



# COMPOSITE HIGGS BOSON THEORIES

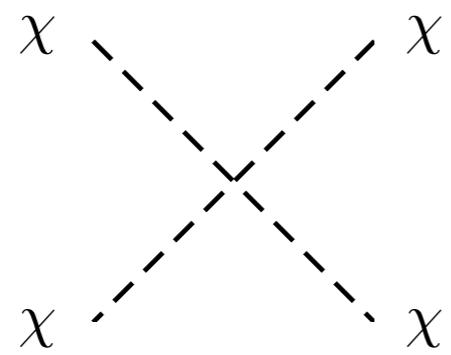
Roberto Contino  
EPFL, Lausanne & CERN

Based on: RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

CLIC Detector and Physics Collaboration Meeting, CERN 1-2 October 2013

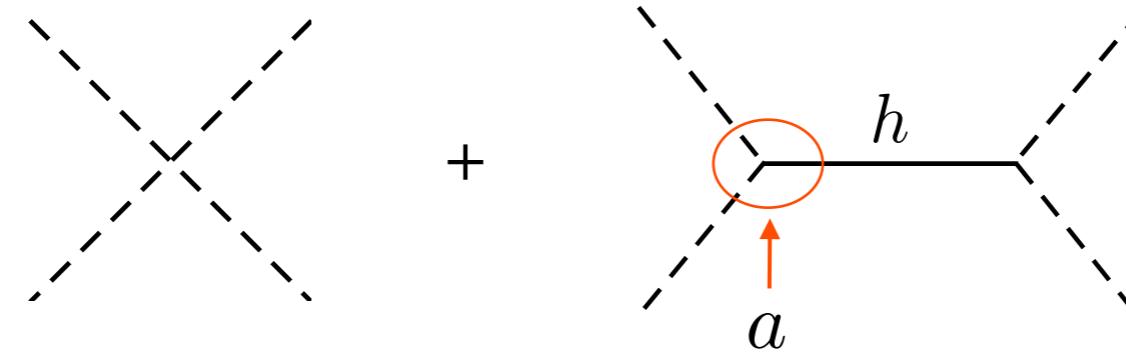
# Strong vs Weak EWSB

In the {SM-H}



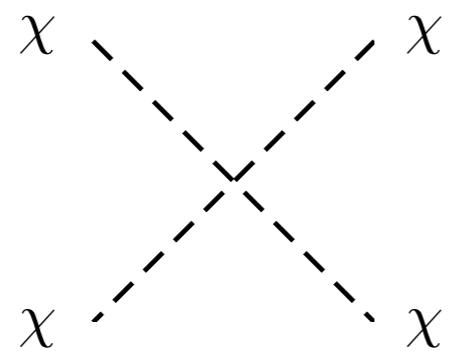
$$A(W_L W_L \rightarrow W_L W_L) = A(\chi\chi \rightarrow \chi\chi) \sim \frac{E^2}{v^2} \equiv g^2(E)$$

In the {SM-H} + H



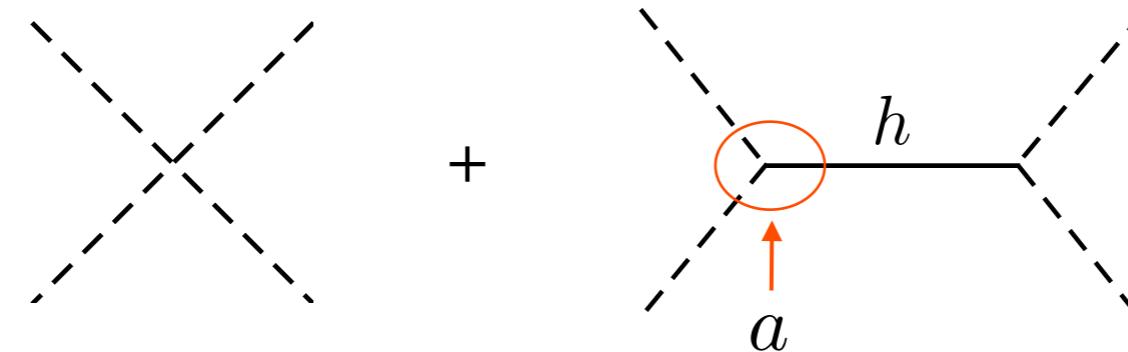
$$A \sim \frac{E^2}{v^2} (1 - a^2) - a^2 \frac{m_h^2}{v^2} \frac{s}{s - m_h^2}$$

In the {SM-H}



$$A(W_L W_L \rightarrow W_L W_L) = A(\chi\chi \rightarrow \chi\chi) \sim \frac{E^2}{v^2} \equiv g^2(E)$$

In the {SM-H} + H



$$A \sim \frac{E^2}{v^2} (1 - a^2) - a^2 \frac{m_h^2}{v^2} \frac{s}{s - m_h^2}$$

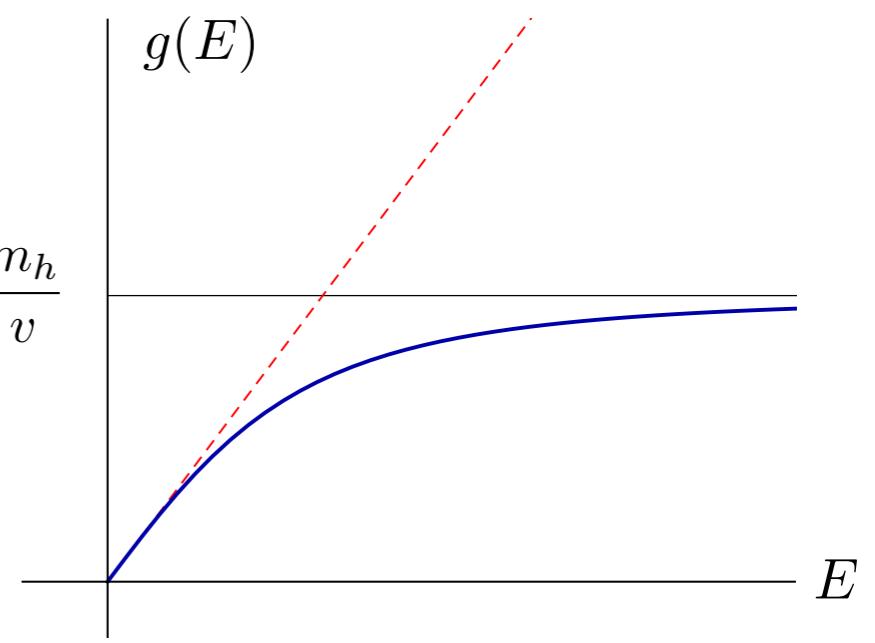
$$= 0$$

Elementary Higgs:

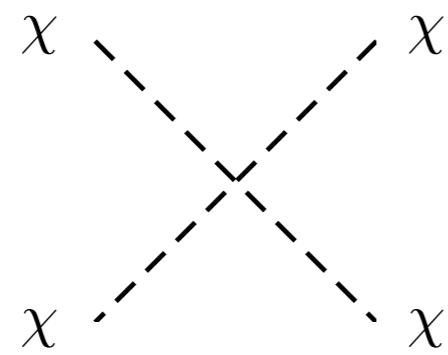
$$a = 1$$

weak  $\rightarrow$

$$\frac{m_h}{v}$$

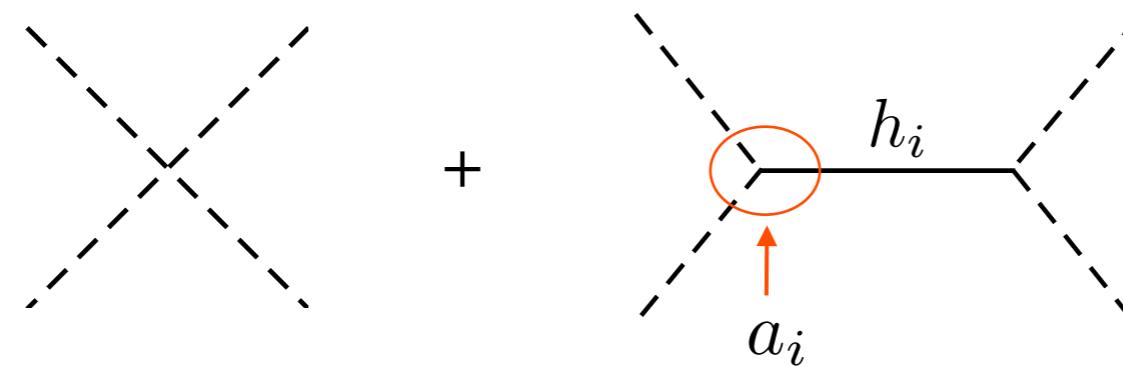


In the  $\{\text{SM-H}\}$



$$A(W_L W_L \rightarrow W_L W_L) = A(\chi\chi \rightarrow \chi\chi) \sim \frac{E^2}{v^2} \equiv g^2(E)$$

In the  $\{\text{SM-H}\} + \text{H}$



$$A \sim \frac{E^2}{v^2} \left( 1 - \sum_i a_i^2 \right) + \dots$$

$\underbrace{\phantom{1 - \sum_i a_i^2}}$

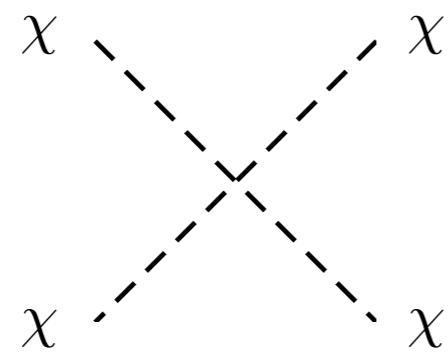
$$= 0$$

Elementary Higgses:  
(more than one)

- $\delta a_i \sim O(1)$  possible

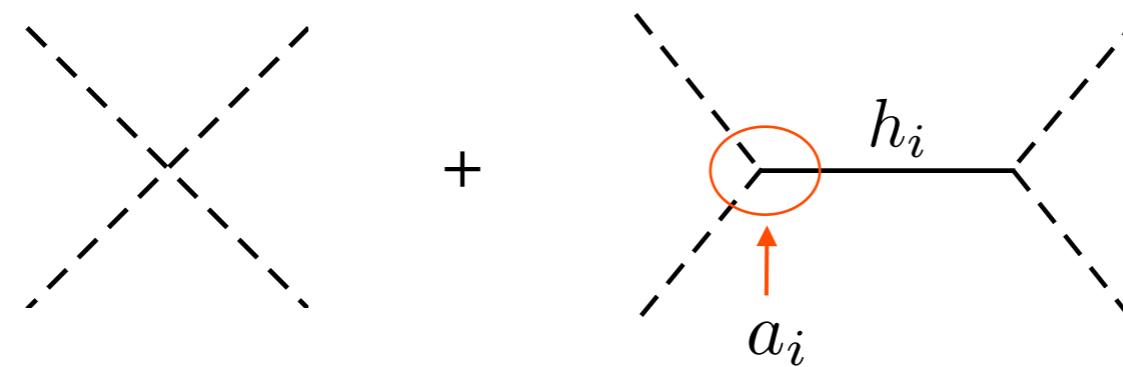
- sum rule:  $\sum_i a_i = 1$

In the {SM-H}



$$A(W_L W_L \rightarrow W_L W_L) = A(\chi\chi \rightarrow \chi\chi) \sim \frac{E^2}{v^2} \equiv g^2(E)$$

In the {SM-H} + H



$$A \sim \frac{E^2}{v^2} \left( 1 - \sum_i a_i^2 \right) + \dots$$

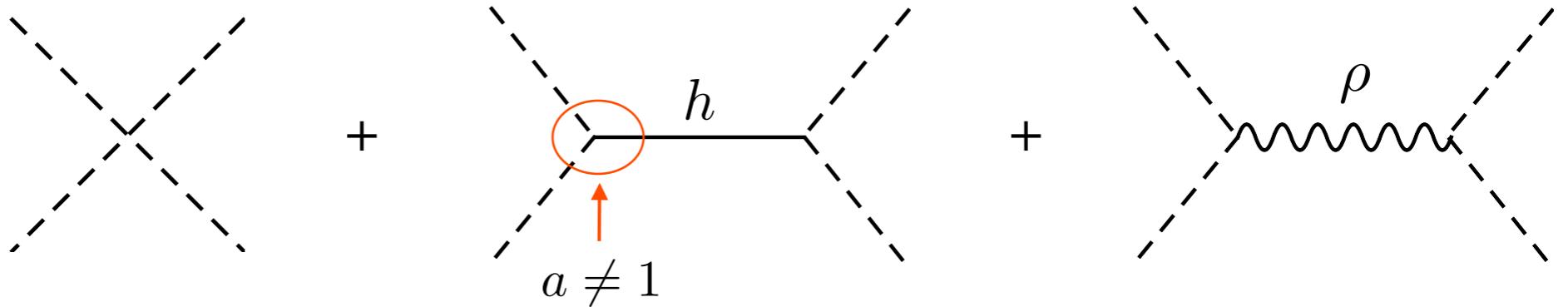
$\underbrace{\phantom{1 - \sum_i a_i^2}}$

$$= 0$$

Elementary Higgses:  
(more than one)

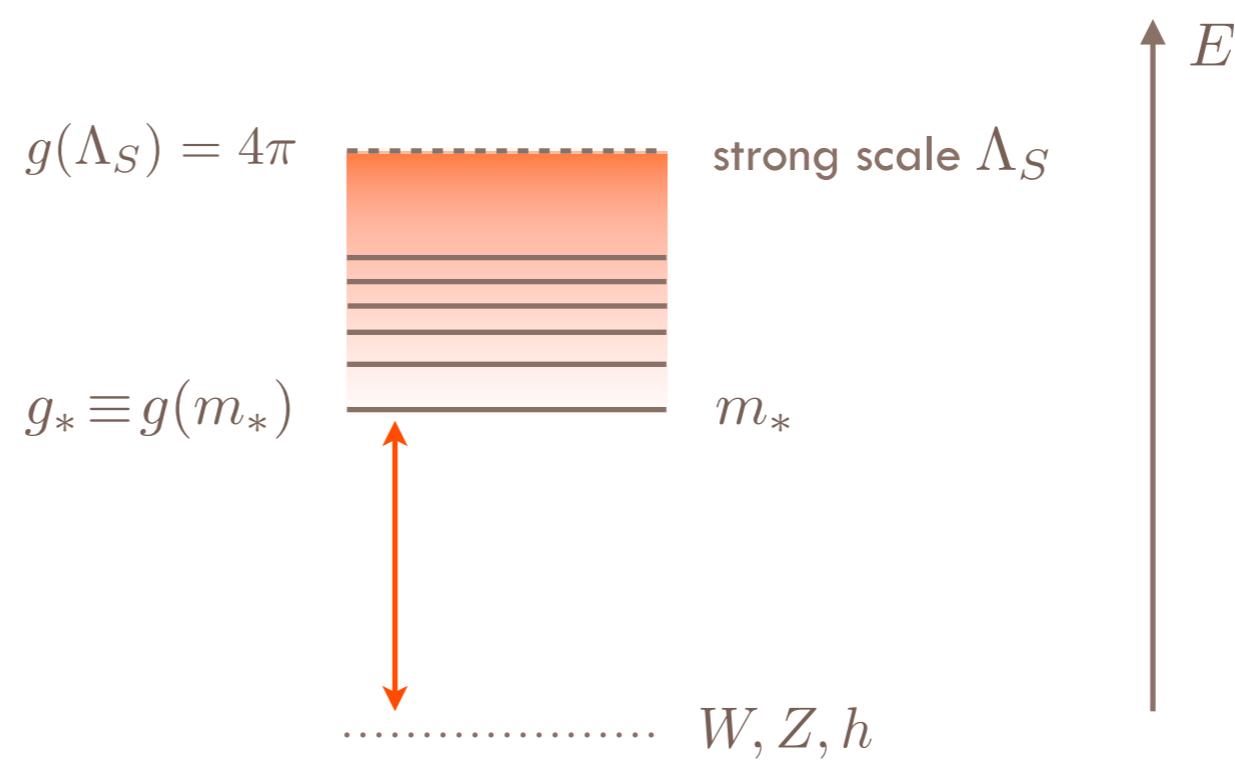
- $\delta a_i \sim O(1)$  possible
- sum rule:  $\sum_i a_i = 1$

Composite Higgs:



coupling strength grows with energy and saturates at  $g_* \lesssim 4\pi$

Energy cartoon:



Analogy with  $\pi\pi$  scattering in QCD:  $h \leftrightarrow \sigma$



Q: why light and narrow ?

Analogy with  $\pi\pi$  scattering in QCD:  $h \leftrightarrow \sigma$



Q: why light and narrow ?

A: the Higgs is itself a (pseudo) NG boson [ Georgi & Kaplan, '80 ]

ex:  $\frac{SO(5)}{SO(4)} \rightarrow$  4 NGBs transforming as a (2,2) of  $SO(4)$

[ Agashe, RC, Pomarol  
NPB 719 (2005) 165 ]

$$f^2 \left| \partial_\mu e^{i\pi/f} \right|^2 = (\partial\pi)^2 + \frac{(\pi\partial\pi)^2}{f^2} + \frac{\pi^2(\pi\partial\pi)^2}{f^4} + \dots$$

Analogy with  $\pi\pi$  scattering in QCD:  $h \leftrightarrow \sigma$



Q: why light and narrow ?

A: the Higgs is itself a (pseudo) NG boson [ Georgi & Kaplan, '80 ]

ex:  $\frac{SO(5)}{SO(4)}$  → 4 NGBs transforming as a (2,2) of SO(4) [ Agashe, RC, Pomarol  
NPB 719 (2005) 165 ]

$$f^2 \left| \partial_\mu e^{i\pi/f} \right|^2 = |D_\mu H|^2 + \frac{c_H}{2f^2} [\partial_\mu (H^\dagger H)]^2 + \frac{c'_H}{2f^4} (H^\dagger H) [\partial_\mu (H^\dagger H)]^2 + \dots$$

[ Giudice et al. JHEP 0706 (2007) 045 ]

Analogy with  $\pi\pi$  scattering in QCD:  $h \leftrightarrow \sigma$



Q: why light and narrow ?

A: the Higgs is itself a (pseudo) NG boson [ Georgi & Kaplan, '80 ]

ex:  $\frac{SO(5)}{SO(4)} \rightarrow$  4 NGBs transforming as a (2,2) of  $SO(4)$  [ Agashe, RC, Pomarol  
NPB 719 (2005) 165 ]

$$f^2 \left| \partial_\mu e^{i\pi/f} \right|^2 = |D_\mu H|^2 + \frac{c_H}{2f^2} [\partial_\mu (H^\dagger H)]^2 + \frac{c'_H}{2f^4} (H^\dagger H) [\partial_\mu (H^\dagger H)]^2 + \dots$$

[ Giudice et al. JHEP 0706 (2007) 045 ]

1.  $O(v^2/f^2)$  shifts in tree-level Higgs couplings. Ex:  $a = 1 - c_H \left(\frac{v}{f}\right)^2 + \dots$

Analogy with  $\pi\pi$  scattering in QCD:  $h \leftrightarrow \sigma$



Q: why light and narrow ?

