



PECULIARITIES OF DARK CONDUCTIVITY IN IRRADIATED SILICON

J. Vaitkus, A. Mekys, V.Rumbauskas, J. Storasta
*Institute of Applied Research & Department of
Semiconductor Physics, Vilnius University*

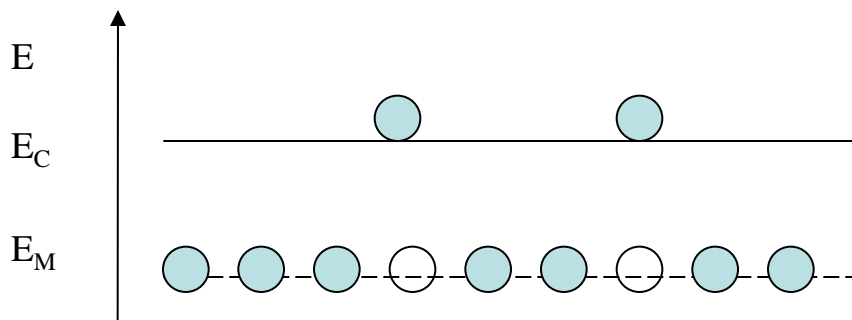


Purpose:

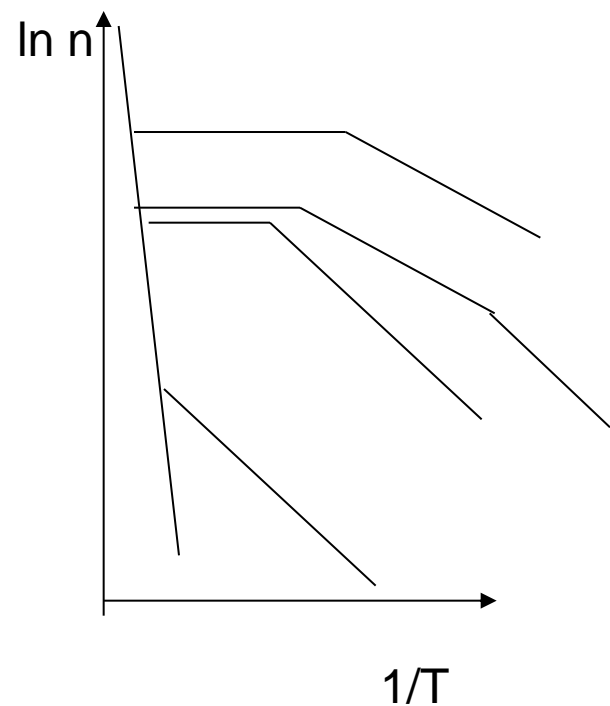
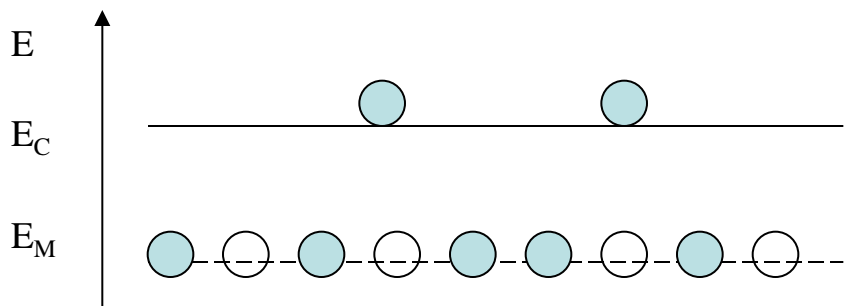
- To point out the main deep levels contributing to the dark conductivity in the irradiated silicon
- To pay attention to the problems seen in the analysis
- To propose a model that can explain the observed non-classical behavior.

Classics:

- $n = N_M / \{1 + \sqrt{[1 + (4\beta N_M / N_C) \exp(\Delta E_M / kT)]}\}$.

