Comparison between Simulation and Experiment on Injection Beam Loss and Other Beam Behaviours in the SPring-8 Storage Ring

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1. Motivation

 In ideal top-up operation, reduction of injection beam loss and suppression of stored beam oscillation by frequent beam injections are critically important.

•Aiming at understanding of mechanism of injection beam loss and its suppression, a precise simulation model has been developed.

Key: Precise simulation of particle motion with a large amplitude (x=~±10mm)

2. Simulation Model 2.1. Hamiltonian $H = \hat{p}_{\sigma} - h \cdot \left\{ (1 + \delta)^2 - (\hat{p}_x - \frac{e}{P_0} A_x)^2 - (\hat{p}_y - \frac{e}{P_0} A_y)^2 \right\}^{1/2}$ $-\frac{e}{P_0}A_s$, (1) where $(x, \hat{p}_x \equiv \frac{p_x}{p_0})$, $(x(s) \cdot e_x, y(s) \cdot e_y)$ $\int_{s} (y, \hat{p}_{y} \equiv \frac{p_{y}}{p_{0}}),$ $r_{0(s)}$ $(\sigma \equiv s - v_0 t, \hat{p}_{\sigma} \equiv \frac{E - E_0}{p_0 v_0}), \delta = \frac{p - p_0}{p_0},$ $h \equiv 1 + \frac{x}{\rho_{\rm r}(s)} + \frac{y}{\rho_{\rm v}(s)} \,.$

Refs. [1] D.P. Barber, G. Ripken, F. Schmidt; DESY 87-036 [2] G. Ripken; DESY 87-036 2.2. Components and their Simplectic Integration

(a) Bending Magnet

Equation (1) is integrated without expanding the square root part.

Analytical solutions derived by E. Forest are used for "rectangular" and "sector" magnet integration.

Ref. [3] E. Forest, M.F. Reusch, D.L. Bruhwiler, A. Amiry; Particle Accel.45(1994)65. 2.2. Component Model and Simplectic Integration (2)

(b) Quadrupole and Sextupole Magnets

In this case, Eq. (1) has a separating form of A(p) + V(x).

4th order explicit integration method is adopted. Sextupoles are treated as thick elements.

Refs. [4] E. Forest and R.D. Ruth; Pysica D 43 (1990)105.
[5] H. Yoshida; Celestial Mechanics and Dynamical Astronomy 56 (1993) 27.

2.2. Component Model and Simplectic Integration (3) (c) Other Magnets (except for IDs) All other magnetic components are treated as thin kicks. Multi-pole up to 20 poles available. (d) RF Cavities (Ref. [1]) Cavities are treated as thin elements. $\frac{e}{p_0}A_s = \frac{-L}{2\pi k} \cdot \frac{e}{p_0} \cdot V(s) \cdot \left[\cos\left(\frac{2\pi k}{L} \cdot \sigma + \varphi\right) + G \cdot \frac{2\pi k}{L} \cdot \sigma \cdot \sin\varphi\right], (2)$ G = 1 without Radiation, G = 0 with radiation.

2.2. Component Model and Simplectic Integration (4)

(e) Radiation Effects (1970).

Ref. [6] M. Sands; Report SLAC-121

Normalized photon energy spectrum + random number ranging from 0 to 1 Radiation from BMs, QMs, SMs and IDs are considered.

(f) Fringe Fields

Lowest order effect is only considered for BMs, QMs and SMs. Bivis, Givis and Orvis. $\frac{1}{2} \cdot y^2 \cdot p_x$ (Ref BM fringe -> $H_{\pm} = \pm K_x \cdot \frac{\frac{1}{2} \cdot y^2 \cdot p_x}{\sqrt{(1+\delta)^2 - p_x^2 - p_y^2}}, K_x = \frac{1}{\rho}$. (3) (Ref. [3])

2.3. Magnetic Errors

(a) Normal and Skew Quadrupole Errors [7] 236 Q-error kicks and 132 SQ-error kicks are considered. These strengths are estimated by fitting measured beam response with 4x4 formalism. (b) Multipole errors Systematic errors lower than 20 poles are considered. 10pole for BM and 12pole for QM.

Ref. [7] J. Safranek; NIMA 388 (1997)27.

2.4. ID Model

$$\frac{e}{p_0}A_x = \frac{\rho_0^{-1}}{kz} \cdot \cos(kx \cdot x) \times \cosh(ky \cdot y) \times \sin(kz \cdot z)$$
$$-\hat{\rho}_0^{-1} \cdot \frac{\hat{k}y}{kz \cdot \hat{k}x} \cdot \sinh(\hat{k}x \cdot x) \times \sin(\hat{k}y \cdot y) \times \sin(kz \cdot z + \varphi)$$
$$\frac{e}{p_0}A_y = \rho_0^{-1} \frac{kx}{kz \cdot ky} \cdot \sin(kx \cdot x) \times \sinh(ky \cdot y) \times \sin(kz \cdot z)$$
(4)

$$-\frac{\widehat{\rho_0}^1}{kz} \cdot \cosh\left(\widehat{kx} \cdot x\right) \times \cosh\left(\widehat{ky} \cdot y\right) \times \sin\left(kz \cdot z + \varphi\right)$$

For ID integration, the square root of Eq.(1) is expanded and the lowest order parts are only integrated by means of a generating function.

Refs. [8] K. Halbach; NIM 187(1981)109. [9] E. Forest and K. Ohmi; KEK Report 92-14 (1992).

3. Comparison between Simulation and Experiments

3.1. Amplitude Dependent Tune Shift



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3.2. Smear of Coherent Oscillation



3.3. Nonlinear Momentum Compaction



Ref. [12] H. Tanaka et.al.; NIMA 431(1999)396.

