

QCD back-scattering photons in heavy ion collisions

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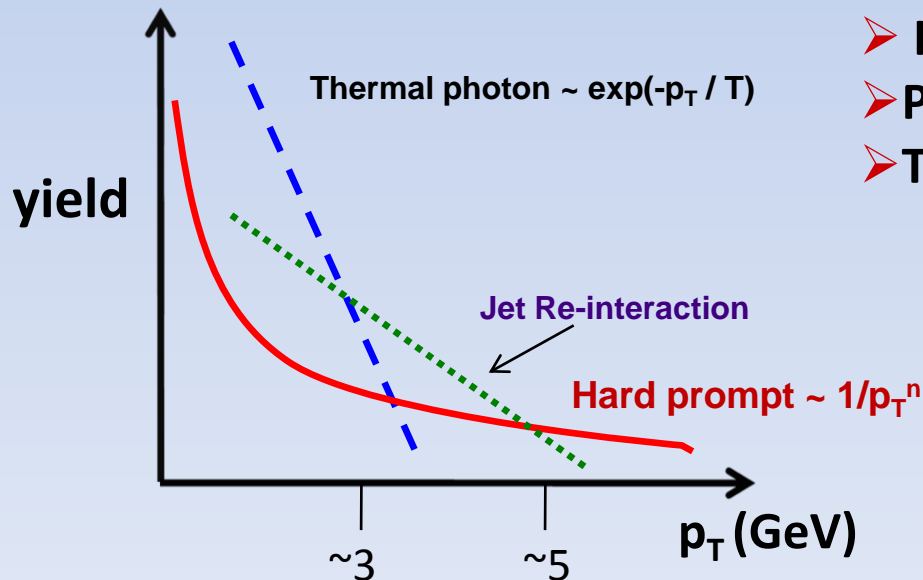
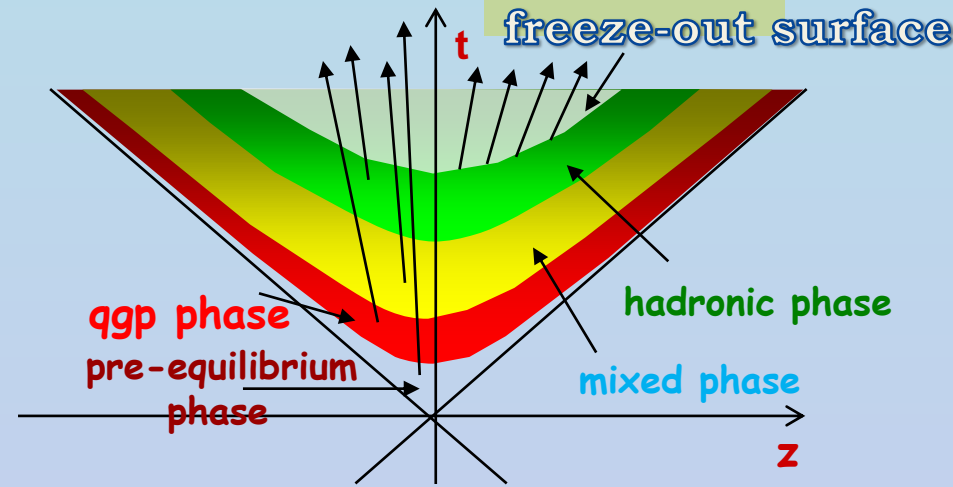
**In collaboration with: Dr. Rainer J Fries, Texas A & M University
Dr. Dinesh K Srivastava, VECC, Kolkata**

Matter at Extreme Conditions: Then & Now ; Bose Institute, Kolkata

Electromagnetic probes in Heavy Ion Collisions:

Direct photons are considered as the cleanest probe in HIC

Being electromagnetic in nature, photons carry information undistorted from each stages of evolution.

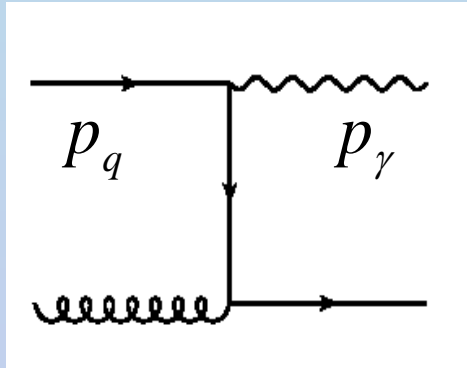


- Hard photons (Direct + Jet fragment)
- Pre-equilibrium and jet-medium photons
- Thermal photons (QGP + Hadron matter)

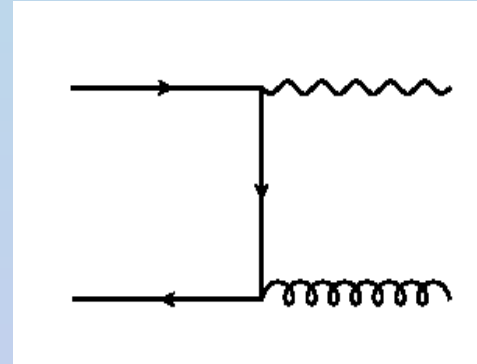
Experimental challenge to separate different sources of direct photon