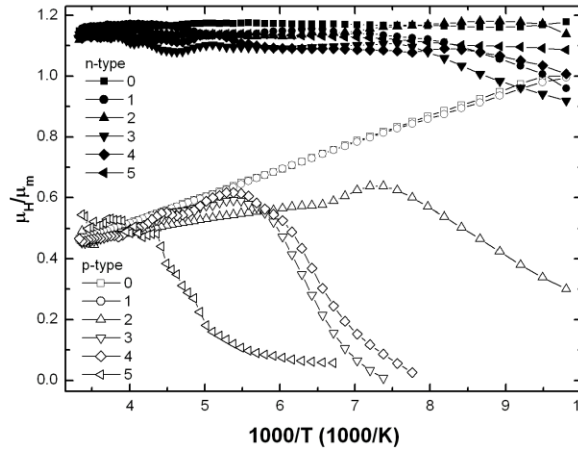
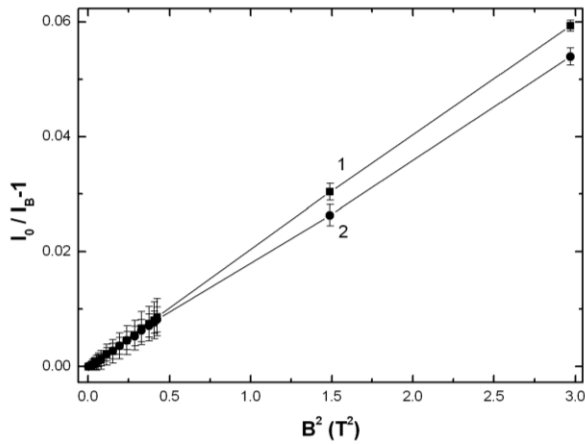


- $n = (N_M - N_K) / [1 + (\beta N_K / N_C) \exp(\Delta E_M / kT)]$
 - $n = \beta N_C [(N_M - N_K) / N_K] \exp(-\Delta E_M / kT)$,



Mobility measurement results

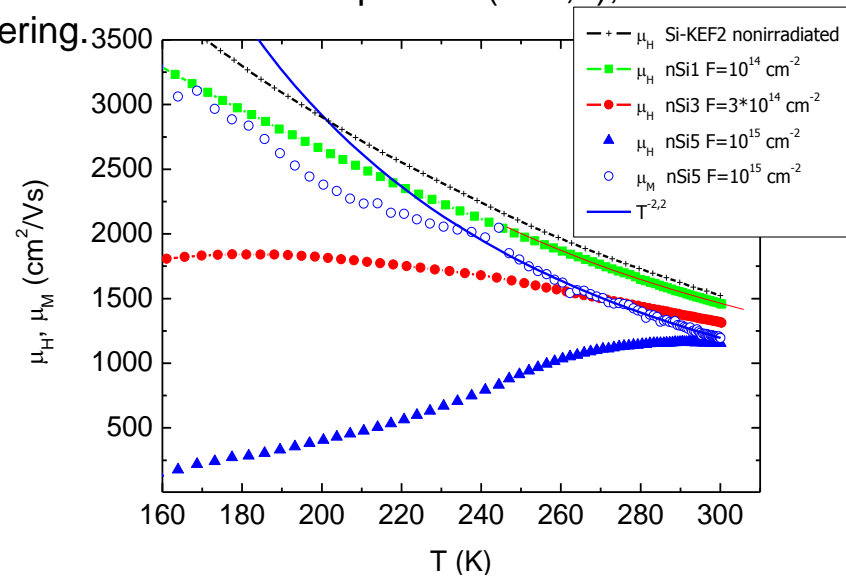
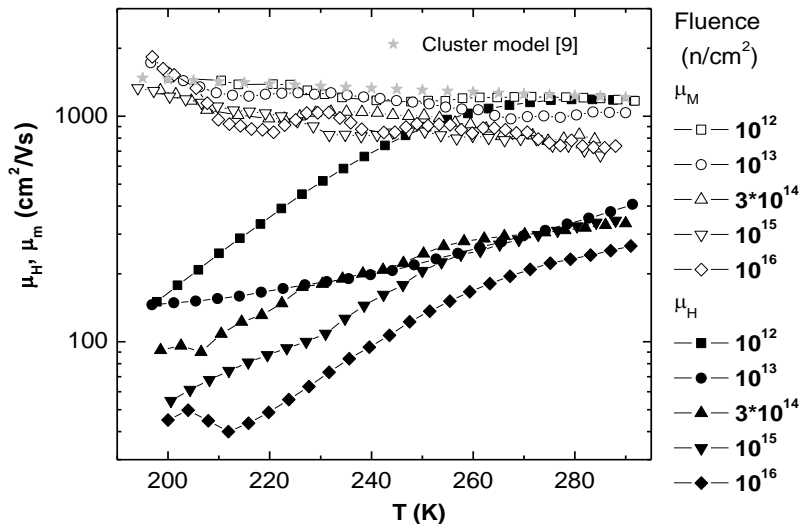
$$\mu_H = \frac{V_H}{BV_X} \quad \mu_M = \frac{1}{B} \sqrt{\frac{\rho_B - \rho_0}{\rho_0}}$$