A: the Higgs is itself a (pseudo) NG boson [ Georgi & Kaplan, '80 ]

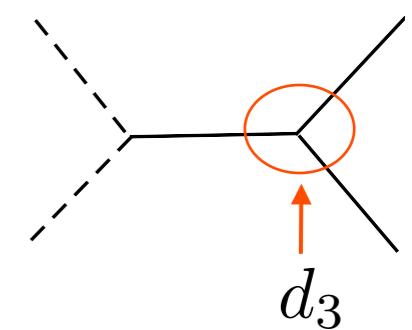
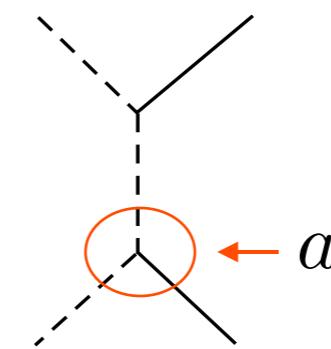
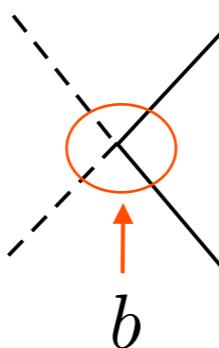
ex:  $\frac{SO(5)}{SO(4)} \rightarrow 4 \text{ NGBs transforming as a (2,2) of } SO(4)$  [ Agashe, RC, Pomarol  
NPB 719 (2005) 165 ]

$$f^2 \left| \partial_\mu e^{i\pi/f} \right|^2 = |D_\mu H|^2 + \frac{c_H}{2f^2} [\partial_\mu (H^\dagger H)]^2 + \frac{c'_H}{2f^4} (H^\dagger H) [\partial_\mu (H^\dagger H)]^2 + \dots$$

[ Giudice et al. JHEP 0706 (2007) 045 ]

## 2. Scatterings involving the Higgs also grow with energy

$$A(WW \rightarrow hh) \sim \frac{s}{v^2} (a^2 - b)$$



## How to test Higgs compositeness

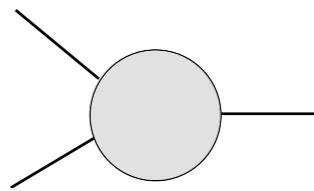
1. **Direct:** Reach energy threshold for direct production of new resonances
2. **Indirect:** Precision measurement of low-energy quantities

## How to test Higgs compositeness

1. **Direct:** Reach energy threshold for direct production of new resonances

2. **Indirect:** Precision measurement of low-energy quantities

i) virtual corrections to single-Higgs processes

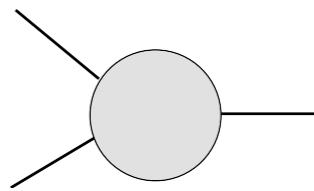


# How to test Higgs compositeness

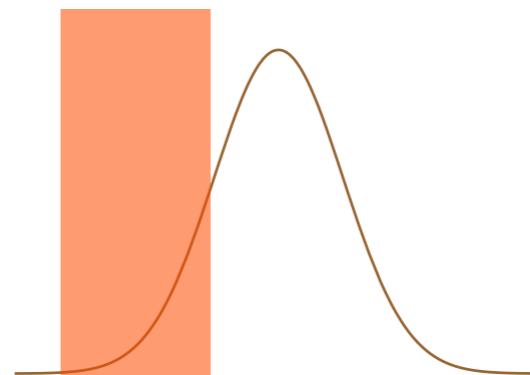
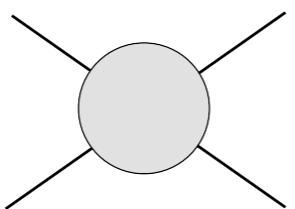
1. **Direct:** Reach energy threshold for direct production of new resonances

2. **Indirect:** Precision measurement of low-energy quantities

i) virtual corrections to single-Higgs processes

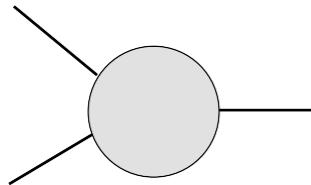


ii) tails in scattering amplitudes

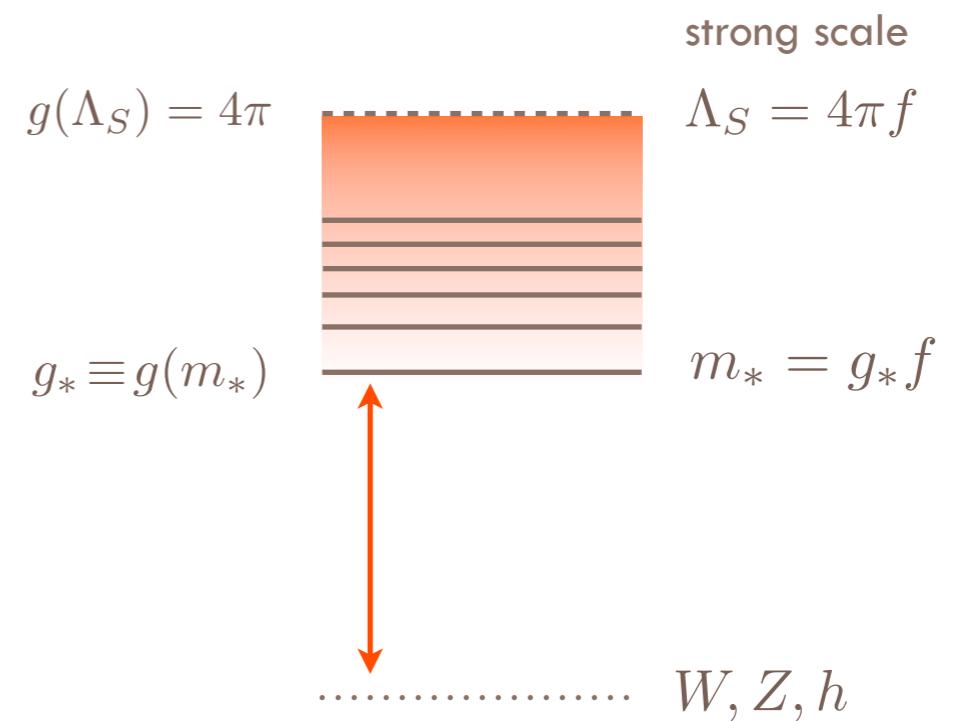


# Corrections to Higgs couplings

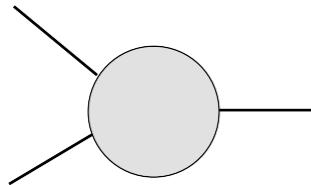
*Precision measurement of Higgs couplings  
can give an appraisal of the strength of  
the underlying interactions*



$$\frac{\delta \mathcal{O}}{\mathcal{O}} \sim \frac{v^2}{f^2} + \frac{g_*^2 v^2}{m_*^2}$$

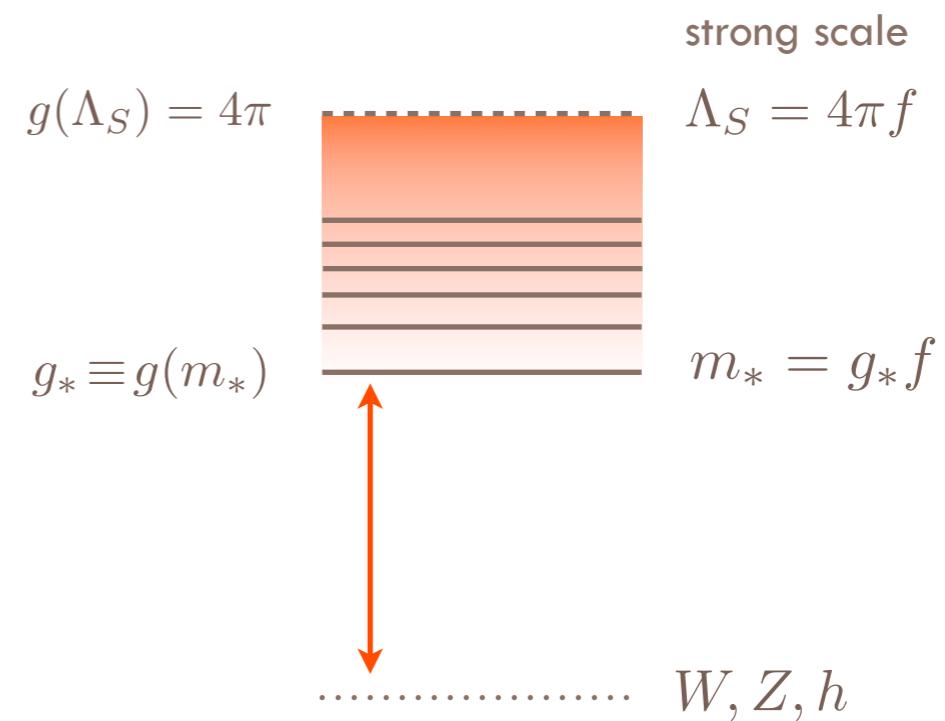


Precision measurement of Higgs couplings can give an appraisal of the strength of the underlying interactions

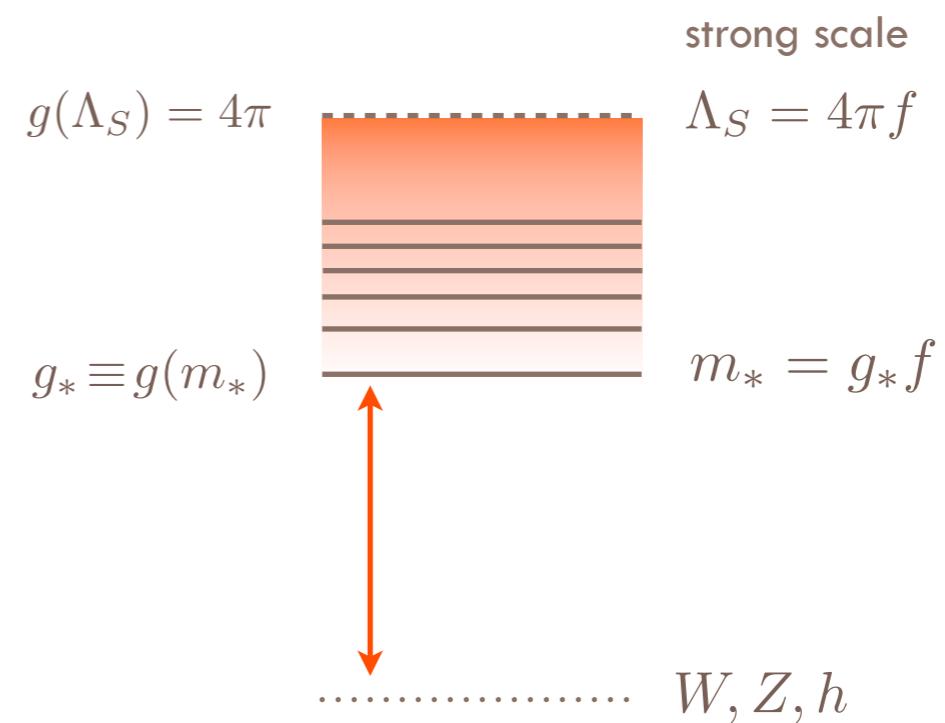
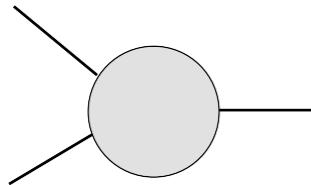


$$\frac{\delta \mathcal{O}}{\mathcal{O}} \sim \frac{v^2}{f^2} + \frac{g_*^2 v^2}{m_*^2}$$

from NL sigma model



Precision measurement of Higgs couplings can give an appraisal of the strength of the underlying interactions

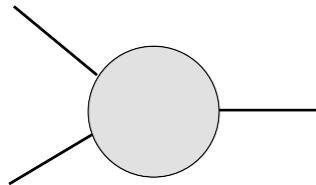


contribution of resonances

$$\frac{\delta \mathcal{O}}{\mathcal{O}} \sim \frac{v^2}{f^2} + \frac{g_*^2 v^2}{m_*^2}$$

from NL sigma model

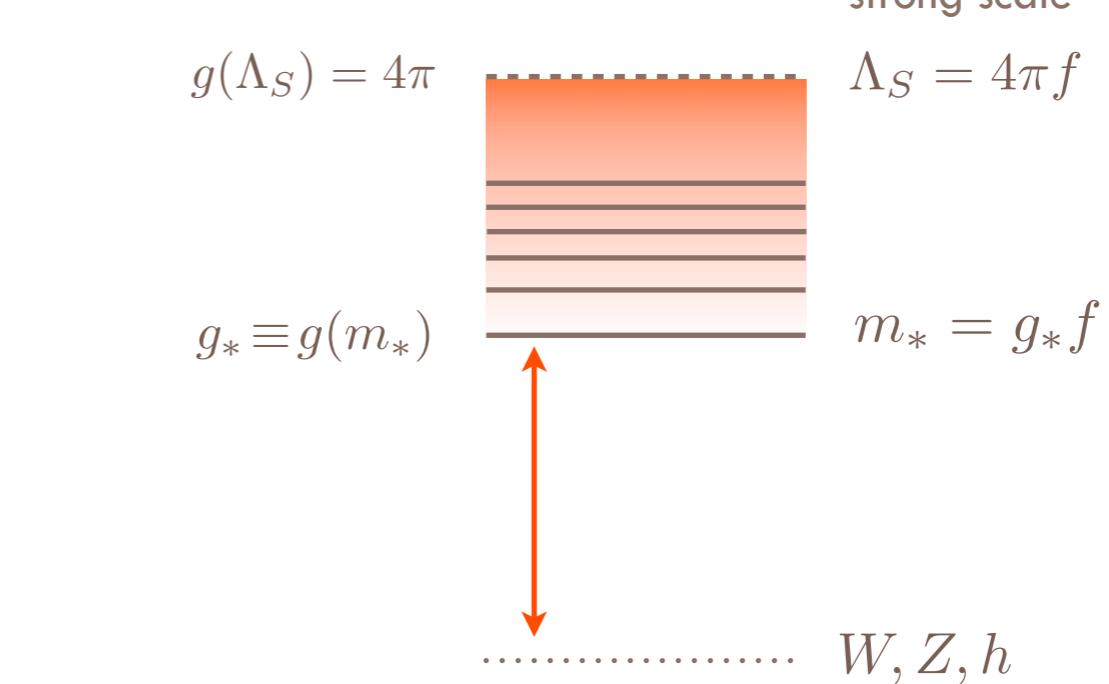
Precision measurement of Higgs couplings can give an appraisal of the strength of the underlying interactions



contribution of resonances

$$\frac{\delta \mathcal{O}}{\mathcal{O}} \sim \left( \frac{v^2}{f^2} + \frac{g_*^2 v^2}{m_*^2} \right)$$

from NL sigma model



$$g(\Lambda_S) = 4\pi$$

strong scale

$$\Lambda_S = 4\pi f$$

$$g_* \equiv g(m_*)$$

$$m_* = g_* f$$

$$W, Z, h$$

Suppose we find:

$$\frac{\delta \mathcal{O}}{\mathcal{O}} \Big|_{exp} = \delta_{\mathcal{O}}^{exp} \quad \rightarrow$$

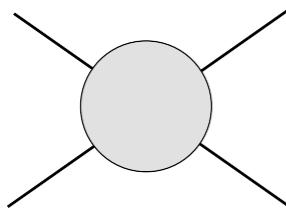
$$g_* > \sqrt{\delta_{\mathcal{O}}^{exp}} \frac{M}{v}$$

$$m_* > M$$

(from direct searches)

# Tails in scattering amplitudes

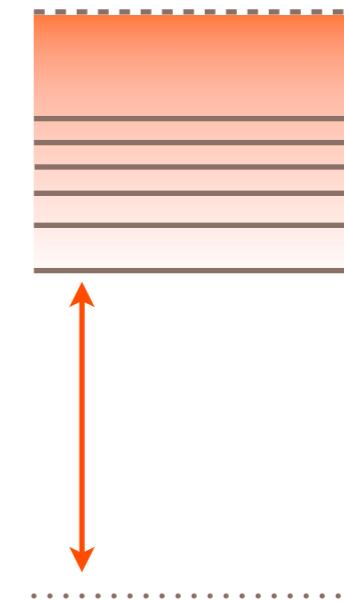
*Precision measurement of scattering amplitudes can give an appraisal of the strength of the underlying interactions*



$$\mathcal{A}(2 \rightarrow 2) = \delta_{hh} \frac{E^2}{v^2} \left( 1 + O\left(\frac{E^2}{m_*^2}\right) \right)$$

$$g(\Lambda_S) = 4\pi$$

$$g_* \equiv g(m_*)$$



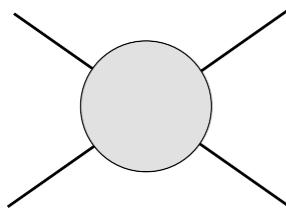
**strong scale**

$$\Lambda_S = 4\pi f$$

$$m_* = g_* f$$

$W, Z, h$

Precision measurement of scattering amplitudes can give an appraisal of the strength of the underlying interactions

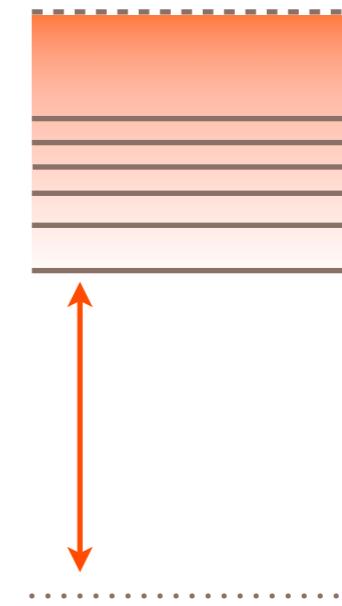


$$\mathcal{A}(2 \rightarrow 2) = \delta_{hh} \frac{E^2}{v^2} \left( 1 + O\left(\frac{E^2}{m_*^2}\right) \right)$$

$\equiv g^2(E)$

$$g(\Lambda_S) = 4\pi$$

$$g_* \equiv g(m_*)$$



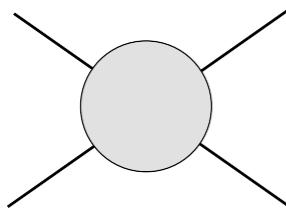
strong scale

$$\Lambda_S = 4\pi f$$

$$m_* = g_* f$$

$W, Z, h$

Precision measurement of scattering amplitudes can give an appraisal of the strength of the underlying interactions



$$\mathcal{A}(2 \rightarrow 2) = \delta_{hh} \frac{E^2}{v^2} \left( 1 + O\left(\frac{E^2}{m_*^2}\right) \right)$$

$\equiv g^2(E)$

Suppose we find:

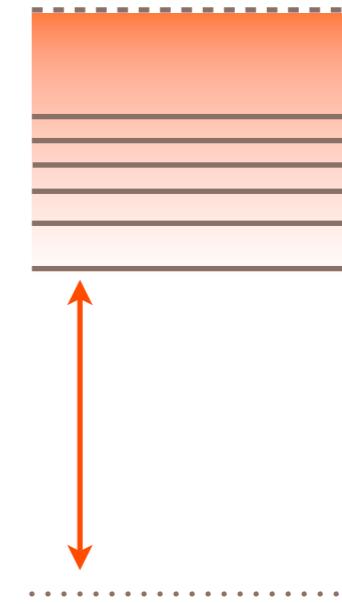
$$\delta_{hh}^{exp} \neq 0$$



$$g_* > g(E) = \sqrt{\delta_{hh}^{exp}} \frac{E}{v}$$

$$g(\Lambda_S) = 4\pi$$

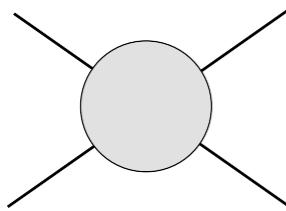
$$g_* \equiv g(m_*)$$



strong scale  
 $\Lambda_S = 4\pi f$

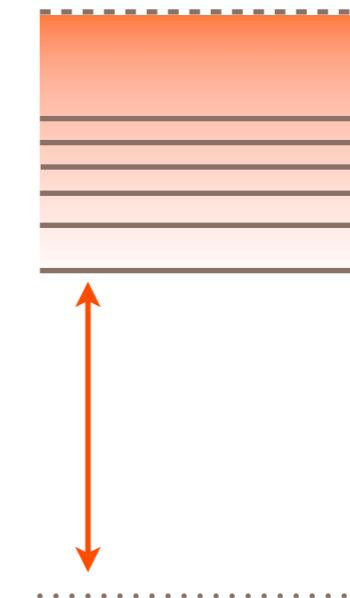
$m_* = g_* f$

Precision measurement of scattering amplitudes can give an appraisal of the strength of the underlying interactions



$$g(\Lambda_S) = 4\pi$$

$$g_* \equiv g(m_*)$$



strong scale

$$\Lambda_S = 4\pi f$$

$$m_* = g_* f$$

$W, Z, h$

Suppose we can bound  $\frac{E^2}{m_*^2} < \epsilon_{hh}$  hence  $m_* > \frac{E}{\sqrt{\epsilon_{hh}}} \equiv M$

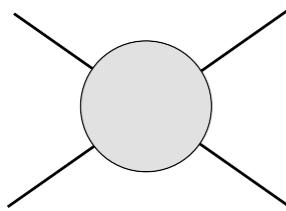
$$\mathcal{A}(2 \rightarrow 2) = \delta_{hh} \frac{E^2}{v^2} \left( 1 + O\left(\frac{E^2}{m_*^2}\right) \right)$$

$$\delta_{hh}^{exp} \neq 0$$



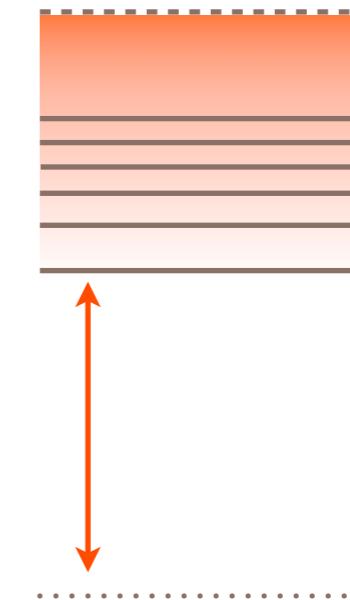
$$g_* > g(E) = \sqrt{\delta_{hh}^{exp}} \frac{E}{v}$$

Precision measurement of scattering amplitudes can give an appraisal of the strength of the underlying interactions



$$g(\Lambda_S) = 4\pi$$

$$g_* \equiv g(m_*)$$



strong scale

$$\Lambda_S = 4\pi f$$

$$m_* = g_* f$$

$W, Z, h$

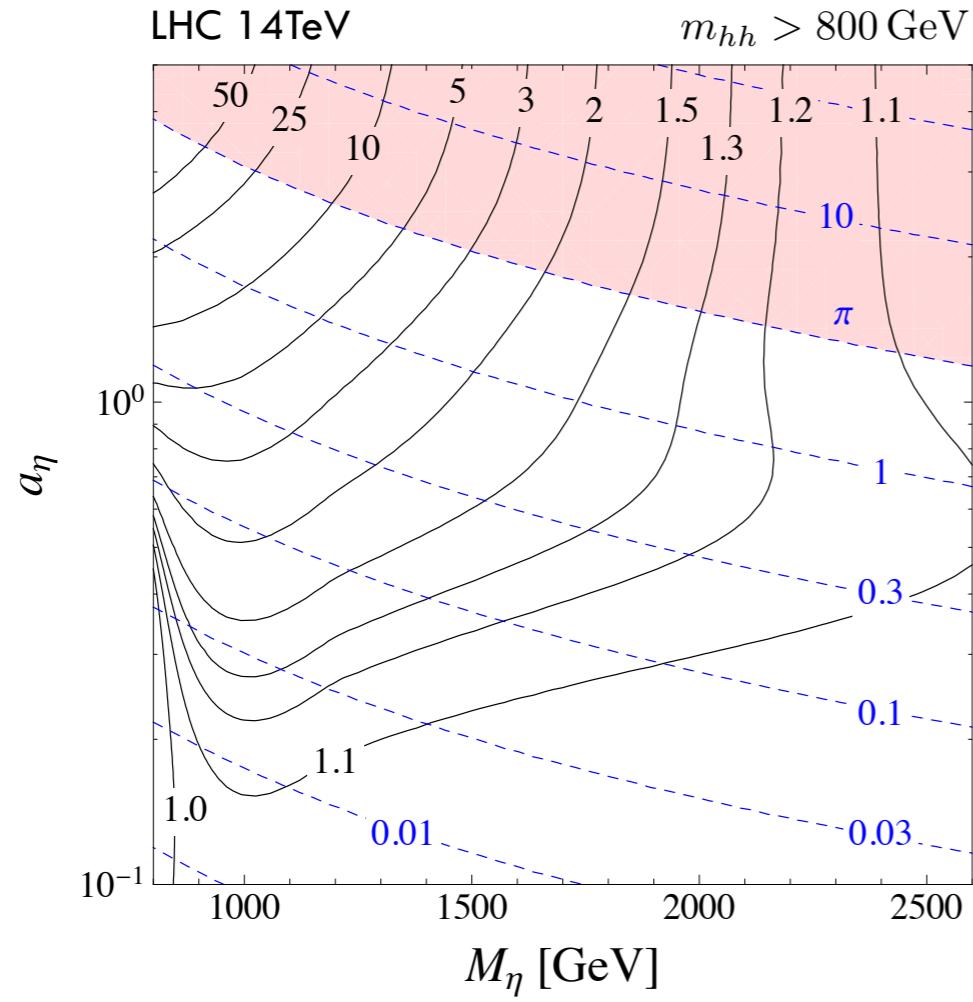
Suppose we can bound  $\frac{E^2}{m_*^2} < \epsilon_{hh}$  hence  $m_* > \frac{E}{\sqrt{\epsilon_{hh}}} \equiv M$

$$\mathcal{A}(2 \rightarrow 2) = \delta_{hh} \frac{E^2}{v^2} \left( 1 + O\left(\frac{E^2}{m_*^2}\right) \right)$$

$\delta_{hh}^{exp} \neq 0 \quad \rightarrow \quad g_* > g(E) = \sqrt{\delta_{hh}^{exp}} \frac{E}{v}$

then we get the stronger limit

$$g_* > g(M) = \sqrt{\frac{\delta_{hh}^{exp}}{\epsilon_{hh}}} \frac{E}{v}$$



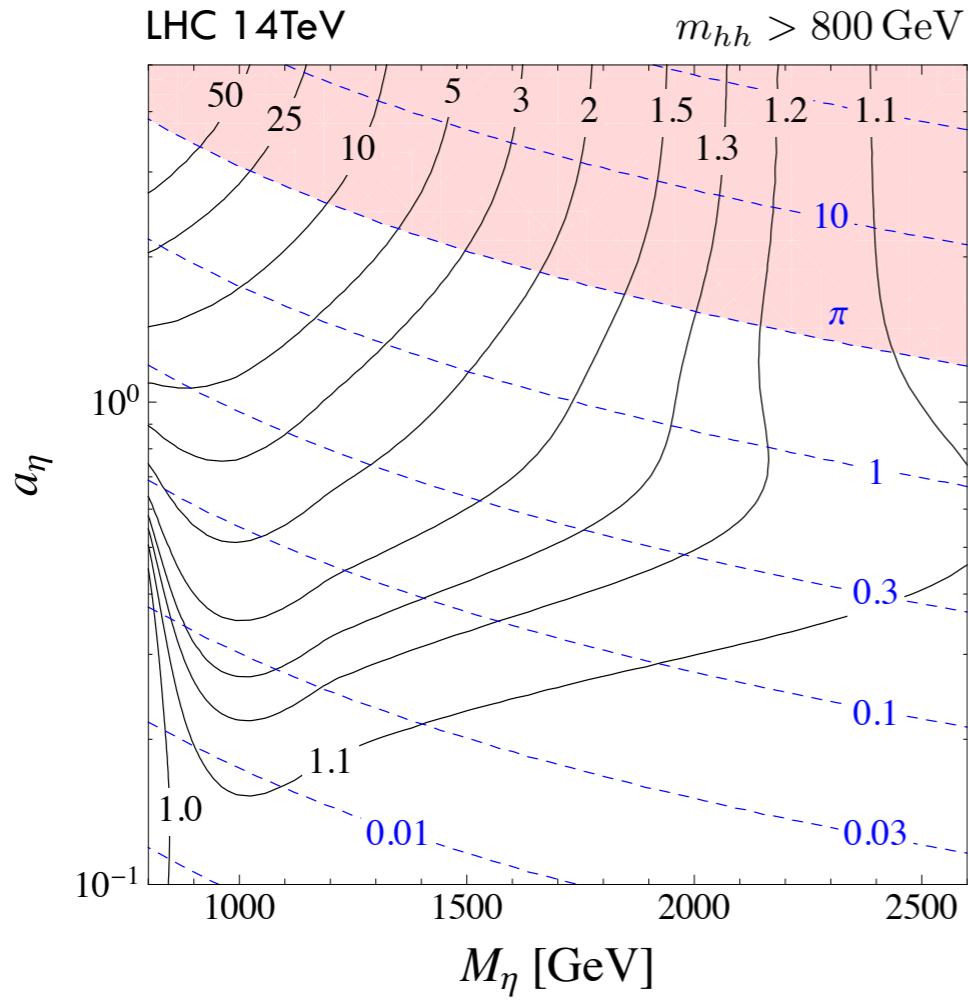
$$R = \frac{\sigma(pp \rightarrow hhjj)}{\sigma(pp \rightarrow hhjj)|_{LET}}$$

$$\mathcal{L} = \frac{a_\eta}{2f} \eta (\partial_\mu \pi)^2 + \dots$$

[ RC, Marzocca, Pappadopulo, Rattazzi, JHEP 1110 (2011) 081 ]

$$\mathcal{A}(2 \rightarrow 2) = \delta_{hh} \frac{E^2}{v^2} \left( 1 + O\left(\frac{E^2}{m_*^2}\right) \right)$$

measurement of resonance effects  
gives direct access to strong dynamics



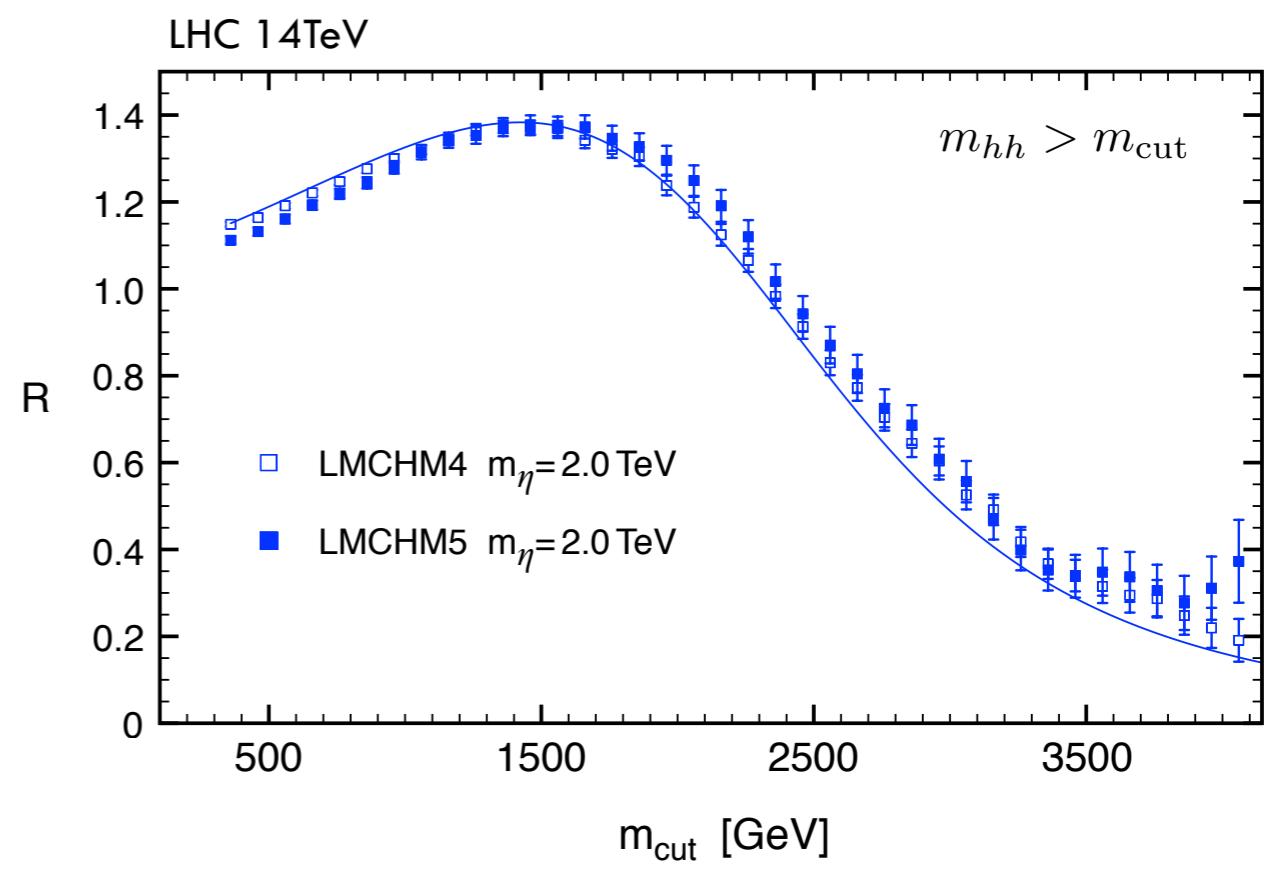
$$\mathcal{A}(2 \rightarrow 2) = \delta_{hh} \frac{E^2}{v^2} \left( 1 + O\left(\frac{E^2}{m_*^2}\right) \right)$$

measurement of resonance effects  
gives direct access to strong dynamics

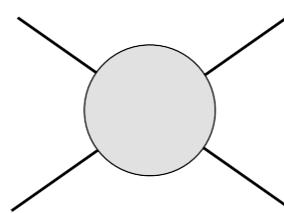
$$R = \frac{\sigma(pp \rightarrow hhjj)}{\sigma(pp \rightarrow hhjj)|_{LET}}$$

$$\mathcal{L} = \frac{a_\eta}{2f} \eta (\partial_\mu \pi)^2 + \dots$$

[ RC, Marzocca, Pappadopulo, Rattazzi, JHEP 1110 (2011) 081 ]



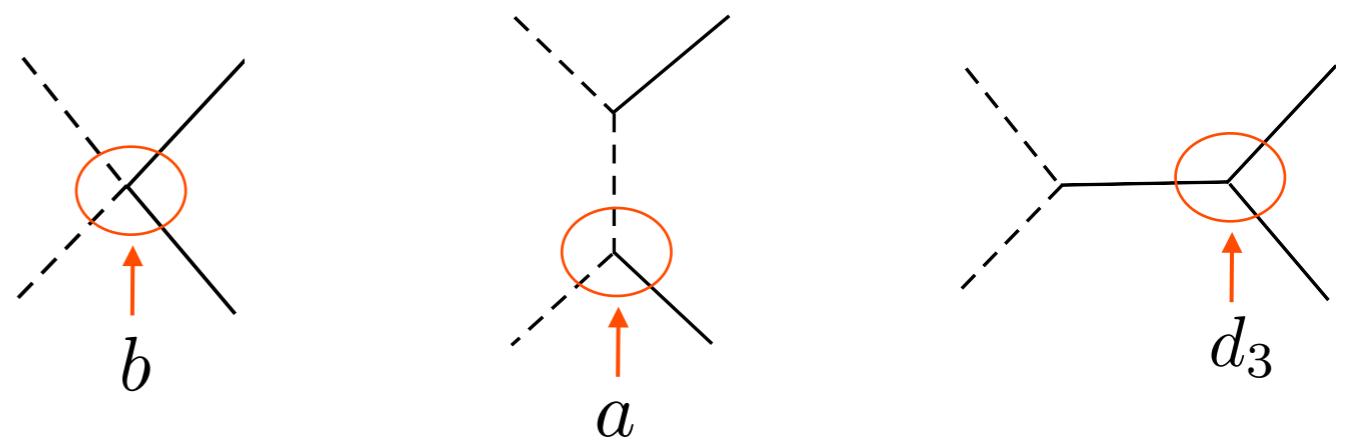
A high-energy  $e^+e^-$  collider  
(such as CLIC 3TeV) can  
provide a clean environment to  
make precision studies of  
scattering amplitudes



[ RC, Grojean, Pappadopulo,  
Rattazzi, Thamm, arXiv:1309.7038 ]

Example:  $WW \rightarrow hh$

$$A(WW \rightarrow hh) \sim \frac{s}{v^2} (a^2 - b)$$



dim 6:  $O_H = \frac{c_H}{2f^2} \partial_\mu |H|^2 \partial^\mu |H|^2$

$$a = 1 - \frac{c_H}{2} \frac{v^2}{f^2} + \left( \frac{3c_H^2}{8} - \frac{c'_H}{4} \right) \frac{v^4}{f^4}$$

dim 8:  $O'_H = \frac{c'_H}{2f^4} |H|^2 \partial_\mu |H|^2 \partial^\mu |H|^2$

$$b = 1 - 2c_H \frac{v^2}{f^2} + \left( 3c_H^2 - \frac{3c'_H}{2} \right) \frac{v^4}{f^4}$$

[ Higgs Effective Lagrangian (SILH basis) ]

Ex:  $\text{SO}(5)/\text{SO}(4)$

In PNGB Higgs theories the whole series in  $H/f$  can be re-summed:

$$a = \sqrt{1 - \xi}$$

$$\xi = \frac{v^2}{f^2}$$

$$b = 1 - 2\xi$$

At dimension-6 level:

$$\Delta b = 2\Delta a^2 (1 + O(\Delta a^2))$$

$$\begin{aligned}\Delta b &\equiv 1 - b \\ \Delta a^2 &\equiv 1 - a^2\end{aligned}$$

Ex:  $\text{SO}(5)/\text{SO}(4)$

In PNGB Higgs theories the whole series in  $H/f$  can be re-summed:

$$a = \sqrt{1 - \xi}$$

$$\xi = \frac{v^2}{f^2}$$

$$b = 1 - 2\xi$$

At dimension-6 level:

$$\Delta b = 2\Delta a^2 (1 + O(\Delta a^2))$$

$$\begin{aligned}\Delta b &\equiv 1 - b \\ \Delta a^2 &\equiv 1 - a^2\end{aligned}$$