Photons from re-scattering of jets in quark gluon plasma

- QCD Compton and Annihilation process :



$$\frac{d\sigma}{dt} = -\frac{\pi\alpha\alpha_s e_q^2}{3s^2} \left(\frac{u}{s} + \frac{s}{u} \right)$$



$$\frac{d\sigma}{dt} = \frac{8\pi\alpha\alpha_s e_q^2}{9s^2} \left(\frac{u}{t} + \frac{t}{u} \right)$$

Cross sections are maximum for small values of t and u

Forward scatt. $\vec{p}_\gamma \approx \vec{p}_q(\vec{p}_{jet})$ Backward scatt. $\vec{p}_\gamma \approx \vec{p}_{\bar{q}}(\vec{p}_{jet})$

Jet-Photon Conversion

The rate of production:

$$E_\gamma \frac{dN^{(A)}}{d^4 x d^3 p_\gamma} = \frac{16E_\gamma}{2(2\pi)^6} \sum_{q=1}^{N_f} f_q(p_\gamma) \\ \times \int d^3 p f_q^-(p) [1 + f_g(p)] \sigma^{(A)}(s) \frac{\sqrt{s(s-4m^2)}}{2E_\gamma E} + (q \leftrightarrow \bar{q})$$

$$E_\gamma \frac{dN^{(C)}}{d^4 x d^3 p_\gamma} = \frac{16E_\gamma}{2(2\pi)^6} \sum_{q=1}^{N_f} f_q(p_\gamma) \\ \times \int d^3 p f_g(p) [1 - f_q(p)] \sigma^{(C)}(s) \frac{s-m^2}{2E_\gamma E} + (q \leftrightarrow \bar{q})$$

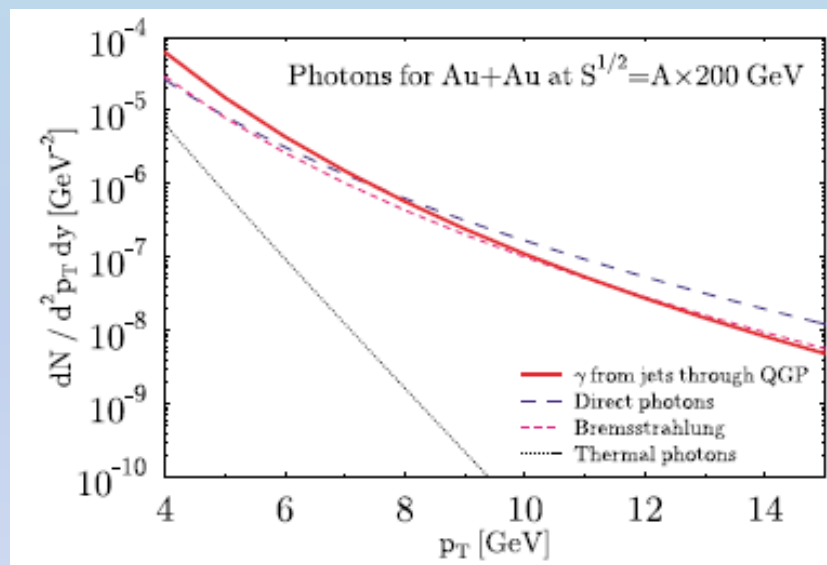
$$f^q(p) = f^{jet}(p) + f^{th}(p)$$

$$dN^\gamma = f_{th} \otimes f_{th} + f_{th} \otimes f_{jet} + \dots \quad \text{Jet-converted photons}$$

Jet-photon contd..

▪ Total inclusive yield:

$$E_\gamma \frac{dN}{d^4 x d^3 p_\gamma} = \frac{\alpha \alpha_s}{4\pi^2} \sum_f \left(\frac{e_f}{e}\right)^2 [f_q(x, p_\gamma) + f_{\bar{q}}(x, p_\gamma)] T^2 \left[\ln \frac{3E_\gamma}{\alpha_s \pi T} + C \right]$$