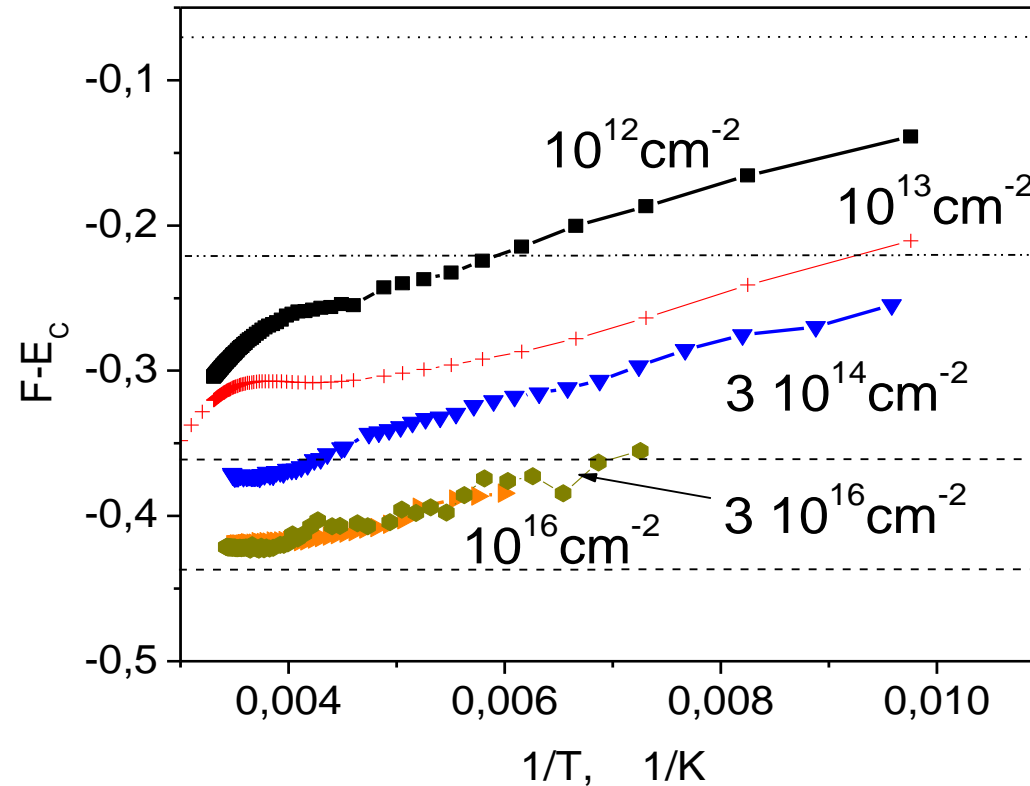
1,2,3,4,5 times 10^{16} e/cm².

Details will be published as a joint paper with L.Makarenko group (Minsk), and it was presented at earlier meeting of RD50

At $\sim T_{\text{Room}}$ the electron μ_H and μ_M are $\sim T^{-\alpha}$, and α is smaller than in a pure Si ($\alpha=2,4$), and it is related to the contribution of impurity (and the cluster) scattering.

