# Dark Forces and New Physics at the Intensity Frontier

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# 38 years rule = new forces of nature are discovered every 38 years for the last 150 yrs

- 1860s first papers of Maxwell on EM. Light is EM excitation.
   E & M unification.
- 1897 Becquerel discovers radioactivity first evidence of weak charged currents (in retrospect).
- 1935 Chadwick gets NP for his discovery of neutron with subsequent checks that there exists strong n-p interaction. Strong force is established.
- 4. 1973 Gargamelle experiment sees the evidence for weak neutral currents in nu-N scattering
- 5. 2011/2012 Discovery of the Higgs, i.e. new Yukawa force.
- 6. *Prediction: Discovery of a new dark force 2050?*

(+/-2 years or so).

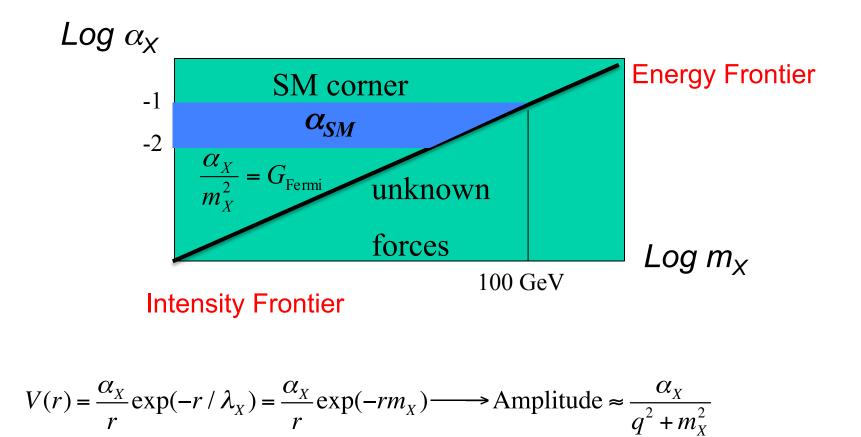
# Outline of the talk

- 1. Energy and Intensity Frontiers. Portals to SM. Implications of the LHC results.
- 2. "Anomalies" and various rationales for dark forces at low energy. Secluded U(1) (= dark photon) model. Possible connection to dark matter. Main features and signatures.
- 3. Selected new results/ideas for secluded sectors:
  - 3a. Fixed target searches of dark photons and light (MeV scale) dark matter
  - 3c. Very very dark photons. Implication for CMB/BBN + constraints from direct detection.

3d. Dark forces and world peace B-modes.

4. Conclusions.

# **Intensity and Energy Frontiers**



LHC can realistically pick up New Physics with  $\alpha_X \sim \alpha_{SM}$ , and  $m_X \sim 1$ TeV, while having no success with  $\alpha_X < 10^{-6}$ , and  $m_X \sim \text{GeV. 4}$ 

# Neutral doors ["portals"] to the SM

Let us *use* these doors, and attach the Dark Matter to the SM  $H^+H(\lambda S^2 + AS)$  Higgs-singlet scalar interactions  $B_{\mu\nu}V_{\mu\nu}$  "Kinetic mixing" with additional U(1)' group (becomes a specific example of  $J_{\mu}^{\ i}A_{\mu}$  extension) neutrino Yukawa coupling, N - RH neutrino LHN  $J_{\mu}^{i}A_{\mu}$  requires gauge invariance and anomaly cancellation It is very likely that the observed neutrino masses indicate that Nature may have used the *LHN* portal...

Dim>4

. . . . . . . . . .

 $J_{\mu}^{A} \partial_{\mu} a / f$  axionic portal

$$\mathcal{L}_{\text{mediation}} = \sum_{k,l,n}^{k+l=n+4} \frac{\mathcal{O}_{\text{med}}^{(k)} \mathcal{O}_{\text{SM}}^{(l)}}{\Lambda^n},$$

### Simplest example of a dark force (Holdom 1986; earlier paper by Okun')

$$\mathcal{L} = -\frac{1}{4}V_{\mu\nu}^2 - \frac{\kappa}{2}V_{\mu\nu}F^{\mu\nu} + |D_{\mu}\phi|^2 - V(\phi),$$

This Lagrangian describes an extra U(1)' group (dark force, hidden photon, secluded gauge boson, shadow boson etc, also known as U-boson, V-boson, A-prime, gamma-prime etc), attached to the SM via a vector portal (kinetic mixing). Mixing angle  $\kappa$  (also known as  $\varepsilon$ ,  $\eta$ ) controls the coupling to the SM. New gauge bosons can be light if the mixing angle is small.