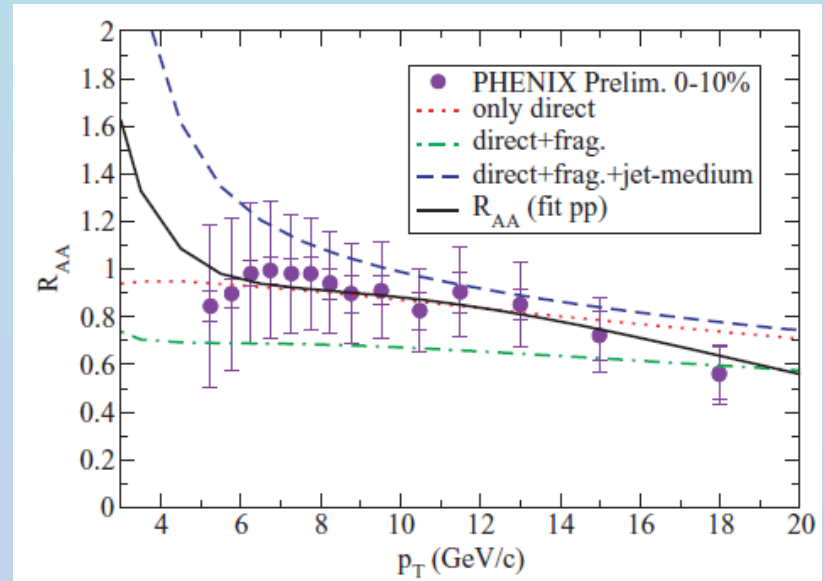
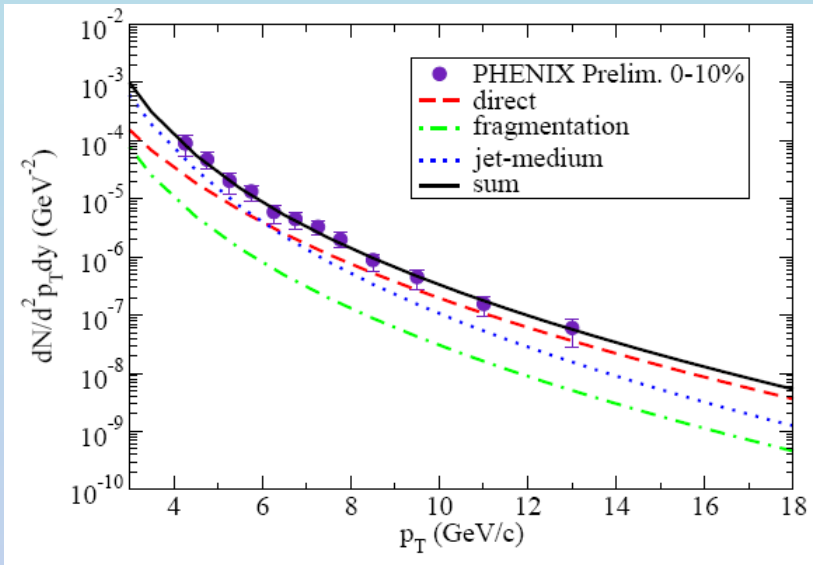
First proposed by :
Fries, Muller, Srivastava (
PRL,90, 132301 (2003))

Shows substantial
contribution for $p_T \leq 6$
GeV at RHIC energies

Experimental measurement of photons :

- Inclusive yield and Nuclear modification of direct photons
- Azimuthal momentum anisotropy coeff. (v_2)

Inclusive yield and R_{AA} :



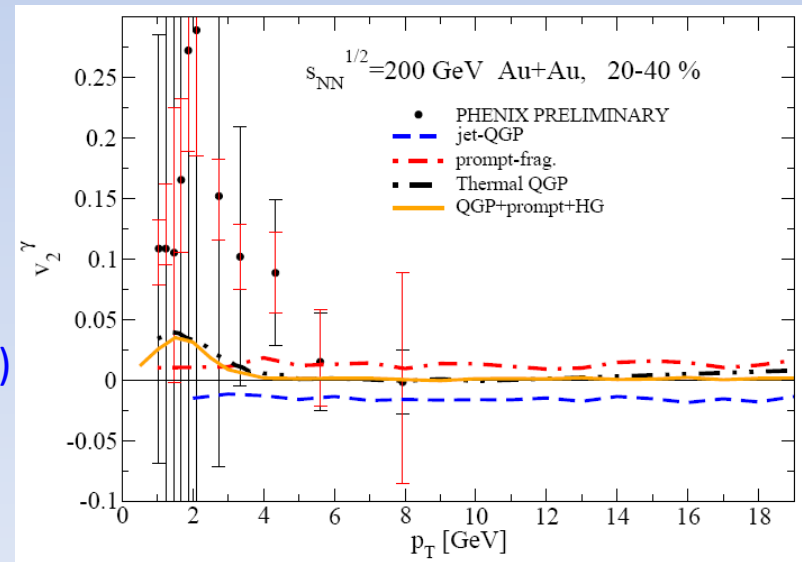
Qin, Ruppert, Gale, Jeon & Moore, PRC 80 (2009)

Azimuthal momentum anisotropy:

- Jet-medium photons shows negative v_2 .
- Theoretical predictions are inconclusive.

Chatterjee, Frodermann, Heinz, Srivastava; PRL 96 (2006)
 Turbide, Gale, Fries; PRL 96 (2006)

Not promising, so far...



Turbide et al. PRC 77 (2008)

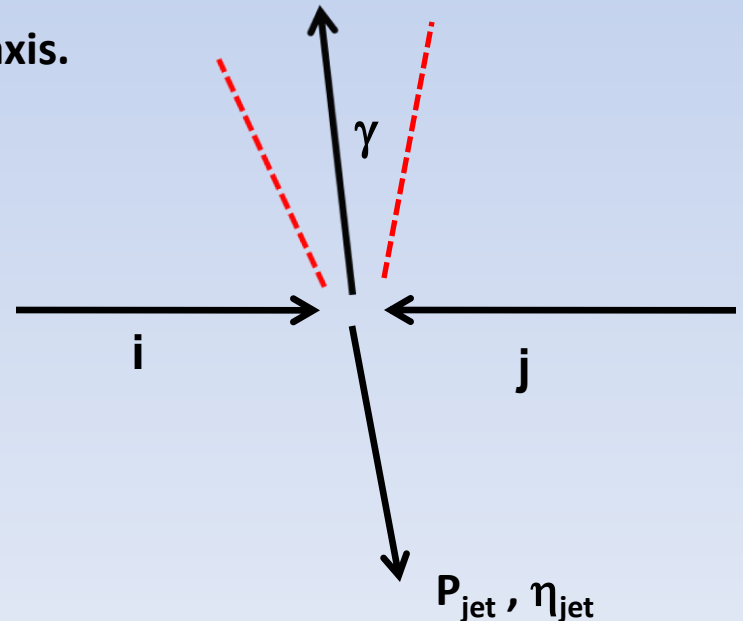
Jet - tagged photon measurement:

▪ Motivation:

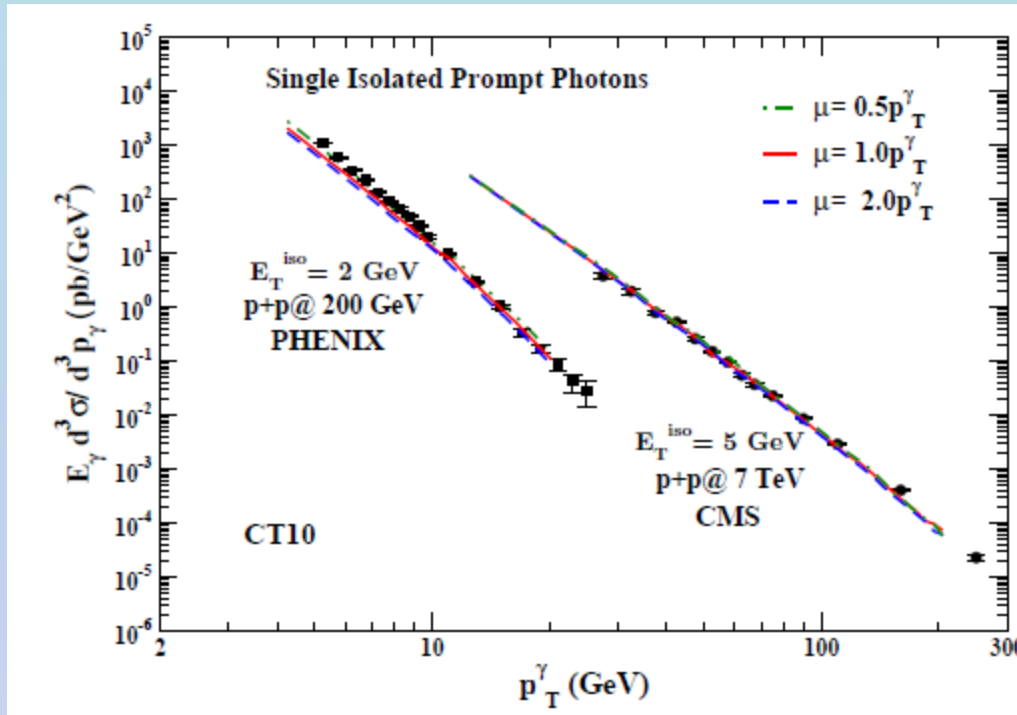
- ❖ The back-scattered photons has strong correlation with the parent jet momentum
- ❖ Also has strong momentum correlation with the away-side leading jet.

▪ Strategy:

- Fix the momentum and rapidity of the away-side leading jet.
- Study the photons, very close to the away-side jet axis.
- The initial hard photons are treated as background.
- Get rid of thermal and pre-equilibrium photons.



Estimation of Background



S D, arXiv: 1305.0624

JETPHOX:
S.Catani, M.
Fontannaz, J. Ph.
Guillet, E. Pilon
JHEP 05 (2002)
028

- Photons from initial hard collision + Fragmentation of jets
- The background is calculated from JETPHOX-ROOT version
- EPS09 nuclear pdf is used for A+A collisions
- No E_T cut for A+A case

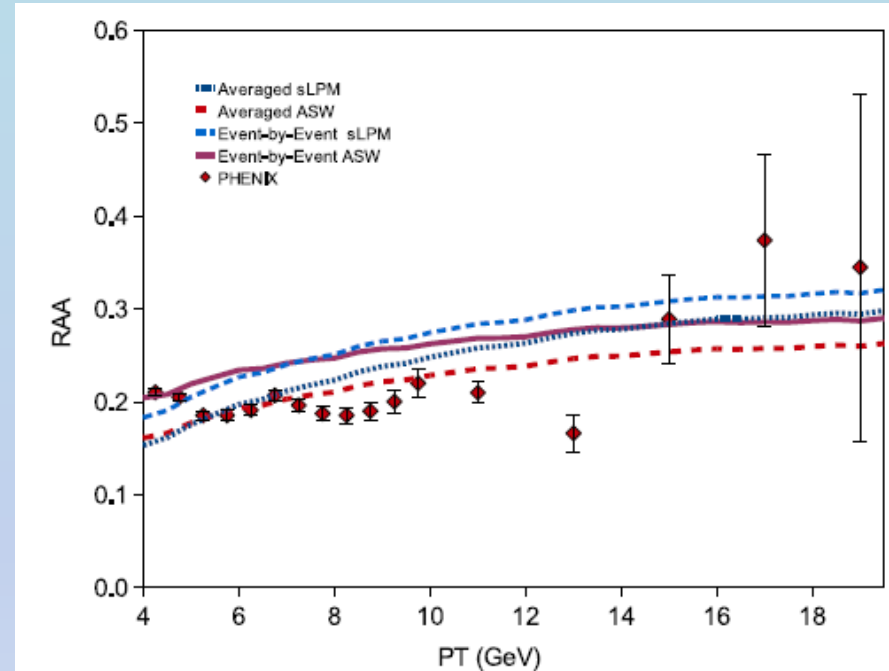
The energy loss model : ppm

We have used the fireball model by:

R Rodriguez, R J Fries, E Ramirez
PLB 693, 108

Path travelled by the jet:

$$I_{\beta}(\vec{r}, \varphi) = \int d\tau \tau^{\beta} \rho(\vec{r} + \tau \vec{e}_{\varphi})$$



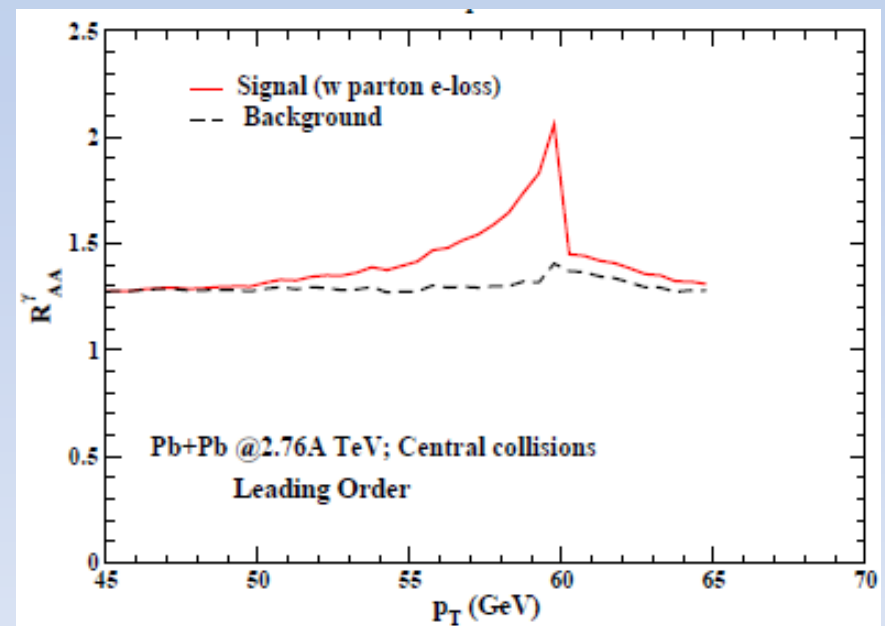
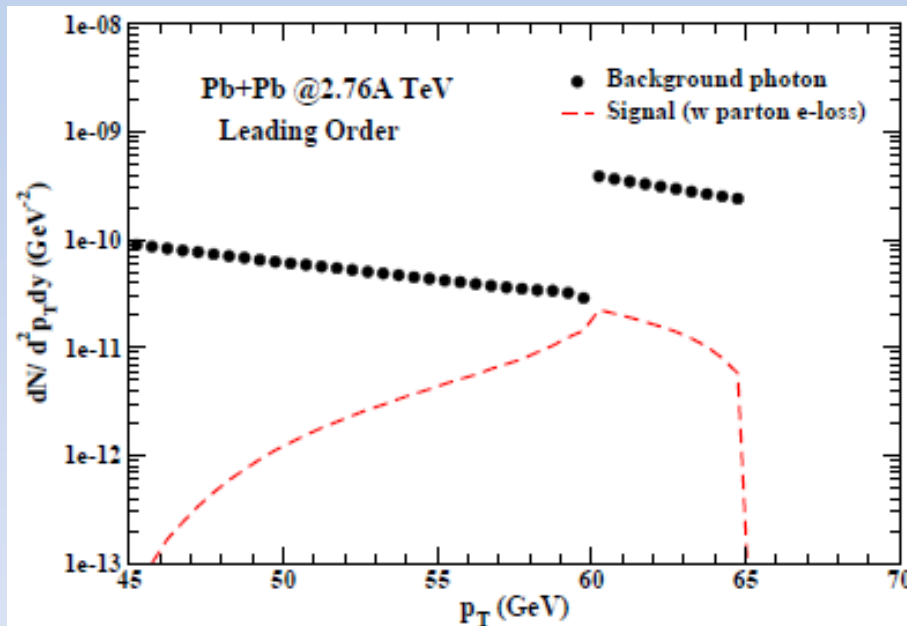
We have used LPM type of energy loss; $\beta=1$ $\Delta E = C_{LPM} I_1(r, \varphi)$

The coeff. C_{LPM} is determined from the fitting of R_{AA} of hadrons

QCD back-scattering photons at LHC : LO

$$R_{AA} : (\text{Signal} + \text{Background})_{AA} / N_{\text{coll}} \times (\text{Background})_{pp}$$