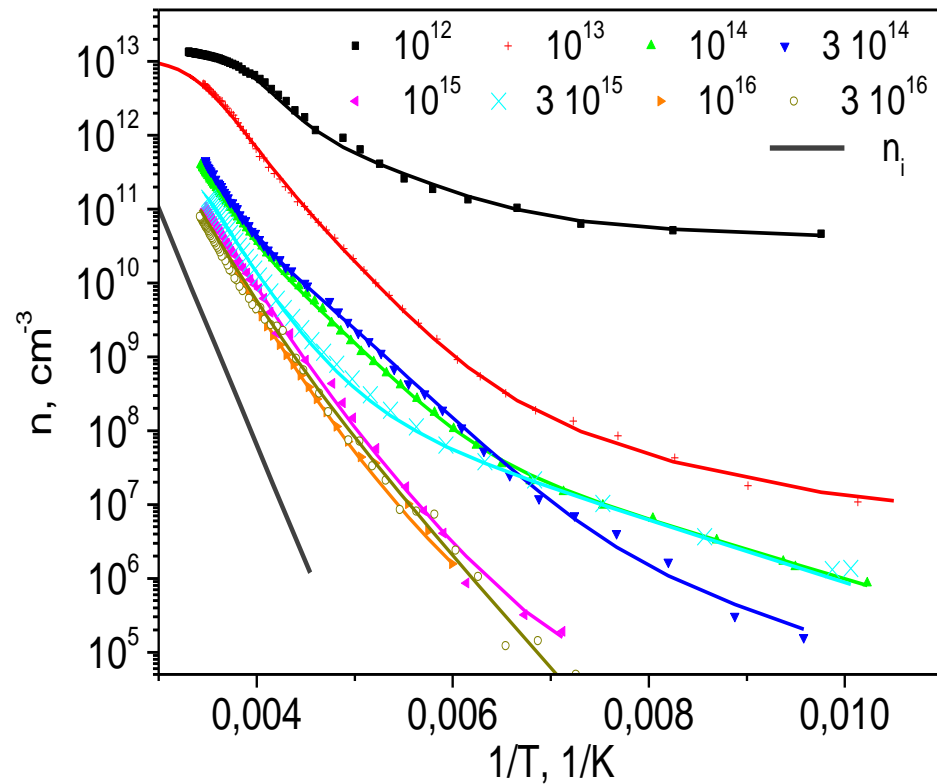


Fermi level = $j(1/T)$ in irradiated Si



- According to this dependence it is possible to think about the model that might fit to the experiment

$\lg N = f(1/T)$ in the irradiated Si



	Deep level 0,40 eV	Deep level 0,36 eV	Deep level 0,22 eV (0.23 eV)	Shallow level 70 meV	Shallow level 46 meV
Fluence 1e12 cm ⁻²					
$(N_M - N_K),$ cm ⁻³		1.3e13	1.3e12 (8e11)	5e10	
βN_M		1.2e13 1	2e13 (4e12)		
Fluence 1e13 cm ⁻²					
$(N_M - N_K),$ cm ⁻³		1.2e13 (1.2e13)	1e11 (1e11)	2e8 (2e8)	3e6 (3e6)
βN_M		8e12 (9e13)	2e13 (1e13)	2e15 (2e15)	
Fluence 1e14 cm ⁻²					
$(N_M - N_K),$ cm ⁻³		1e14 (1e14)	1e10 (1e10)	1e8	$\leq 4e5$
βN_M		2.4e14 (2.4e14)	2e13 (1e13)	1e16	
Fluence 3e14 cm ⁻²					
$(N_M - N_K),$ cm ⁻³	5e12 (6e12)		2e10 (3e10)	3e7 (3e7)	
βN_M	6e12 (2e12)		2e13 (2e13)	2e16 (2e16)	

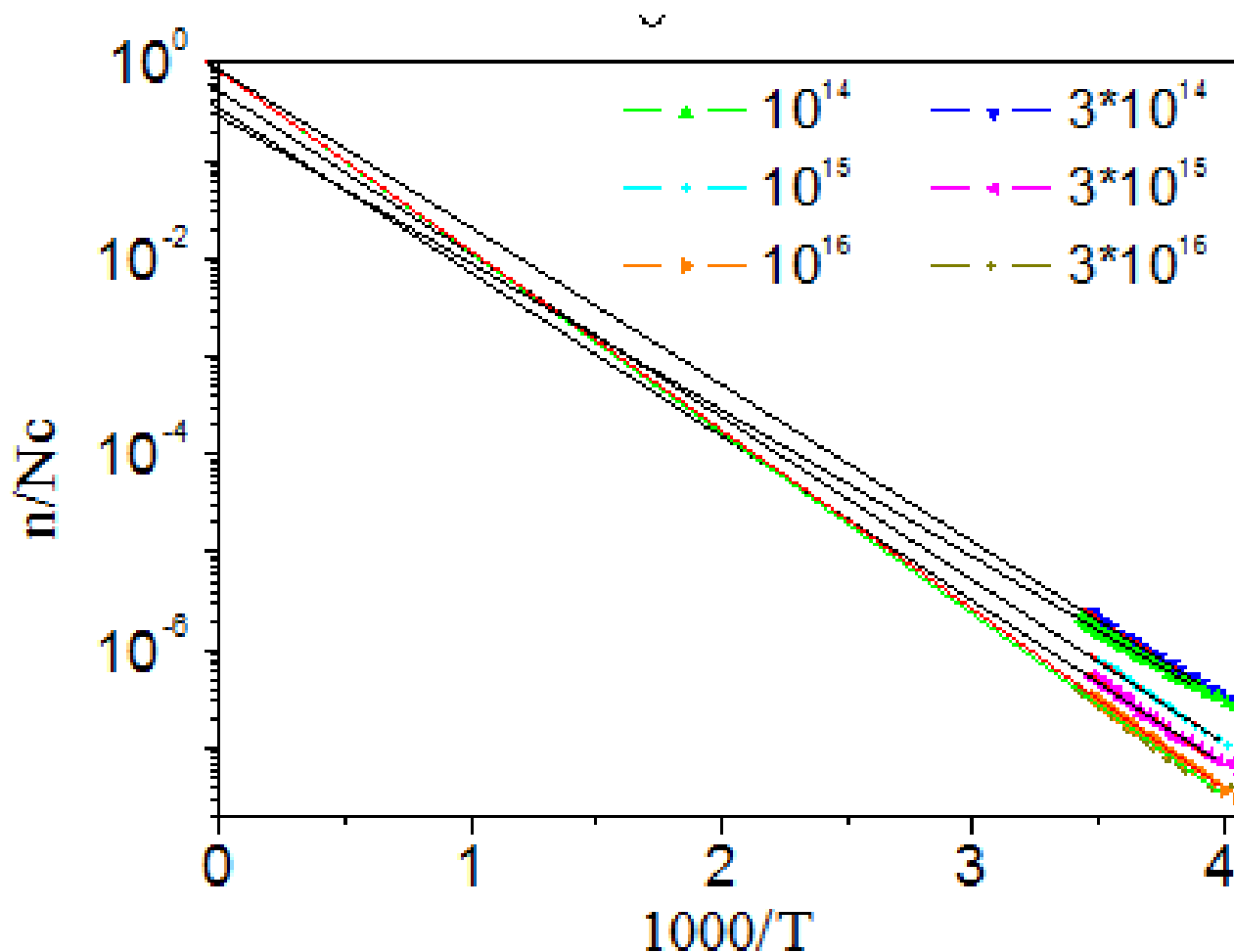
V_2 and V_3 contribution is enough to fit the experimental data in the samples irradiated to 1e12-1e14 cm⁻²