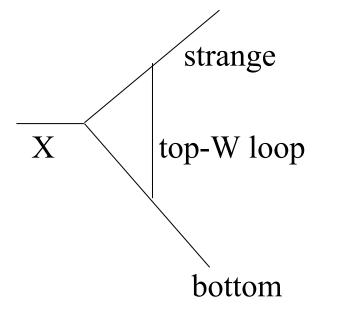
Low-energy content: Additional massive photon-like vector V, and a new light Higgs h', both with small couplings.

Well over 100 theory papers have been written with the use of this model in some form in the last four years.

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# Why EM or baryonic currents are "safe" from flavor constraints

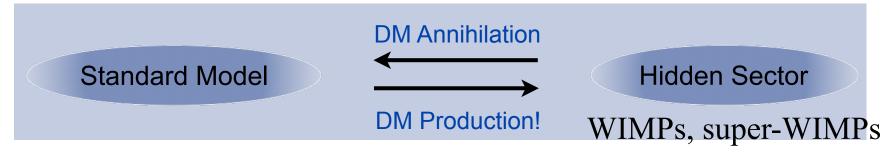
Conserved vector currents are uniquely positioned to avoid very strong flavor constraints. Axial vector portals, Higgs portals are potentially liable to very strong flavor constraints. Consider generic FCNC penguin-type loop correction.



For a conserved vector current,  $G_F q^2$ For axial vector current,  $G_F m_t^2$ 

There is extremely strong sensitivity to new scalars, pseudoscalars axial-vectors in rare K and B decays.

# Possible connection to WIMP-y dark matter



Mediators (SM Z, h etc or dark force)

Heavy WIMP/heavy mediators: - "mainstream" literature Light WIMPs/light mediators: Boehm et al; Fayet; MP, Ritz, Voloshin; Hooper, Zurek; others

Heavy WIMPs/light mediators: Finkbeiner, Weiner; Pospelov, Ritz, Voloshin (secluded DM); Arkani-Hamed et al., many others

Light WIMPs/heavy mediators: does not work. (Except for super-WIMPs; or non-standard thermal history)

### "Non-decoupling" of secluded U(1) Theoretical expectations for masses and mixing

Suppose that the SM particles are not charged under new  $U_s(1)$ , and communicate with it only via extremely heavy particles of mass scale  $\Lambda$  (however heavy!, e.g. 100000 TeV) charged under the SM  $U_{\rm v}(1)$  and  $U_{\rm s}(1)$ (B. Holdom, 1986) Λ  $U_{\rm v}(1)$  $U_{\rm V}(1)$  does not decouple! Diagram A mixing term is induced,  $\kappa F_{\mu\nu}^{\gamma} F_{\mu\nu}^{S}$ , With  $\kappa$  having only the log dependence on mass scale  $\Lambda$  $\kappa \sim (\alpha \alpha')^{1/2} (3\pi)^{-1} \log(\Lambda_{UV}/\Lambda) \sim 10^{-3}$  $M_V \sim e' \kappa M_{FW} (M_Z \text{ or TeV}) \sim \text{MeV} - \text{GeV}$ This is very "realistic" in terms of experimental sensitivity range of parameters.

### Some specific motivations for new states/ new forces below GeV

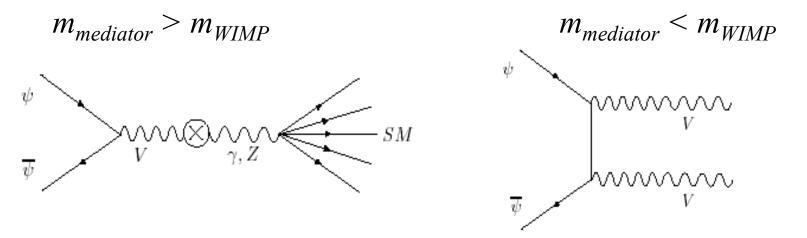
- 1. Theoretical motivation to look for an extra U(1) gauge group.
- 2. Recent intriguing results in astrophysics. 511 keV line, PAMELA positron rise.
- 3. A decade old discrepancy of the muon g-2.
- 4. New discrepancy of the muonic hydrogen Lamb shift.
- 5. Other motivations.

• • • •

### Secluded WIMP idea – heavy WIMPs, light mediators

 $\mathcal{L}_{\text{WIMP+mediator}} = -\frac{1}{4}V_{\mu\nu}^2 - \frac{\kappa}{2}V_{\mu\nu}B_{\mu\nu} - |D_{\mu}\phi|^2 - U(\phi\phi^*) + \bar{\psi}(iD_{\mu}\gamma_{\mu} - m_{\psi})\psi.$ 

 $\psi$  – weak scale Dark Matter; V – mediator particle.