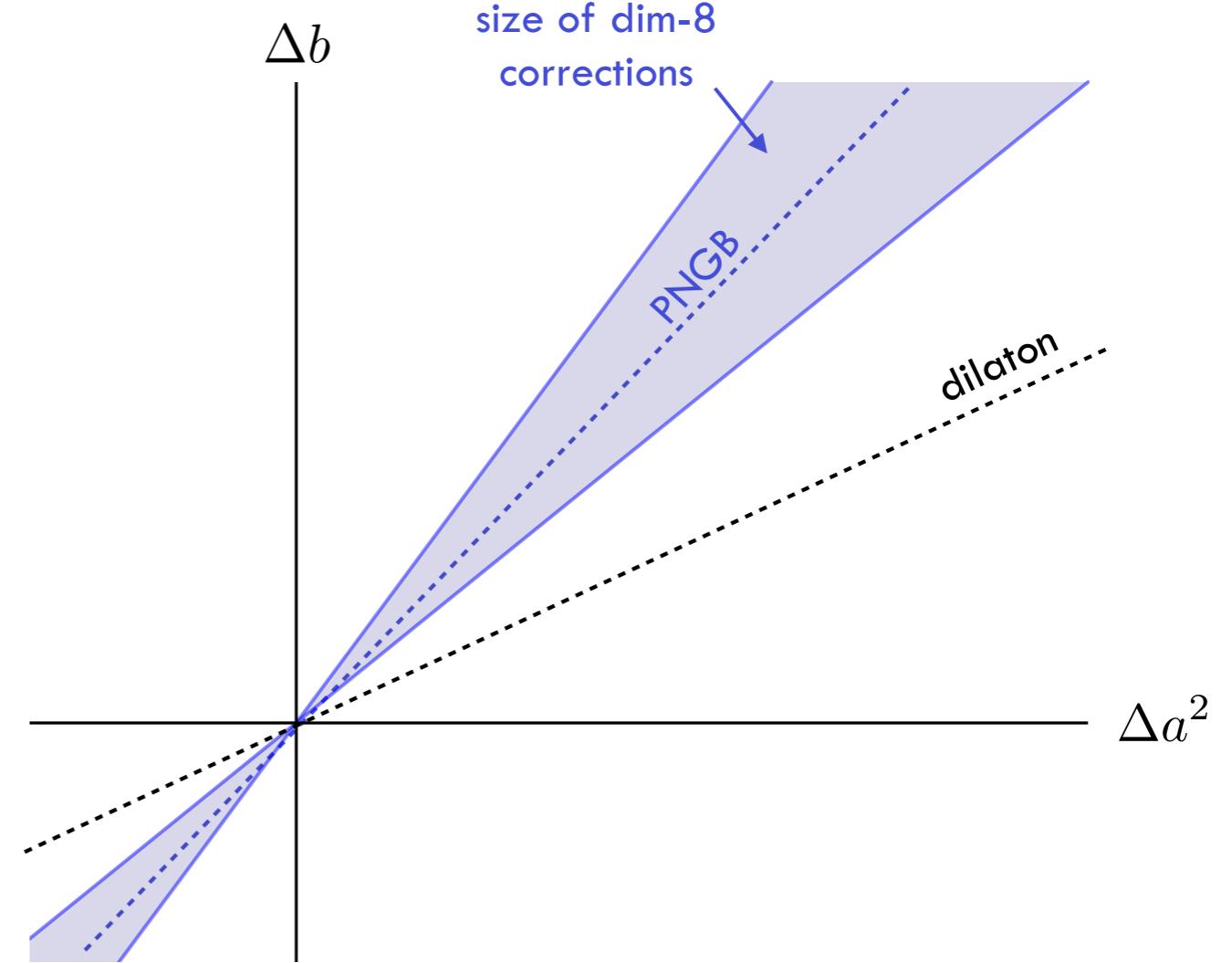
Scenario 1:

$$\Delta a^2 \sim \Delta b \sim 10\%$$

$$\text{Exp. precision} \sim 1\%$$



Test dim-8 operators



Ex:  $\text{SO}(5)/\text{SO}(4)$

In PNGB Higgs theories the whole series in  $H/f$  can be re-summed:

$$a = \sqrt{1 - \xi}$$

$$\xi = \frac{v^2}{f^2}$$

$$b = 1 - 2\xi$$

At dimension-6 level:

$$\Delta b = 2\Delta a^2 (1 + O(\Delta a^2))$$

$$\begin{aligned}\Delta b &\equiv 1 - b \\ \Delta a^2 &\equiv 1 - a^2\end{aligned}$$

Scenario 1:

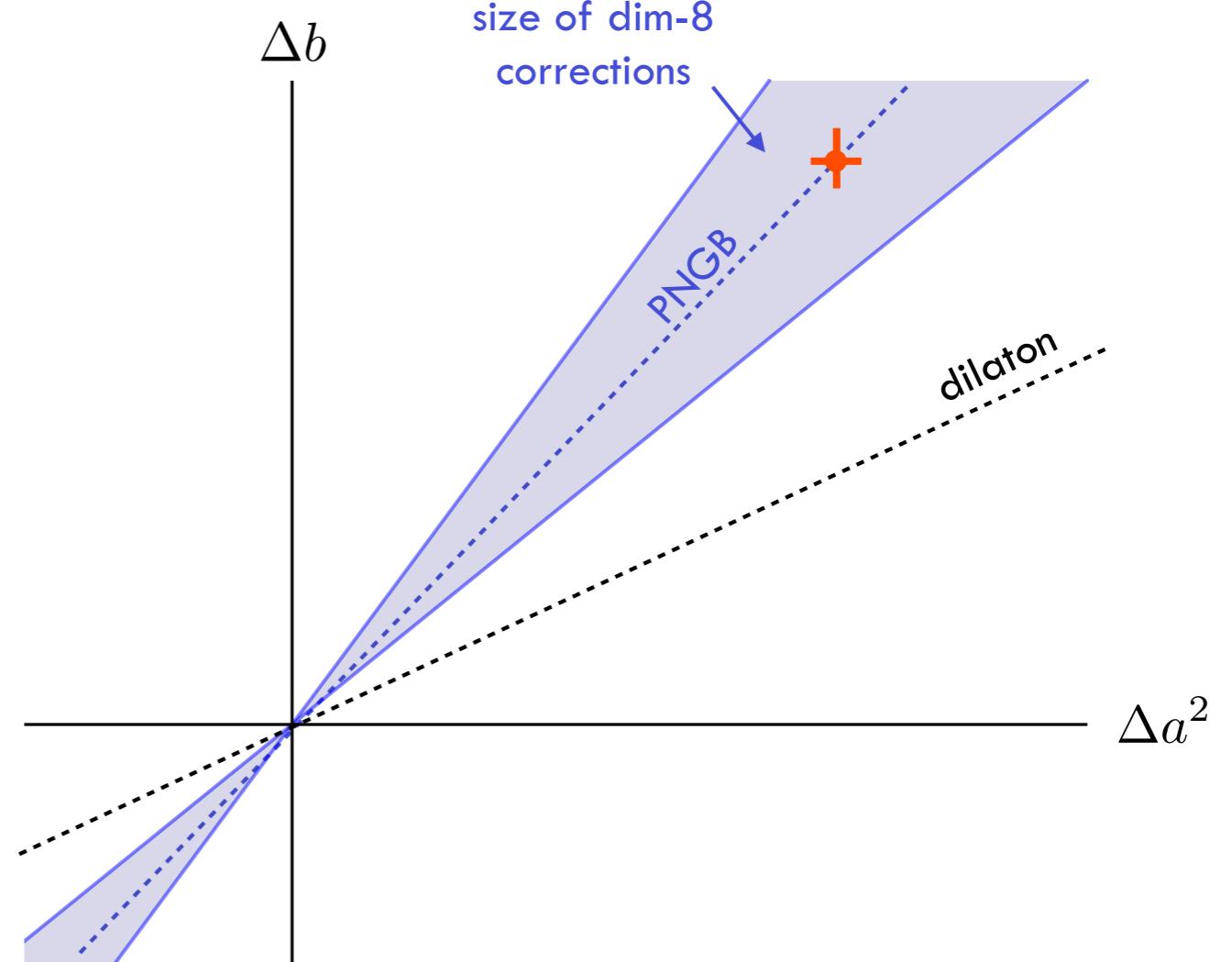
$$\Delta a^2 \sim \Delta b \sim 10\%$$

$$\text{Exp. precision} \sim 1\%$$



Test dim-8 operators

1. PNGB (and specific coset) proved



Ex:  $\text{SO}(5)/\text{SO}(4)$

In PNGB Higgs theories the whole series in  $H/f$  can be re-summed:

$$a = \sqrt{1 - \xi}$$

$$\xi = \frac{v^2}{f^2}$$

$$b = 1 - 2\xi$$

At dimension-6 level:

$$\Delta b = 2\Delta a^2 (1 + O(\Delta a^2))$$

$$\begin{aligned}\Delta b &\equiv 1 - b \\ \Delta a^2 &\equiv 1 - a^2\end{aligned}$$

Scenario 1:

$$\Delta a^2 \sim \Delta b \sim 10\%$$

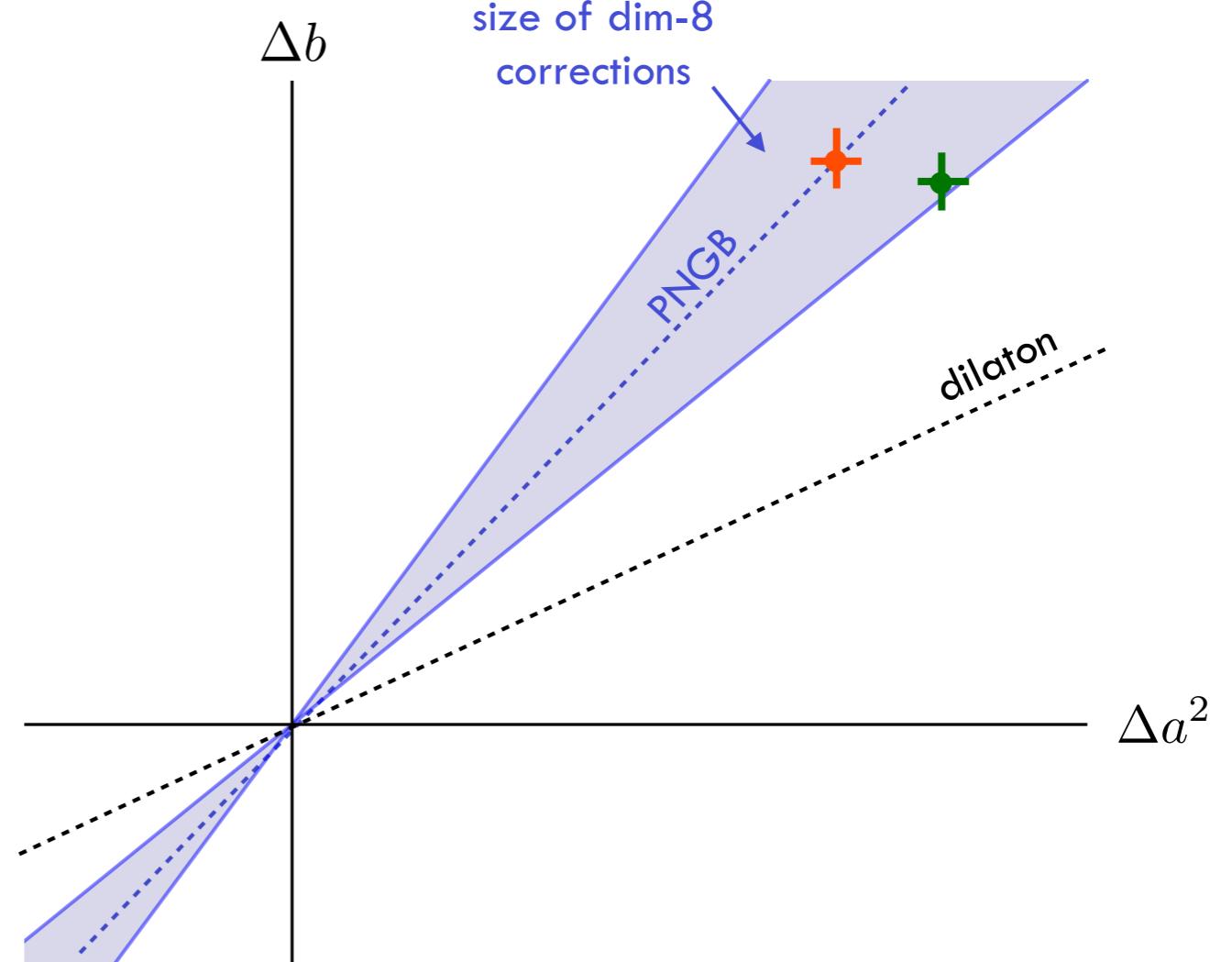
$$\text{Exp. precision} \sim 1\%$$



Test dim-8 operators

1. PNGB (and specific coset) proved

2. SILH proved, PNGB disproved



Ex:  $\text{SO}(5)/\text{SO}(4)$

In PNGB Higgs theories the whole series in  $H/f$  can be re-summed:

$$a = \sqrt{1 - \xi}$$

$$\xi = \frac{v^2}{f^2}$$

$$b = 1 - 2\xi$$

At dimension-6 level:

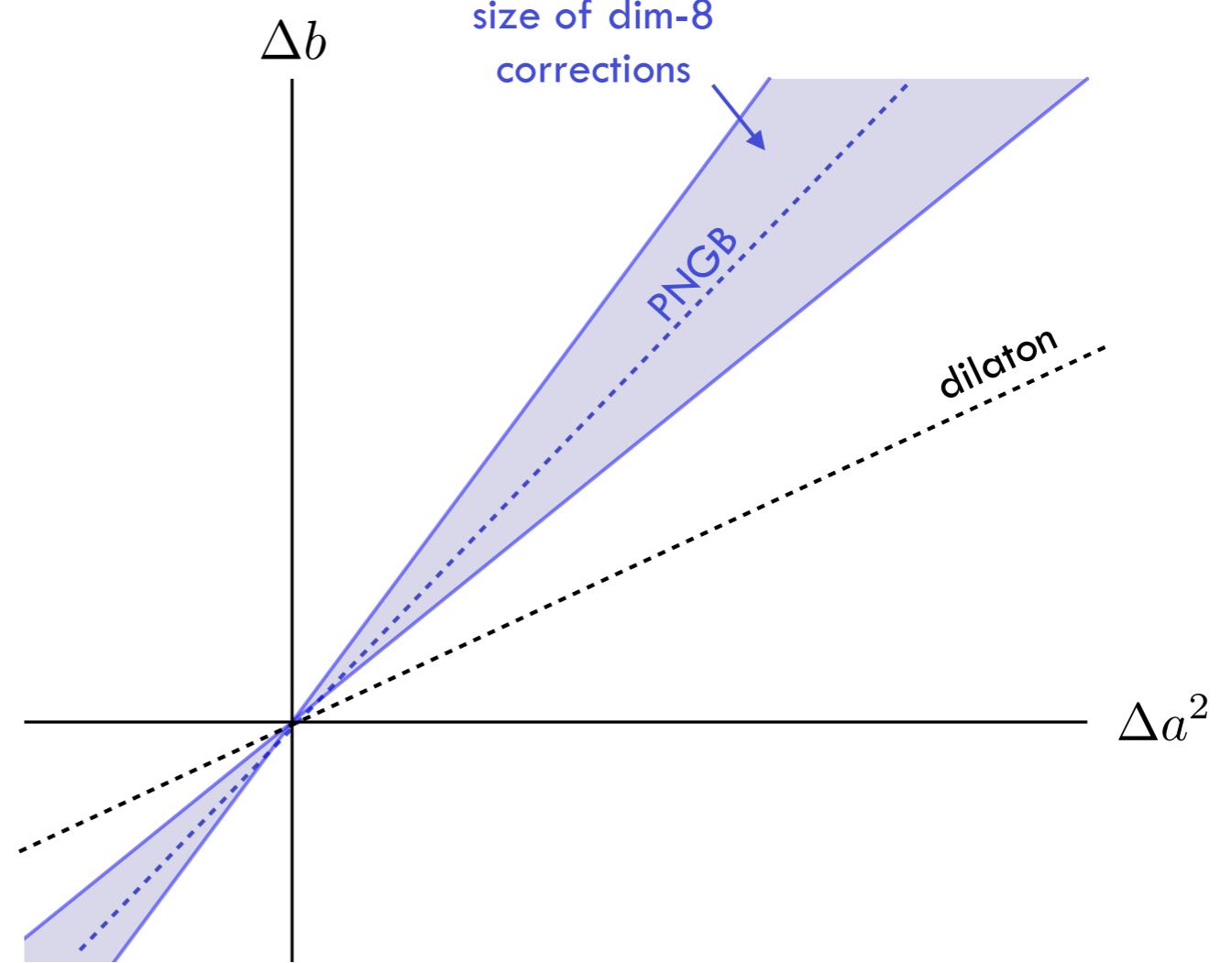
$$\Delta b = 2\Delta a^2 (1 + O(\Delta a^2))$$

$$\begin{aligned}\Delta b &\equiv 1 - b \\ \Delta a^2 &\equiv 1 - a^2\end{aligned}$$