Level activation energy	0,44 eV	0.43 eV	0.41 eV	0.28 eV	0.27 eV	0,22 eV	70 meV
Fluence 1e15 cm ⁻²							
$(N_M - N_K), \text{cm}^{-3}$			3.2e12			3e10	3e8
βN_M			2e12			3.3e13	3e16
Fluence 3e15 cm ⁻²							
$(N_M - N_K), \text{cm}^{-3}$		5e12			3e11		2e6
βN_M		2e12			5e14		3e16
Fluence 1e16 cm ⁻²							
$(N_M - N_K), \text{cm}^{-3}$	4e12			2e13			3e6
βN_M	2e12			4e16			3e16
Fluence 3e16 cm ⁻²							
$(N_M - N_K), \text{cm}^{-3}$	4.2e12			3e11			
βN_M	2e12			3.3e14			

The main problem: why activation energy depends on the irradiation fluence???

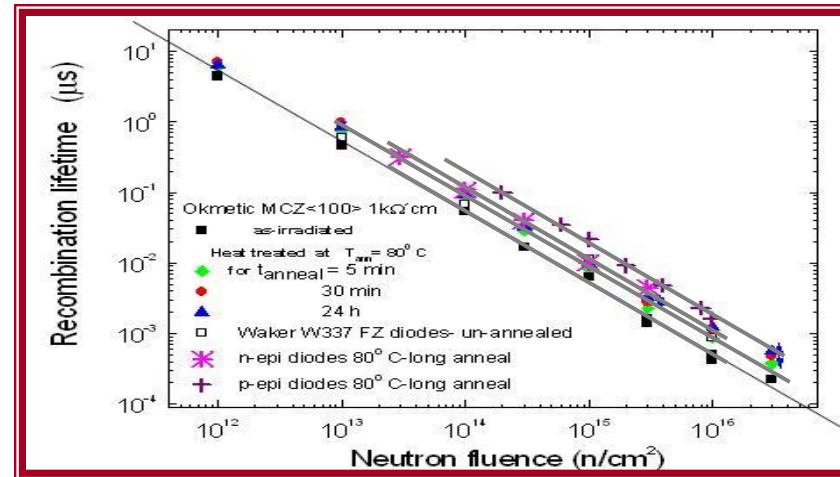
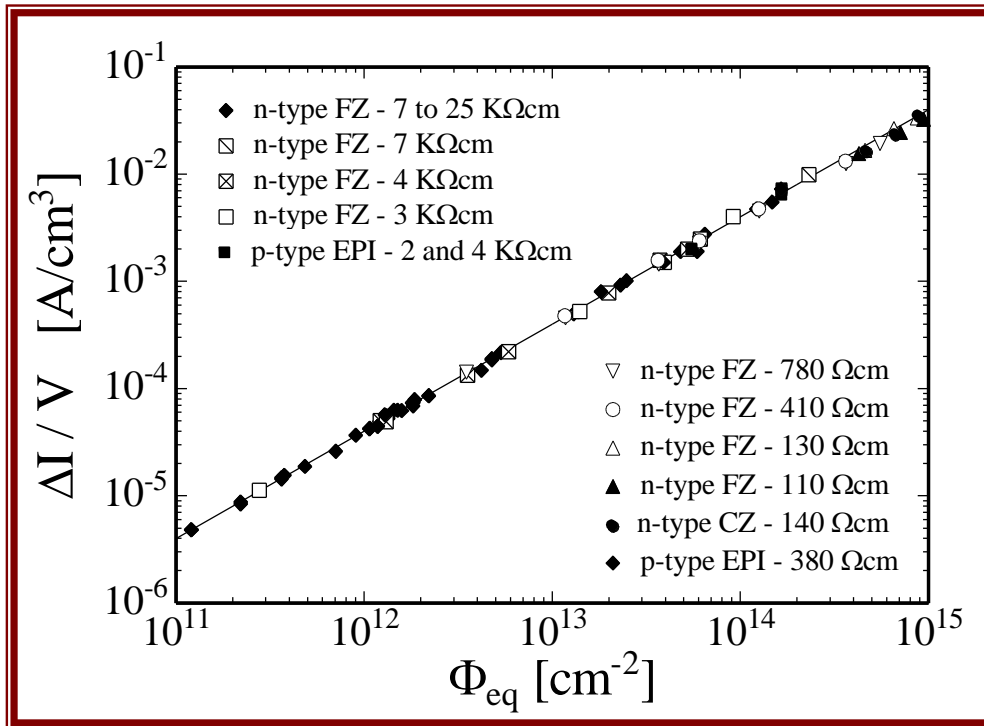
$\lg N = f(1/T)$ in the irradiated Si



A fan type dependencies.

Earlier investigations show: the increase of irradiation fluence does not introduce evident nonlinearity:

