the «death line» in the form  $\dot{P}/P^3 = const$

# Our extension of the model

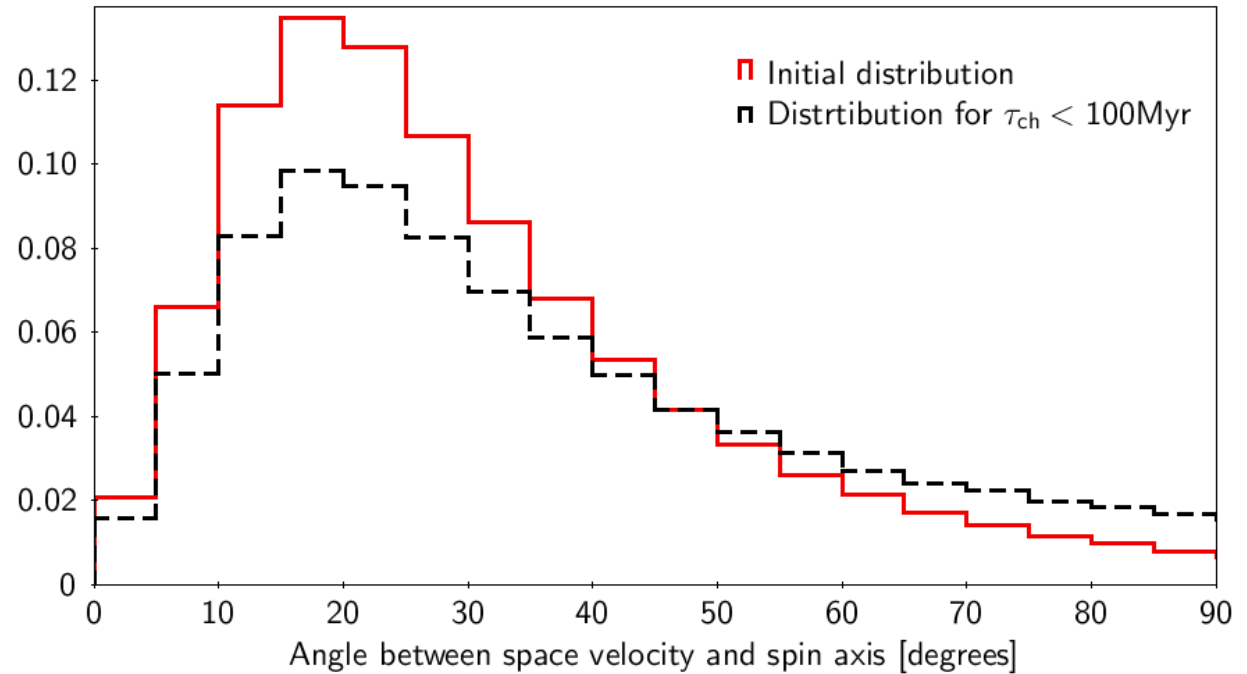
- Spin axis  $\Omega$  is directed precisely along  $\mathbf{V}_{\text{kick}}$
- Initial magnetic angles  $\chi_0$  are distributed isotropically, i.e.  $PDF(\chi_0) = \sin \chi_0/2$
- Pulsar beam has cone-like geometry with opening angle  $2\rho = 0.2/\text{sqrt}(P [\text{sec}])$  radians
- Velocity of the Sun's LSR is 220 km/s
- Spindown law in generalized form:  $\dot{\Omega} = -\alpha B_0^2 \beta_0^2 \cdot \exp\left(-\frac{2t}{T_\Omega}\right) \cdot \Omega^3$
- Magnetic angle evolution in generalized form:  $F(\chi) = F(\chi_0) \cdot \exp\left(-\frac{t}{T_\chi}\right)$