Scenario 2:

$$\Delta a^2 \sim \Delta b \sim 1\%$$

$$\text{Exp. precision} \sim 1\%$$



Ex:  $\text{SO}(5)/\text{SO}(4)$

In PNGB Higgs theories the whole series in  $H/f$  can be re-summed:

$$a = \sqrt{1 - \xi}$$

$$\xi = \frac{v^2}{f^2}$$

$$b = 1 - 2\xi$$

At dimension-6 level:

$$\Delta b = 2\Delta a^2 (1 + O(\Delta a^2))$$

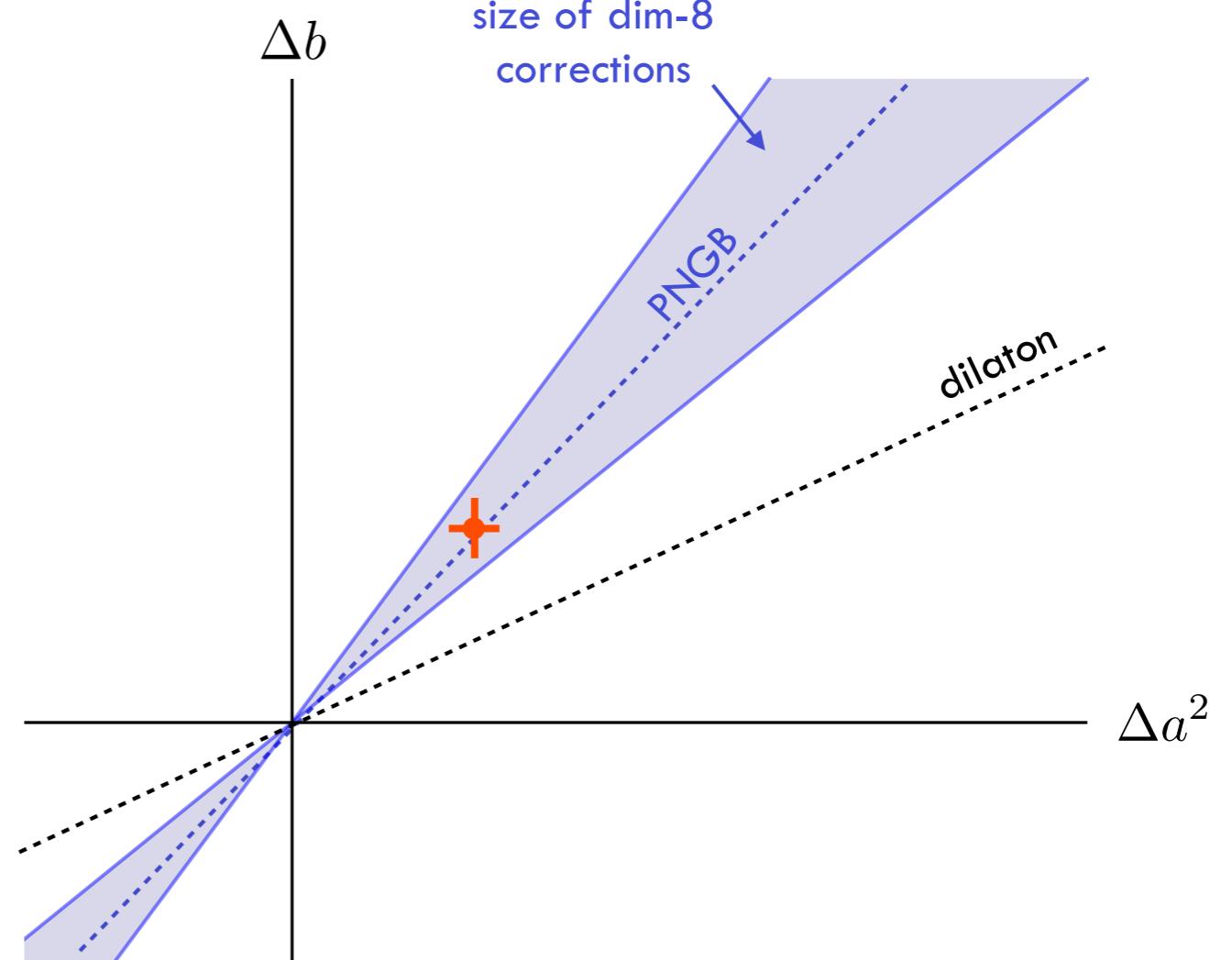
$$\begin{aligned}\Delta b &\equiv 1 - b \\ \Delta a^2 &\equiv 1 - a^2\end{aligned}$$

Scenario 2:

$$\Delta a^2 \sim \Delta b \sim 1\%$$

Exp. precision  $\sim 1\%$

1. SILH proved



Ex:  $\text{SO}(5)/\text{SO}(4)$

In PNGB Higgs theories the whole series in  $H/f$  can be re-summed:

$$a = \sqrt{1 - \xi}$$

$$\xi = \frac{v^2}{f^2}$$

$$b = 1 - 2\xi$$

At dimension-6 level:

$$\Delta b = 2\Delta a^2 (1 + O(\Delta a^2))$$

$$\begin{aligned}\Delta b &\equiv 1 - b \\ \Delta a^2 &\equiv 1 - a^2\end{aligned}$$

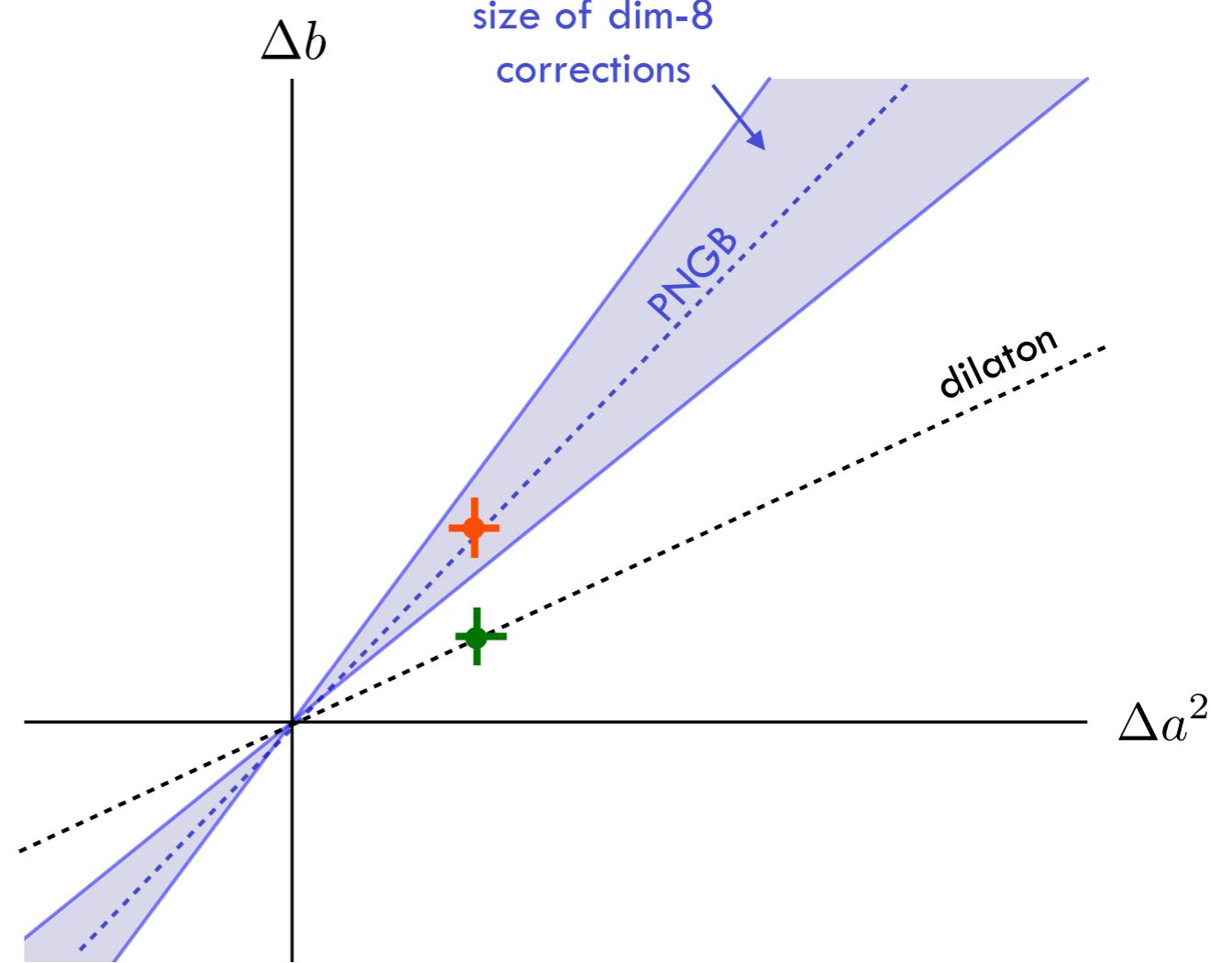
Scenario 2:

$$\Delta a^2 \sim \Delta b \sim 1\%$$

Exp. precision  $\sim 1\%$

1. SILH proved

2. SILH (i.e. Higgs doublet) disproved



# Extracting $b, d_3$ at CLIC with $\sqrt{s} = 3 \text{ TeV}$ through $e^+e^- \rightarrow \nu\bar{\nu} hh \rightarrow \nu\bar{\nu} b\bar{b}b\bar{b}$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

$$\sigma = a^4 \sigma_{SM} (1 + A \delta_b + B \delta_{d_3} + C \delta_b \delta_{d_3} + D \delta_b^2 + E \delta_{d_3}^2)$$

$$\delta_b = 1 - b/a^2$$

$$\delta_{d_3} = 1 - d_3/a$$

# Extracting $b, d_3$ at CLIC with $\sqrt{s} = 3 \text{ TeV}$ through $e^+e^- \rightarrow \nu\bar{\nu} hh \rightarrow \nu\bar{\nu} b\bar{b}b\bar{b}$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

$$\sigma = a^4 \sigma_{SM} (1 + A \delta_b + B \delta_{d_3} + C \delta_b \delta_{d_3} + D \delta_b^2 + E \delta_{d_3}^2)$$

$$\delta_b = 1 - b/a^2$$

$$\delta_{d_3} = 1 - d_3/a$$

The analysis in a nutshell

Parton level + Gaussian smearing of jet energy  $\Delta E/E = 5\%$

Selected events with: 4 jets

Basic acceptance cuts:  $E_j > 20 \text{ GeV}$      $|\eta_j| < 2$      $\Delta R_{jj} > 0.4$

$E_l > 5 \text{ GeV}$      $|\eta_l| < 2$      $\Delta R_{jl} > 0.4$

Higgs reconstruction: 1) choose pairing  $(j_1 j_2, j_3 j_4)$  that minimizes  $\chi^2 = (m_{j_1 j_2} - m_h)^2 + (m_{j_3 j_4} - m_h)^2$

2) each candidate must satisfy  $|m_{jj} - m_h| < 15 \text{ GeV}$

3) keep events with at least 3 b-tags

# Extracting $b, d_3$ at CLIC with $\sqrt{s} = 3 \text{ TeV}$ through $e^+e^- \rightarrow \nu\bar{\nu} hh \rightarrow \nu\bar{\nu} b\bar{b}b\bar{b}$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

$$\sigma = a^4 \sigma_{SM} (1 + A \delta_b + B \delta_{d_3} + C \delta_b \delta_{d_3} + D \delta_b^2 + E \delta_{d_3}^2)$$

$$\delta_b = 1 - b/a^2$$

$$\delta_{d_3} = 1 - d_3/a$$

**Background negligible (requires good mass resolution h vs Z)**  
 (largest processes:  $hZ\nu\bar{\nu}$ ,  $ZZ\nu\bar{\nu}$ ,  $ZZe^+e^-$ )

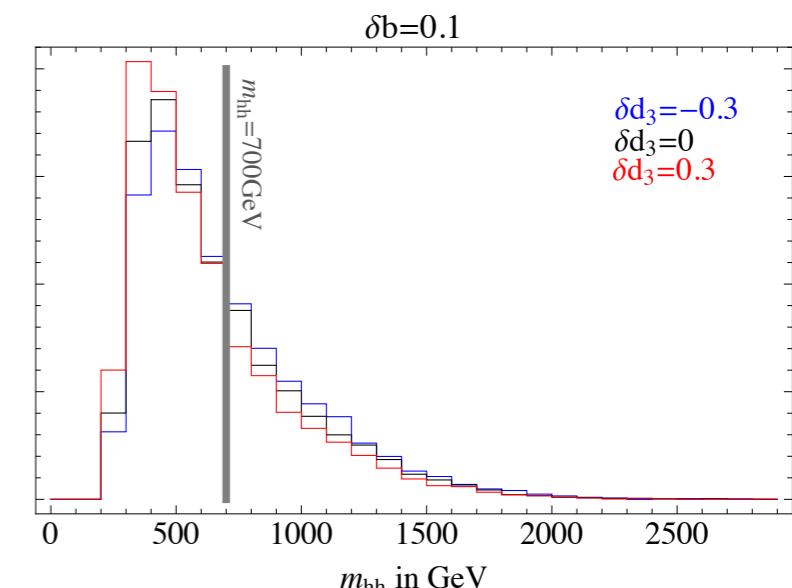
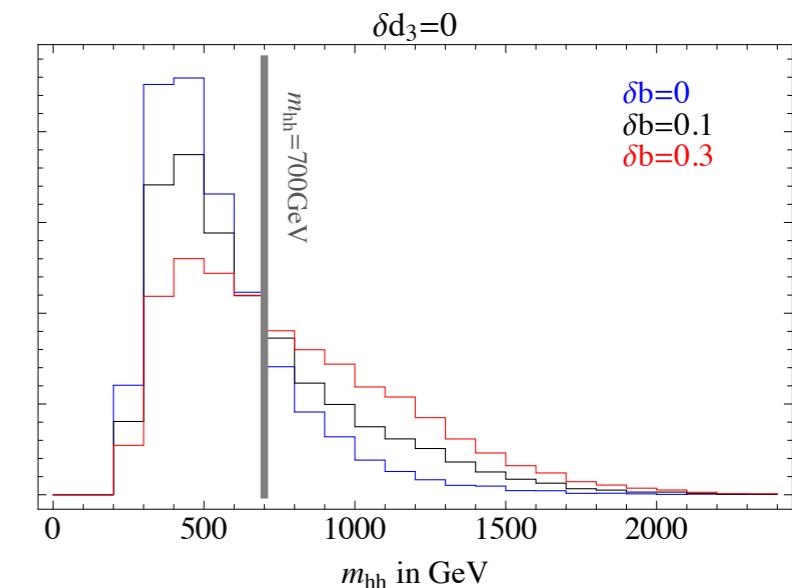
**Signal analyzed in 4 kinematic regions to enhance sensitivity on Higgs couplings:**

I :  $m_{hh} > 700 \text{ GeV}$  and  $H_T > 400 \text{ GeV}$

II :  $m_{hh} > 700 \text{ GeV}$  and  $H_T < 400 \text{ GeV}$

III :  $m_{hh} < 700 \text{ GeV}$  and  $H_T > 400 \text{ GeV}$

IV :  $m_{hh} < 700 \text{ GeV}$  and  $H_T < 400 \text{ GeV}$



# Extracting $b, d_3$ at CLIC with $\sqrt{s} = 3 \text{ TeV}$ through $e^+e^- \rightarrow \nu\bar{\nu} hh \rightarrow \nu\bar{\nu} b\bar{b}b\bar{b}$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

$$\sigma = a^4 \sigma_{SM} (1 + A \delta_b + B \delta_{d_3} + C \delta_b \delta_{d_3} + D \delta_b^2 + E \delta_{d_3}^2)$$

$$\delta_b = 1 - b/a^2$$

$$\delta_{d_3} = 1 - d_3/a$$

**Expected precision on  $\delta_b$  with  $L = 1 \text{ ab}^{-1}/(a^2 BR(b\bar{b})/BR(b\bar{b})_{SM})^2$**

measured $\delta_b$	$\bar{\delta}_{d_3}$							
	-0.5	-0.3	-0.1	0	0.1	0.3	0.5	
0	$-0.045^{+0.060}_{-0.025}$	$0.015^{+0.020}_{-0.040}$	$0.010^{+0.070}_{-0.045}$	$0.00^{+0.05}_{-0.05}$	$0.00^{+0.03}_{-0.03}$	$0.00^{+0.03}_{-0.03}$	$0.00^{+0.03}_{-0.03}$	
0.01	$-0.055^{+0.070}_{-0.020}$	$0.030^{+0.030}_{-0.045}$	$0.020^{+0.080}_{-0.035}$	$0.015^{+0.030}_{-0.035}$	$0.010^{+0.020}_{-0.030}$	$0.010^{+0.025}_{-0.025}$	$0.010^{+0.025}_{-0.025}$	
$\bar{\delta}_b$	$0.02^{+0.030}_{-0.035}$	$0.040^{+0.040}_{-0.050}$	$0.025^{+0.075}_{-0.020}$	$0.020^{+0.030}_{-0.035}$	$0.020^{+0.025}_{-0.025}$	$0.020^{+0.025}_{-0.025}$	$0.020^{+0.025}_{-0.025}$	
0.03	$0.03^{+0.030}_{-0.035}$	$0.050^{+0.040}_{-0.050}$	$0.035^{+0.030}_{-0.020}$	$0.030^{+0.025}_{-0.025}$	$0.030^{+0.025}_{-0.025}$	$0.030^{+0.025}_{-0.025}$	$0.030^{+0.020}_{-0.020}$	
0.05	$0.05^{+0.030}_{-0.035}$	$0.080^{+0.020}_{-0.040}$	$0.055^{+0.025}_{-0.020}$	$0.050^{+0.025}_{-0.020}$	$0.050^{+0.025}_{-0.025}$	$0.050^{+0.025}_{-0.025}$	$0.050^{+0.020}_{-0.020}$	
0.1	$0.12^{+0.025}_{-0.030}$	$0.10^{+0.03}_{-0.02}$	$0.10^{+0.03}_{-0.03}$	$0.10^{+0.02}_{-0.03}$	$0.10^{+0.02}_{-0.02}$	$0.10^{+0.02}_{-0.02}$	$0.10^{+0.02}_{-0.02}$	
0.3	$0.30^{+0.02}_{-0.02}$	$0.30^{+0.02}_{-0.02}$	$0.30^{+0.02}_{-0.02}$	$0.30^{+0.02}_{-0.02}$	$0.30^{+0.02}_{-0.02}$	$0.30^{+0.02}_{-0.02}$	$0.30^{+0.02}_{-0.02}$	
0.5	$0.50^{+0.02}_{-0.02}$	$0.50^{+0.02}_{-0.02}$	$0.50^{+0.02}_{-0.02}$	$0.50^{+0.02}_{-0.02}$	$0.50^{+0.02}_{-0.02}$	$0.50^{+0.02}_{-0.02}$	$0.50^{+0.02}_{-0.02}$	