The reverse current [Lindstrom G., Moll M., Fretwurst E. Radiation hardness of silicon detectors - a challenge from high-energy physics. NIMA 426, 1999, 1-13 p.] and free carrier lifetime [E. Gaubas, T. Čeponis, A. Uleckas, J. Vaitkus. Anneal dependent variations of recombination and generation lifetime in neutron irradiated MCZ Si. NIMA 612, 2010, 563-565 p.] linearly depend on the fluence in a wide range of fluence.



$$\frac{1}{\tau} = \gamma P_R = aF$$



A summary of non-classics:

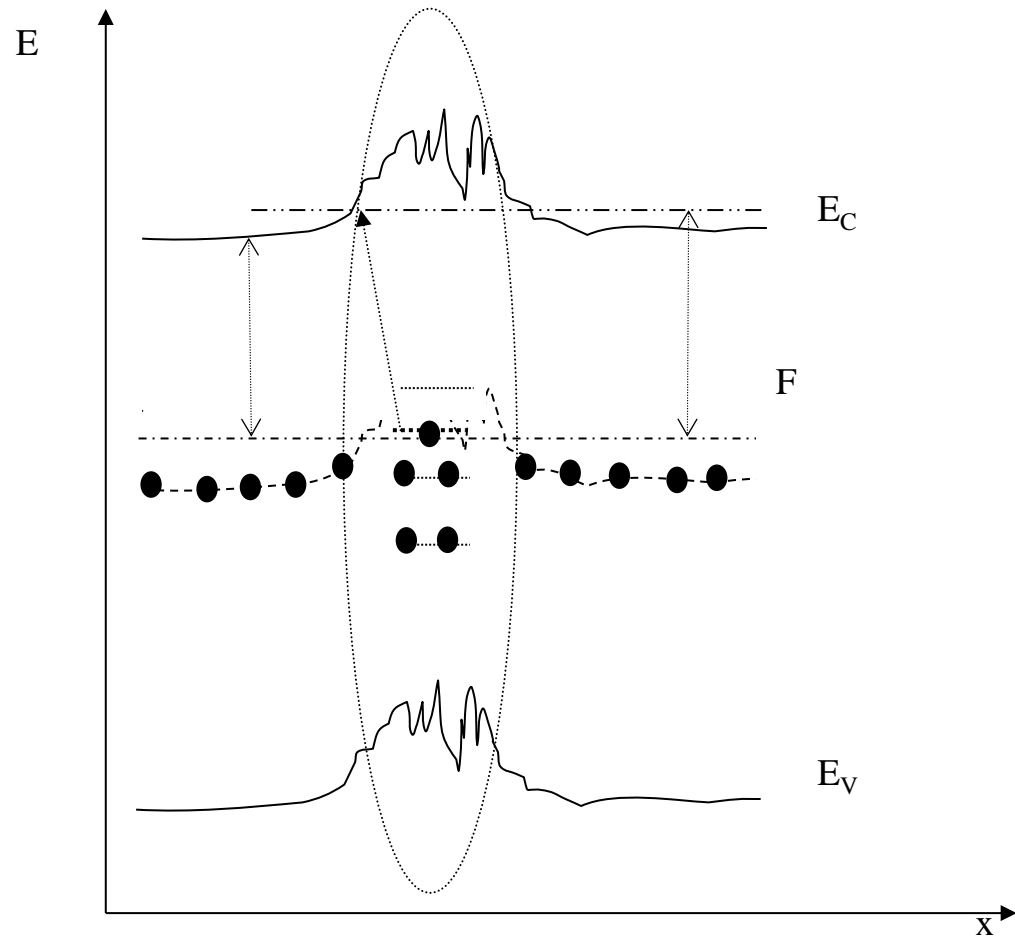
- The deepest level responsible for the dark conductivity is partly filled by electrons;
- The activation energy depends on the compensation of conductivity (that directly related to the fluence).