Mod	Description	$\beta_0$	$T_\Omega$	$T_\chi$	$F(\chi)$
A	FG&K 2006 spindown: $\chi = \text{const}, \dot{\Omega} = -\alpha B_0^2 \cdot \Omega^3$	1	<i>inf</i>	<i>inf</i>	$\chi$
B	Magnetic alignment: $\sin \chi = \sin \chi_0 \cdot e^{-t/T_\chi}$ , $\dot{\Omega} = -\alpha B_0^2 \cdot \Omega^3$	1	<i>inf</i>	$10^7$ yr	$\sin \chi$
C	Magnetic disalignment: $\cos \chi = \cos \chi_0 \cdot e^{-t/T_\chi}$ , $\dot{\Omega} = -\alpha B_0^2 \cdot \Omega^3$	1	<i>inf</i>	$10^7$ yr	$\cos \chi$
D	$\chi = \text{const}, B(t) = B_0 \cdot e^{-t/T_\Omega}$ , $\dot{\Omega} = -\alpha B^2 \cdot \Omega^3$	1	$10^7$ yr	<i>inf</i>	$\chi$
E	Classical MD spindown-like law with magnetic alignment	$\sin \chi_0$	$10^7$ yr	$10^7$ yr	$\sin \chi$
F	Electric current spindown-like law with magnetic disalignment (Beskin et al. 1993)	$\cos \chi_0$	$10^7$ yr	$10^7$ yr	$\cos \chi$

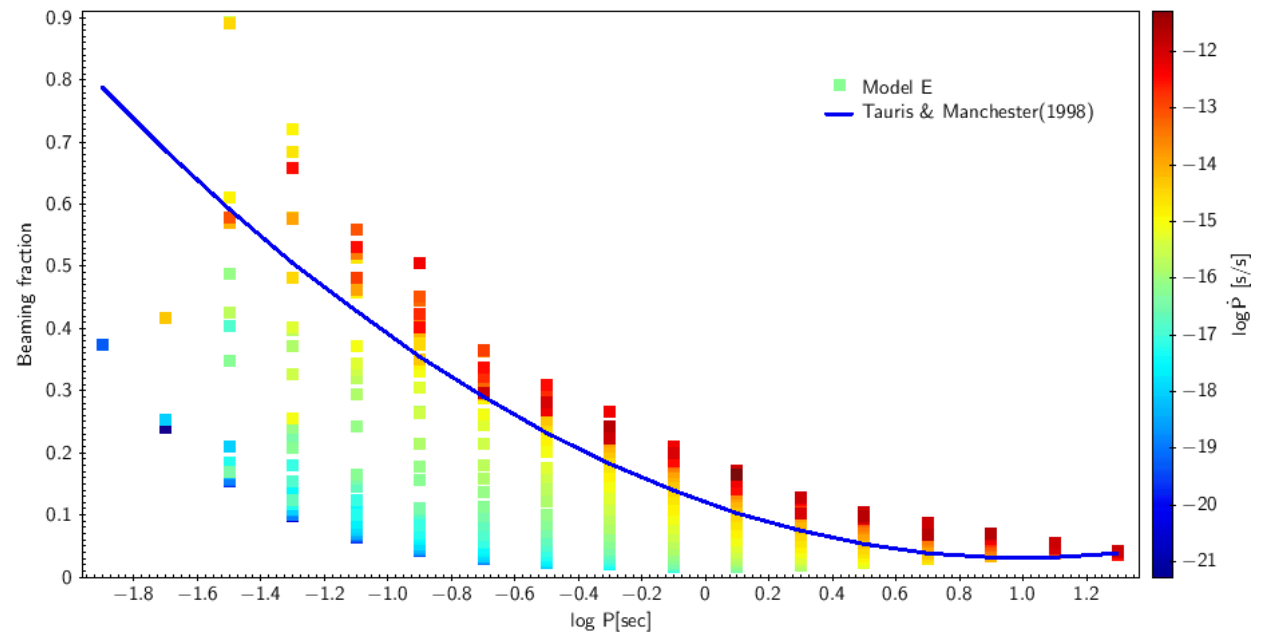


# Results: spin-velocity alignment & average beaming fraction

Orientation of space velocity vector is always partially uncorrelated with that of spin axis due to LSR velocity



We have reproduced the empirical dependence of Tauris & Manchester (1998) for beaming fraction.