3.4. Nonlinear Dispersion Ref. [12] $\eta(\delta) = \eta_1 + \eta_2 \delta + \eta_3 \delta^2 + \eta_4 \delta^3 + \eta_5 \delta^4 \cdots$



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3.5. 3D Closed Orbit



3.6. Injection Beam Loss 3.6.1. Parameters of Injection Beam and Injection scheme

Injection Beam Parameters

 ϵ_x =220 nmrad, (ϵ_y/ϵ_x) =0.002 σ_{δ} =0.0013, σ_{τ} =63 psec Optical Functions at IP β_x =~13.5m, α_x =-0.11 β_y =~13.7m, α_y =-0.20

Injection Bump System

Injection Point: Δx =-24.5mm Bump height=~14.5mm Bump Pulse Width=~8usec (Revolution=~4.7usec)

Aperture Limit

Septum Inner wall: -19.5mm Chamber Aperture 70(H) / 40(V) mm In-vacuum ID Vert. Lim.: min.7mm

Ring Parameters

 $\varepsilon_x = 2.8 \sim 6.6 \text{ nmrad}, \text{ Coupling}(\varepsilon_v / \varepsilon_x) = 0.002$ $\sigma_8 = 0.0011, \sigma_7 = 13 \text{ psec}$ v_x=40.15, v_v=18.35 **Optical Functions at IP** $\beta_x = -23m, \alpha_x = -0.2$ $\beta_v = ~7.5 \text{m}, \alpha_v = -0.45$ 1.532 Supersion Function 24 and β_{y} [m] 16 0.5 8 0 0.5 ਤੁ -16

3.6.2. ID Parameters(1)

#	Min.Gap	Kx	Ky	Power	Туре
	[mm]			[kW]	
19	12		1.8	35	In-Vacuum Long
20	7		2.1	14	In-Vacuum Hybrid
9	9.6		2.1	8.8	In-Vacuum Standard
10	9.6		2.2	9.6	In-Vacuum Standard
11	9.6		2.2	9.9	In-Vacuum Standard
12	9.6		2.2	9.6	In-Vacuum Standard
13	9.6		2.2	9.8	In-Vacuum Standard
16	12.9		2.4	7.4	In-Vacuum
22	9.98		2.9	12.3	In-Vacuum
24	9.6	1.3	1.4	4	In-Vacuum Figure-8
29	8.8		2.4	11.2	In-Vacuum Standard
35	9		2.3	10.8	In-Vacuum Standard
37	8		2.6	13.3	In-Vacuum Standard
39	8.6		2.3	10.9	In-Vacuum Standard
40	8.3	1.1	1.0	3.5	In-Vacuum Helical
41	9.6		2.0	8.1	In-Vacuum Standard
44	9		2.3	10.9	In-Vacuum Standard
45	8	1.7		1.4	In-Vacuum Vertical tanden
46	8		1.5	8.1	In-Vacuum Hybrid
47	9.6		2.1	9.1	In-Vacuum Standard

3.6.2. ID Parameters(2)

#	Min.Gap	Kx	Ky	Power	Туре
	[mm]			[kW]	
8	25.5	1.1	11.2	18.04	Out-Vacuum Elliptical Wiggler
15	20.0		2.2	5.4	Out-Vacuum Revolver (Linear)
	20.0	3.4	3.4	5.6	Out-Vacuum Revolver (Helical)
25	30.0	4.8	4.6	2	Out-Vacuum Helical tandem
27	37.0	3.9	5.4	6.6	Out-Vacuum Figure-8
23	36.0	3.4	3.4		Out-Vacuum APPLE-II (Circular)
	36.0		5.8		Out-Vacuum APPLE-II (Linear)
	36.0	4.2			Out-Vacuum APPLE-II (Vertical)
17*	20.0				Installed but under commissioning

* ID 17 was recently installed in the ring. Gap of ID 17 is usually closed. The phase is shifted so that the magnetic field is cancelled out at the center. The field is no zero and nonlinear at the off-center.

3.6.3. Behaviour of Injection Beam (1-Calc)



3.6.3. Behaviour of Injection Beam (2-Calc)



3.6.4. Beam Loss v.s. Gap of ID47(1)



3.6.4. Beam Loss v.s. Gap of ID47(2)



3.6.5. Beam Loss v.s. Symmetry restoration of Optics



3.6.6. Beam Loss v.s.Initial Condition (Calc)





Lost particles have clear correlation with horizontal emittance .

3.6.7. Beam Loss v.s. Beam Collimation (1)



Planner type IDs are only considered in the simulation.

3.6.7. Beam Loss v.s. Beam Collimation (2)



Planner type IDs are only considered in the simulation.

4. Summary(1)

•Developed simulator well describes nonlinear particle motion in a storage ring. However, the simulator can not explain injection beam loss quantitatively.

•Possible causes are:

- •Ambiguity in injection beam distribution
- Incorrect ID Modeling at especially large amplitude
- •Nonlinearity not included in the model
- •Some mistake in calculation, etc

4. Summary(2)

•Frequency Map Analysis (FMA) could be a powerful tool to investigate mechanism of injection beam loss.

•To use FMA effectively, we need " a precise ring model".

5. Tentative FMA Results Tune modulation in damping of injection beam



FMA of Normal Optics w/o Errors Chromaticity (+8, +8)



Stability Map of Normal Optics w/o Errors Chromaticity (+8, +8)



FMA of Normal Optics with Errors Chromaticity (+8, +8)



Stability Map of Normal Optics with Errors Chromaticity (+8, +8)