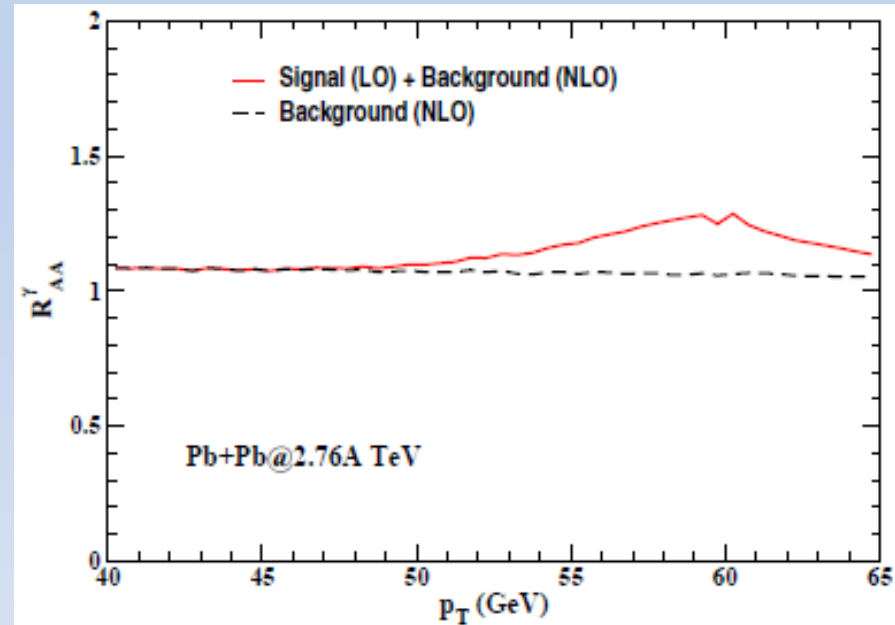
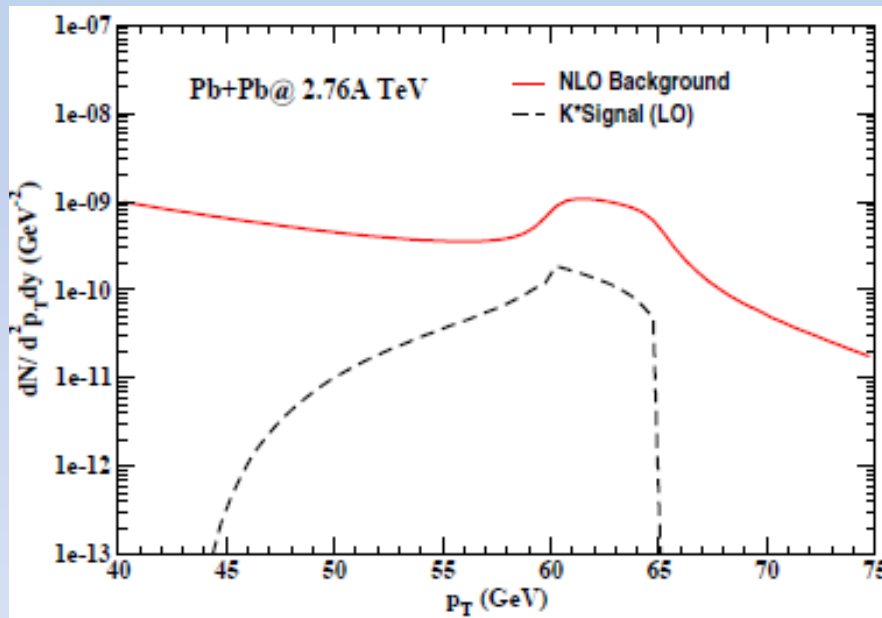
- For central Pb+Pb collisions at 2.76 TeV at mid rapidity
- Photons opposite to the 60-65 GeV jet within ± 15 degrees
- The quarks suffer energy loss before conversion



A clear back scattering peak is seen, despite finite trigger interval

QCD back-scattering photons at LHC : NLO

- Back ground is calculated in the Next-to Leading order
- Kinematics of jet-conversion is still leading order
- Parton energy loss is accounted