# Extracting $b, d_3$ at CLIC with $\sqrt{s} = 3 \text{ TeV}$ through $e^+e^- \rightarrow \nu\bar{\nu} hh \rightarrow \nu\bar{\nu} b\bar{b}b\bar{b}$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

$$\sigma = a^4 \sigma_{SM} (1 + A \delta_b + B \delta_{d_3} + C \delta_b \delta_{d_3} + D \delta_b^2 + E \delta_{d_3}^2)$$

$$\delta_b = 1 - b/a^2$$

$$\delta_{d_3} = 1 - d_3/a$$

Expected precision on  $\delta_b$  with  $L = 1 \text{ ab}^{-1}/(a^2 BR(b\bar{b})/BR(b\bar{b})_{SM})^2$

measured $\delta_b$	$\bar{\delta}_{d_3}$							
	-0.5	-0.3	-0.1	0	0.1	0.3	0.5	
0	-0.045 <sup>+0.060</sup> <sub>-0.025</sub>	0.015 <sup>+0.020</sup> <sub>-0.040</sub>	0.010 <sup>+0.070</sup> <sub>-0.045</sub>	0.00 <sup>+0.05</sup> <sub>-0.05</sub>	+0.03 -0.03	0.00 <sup>+0.03</sup> <sub>-0.03</sub>	0.00 <sup>+0.03</sup> <sub>-0.03</sub>	
0.01	-0.055 <sup>+0.070</sup> <sub>-0.020</sub>	0.030 <sup>+0.030</sup> <sub>-0.045</sub>	0.020 <sup>+0.080</sup> <sub>-0.035</sub>	0.015 <sup>+0.030</sup> <sub>-0.035</sub>	0.010 <sup>+0.020</sup> <sub>-0.030</sub>	0.010 <sup>+0.025</sup> <sub>-0.025</sub>	0.010 <sup>+0.025</sup> <sub>-0.025</sub>	
$\bar{\delta}_b$	0.02	0.02 <sup>+0.030</sup> <sub>-0.035</sub>	0.040 <sup>+0.040</sup> <sub>-0.050</sub>	0.025 <sup>+0.075</sup> <sub>-0.020</sub>	0.020 <sup>+0.030</sup> <sub>-0.035</sub>	0.020 <sup>+0.025</sup> <sub>-0.025</sub>	0.020 <sup>+0.025</sup> <sub>-0.025</sub>	
0.03	0.03 <sup>+0.030</sup> <sub>-0.035</sub>	0.050 <sup>+0.040</sup> <sub>-0.050</sub>	0.035 <sup>+0.030</sup> <sub>-0.020</sub>	0.030 <sup>+0.025</sup> <sub>-0.025</sub>	0.030 <sup>+0.025</sup> <sub>-0.025</sub>	0.030 <sup>+0.025</sup> <sub>-0.025</sub>	0.030 <sup>+0.020</sup> <sub>-0.020</sub>	
0.05	0.05 <sup>+0.030</sup> <sub>-0.035</sub>	0.080 <sup>+0.020</sup> <sub>-0.040</sub>	0.055 <sup>+0.025</sup> <sub>-0.020</sub>	0.050 <sup>+0.025</sup> <sub>-0.020</sub>	0.050 <sup>+0.025</sup> <sub>-0.025</sub>	0.050 <sup>+0.025</sup> <sub>-0.025</sub>	0.050 <sup>+0.020</sup> <sub>-0.020</sub>	
0.1	0.12 <sup>+0.025</sup> <sub>-0.030</sub>	0.10 <sup>+0.03</sup> <sub>-0.02</sub>	0.10 <sup>+0.03</sup> <sub>-0.03</sub>	0.10 <sup>+0.02</sup> <sub>-0.03</sub>	0.10 <sup>+0.02</sup> <sub>-0.02</sub>	0.10 <sup>+0.02</sup> <sub>-0.02</sub>	0.10 <sup>+0.02</sup> <sub>-0.02</sub>	
0.3	0.30 <sup>+0.02</sup> <sub>-0.02</sub>	0.30 <sup>+0.02</sup> <sub>-0.02</sub>	0.30 <sup>+0.02</sup> <sub>-0.02</sub>	0.30 <sup>+0.02</sup> <sub>-0.02</sub>	0.30 <sup>+0.02</sup> <sub>-0.02</sub>	0.30 <sup>+0.02</sup> <sub>-0.02</sub>	0.30 <sup>+0.02</sup> <sub>-0.02</sub>	
0.5	0.50 <sup>+0.02</sup> <sub>-0.02</sub>	0.50 <sup>+0.02</sup> <sub>-0.02</sub>	0.50 <sup>+0.02</sup> <sub>-0.02</sub>	0.50 <sup>+0.02</sup> <sub>-0.02</sub>	0.50 <sup>+0.02</sup> <sub>-0.02</sub>	0.50 <sup>+0.02</sup> <sub>-0.02</sub>	0.50 <sup>+0.02</sup> <sub>-0.02</sub>	

For injected SM: precision on  $\delta_b$  is 5%

# Extracting $b, d_3$ at CLIC with $\sqrt{s} = 3 \text{ TeV}$ through $e^+e^- \rightarrow \nu\bar{\nu} hh \rightarrow \nu\bar{\nu} b\bar{b}b\bar{b}$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

$$\sigma = a^4 \sigma_{SM} (1 + A \delta_b + B \delta_{d_3} + C \delta_b \delta_{d_3} + D \delta_b^2 + E \delta_{d_3}^2)$$

$$\delta_b = 1 - b/a^2$$

$$\delta_{d_3} = 1 - d_3/a$$

**Expected precision on  $\delta_{d_3}$  with  $L = 1 \text{ ab}^{-1}/(a^2 BR(b\bar{b})/BR(b\bar{b})_{SM})^2$**

measured $\delta_{d_3}$	$\bar{\delta}_{d_3}$							
	-0.5	-0.3	-0.1	0	0.1	0.3	0.5	
0	$-0.50^{+0.35}_{-0.25}$	$-0.25^{+0.20}_{-0.50}$	$0.00^{+0.25}_{-0.40}$	$0.05^{+0.30}_{-0.30}$	$0.10^{+0.25}_{-0.20}$	$0.30^{+0.20}_{-0.15}$	$0.50^{+0.15}_{-0.15}$	
0.01	$-0.45^{+0.35}_{-0.30}$	$-0.20^{+0.30}_{-0.55}$	$-0.05^{+0.30}_{-0.30}$	$0.00^{+0.25}_{-0.25}$	$0.10^{+0.20}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$	
$\bar{\delta}_b$	0.02	$-0.35^{+0.30}_{-0.35}$	$-0.25^{+0.25}_{-0.60}$	$-0.10^{+0.25}_{-0.30}$	$0.00^{+0.20}_{-0.25}$	$0.10^{+0.15}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
0.03		$-0.40^{+0.30}_{-0.35}$	$-0.25^{+0.20}_{-0.70}$	$-0.10^{+0.20}_{-0.25}$	$0.00^{+0.15}_{-0.20}$	$0.10^{+0.15}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
0.05		$-0.55^{+0.30}_{-0.40}$	$-0.30^{+0.20}_{-0.30}$	$-0.10^{+0.20}_{-0.20}$	$0.00^{+0.15}_{-0.20}$	$0.10^{+0.15}_{-0.15}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.10}_{-0.10}$
0.1		$-0.50^{+0.15}_{-0.25}$	$-0.30^{+0.15}_{-0.20}$	$-0.10^{+0.20}_{-0.20}$	$0.00^{+0.15}_{-0.15}$	$0.10^{+0.15}_{-0.15}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$
0.3		$-0.50^{+0.15}_{-0.15}$	$-0.30^{+0.15}_{-0.15}$	$-0.10^{+0.10}_{-0.10}$	$0.00^{+0.10}_{-0.10}$	$0.10^{+0.10}_{-0.10}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$
0.5		$-0.50^{+0.15}_{-0.10}$	$-0.30^{+0.10}_{-0.10}$	$-0.10^{+0.10}_{-0.10}$	$0.00^{+0.10}_{-0.10}$	$0.10^{+0.10}_{-0.10}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$

# Extracting $b, d_3$ at CLIC with $\sqrt{s} = 3 \text{ TeV}$ through $e^+e^- \rightarrow \nu\bar{\nu} hh \rightarrow \nu\bar{\nu} b\bar{b}b\bar{b}$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

$$\sigma = a^4 \sigma_{SM} (1 + A \delta_b + B \delta_{d_3} + C \delta_b \delta_{d_3} + D \delta_b^2 + E \delta_{d_3}^2)$$

$$\delta_b = 1 - b/a^2$$

$$\delta_{d_3} = 1 - d_3/a$$

Expected precision on  $\delta_{d_3}$  with  $L = 1 \text{ ab}^{-1}/(a^2 BR(b\bar{b})/BR(b\bar{b})_{SM})^2$

measured $\delta_{d_3}$	$\bar{\delta}_{d_3}$							
	-0.5	-0.3	-0.1	0	0.1	0.3	0.5	
$\bar{\delta}_b$	0	$-0.50^{+0.35}_{-0.25}$	$-0.25^{+0.20}_{-0.50}$	$0.00^{+0.25}_{-0.40}$	$0.05^{+0.30}_{-0.30}$	$0.10^{+0.25}_{-0.20}$	$0.30^{+0.20}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
	0.01	$-0.45^{+0.35}_{-0.30}$	$-0.20^{+0.30}_{-0.55}$	$-0.05^{+0.30}_{-0.30}$	$0.00^{+0.25}_{-0.25}$	$0.10^{+0.20}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
	0.02	$-0.35^{+0.30}_{-0.35}$	$-0.25^{+0.25}_{-0.60}$	$-0.10^{+0.25}_{-0.30}$	$0.00^{+0.20}_{-0.25}$	$0.10^{+0.15}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
	0.03	$-0.40^{+0.30}_{-0.35}$	$-0.25^{+0.20}_{-0.70}$	$-0.10^{+0.20}_{-0.25}$	$0.00^{+0.15}_{-0.20}$	$0.10^{+0.15}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
	0.05	$-0.55^{+0.30}_{-0.40}$	$-0.30^{+0.20}_{-0.30}$	$-0.10^{+0.20}_{-0.20}$	$0.00^{+0.15}_{-0.20}$	$0.10^{+0.15}_{-0.15}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.10}_{-0.10}$
	0.1	$-0.50^{+0.15}_{-0.25}$	$-0.30^{+0.15}_{-0.20}$	$-0.10^{+0.20}_{-0.20}$	$0.00^{+0.15}_{-0.15}$	$0.10^{+0.15}_{-0.15}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$
	0.3	$-0.50^{+0.15}_{-0.15}$	$-0.30^{+0.15}_{-0.15}$	$-0.10^{+0.10}_{-0.10}$	$0.00^{+0.10}_{-0.10}$	$0.10^{+0.10}_{-0.10}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$
	0.5	$-0.50^{+0.15}_{-0.10}$	$-0.30^{+0.10}_{-0.10}$	$-0.10^{+0.10}_{-0.10}$	$0.00^{+0.10}_{-0.10}$	$0.10^{+0.10}_{-0.10}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$

For injected SM: precision on  $\delta_{d_3}$  is 30%

# Extracting $b, d_3$ at CLIC with $\sqrt{s} = 3 \text{ TeV}$ through $e^+e^- \rightarrow \nu\bar{\nu} hh \rightarrow \nu\bar{\nu} b\bar{b}b\bar{b}$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

$$\sigma = a^4 \sigma_{SM} (1 + A \delta_b + B \delta_{d_3} + C \delta_b \delta_{d_3} + D \delta_b^2 + E \delta_{d_3}^2)$$

$$\delta_b = 1 - b/a^2$$

$$\delta_{d_3} = 1 - d_3/a$$

Expected precision on  $\delta_{d_3}$  with  $L = 1 \text{ ab}^{-1}/(a^2 BR(b\bar{b})/BR(b\bar{b})_{SM})^2$

measured $\delta_{d_3}$	$\bar{\delta}_{d_3}$							
	-0.5	-0.3	-0.1	0	0.1	0.3	0.5	
$\bar{\delta}_b$	0	$-0.50^{+0.35}_{-0.25}$	$-0.25^{+0.20}_{-0.50}$	$0.00^{+0.25}_{-0.40}$	$0.05^{+0.30}_{-0.30}$	$0.10^{+0.25}_{-0.20}$	$0.30^{+0.20}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
	0.01	$-0.45^{+0.35}_{-0.30}$	$-0.20^{+0.30}_{-0.55}$	$-0.05^{+0.30}_{-0.30}$	$0.00^{+0.25}_{-0.25}$	$0.10^{+0.20}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
	0.02	$-0.35^{+0.30}_{-0.35}$	$-0.25^{+0.25}_{-0.60}$	$-0.10^{+0.25}_{-0.30}$	$0.00^{+0.20}_{-0.25}$	$0.10^{+0.15}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
	0.03	$-0.40^{+0.30}_{-0.35}$	$-0.25^{+0.20}_{-0.70}$	$-0.10^{+0.20}_{-0.25}$	$0.00^{+0.15}_{-0.20}$	$0.10^{+0.15}_{-0.20}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.15}_{-0.15}$
	0.05	$-0.55^{+0.30}_{-0.40}$	$-0.30^{+0.20}_{-0.30}$	$-0.10^{+0.20}_{-0.20}$	$0.00^{+0.15}_{-0.20}$	$0.10^{+0.15}_{-0.15}$	$0.30^{+0.15}_{-0.15}$	$0.50^{+0.10}_{-0.10}$
	0.1	$-0.50^{+0.15}_{-0.25}$	$-0.30^{+0.15}_{-0.20}$	$-0.10^{+0.20}_{-0.20}$	$0.00^{+0.15}_{-0.15}$	$0.10^{+0.15}_{-0.15}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$
	0.3	$-0.50^{+0.15}_{-0.15}$	$-0.30^{+0.15}_{-0.15}$	$-0.10^{+0.10}_{-0.10}$	$0.00^{+0.10}_{-0.10}$	$0.10^{+0.10}_{-0.10}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$
	0.5	$-0.50^{+0.15}_{-0.10}$	$-0.30^{+0.10}_{-0.10}$	$-0.10^{+0.10}_{-0.10}$	$0.00^{+0.10}_{-0.10}$	$0.10^{+0.10}_{-0.10}$	$0.30^{+0.10}_{-0.10}$	$0.50^{+0.10}_{-0.10}$

For injected SM: precision on  $\delta_{d_3}$  is 30%  $\Rightarrow$

Much stronger sensitivity on  $b$  than on  $d_3$

# Reach on compositeness scale

	$\xi = (v/f)^2$	$\Lambda = 4\pi f$
LHC 14 TeV $L = 300 \text{ fb}^{-1}$	0.5 (double Higgs [1, 2])	4.5 TeV
	0.1 (single Higgs [3, 4])	10 TeV
ILC 250 GeV $L = 250 \text{ fb}^{-1}$ + 500 GeV $L = 500 \text{ fb}^{-1}$	$0.6-1.2 \times 10^{-2}$ (single Higgs [5, 6])	30-40 TeV
CLIC 3 TeV $L = 1 \text{ ab}^{-1}$	$2-5 \times 10^{-2}$ (double Higgs [7])	15-20 TeV
CLIC 350 GeV $L = 500 \text{ fb}^{-1}$ + 1.4 TeV $L = 1.5 \text{ ab}^{-1}$ + 3.0 TeV $L = 2 \text{ ab}^{-1}$	$1.1-2.4 \times 10^{-3}$ (single Higgs [8])	60-90 TeV

- [1] G. F. Giudice, C. Grojean, A. Pomarol and R. Rattazzi, JHEP **0706** (2007) 045 [arXiv:hep-ph/0703164].
- [2] R. Contino, C. Grojean, M. Moretti, F. Piccinini and R. Rattazzi, JHEP **1005** (2010) 089 [arXiv:1002.1011 [hep-ph]].
- [3] CMS Collaboration, 2012, CMS NOTE-2012/006.
- [4] ATLAS Collaboration, 2012, ATL-PHYS-PUB-2012-004.
- [5] H. Baer, T. Barklow, K. Fujii, Y. Gao, A. Hoang, S. Kanemura, J. List and H. E. Logan *et al.*, “The International Linear Collider Technical Design Report - Volume 2: Physics,” [arXiv:1306.6352 [hep-ph]].
- [6] S. Dutta, K. Hagiwara and Y. Matsumoto, Phys. Rev. D **78** (2008) 115016 [arXiv:0808.0477 [hep-ph]]. M. E. Peskin, [arXiv:1207.2516 [hep-ph]]. M. Klute, R. Lafaye, T. Plehn, M. Rauch and D. Zerwas, Europhys. Lett. **101** (2013) 51001 [arXiv:1301.1322 [hep-ph]].
- [7] R. Contino, C. Grojean, D. Pappadopulo, R. Rattazzi and A. Thamm, arXiv:1309.7038 [hep-ph].
- [8] H. Abramowicz *et al.* [CLIC Detector and Physics Study Collaboration], “Physics at the CLIC e+e- Linear Collider – Input to the Snowmass process 2013,” [arXiv:1307.5288 [hep-ex]].

# Conclusions

# Conclusions

- Tests of Higgs compositeness (i.e. strong EWSB) can be done by precisely measuring low-energy quantities (e.g. tails of scattering amplitudes)

# Conclusions

- Tests of Higgs compositeness (i.e. strong EWSB) can be done by precisely measuring low-energy quantities (e.g. tails of scattering amplitudes)
  
- From study of  $W \rightarrow hh$  at CLIC with  $\sqrt{s} = 3 \text{ TeV}$  :
  - precision:  $Whh$  at 5% ;  $hhh$  at 30%
  - tests of Higgs effective Lagrangian at dim-8 level: PNBG vs SILH

# Conclusions

- Tests of Higgs compositeness (i.e. strong EWSB) can be done by precisely measuring low-energy quantities (e.g. tails of scattering amplitudes)
- From study of  $W \rightarrow hh$  at CLIC with  $\sqrt{s} = 3 \text{ TeV}$ :
  - precision:  $Whh$  at 5% ;  $hhh$  at 30%
  - tests of Higgs effective Lagrangian at dim-8 level: PNBG vs SILH
- Take the right approach to look for NP  
Ex: much more information from  $VV \rightarrow hh$  than just the trilinear coupling

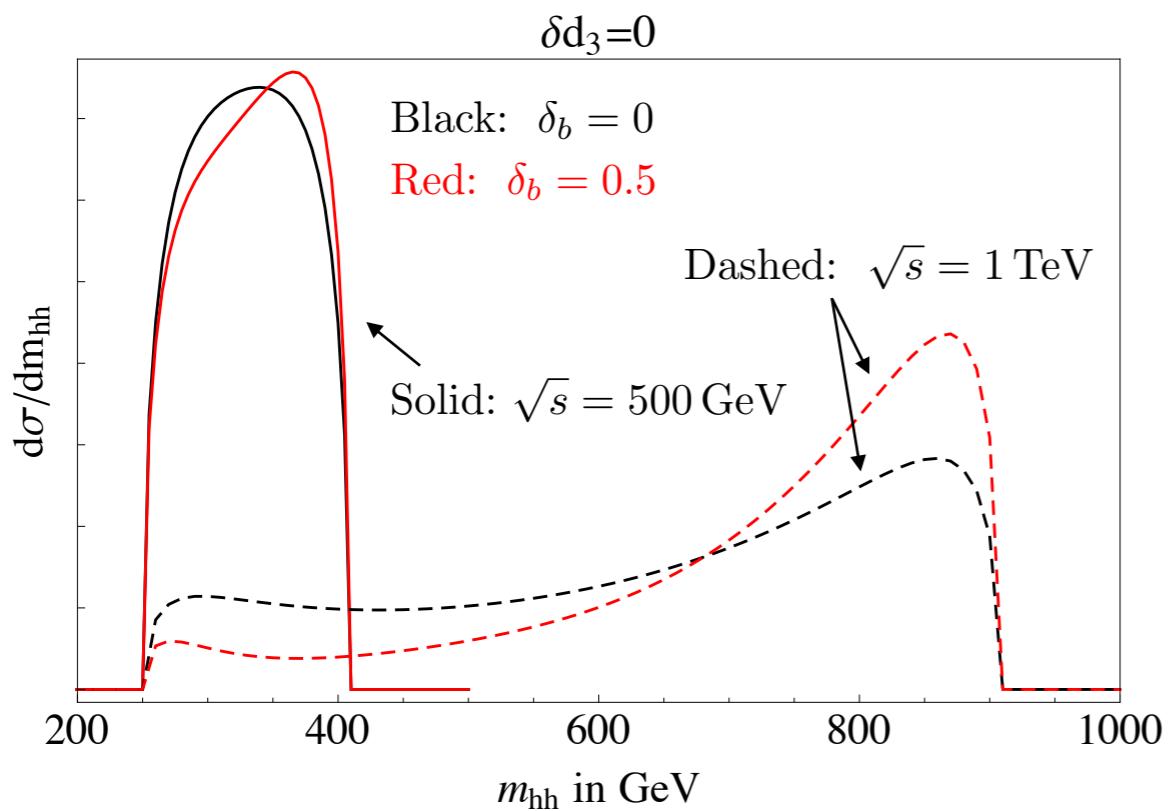
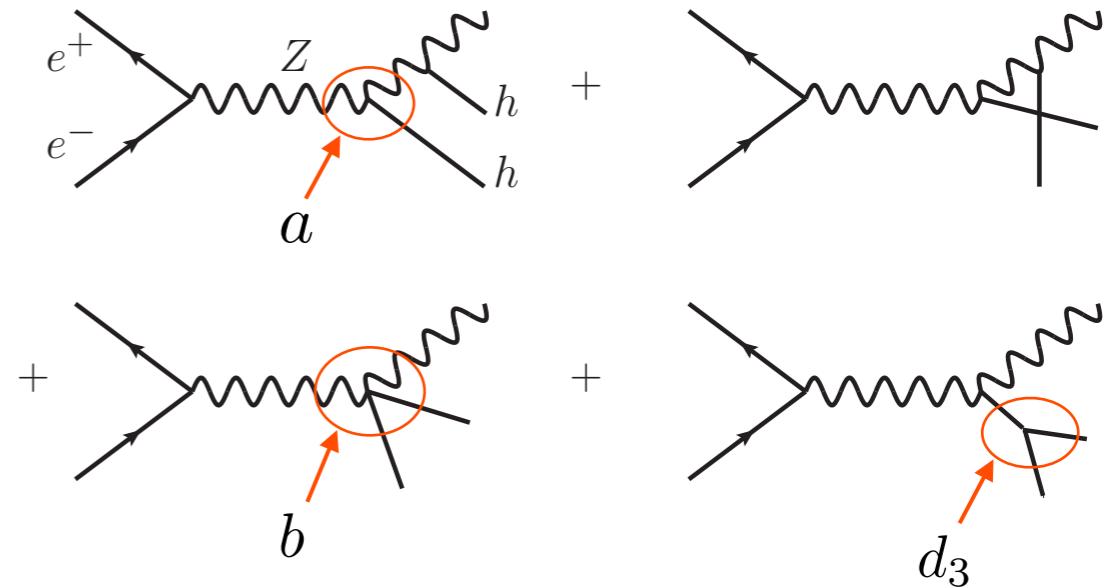
# Extra Slides