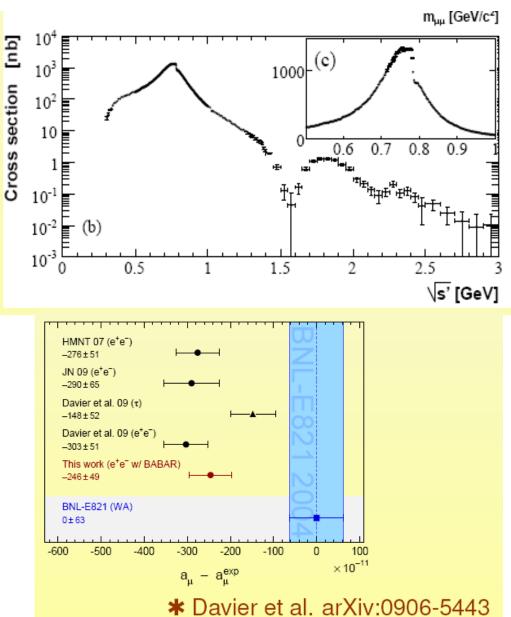


Second regime of annihilation into on-shell mediators (called *secluded*) does not have any restrictions on the size of mixing angle  $\kappa$ .

It turns out *this helps* to tie PAMELA positron rise and WIMP idea together.

# g-2 of muon

BaBar contribution to the "hadronic piece" of VP diagram



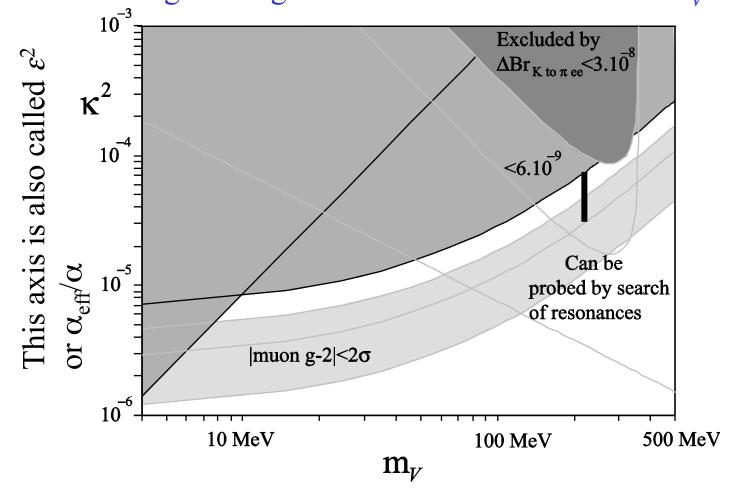
More than 3 sigma discrepancy for most of the analyses. Possibly a sign of new physics, but some complicated strong interaction dynamics could still be at play.

Supersymmetric models with large-ish  $tan\beta$ ; light-ish sleptons, and right sign of  $\mu$  parameter can account for the discrepancy.

Sub-GeV scale vectors can also be at play. 12

### $\kappa$ - $m_V$ parameter space

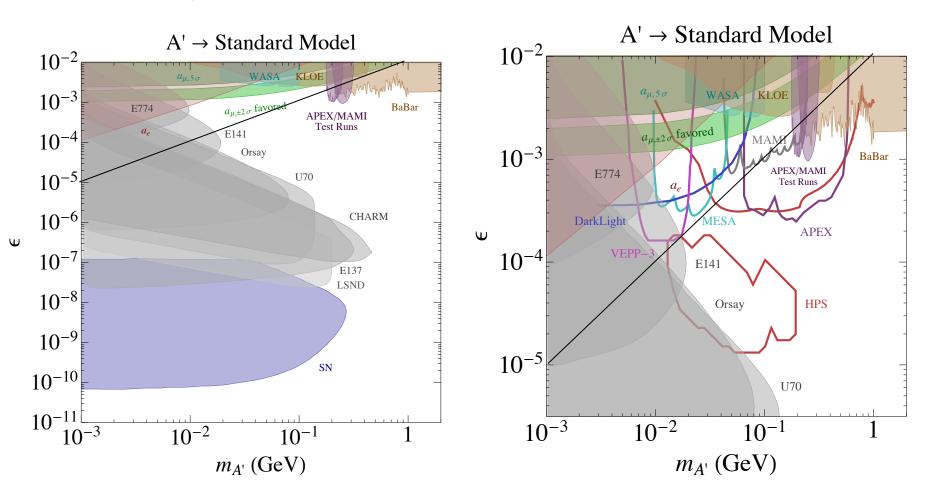
If g-2 discrepancy taken seriously, a new vector force can account for deficit. (Krasnikov, Gninenko; Fayet; Pospelov) E.g. mixing of order few 0.001 and mass  $m_V \sim m_\mu$ 