Signal weakens but survives at NLO

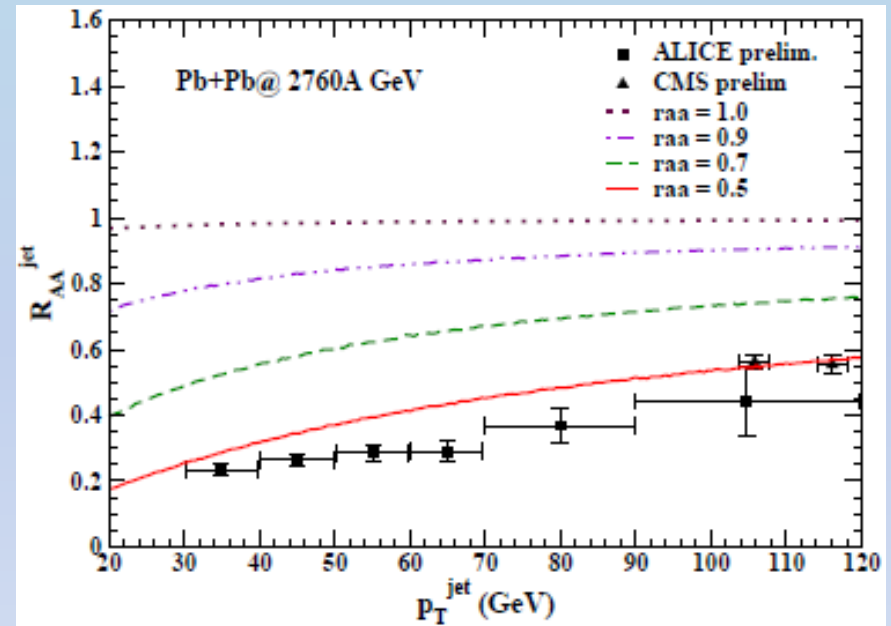
Effect of trigger jet energy loss

❖ Recent measurements at LHC (arXiv:1304.5945) suggest a strong suppression of trigger jets.

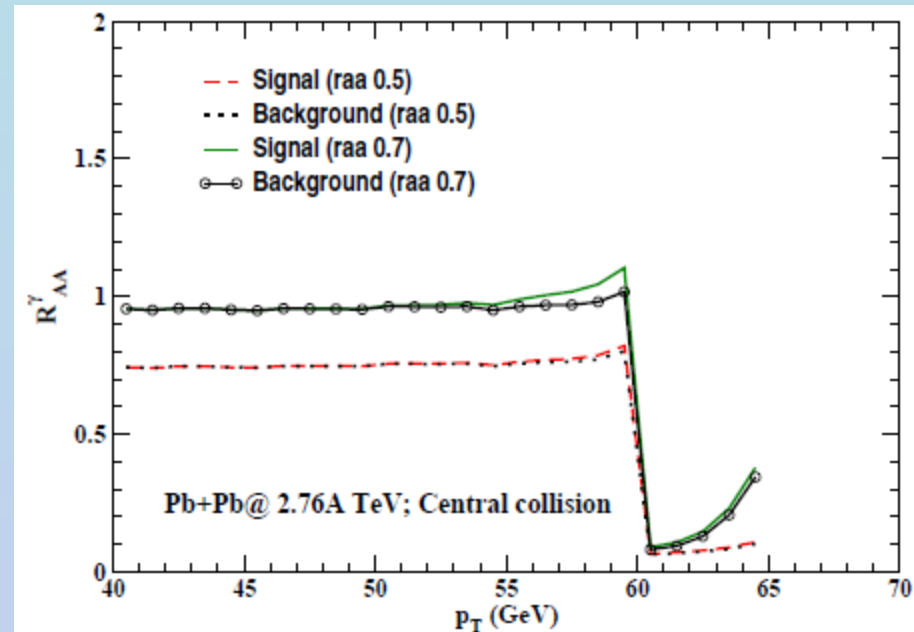
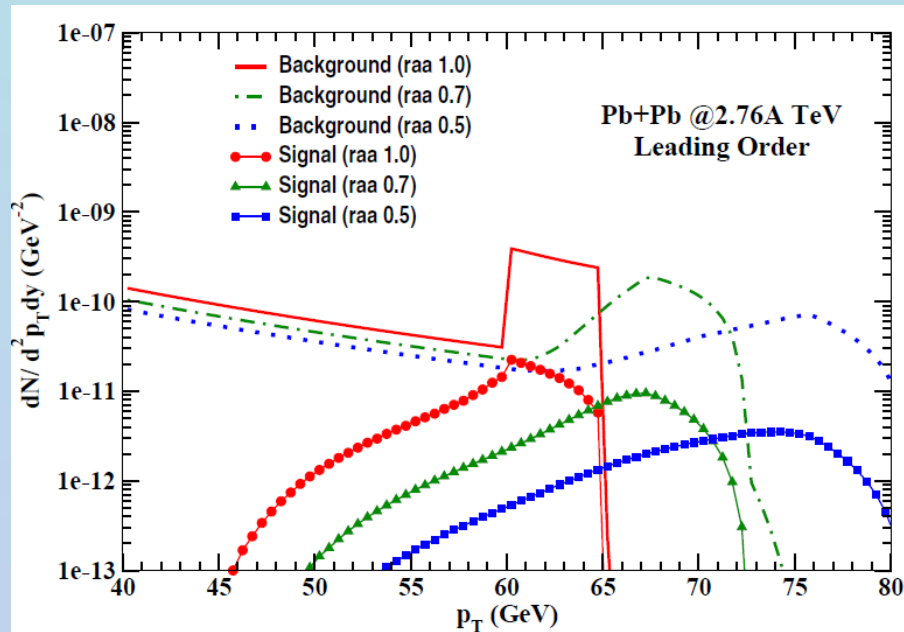
We model the trigger jet energy loss as:

$$\frac{dE_T}{d\tau} = -\hat{r} \ln\left(\frac{E_T}{\Lambda}\right)$$

\hat{r} is proportional to local entropy density, $\Lambda = 0.2$ GeV



➤ Trigger jet energy loss affects the Background (direct and fragmentation photon) and as well as the Signal



- Signal and Background are calculated for Leading order kinematics
- Both trigger jet and parton energy loss are taken into account
- Trigger jet energy loss tends to wash out the strong correlation with parent jet