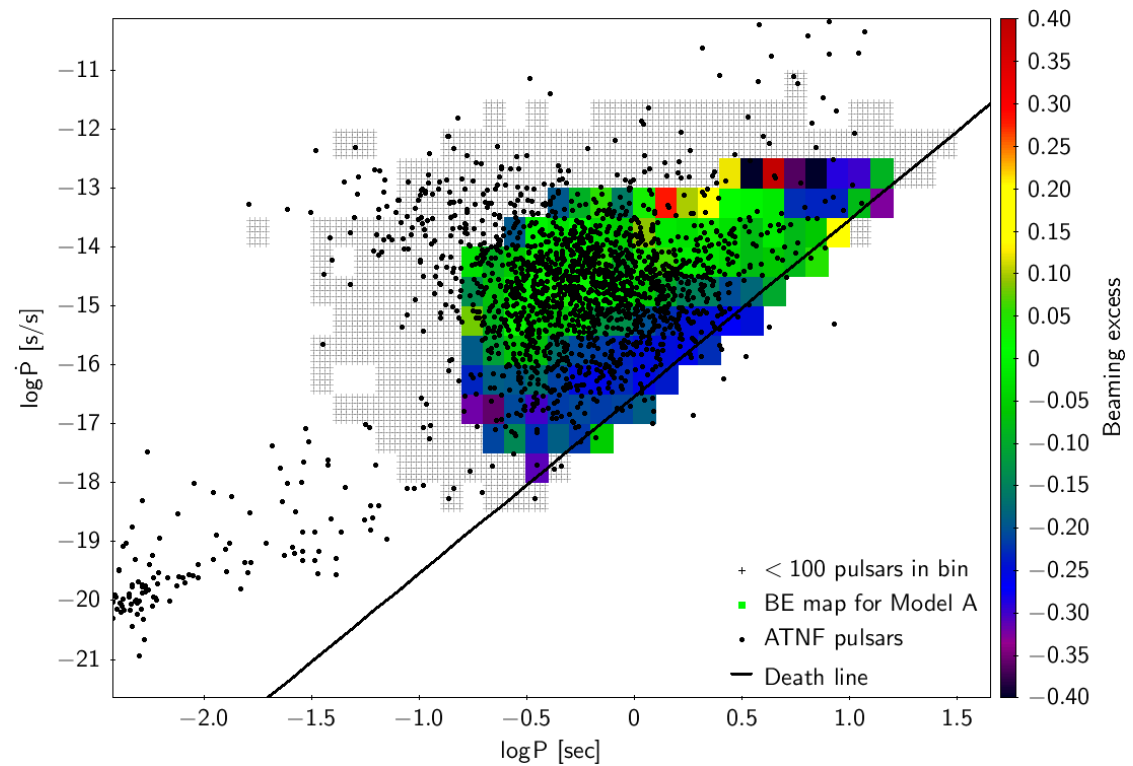
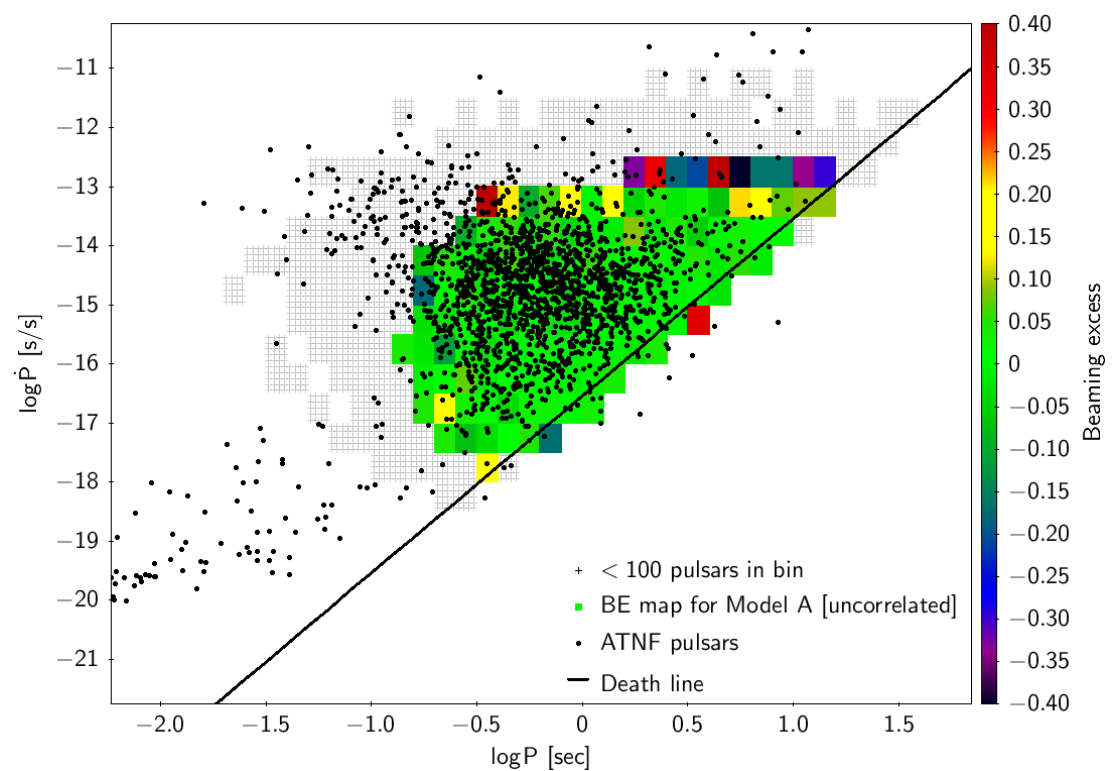
# Results: FG&K model