Since 2008 a lot more of parameter space got constrained

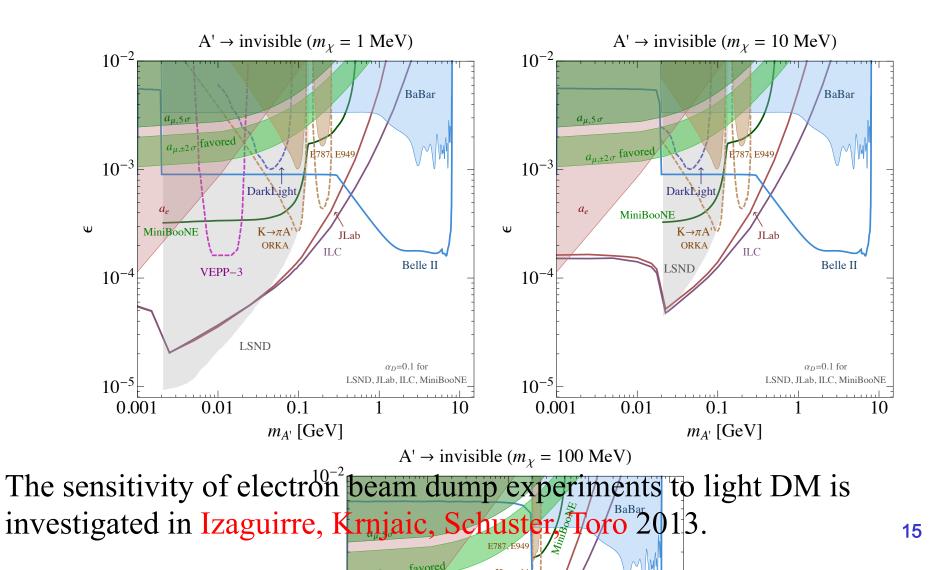
MP, 2008

### $\kappa$ - $m_V$ parameter space, Essig et al 2013

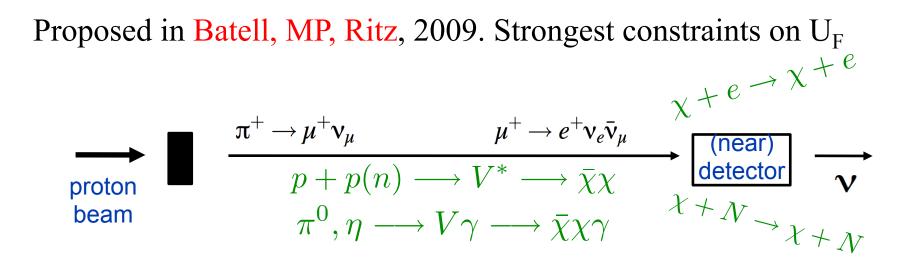


Dark photon models with mass under 1 GeV, and mixing angles ~  $10^{-3}$  represent a "window of opportunity" for the high-intensity experiments, and soon the g - 2 ROI will be completely covered. *But what if dark photon decays to light dark matter?* 

# Compilation of current constraints on dark photons decaying to light DM



### Fixed target probes - Neutrino Beams



We can use the neutrino (near) detector as a dark matter detector, looking for recoil, but now from a relativistic beam. E.g.

T2K 30 GeV protons (IIIII) ~5x10<sup>21</sup> POT) 280m to on- and offaxis detectors

### MINOS

120 GeV protons 10<sup>21</sup> POT 1km to (~27ton) segmented detector MiniBooNE 8.9 GeV protons 10<sup>21</sup> POT 540m to (~650ton) mineral oil detector

#### Light Mass WIMP Searches with a Neutrino Experiment: A Request for Further MiniBooNE Running

September 19, 2012

#### The MiniBooNE Collaboration

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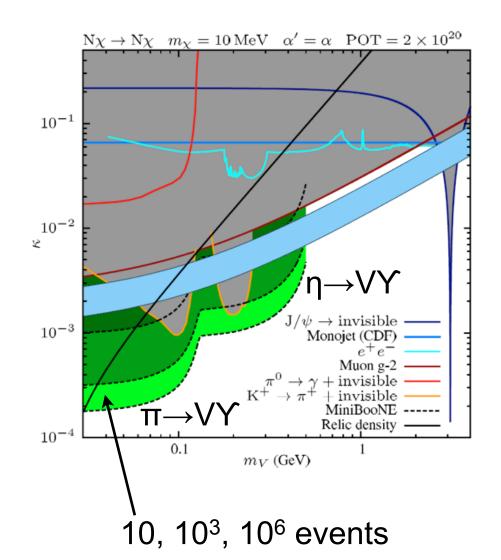
#### The Theory Collaboration

