

Primordial black holes as dark matter candidates: closing the remaining mass window

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Capela, Pshirkov, PT, PRD87 (2013) 023507

Capela, Pshirkov, PT, PRD87 (2013) 123524

Capela, Pshirkov, PT, arXiv:1403.7098

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Outline

- 1 Introduction
- 2 Capture of PBH in stars
- 3 Constraints on PBH
- 4 Summary

INTRODUCTION

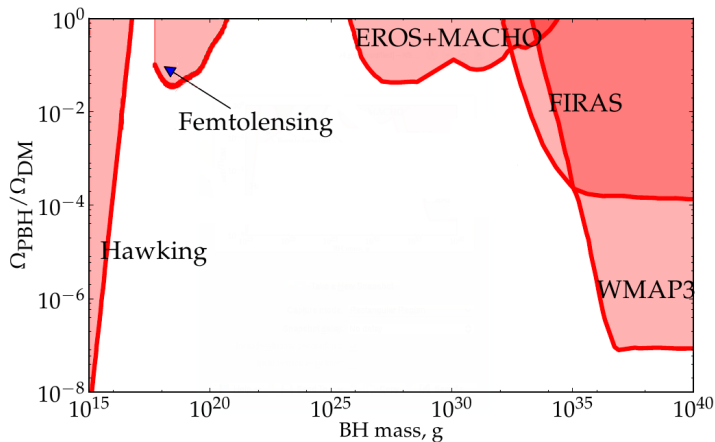
- Many (indirect) arguments suggest the existence of dark matter with $\Omega_{\text{DM}} \simeq 0.26$
 - Rotation curves of galaxies
 - Gas temperature in clusters
 - Gravitational lensing
 - Structure formation
- The DM is often assumed to be a new stable particle. There are many different DM candidates: axion-like particles, sterile neutrinos, WIMPs, ...
- We will consider an alternative option — that DM is composed of primordial black holes (PBH)
 - PBH may have been produced in the early Universe. Many mechanisms have been considered (e.g., at phase transitions).
 - Attractive feature of this scenario is that no new particles are required.
 - From theory, the PBH masses are not constrained.

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- A large part of the parameter space is already constrained from various arguments



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- In the remaining mass window $10^{16} - 10^{26}$ g, the abundance of PBH may be constrained from observations of **compact stars — WD and NS**
- Compact stars are special because if a PBH gets inside such a star, the star gets destroyed. Requiring the probability of such event is $\ll 1$ imposes constraints on PBH abundance.
- DM capture in stars has been considered before
 - Press, Spergel Astrophys.J. 296 (1985) 679-684;*
 - Goldman, Nussinov Phys. Rev. D40, 3221 (1989);*
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 - Sadin, Ciarcelluti, Astropart. Phys. 32 (2009) 278-284;*
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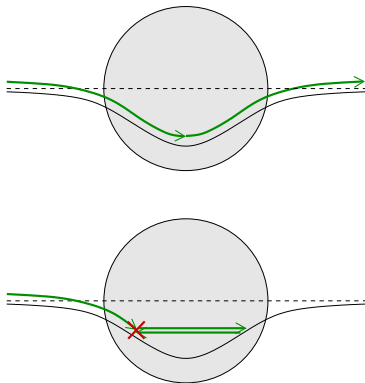
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Two ways to capture a PBH: (I) during star lifetime



- For capture the energy loss is required
- In case of PBH, the energy loss occurs due to dynamical friction and accretion of star matter onto the PBH
- Total energy loss is

$$E_{loss} = \frac{4m_{BH}^2 M_*}{M_{Pl}^4 R_*^2} \left\langle \frac{\ln \Lambda}{v^2} \right\rangle$$

- Averaging with Maxwellian distribution, the capture rate is

$$F = \sqrt{6\pi} \frac{\rho_{DM}}{v_{\infty} m_{BH}} \frac{R_g R_*}{1 - R_g/R_*} \left[1 - \exp\left(-\frac{3E_{\text{loss}}}{m_{BH} v_{\infty}^2}\right) \right]$$

$$\simeq 3\sqrt{6\pi} \frac{\rho_{DM}}{v_{\infty}^3} \frac{R_g R_*}{m_{DM}^2} E_{\text{loss}} \quad \text{at } E_{\text{loss}} \ll m_{DM} v_{\infty}^2$$

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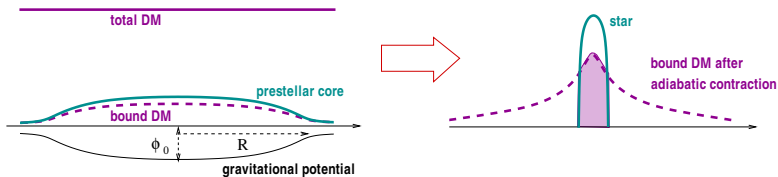
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Two ways to capture a PBH: (II) at star formation

- The stars are formed in the collapse of baryonic matter in giant molecular clouds. These clouds have some DM density gravitationally bound to them.
- Collapsing baryons gravitationally drag the DM along by adiabatic contraction, so some PBHs end up inside the star
- When the star evolves into a compact remnant (NS or WD), some of these PBHs may be inherited by the latter.