BE distribution in case when  $\Omega$  is distributed isotropically and independent of the kick direction.

Spin-velocity alignment included into the model leads to negative beaming excess.

Integral (over P-Pdot distribution) value:  
 $-0.173 \pm 0.003$

I.e. the fraction of pulsars directed to an observer is 17.3% less than their average beaming factor.

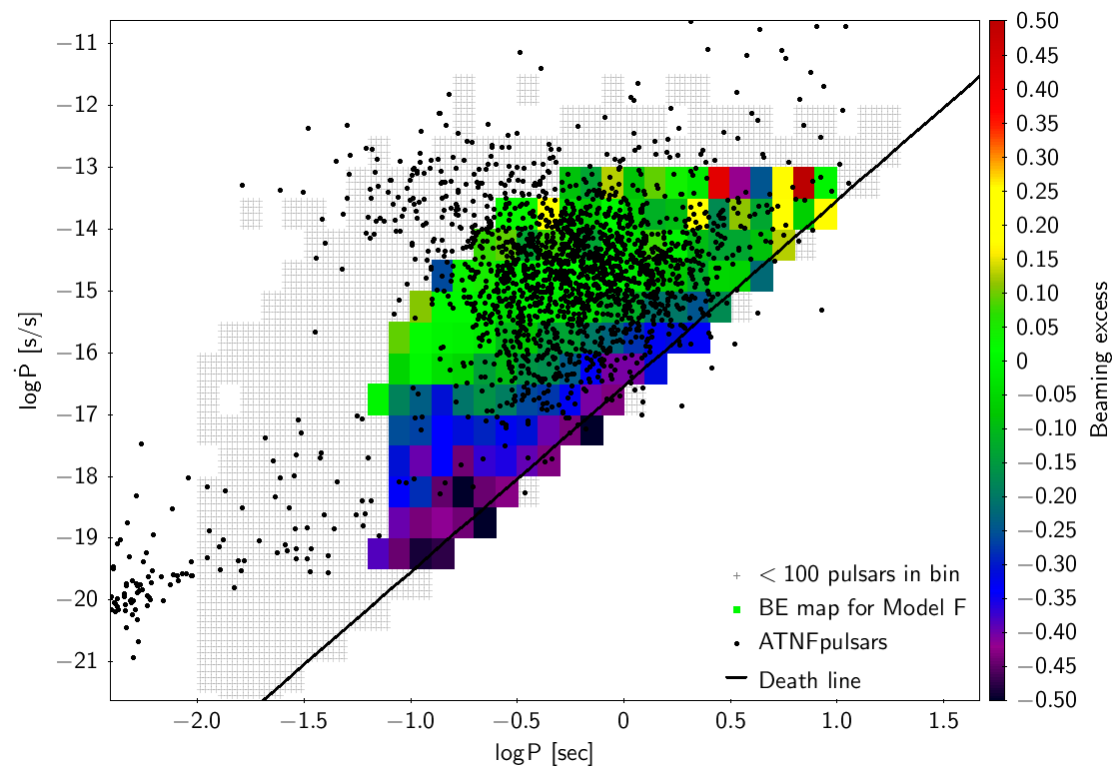
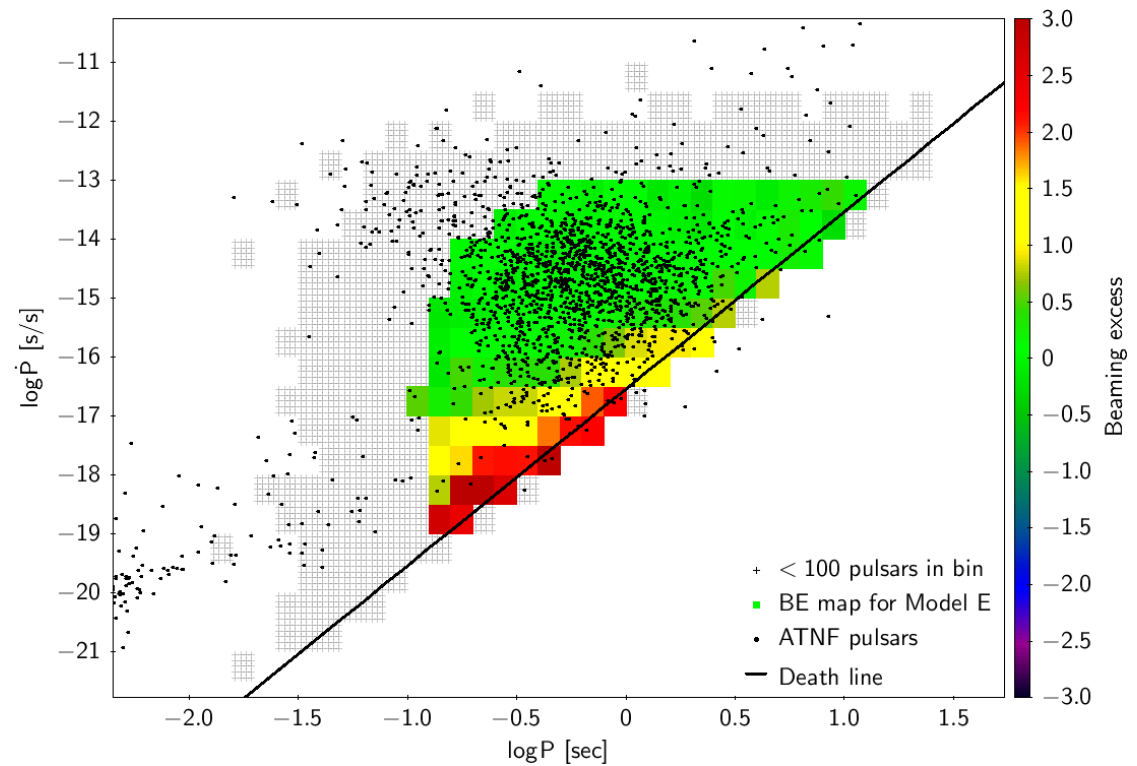


# Results: alignment vs. disalignment

Integral BE for Model E (alignment):  
 $+0.31 \pm 0.01$

Magnetic alignment leads to strong positive BE, while disalignment leads to the strong negative one.