B. Batell University of Chicago, IL 60615

P. deNiverville , D. McKeen, M. Pospelov, & A. Ritz University of Victoria, Victoria, BC, V8N-1M5

# Secured several months of running in 2013-2014!

### Batell, deNiverville, McKeen, MP, Ritz MiniBooNE sensitivity – quite a bit of new ground can be covered



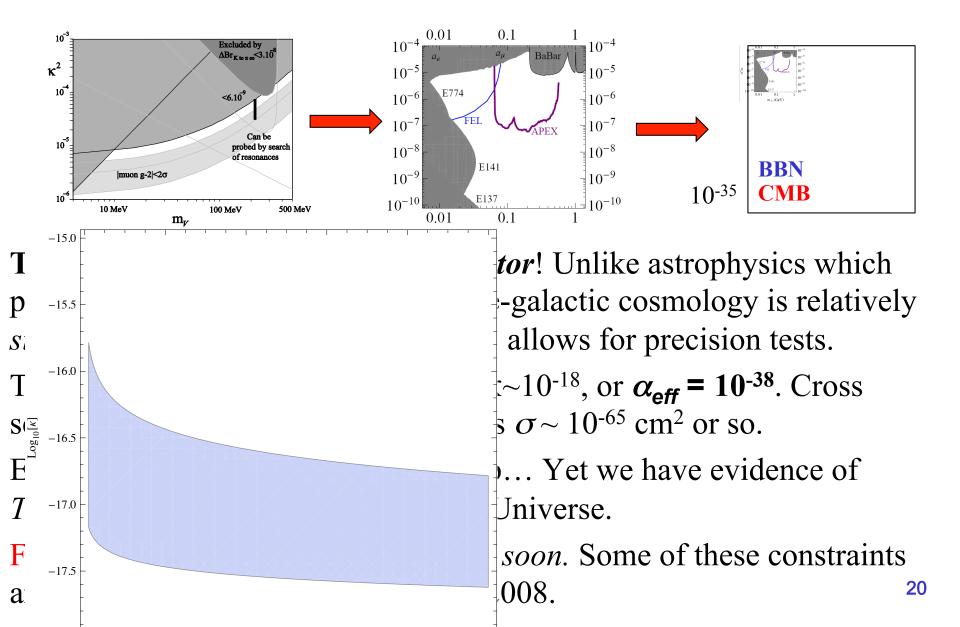
### Other type of dark vector force?

The work on other type of "dark vector forces" is on-going.

- Batell et al. investigate the sensitivity of MiniBooNE to the "baryonic" portal, or gauged baryon number – this portal is also "flavor-safe".
- 2. Altmannshofer et al. investigate dark forces coupled to gauged lepton numbers (such as  $L_{\mu}-L_{\tau}$ ). It turns out that the most sensitive probe of such models is the well-forgotten (25 yr old) observation of muon pair-production by the muon neutrinos,  $Z+\nu \rightarrow Z + \mu^+ \mu^- + \nu$ , the so-called muonic trident.

**Mini-conclusions**: systematic studies of all possibilities for the intensity frontier physics, searching for light weakly coupled states, is gathering momentum. There is still plenty of "low hanging fruits", and in many instances sensitivity reach can be advanced by many orders of magnitude by a relatively modest theoretical and experimental investment.

# Very [very] dark photons



### **Produced early – decays late**

- The production cross section is ridiculously small, but in the early Universe at  $T > m_V$ , in fact, *every colliding pair of particles can produce such Vectors*, and there is a lot of time available for this.
- Once produced such particles *live for a very long time*, and decay in the "quiet" Universe, depositing non-thermal amounts of energy and changing physics of primordial matter after recombination.
- Precision determination of optical depth during the CMB, position of Doppler peaks and the slope of the Silk diffusion tale provide tight restrictions on the amount of energy injected.
- Due to BBN we also have a pretty good evidence that the Universe in fact once was at least T ~ a few MeV hot.....

# Filling out details....

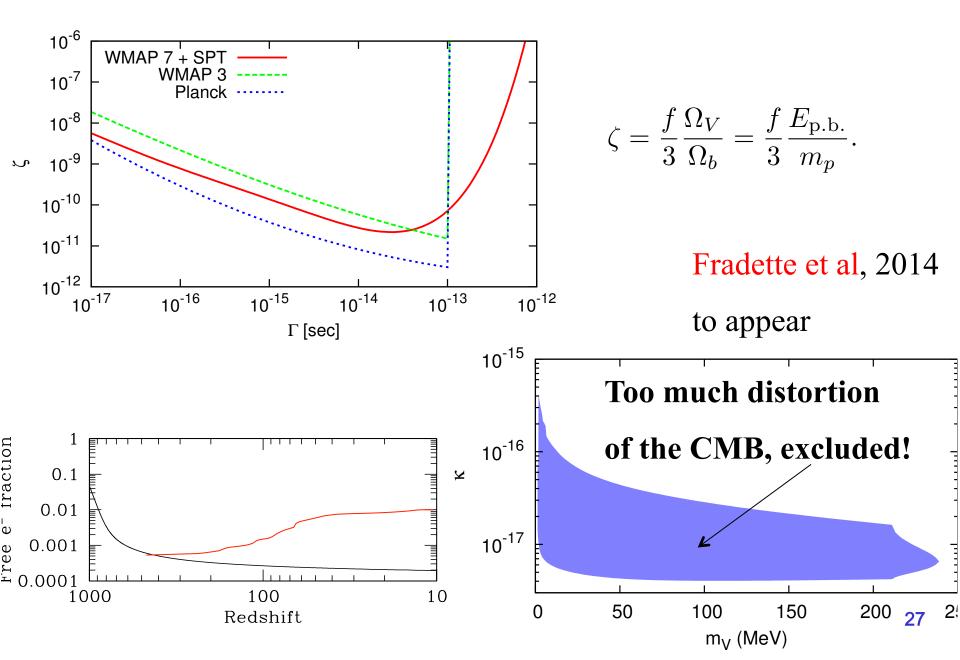
Lifetime against the decay to electron-positron pairs