A summary of predicted cluster properties:

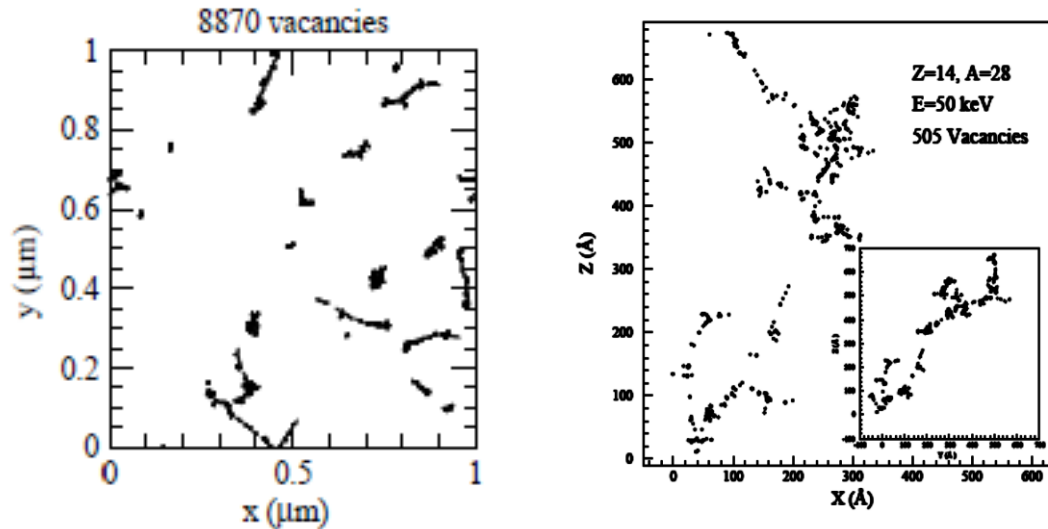
- The cluster has a randomized structure including various complexes of vacancies;
- The crystal structure around the cluster is random and it can introduce the disordered region properties.

A model of cluster following from the dark conductivity measurement:

- The cluster barrier causes a partly filling the centers inside the cluster or at its boundary.
- The compensation increase (an increase of deep acceptor concentration) the Debye length therefore causes the increase of the energy required to excite the electron from the deep level.



Modeling of defect generation



Vacancies generated by 1 MeV neutrons.