For Model F (disalignment):  
 $-0.199 \pm 0.004$



# Observed birthrates of radiopulsars

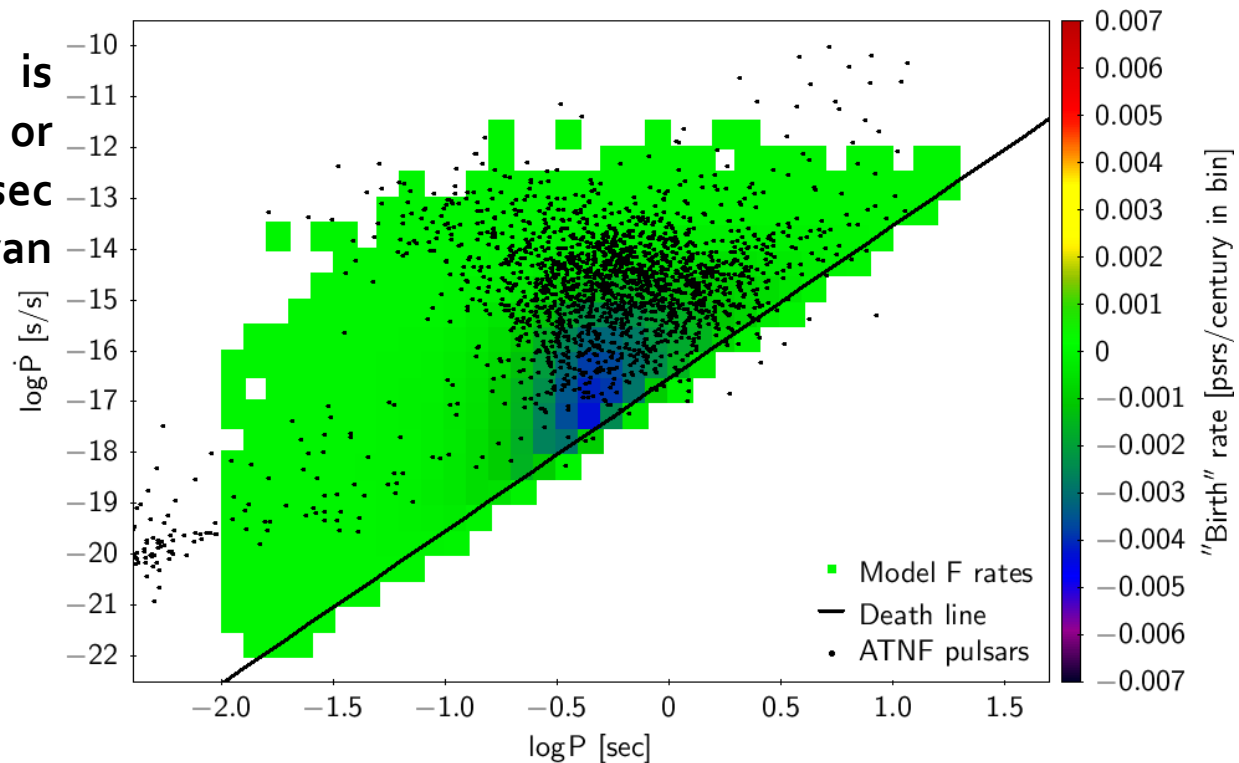
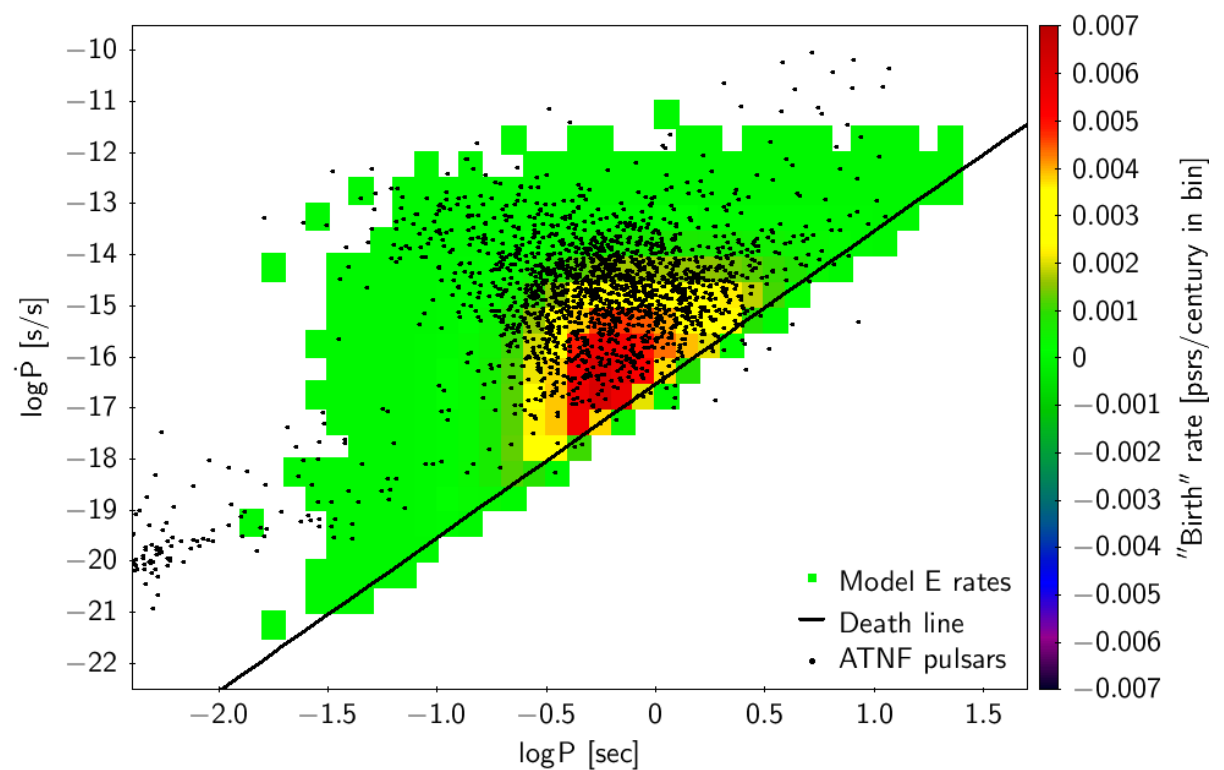
Integral effect for Model E (alignment):