$$\tau_V \simeq \frac{3}{\alpha_{\rm eff} m_V} = 0.6 \text{ mln yr} \times \frac{10 \,\mathrm{MeV}}{m_V} \times \frac{10^{-35}}{\alpha_{\rm eff}}$$

•  $e^+e^- \rightarrow V$  in the early Universe leads to the energy stored per baryon  $E_{p.b.} \sim \frac{m_V \Gamma_{\text{prod}} H_{T=m_V}^{-1}}{n_{b,T=m_V}} \sim \frac{0.1 \alpha_{\text{eff}} M_{\text{Pl}}}{\eta_b} \sim \alpha_{\text{eff}} \times 10^{36} \text{ eV}$  $\int_{0.01}^{0} \int_{0.01}^{0} \int_{0.01}^{0$ 

Once injected back to the medium via V→e<sup>+</sup>e<sup>-</sup> ~ 1/3 of the stored energy leads to ionization. E.g. 1 eV p.b. recreates X<sub>e</sub> ~ few 10<sup>-2</sup>.<sup>26</sup>

### **Constraints from WMAP**



# <BB> = T or "T-like" modes ?

- 1. Every big discovery follows by the period of trying to understand the *result*. E.g. excess of events around 125 GeV  $\rightarrow$  Evidence of a new resonance  $\rightarrow$  Higgs-like properties of the resonance  $\rightarrow$  dropping "-like" after lots of tests. In the process you rule out competitors such as KK-graviton, techni-pion, etc [no matter how creepy they are]. Same process should occur with the discovery of B-modes, but is not happening yet to the fullest.
- 2. The minimal interpretation of B-modes are tensor perturbations, the remnants of inflation that occurred with  $H_{infl}=10^{14}$  GeV. Well, it poses a lot of questions to anyone who tries to play with some physics that has fundamental scale below10<sup>14</sup> GeV.
- 3. One can provide new mechanisms of B-mode generation with a low inflationary scale, e.g.  $H_{infl} \sim 10^{10}$  GeV (MP, Ritz, Skordis, 2008). View it as a competitive explanation of Bicep observations, and try to rule it out from data!

### **Two-axion model**

 Two-axion model is like that. One axion becomes a QCD axion, and the other one remains massless,

$$\left(\frac{a_1}{2g_1} + \frac{a_2}{2g_2}\right) G_{\mu\nu} \tilde{G}^{\mu\nu} + \left(\frac{a_1}{2f_1} + \frac{a_2}{2f_2}\right) F_{\mu\nu} \tilde{F}^{\mu\nu} \rightarrow \mathcal{L}_{\text{QCDa}} + \frac{a}{2f_a} F_{\mu\nu} \tilde{F}^{\mu\nu},$$

Coupling constant is given by

$$f_a^{-1} = (g_2/f_2 - g_1/f_1)/\sqrt{g_1^2 + g_2^2}.$$

### Fluctuating pseudoscalar driven by inflation

The model:

$$\mathcal{L}_{everything} = \mathcal{L}_{SM+gravity} + \mathcal{L}_{inflation} + \frac{1}{2} (\partial_{\mu}a)^2 + \frac{a}{2f_a} F_{\mu\nu} \tilde{F}_{\mu\nu}$$

[Can be viewed as a generic consequence of two QCD axions.

