# The Standard Model (SM) in the Sky (A new consequence of Inflation)

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arXiv:1402.1345

PONT 2014, Avignon, April 14-18, 2014.

I shall be talking about ...

#### 1. Parametrically Excited Fermions $\Rightarrow$ GWs

#### 2. INFLATION $\Rightarrow$ SM Higgs Condensate $\Rightarrow$ SM Fermions $\Rightarrow$ GWs

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#### Scalar field (condensate) after Inflation:

Coherent Oscillations:  $\phi(t) \approx \Phi(t)f(t)$ , f(t+T) = f(t)



$$\psi(\mathbf{x},t) = \int \frac{d\mathbf{k}}{(2\pi)^3} e^{-i\mathbf{k}\cdot\mathbf{x}} \left[ \hat{a}_{\mathbf{k},r} \mathbf{u}_{\mathbf{k},r}(t) + \hat{b}_{-\mathbf{k},r}^{\dagger} \mathbf{v}_{\mathbf{k},r}(t) \right],$$



$$\psi(\mathbf{x},t) = \int \frac{d\mathbf{k}}{(2\pi)^3} e^{-i\mathbf{k}\cdot\mathbf{x}} \Big[ \hat{a}_{\mathbf{k},r} \begin{pmatrix} u_{\mathbf{k},+}(t)S_r \\ u_{\mathbf{k},-}(t)S_r \end{pmatrix} + \hat{b}_{-\mathbf{k},r}^{\dagger} \begin{pmatrix} v_{\mathbf{k},+}(t)S_r \\ v_{\mathbf{k},-}(t)S_r \end{pmatrix} \Big],$$



$$\frac{d^2}{dt^2} u_{\mathbf{k},\pm} + \left(\omega_{\mathbf{k}}^2(t) \pm i \frac{d(am_{\psi})}{dt}\right) u_{\mathbf{k},\pm}(t) = 0, \quad \omega_{\mathbf{k}}^2(t) = k^2 + a^2(t) \frac{m_{\psi}^2(t)}{dt}$$



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$$\frac{d^2}{dt^2}u_{\mathbf{k},\pm} + \left(\omega_{\mathbf{k}}^2(t) \pm i\frac{d(am_{\psi})}{dt}\right)u_{\mathbf{k},\pm}(t) = 0\,, \quad \frac{d}{dt}\omega_{\mathbf{k}} \gg \omega_{\mathbf{k}}^2(t)$$



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#### 1.2 GWs from Fermions

Fermi-Sphere:

phere: 
$$\begin{cases} n_k(k \lesssim k_*) \lesssim 1, \\ n_k(k \gg k_*) \to 0 \end{cases} \Rightarrow T^{(\psi)}_{\mu\nu} \sim \left(\bar{\psi}\gamma_{(\mu}D_{\nu)}\psi\right) [u_{k,\pm}(t)] \\ \left[ \left(T^{(\psi)}_{\mu\nu}\right)^{\mathrm{TT}} \to \text{ GW Source }! \right] \end{cases}$$

**GWs**: 
$$\frac{d\rho_{\text{GW}}}{d\log k}(k,t) = \# \frac{Gk^3}{a^4(t)} \int_0^t \int_0^t dt_1 dt_2 \,\mathcal{G}(k,t_2-t_1) \,\Pi^2(k,t_1,t_2)$$

 $\mathsf{UTC:}\;\left\langle T_{ij}^{\mathrm{TT}}(\mathbf{k},t_{1})\,T_{ij}^{\mathrm{TT}}(\mathbf{k}',t_{2})\right\rangle \equiv (2\pi)^{3}\,\Pi^{2}(k,t_{1},t_{2})\,\delta^{(3)}(\mathbf{k}-\mathbf{k}')$ 

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Enqvist, DGF, Meriniemi PRD'12, JHEP'13

Inflation:  $dS(H_*)$ ,  $(H_* \gg v \equiv 246 \text{ GeV})$ 

SM Higgs: 
$$\Phi = \frac{\varphi}{\sqrt{2}} \rightarrow V(\varphi) = \frac{\lambda(\mu)}{4} \varphi^4, \ \mu = \varphi \gg v$$

 $\begin{array}{l} \textbf{Prob. Dist: } \varphi \text{ light } \left( |V''| < H_*^2 \right) \Rightarrow \left\{ \begin{array}{l} \text{Random Walk } \left( k < aH_* \right) \\ \\ P_{\text{eq}}(\varphi) \propto \text{Exp}\{-c\lambda_*(\varphi/H_*)^4\} \end{array} \right. \end{array} \right.$ 

End of Inflation:  $\varphi_* = \alpha H_* / \lambda_*^{1/4}$   $\alpha \in [0.01, 1]$  (98 %)

Inflation: 
$$dS(H_*)$$
,  $(H_*|_{BICEP2} \sim \mathcal{O}(10^{14}) \text{ GeV})$ 

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**End of Inflation**:  $\varphi_* \neq 0$   $(V \propto \varphi^4) \Rightarrow$  Higgs Oscillations (!)