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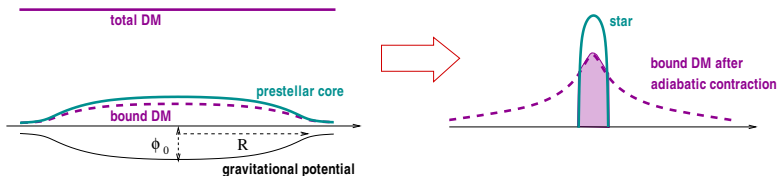
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- The density of bound DM, assuming Maxwellian parent distribution with \bar{v} :

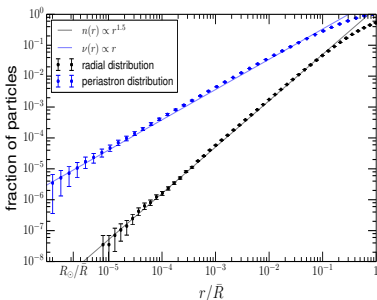
$$\rho_{\text{bound}} \sim \bar{\rho}_{DM} \left(\frac{\phi_0}{\bar{v}^2} \right)^{3/2} = \text{const} \cdot \frac{\bar{\rho}_{DM}}{\bar{v}^3}$$

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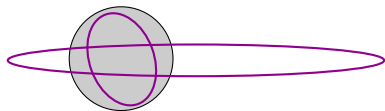


- Number of particles within r ,

$$n(r) \propto r^{3/2}$$

- Number of particles with periastron $< r$,

$$\nu(r) \propto r$$



Time scales

- Two stages
 - When PBH is mostly outside the star

$$\tau_1 \simeq \frac{\sqrt{r_{\max}} R_*^2 v_{\text{esc}}^2}{G \sqrt{R_g} m_{\text{BH}} \ln \Lambda} \sim 2 \times 10^8 \text{ yr} \left(\frac{10^{22} \text{ g}}{m_{\text{BH}}} \right)$$

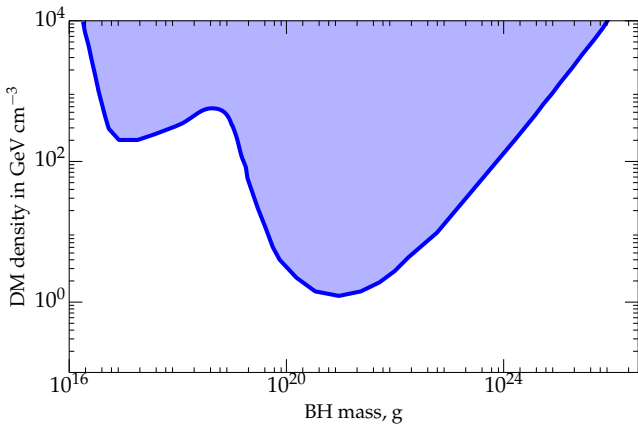
- When PBH is completely inside the star: numerical calculation in a realistic density profile. Rough estimate:

$$\tau_2 \sim 10^2 \frac{M_*^{3/2}}{2\pi \sqrt{G} \rho(0) m_{\text{BH}} R_*^{3/2} \ln \Lambda}$$

- \Rightarrow Insufficient time at small m_{BH}

RESULTING CONSTRAINTS

Assuming DM velocity dispersion $v = 7$ km/s



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Where to look?

- Best constraints come from sites where the DM density is largest and the DM velocity is smallest
 - One such site could be Globular Clusters (GC) — bound compact systems containing $10^4 - 10^7$ stars, very old $\gtrsim 10$ Gyr
 - There are two suggested mechanisms of the GC formation: primordial and 'recent'
 - recently formed GC carry little DM — not enough for constraints
 - primordial GCs should have DM cores with $\rho_D \sim 2 \times 10^3 \text{ GeV cm}^{-3}$.
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- Another candidate — dwarf spheroidals
 - similar to GC in size; DM-dominated: densities $\sim 200 \text{ GeV/cm}^3$ have been inferred from modeling
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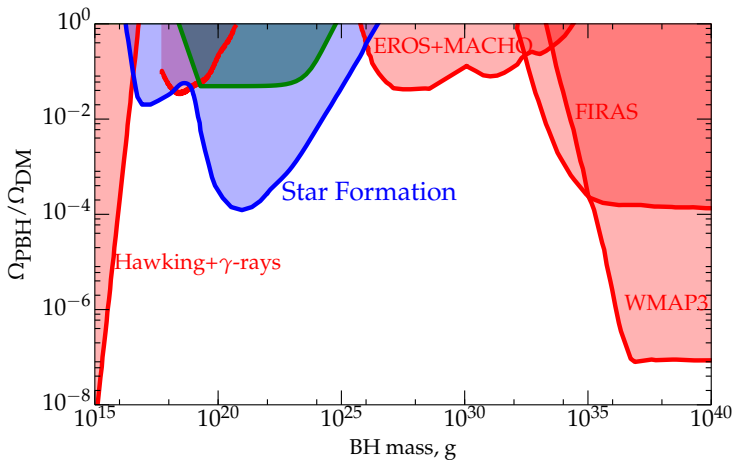
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Assuming $\rho_D = 10^4 \text{ GeV/cm}^3$ and $v = 7\text{ km/s}$
as could be in the cores of GC



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- To close the remaining mass window, one may either show the presence of DM cores in GC (or their primordial origin) or observe pulsars in dSph (feasible, first hints exist)

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