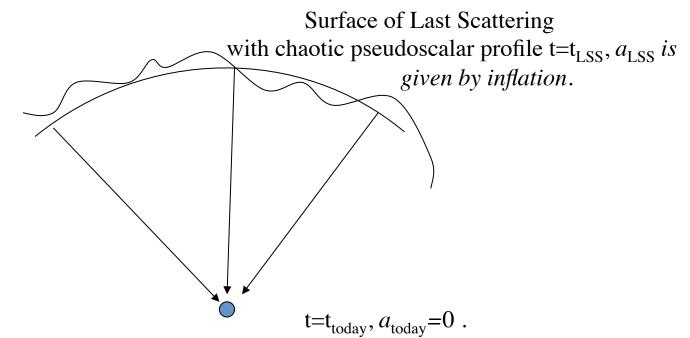
Massless field *a* receives [random, Gaussian, nearly flat-spectrum] fluctuations during inflation,  $\delta a \sim H_{infl}/(2\pi)$ .

Rotation of polarization plane after travelling from point 1 to point 2 is (Harrari, Sikivie; Carroll; Lue, Wang, Kamionkowski...)

$$\psi = \frac{a_1 - a_2}{f_a}$$
$$\langle EE \rangle \to \langle BB \rangle; \qquad \langle TB \rangle = \langle EB \rangle = 0$$

The measure of the r.m.s. angular rotation is  $\delta a \sim H_{infl}/(2\pi f_a) \log z$ 

# Propagation of CMB from the LSS



Polarization of arriving to us CMB photons is randomly rotated by  $\Delta \psi(n) = A_{\text{LSS}}(n) = a_{\text{LSS}}(n) / f_{a}$ . For convenience, we introduce  $c_a$ 

$$c_a = \left(\frac{H}{2\pi f_a}\right)^2, \quad |\Delta\psi| \sim \sqrt{c_a}.$$

Since  $f_a > 10^{11}$  GeV is a mild constraint,  $H \sim 10^{10}$  GeV or below can generate BB

### Master formula for <BB> calculation MP, Ritz, Skordis, 2008

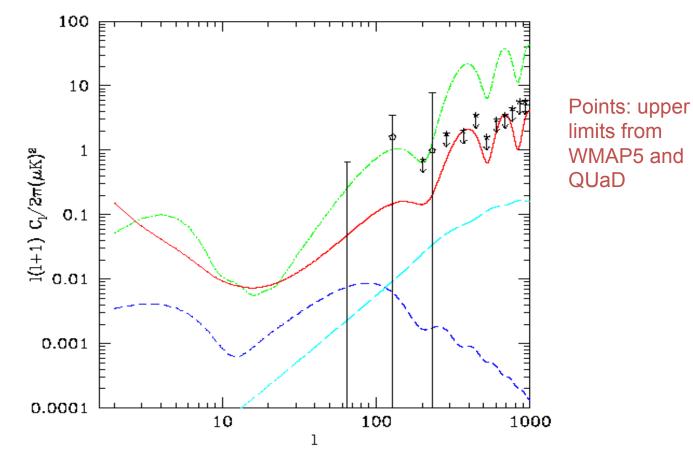
$$C_{Bl} = \frac{1}{2l+1} \sum_{m} \langle a_{Blm}^* a_{Blm} \rangle = \frac{4(4\pi)^3}{2l+1} \frac{(l-2)!}{(l+2)!} \\ \times \sum_{m,l_1,l_2} (2l_1+1)(2l_2+1) \left( \begin{array}{cc} l & l_1 & l_2 \\ 0 & 0 & 0 \end{array} \right)^2 \\ \times \int k^2 \underline{P_{\Phi}} q^2 \underline{P_A} dk dq |\Delta_{l_1 l_2 m}(k,q)|^2,$$

with the generalized transfer function,

$$\Delta_{l_1 l_2 m}(k,q) = \frac{3}{4} \int_0^{\tau_0} d\tau g(\tau) j_{l_1}(x) j_{l_2}(y) \\ \times \left( \frac{(l_1+2)!}{(l_1-2)!} \frac{1}{x^2} - m^2 \right) \Delta_A(\tau,q) \Pi(\tau,k).$$

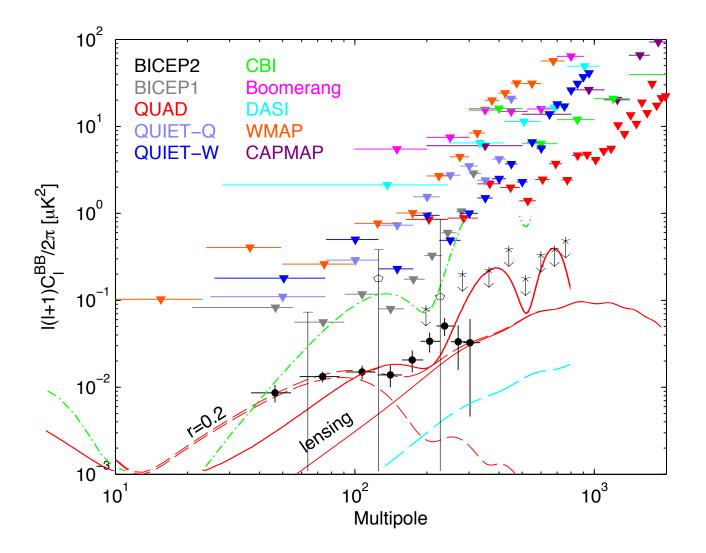
### 2008 limits

### Numerical Results and comparison with experiment



Green: EE; Red: BB with  $c_a = 0.004$ ; Dark blue: BB from gravity waves with r=0.14; light blue: BB lensing background . CAN WE MATCH IT ONTO BICEP???