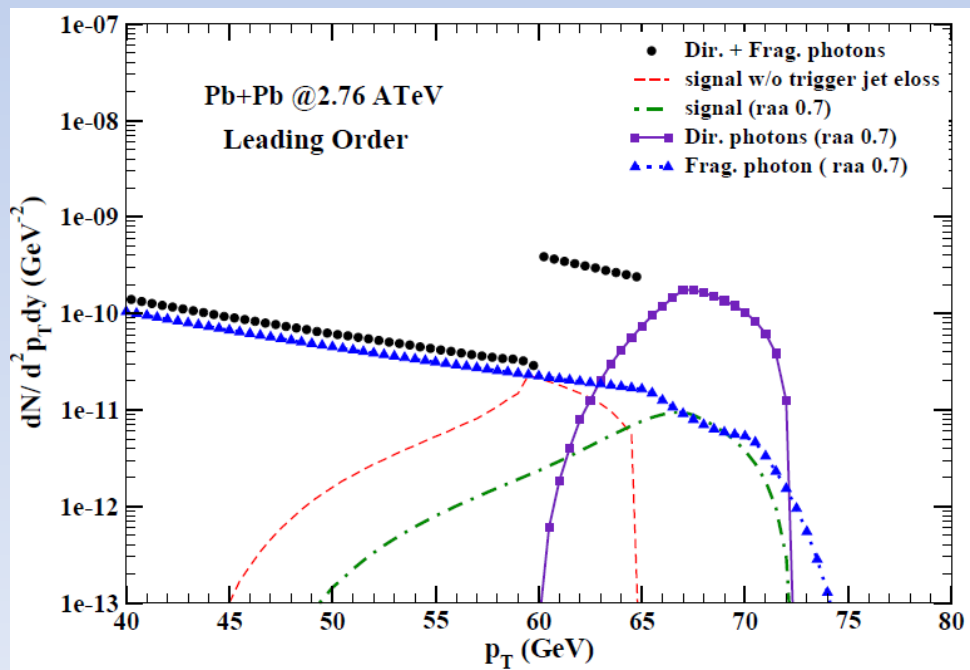
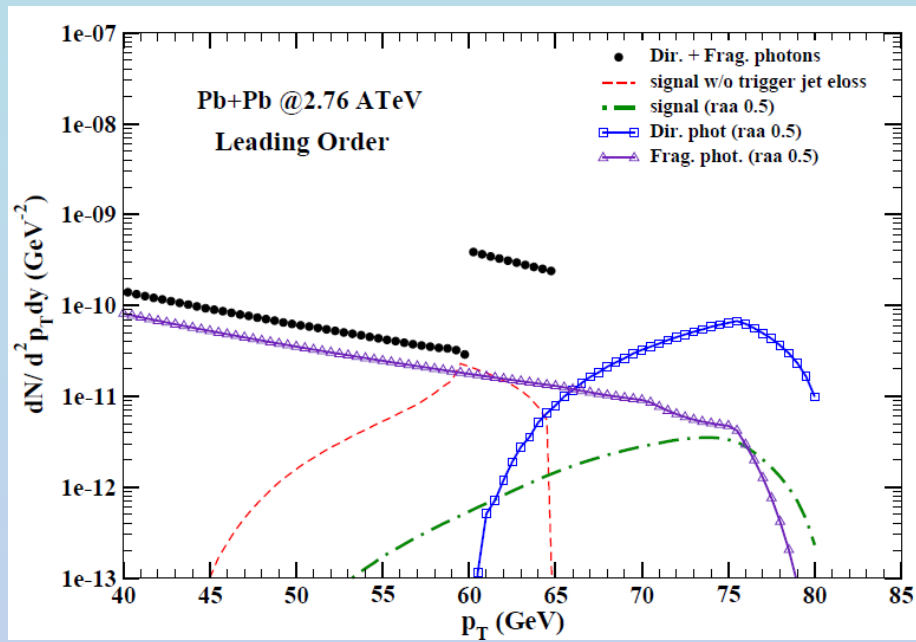
Summary & Conclusion:

- Jet-medium back scattering photon is an important signature of thermalized matter created in heavy ion collisions.
- We propose the use of trigger jet to identify this particular component from the back ground of hard photons.
- Jet-medium photons shows characteristic enhancement in the nuclear modification factor, just below the trigger jet energy.
- The enhancement is clearly visible in Leading order but weakens for the radiative corrections to the process and trigger jet energy loss.
- The shift of the peak from the trigger jet window is sensitive to the parton and trigger jet energy loss in the matter.
- Separation of this signal from other photon sources depends crucially on the initial trigger jet energy estimation.



Thank you





$$f_q^{\mathcal{T}_j}(\mathbf{p}_q, x) = \frac{(2\pi)^3}{g_q \tau p_T} \delta(y - \eta) \rho(\tau, \mathbf{x}_\perp^0) \\ \times \int_{\mathcal{T}_j} dE_T dy_j d\phi_j E_q \frac{dN}{d^3 p_q dE_T dy_j d\phi_j} \Big|_{\substack{\mathbf{p}_q^0 = \mathbf{p}_q + \Delta \mathbf{p}_q \\ E_T^0 = E_T + \Delta E_T}}$$