# Extracting $b, d_3$ at the ILC

with  $\sqrt{s} = 500 \text{ GeV} - 1 \text{ TeV}$

## through double Higgs-strahlung

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

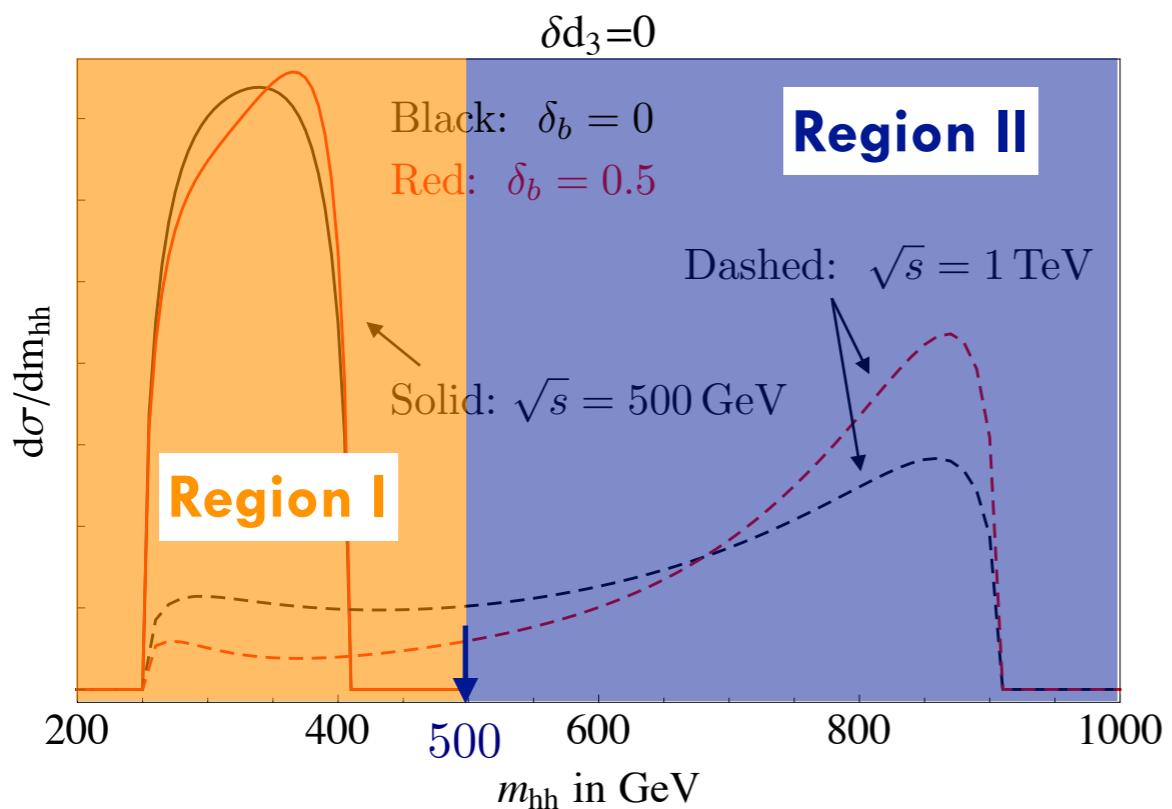
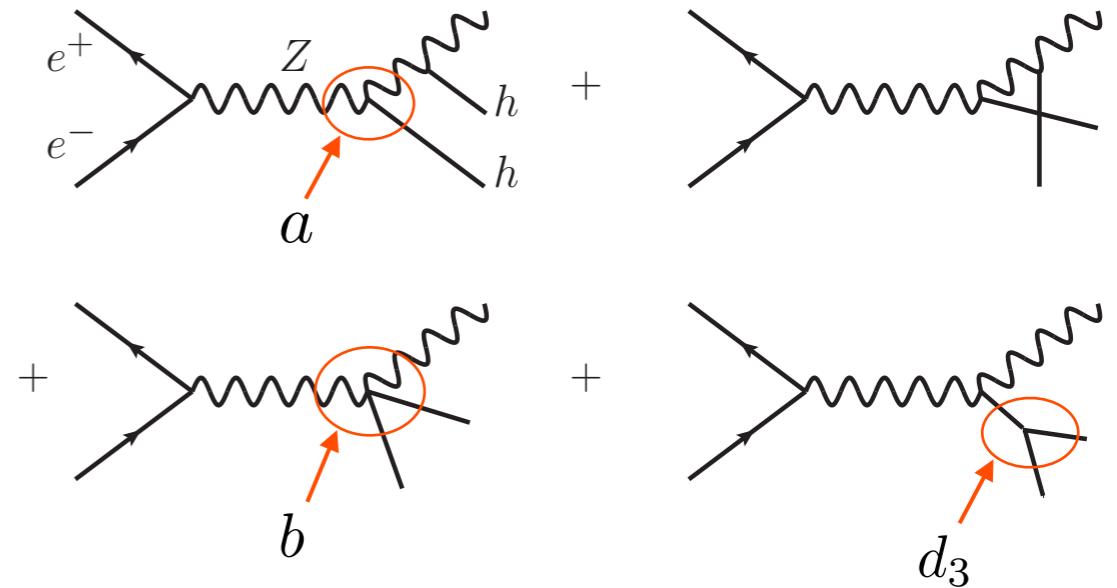


# Extracting $b, d_3$ at the ILC

with  $\sqrt{s} = 500 \text{ GeV} - 1 \text{ TeV}$

## through double Higgs-strahlung

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038



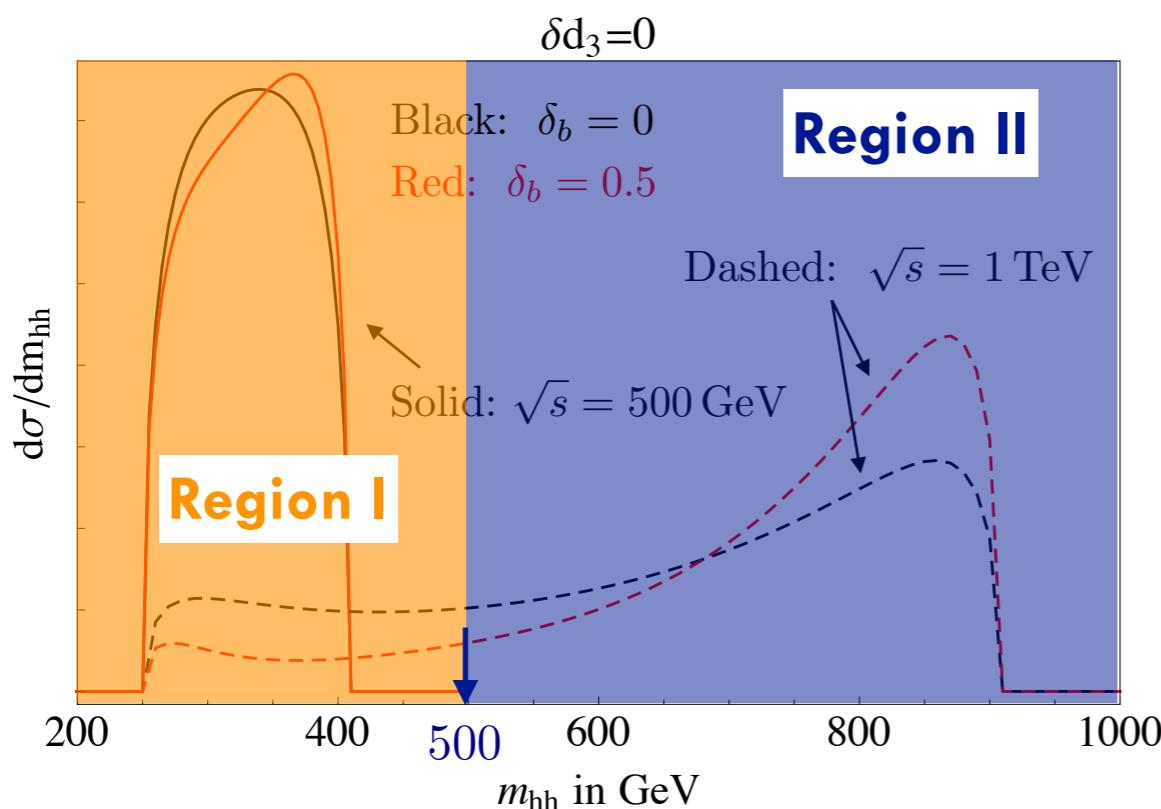
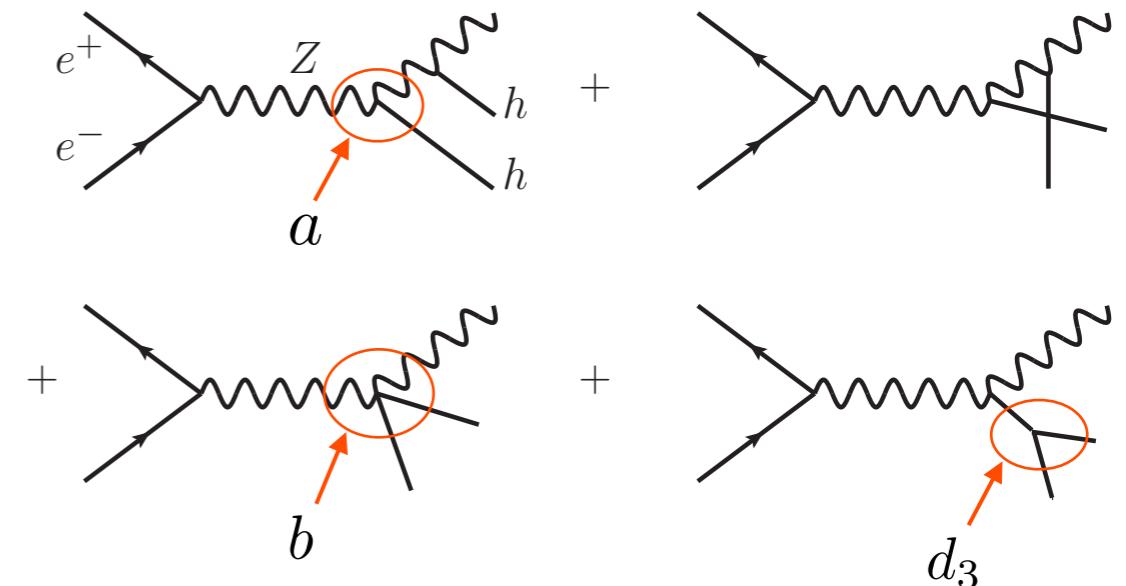
Cut on  $m_{hh}$  useful at  $\sqrt{s} = 1 \text{ TeV}$

# Extracting $b, d_3$ at the ILC

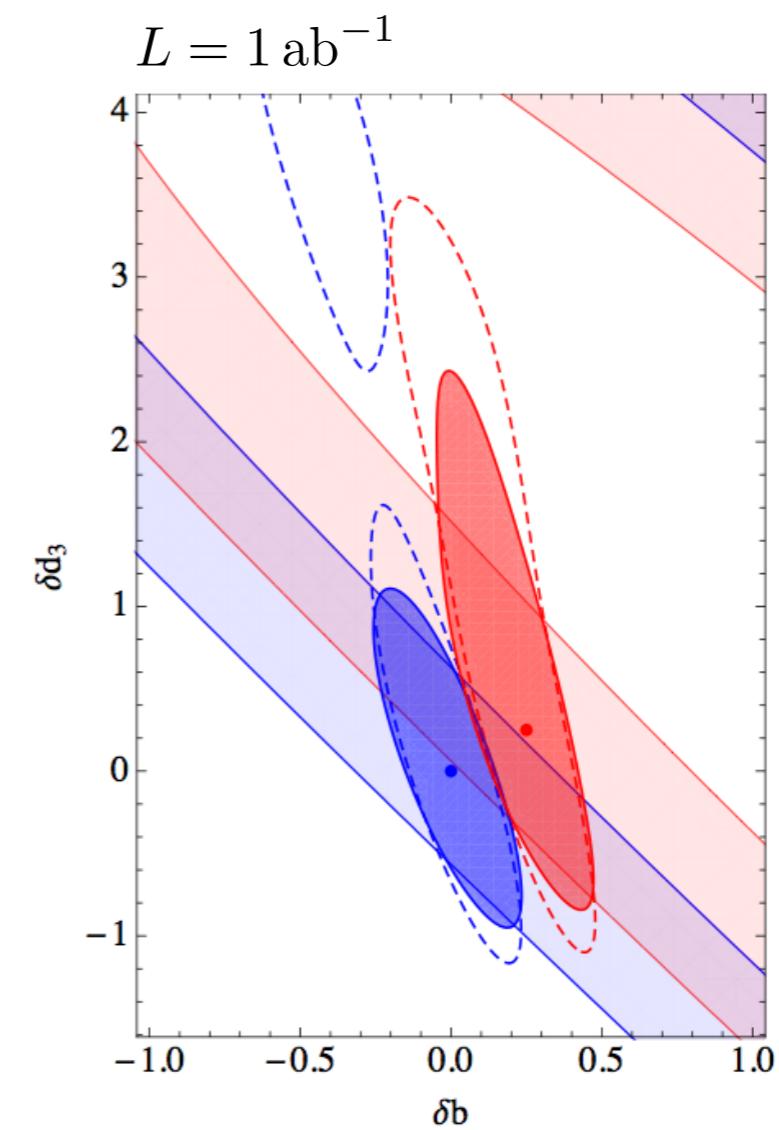
with  $\sqrt{s} = 500 \text{ GeV} - 1 \text{ TeV}$

through double Higgs-strahlung

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038



Cut on  $m_{hh}$  useful at  $\sqrt{s} = 1 \text{ TeV}$



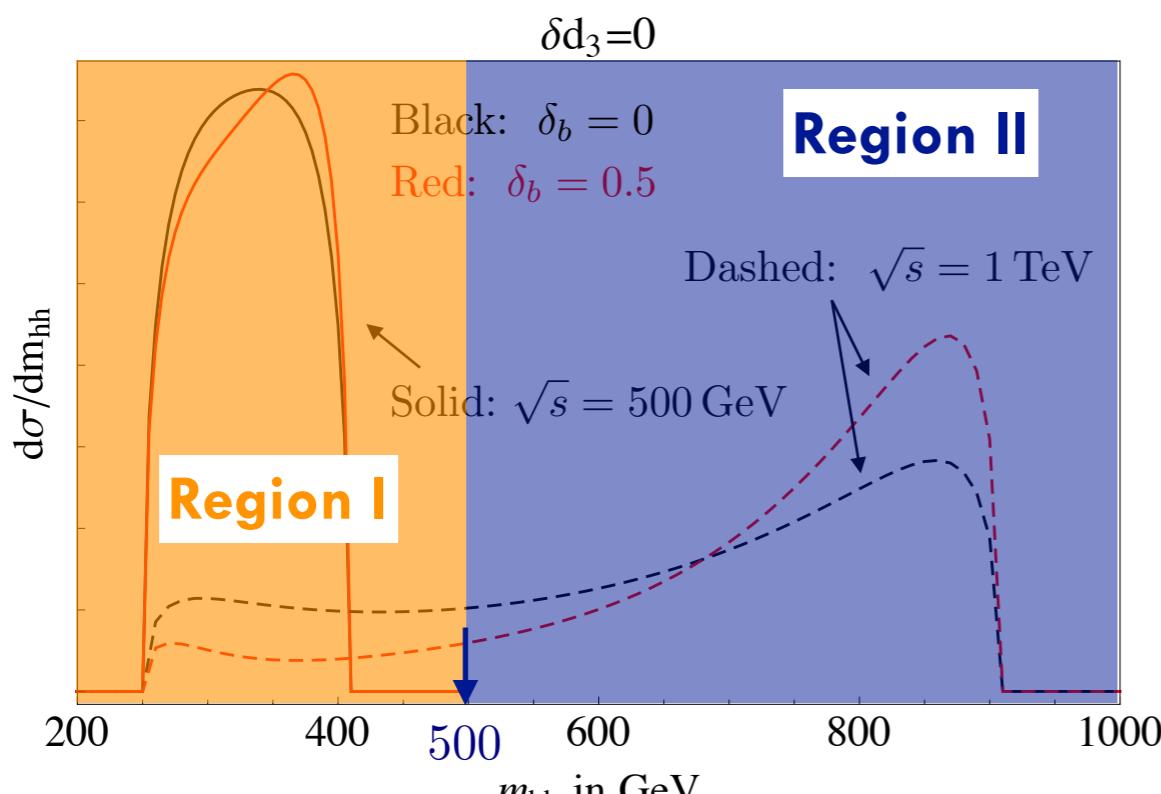
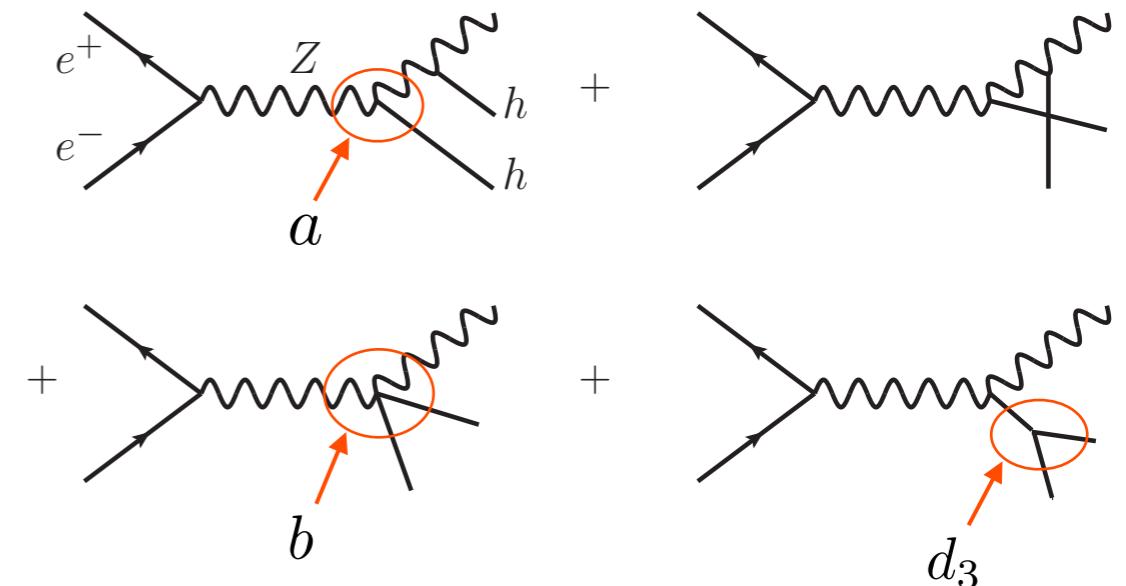
$$a^2(BR(b\bar{b})/BR(b\bar{b})_{SM}) = 1$$