Fluence: 10^{14} cm^{-2} . [M.Huhtinen. NIMA 491 (2002) 194–215]

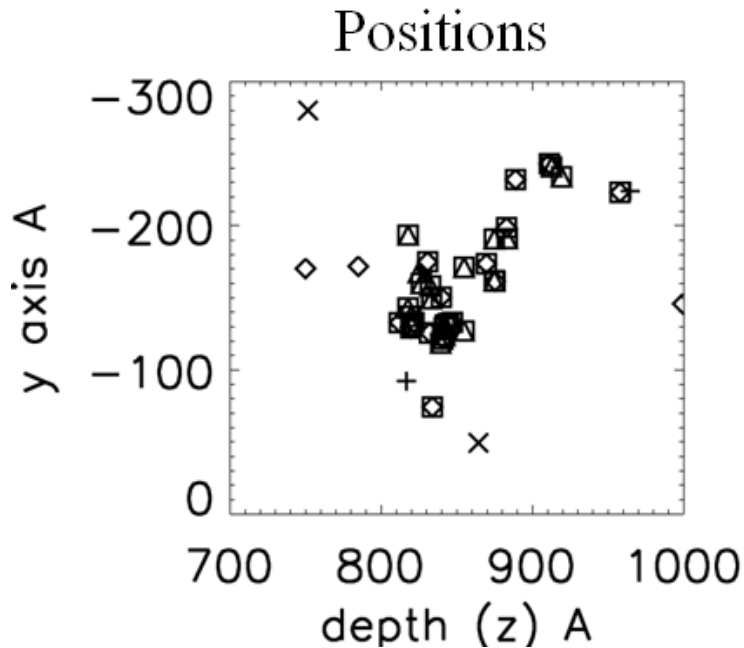
General points: TRIM calculation. A 50 keV ion creates a cluster about 100 nm length in a time of 20 fs. (one vibrational period). The energy deposited gives a high temperature (up to $\sim 1000 \text{ K}$). If $[O] = 10^{18} \text{ cm}^{-3}$, may have 1 O atom in the cluster.

Following the WODEAN (Bucharest) 1.

After random walk (Time nano- to milli-seconds)

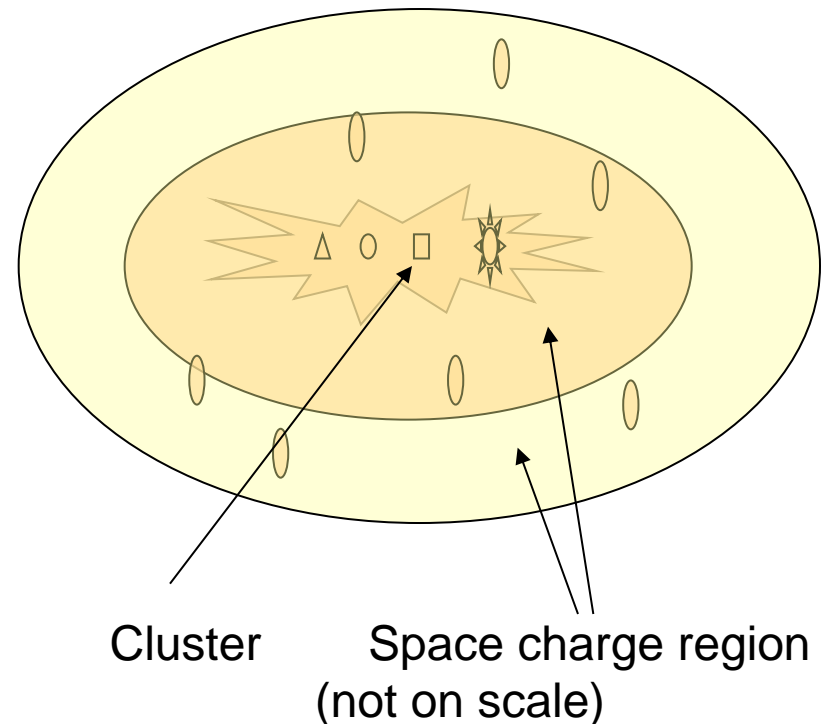
$I (X)$ $V (+)$ $I_2 (\blacklozenge)$ $V_2 (\Delta)$

I-clusters (\blacklozenge) and V-clusters (Δ) in squares



Reconstruction of the cluster in a cube.
G.Davies, RD50 WODEAN conference,
Bucharest:

Approximation:



Cluster

Space charge region
(not on scale)

Conclusions:

- The main impurities responsible for the dark conductivity are similar to bi- and tri- Si vacancies (V_2 and V_3).
- At T near to room T the conductivity is caused by the deep levels in the cluster or its environment, and they are partly filled in the highly irradiated Si.
- The increase of the free carrier activation energy in the highly irradiated samples allows to propose the quenching of thermal generation of free carriers due to the compensation of conductivity.
- The proton irradiation could allow to test the model. (*Who like to collaborate?*)



Thanks to Lithuanian Academy of Sciences grant

THANK YOU FOR YOUR ATTENTION!