$$(+0.230 \pm 0.001) \times \left( \frac{N_{active}}{10^6 \text{ pulsars}} \right) \left[ \frac{psrs}{century} \right]$$

The effect of beaming excess is accompanied by seeming «birth» or «death» of pulsars near  $P \sim 0.5$  sec (discussed by Vivekanand & Narayan 1981; Vranesevic et al. 2004)

For Model F (disalignment):

$$(-0.407 \pm 0.003) \times \left( \frac{N_{active}}{10^6 \text{ pulsars}} \right) \left[ \frac{psrs}{century} \right]$$



# Results: summary

Mod	Description	Total Beaming Excess	Additional «birth» rate [psrs/century] for $N_{\text{total}} = 10^6$
A	FG-K 2006 spindown: $\chi = \text{const}$ , $\dot{\Omega} = -\alpha B_0^2 \cdot \Omega^3$	$-0.173 \pm 0.003$	$-0.042 \pm 0.001$
B	Magnetic alignment: $\sin \chi = \sin \chi_0 \cdot e^{-t/T_\chi}$ , $\dot{\Omega} = -\alpha B_0^2 \cdot \Omega^3$	$+0.79 \pm 0.01$	$+0.050 \pm 0.001$
C	Magnetic disalignment: $\cos \chi = \cos \chi_0 \cdot e^{-t/T_\chi}$ , $\dot{\Omega} = -\alpha B_0^2 \cdot \Omega^3$	$-0.370 \pm 0.003$	$-0.09 \pm 0.01$
D	$\chi = \text{const}$ , $B(t) = B_0 \cdot e^{-t/T_\Omega}$ , $\dot{\Omega} = -\alpha B^2 \cdot \Omega^3$	$-0.065 \pm 0.005$	$-0.108 \pm 0.002$
E	Classical MD spindown-like law with magnetic alignment	$+0.31 \pm 0.01$	$+0.230 \pm 0.001$
F	Electric current spindown-like law with magnetic disalignment	$-0.199 \pm 0.004$	$-0.407 \pm 0.004$

# Conclusions

- Spin-velocity alignment introduces strong anisotropy into the distribution of spin-observer angles of active pulsars.
- This asymmetry leads to the beaming excess – the difference between the observed fraction and average beaming fraction of a pulsars group. Integral value of BE may reach tens of percents. Thus this effect has to be taken into account within a population synthesis.
- The effect of beaming excess is accompanied by the bias in the observed pulsars birthrate and causes a seeming source of newborn pulsars on the P-Pdot diagram.
- BE strongly depends on the model of pulsar evolution. It probably may help to distinguish observationally magnetic alignment and disalignment of radiopulsars.