### The moment of truth



**Inflationary pseudoscalar fluctuations do not give a good fit to Bicep data( too low l<100) !!!** So, it more "T-like" and not at all "*a*-like".

# Conclusion

- 1. Intensity frontier "orthogonal" to energy frontier direction can and should look for light weakly coupled states.
- Search for "dark photons" the simplest model of an additional U(1) has been intensified over the years, fueled by its possible connection to several "anomalies" in particle physics and cosmology. New results brought new constraints, but no independent hints on a signal. "g-2" region is [almost] covered.
- 3. "Very dark photons" with mixing angles  $\sim 10^{-17}$  is an example of unique sensitivity the precision CMB brings to our field.
- 4. Recent discovery of the B-modes if confirmed as coming from the tensor perturbations generated during inflation with  $H_{infl} \sim 10^{14}$  GeV will limit any massless pseudoscalar field to photons as  $f_{a\gamma\gamma} > 10^{15}$  GeV (compare with direct CAST bound of  $10^{11}$  GeV). It is also important to rule chaotic rotation of E to B as the sole source of  $C_{BB}^{l}$ .

# Astrophysical motivations: 511 keV line

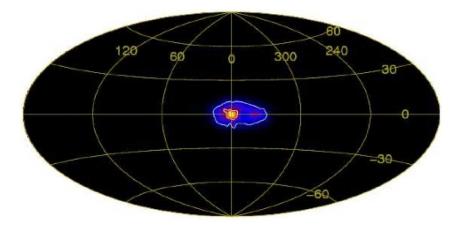


FIG. 4 511 keV line map derived from 5 years of INTE-GRAL/SPI data (from Weidenspointner  $et \ al.$ , 2008a).

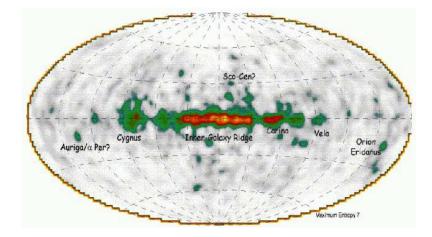
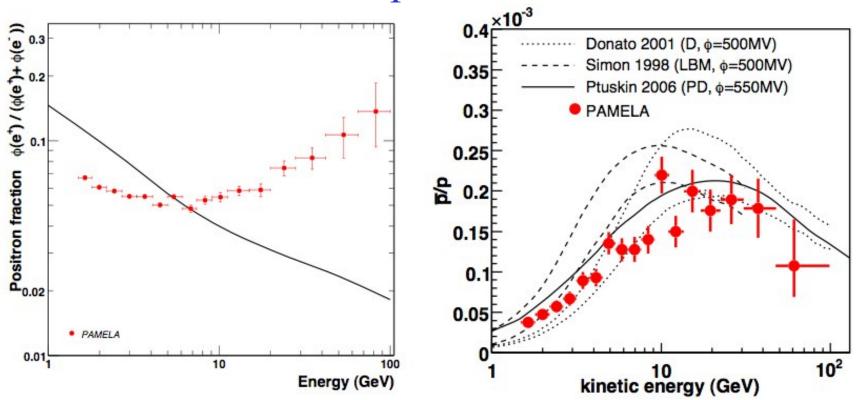


FIG. 7 Map of Galactic  $^{26}\mathrm{Al}$   $\gamma\text{-ray}$  emission after 9-year observations with COMPTEL/CGRO (from Plüschke *et al.*, 2001).

There is a lot more positrons coming from the Galactic Center and the bulge that expected. The emission seems to be diffuse.

- 1. Positrons transported into GC by B-fields?
- 2. Positrons are created by episodic violent events near central BH?
- 3. Positrons being produced by DM? Either annihilation or decay? <sup>32</sup>

### **PAMELA** positron fraction



No surprises with antiprotons, but there is seemingly a need for a new source of positrons!

This is a "boost" factor of 100-1000 "needed" for the WIMP interpretation of PAMELA signal. E.g. SUSY neutralinos would not work, because  $\langle \sigma v \rangle$  is too small. Enhancing it "by hand" does not work because WIMP abundance goes down. Dark forces allow bridging this gap due to the late time enhancement by Coulomb (Sommerfeld)<sup>33</sup>.