$$y_j \varphi \bar{\psi}_j \psi_j : \quad \frac{d^2}{d\tau^2} u_{k,\pm}^{(j)} + \left( \kappa^2 + q_j (a\varphi)^2 \pm i \sqrt{q_j} \frac{d}{d\tau} (a\varphi) \right) u_{k,\pm}^{(j)} = 0 \,, \quad q_j \equiv \frac{y_j^2}{\lambda_1}$$

$$j = \left\{ \begin{array}{c} \{t, b, c, s, u, d\} \\ \{e, \mu, \tau\} \end{array} \right\} \quad \Rightarrow \quad \Omega_{\mathrm{GW}}^{(j)}(k) \equiv \frac{1}{\rho_c} \frac{d\rho_{\mathrm{GW}}}{d\log k}(k; q_j) \,, \quad q_j \equiv \frac{y_j^2}{\lambda_{\mathrm{I}}}$$

$$n_k^{(j)}(k \lesssim k_*^{(j)}) \rightarrow \Omega_{\rm GW}^{(j)}(k) : \left\{ \begin{array}{l} k_p \sim k_*^{(j)} \; ({\rm Max.}) \\ \propto k^3, \; k \ll k_p \\ \propto k^{-1.5}, \; k \gg k_p \end{array} \right\}$$

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$$(\delta \ll 1)$$

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#### SCALING (Universal Shape) :

 $\Omega_{\rm GW}(k; \mathbf{q}_j) = (H_{\rm I}/M_p)^4 \, (a_{\rm I}/a_{\rm F})^{1-3w} \times \mathbf{q}_j^{1.55} \, \mathcal{U}(k/k_p)$ 

$$\mathcal{U}(x) \equiv \mathcal{U}_1 \frac{x^3}{(\alpha + \beta x^{4.5})}, \begin{cases} \mathcal{U}_1 \equiv \mathcal{U}(1) \ [\sim 10^{-5} (\text{RD}), \sim 10^{-6} (\text{MD})] \\ \alpha + \beta = 1 \ [\alpha = 0.25, \beta = 0.75 \ (\text{RD}, \text{MD})] \end{cases}$$



Total GWs :

$$h^2 \Omega_{\rm GW}^{(0)}(f) \simeq \epsilon_{\rm I} \, 10^{-6} \, (H_{\rm I}/M_p)^4 \, \sum_j q_j^{1.55} \, \mathcal{U}(q_j^{-1/4}(k/H_{\rm I}))$$

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SM Yukawa Couplings :  $y_t > y_b > y_\tau > y_c > y_\mu \gtrsim y_s > y_d > y_u > y_e$ 

 $h^2 \Omega_{\rm GW}^{(0)}(f) \propto q^{3/2} \propto y^3 \Rightarrow \text{Top Quark dominates (!)}$ 

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Top Quark GW Peak Today:  $(H_* \sim 10^{14} \text{ GeV}, y_t \sim 0.5)$ 

<u>Today</u>:  $f_p^{(t)} \sim 10^7 \, \text{Hz}$ ,  $h^2 \Omega_{\text{GW}}^{(\text{p})} \Big|_{\text{t}} \sim 10^{-30} \, \lambda_{\text{I}}^{-1.55}$ 

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 $\lambda_{\rm I} \lesssim 10^{-7}, 10^{-10}, 10^{-13} \Rightarrow h^2 \Omega_{\rm GW}^{(\rm p)}|_{\rm t} \gtrsim 10^{-20}, 10^{-15}, 10^{-10}$ 



#### • Inflation, $dS(H_*) \Rightarrow$ Higgs Osc. $\Rightarrow \Psi's$ (Param. Excitation).

- $\ \, \textcircled{\ } \{\Psi_a\} \ \, \to \ \, h^2\Omega^{(0)}_{\rm GW}(f) \propto q_a^{3/2} \propto y_a^3 \ \, \Rightarrow {\rm Top \ Quark \ dominates \ (!)}$
- 0 < λ<sub>I</sub> ≪ 1, H<sub>\*</sub> ~ 10<sup>14</sup> GeV, Peak's Frequency: f<sub>\*</sub> ~ 10<sup>7</sup> Hz h<sup>2</sup>Ω<sup>(p)</sup><sub>GW</sub>|<sub>t</sub> ≥ 10<sup>-20</sup>, 10<sup>-15</sup>, 10<sup>-10</sup> (λ<sub>I</sub> ≤ 10<sup>-7</sup>, 10<sup>-10</sup>, 10<sup>-13</sup>).
- Similar Conclusions for Higgs Inflation ! Also expected in BSM !
- Ouriversal Effect: Inflation + SM Higgs ⇒ GWs: Spectroscopy of Particle Physics: Probing the Most Strongly Interacting Particle (!)

- Inflation,  $dS(H_*) \Rightarrow$  Higgs Osc. $\Rightarrow \Psi's$  (Param. Excitation).
- $\{\Psi_a\} \rightarrow h^2 \Omega_{GW}^{(0)}(f) \propto q_a^{3/2} \propto y_a^3 \Rightarrow \text{Top Quark dominates (!)}$
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