# Extracting $b, d_3$ at the ILC

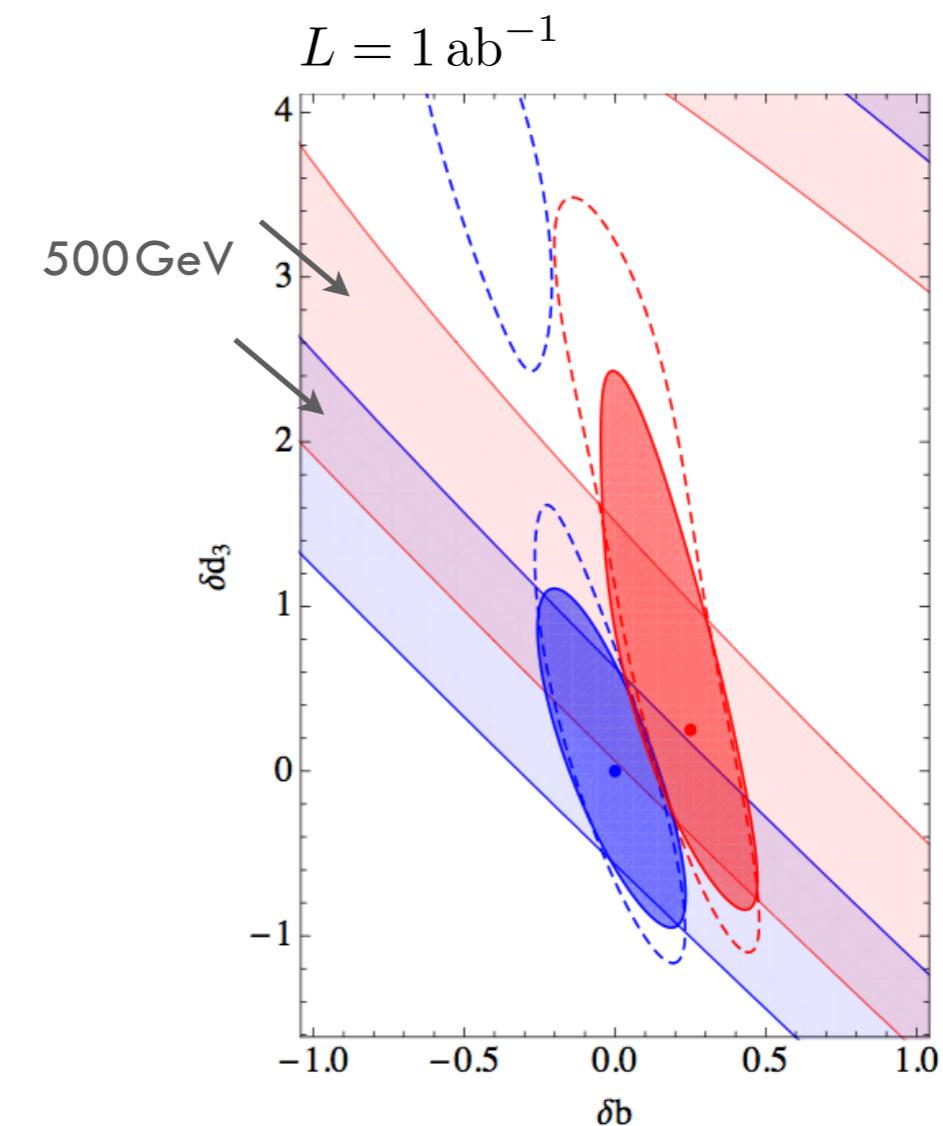
with  $\sqrt{s} = 500 \text{ GeV} - 1 \text{ TeV}$

through double Higgs-strahlung

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038



Cut on  $m_{hh}$  useful at  $\sqrt{s} = 1 \text{ TeV}$



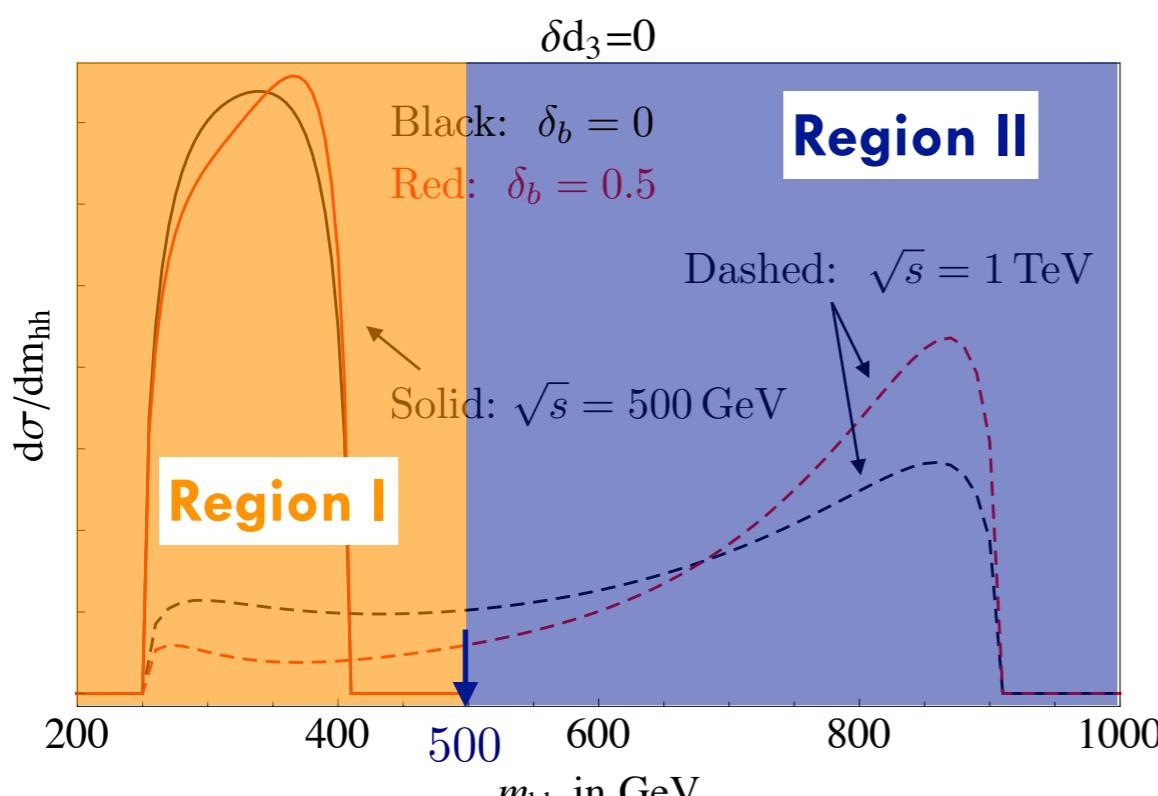
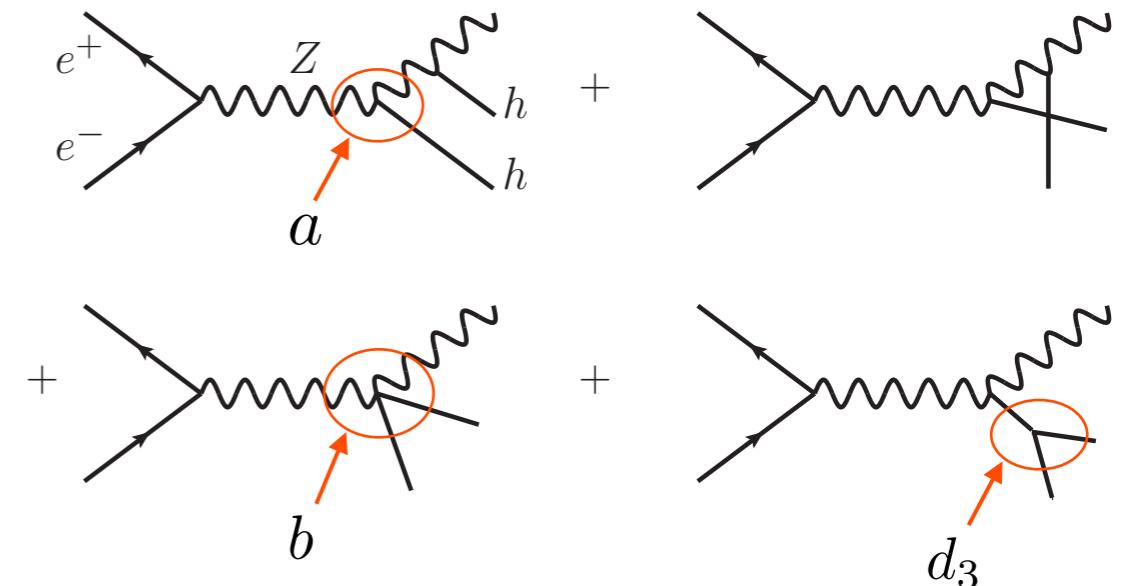
$$a^2(BR(b\bar{b})/BR(b\bar{b})_{SM}) = 1$$

# Extracting $b, d_3$ at the ILC

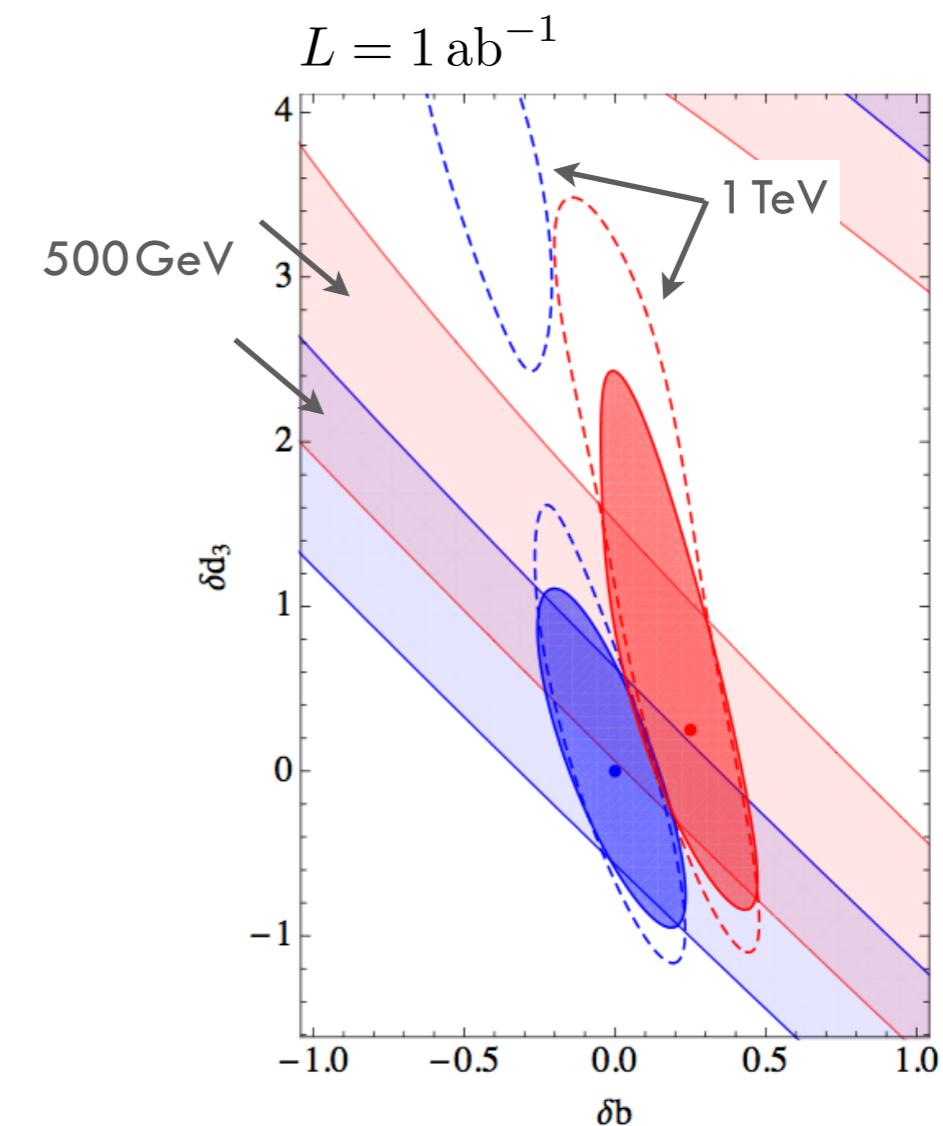
with  $\sqrt{s} = 500 \text{ GeV} - 1 \text{ TeV}$

through double Higgs-strahlung

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038



Cut on  $m_{hh}$  useful at  $\sqrt{s} = 1 \text{ TeV}$



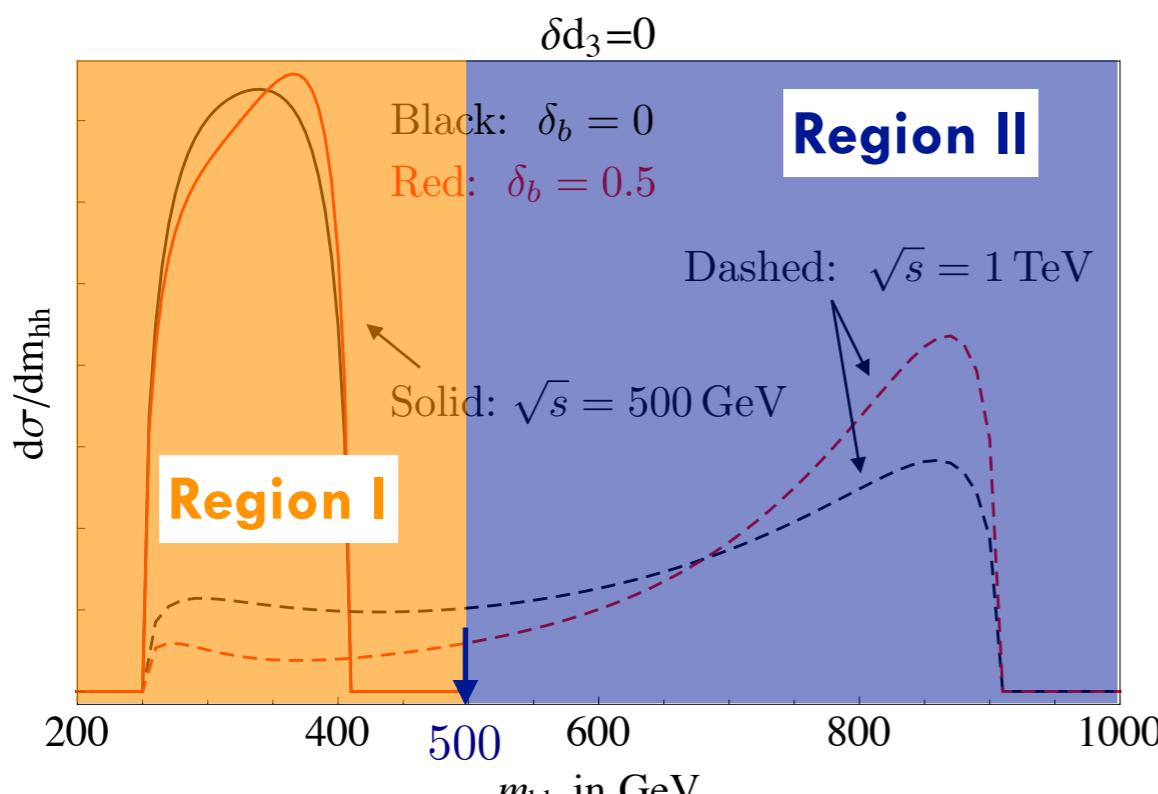
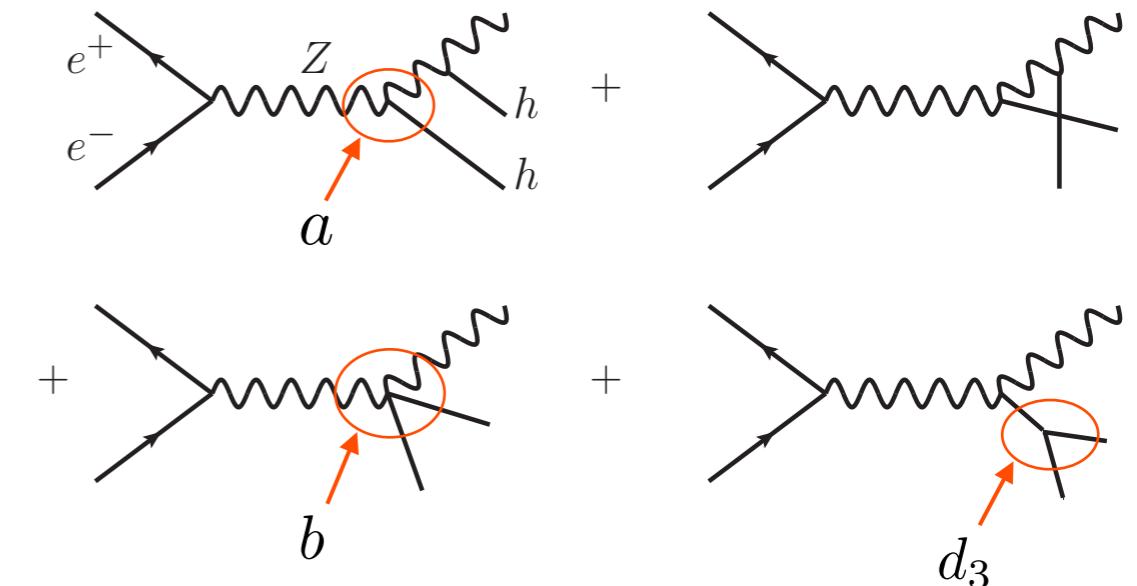
$$a^2(BR(b\bar{b})/BR(b\bar{b})_{SM}) = 1$$

# Extracting $b, d_3$ at the ILC

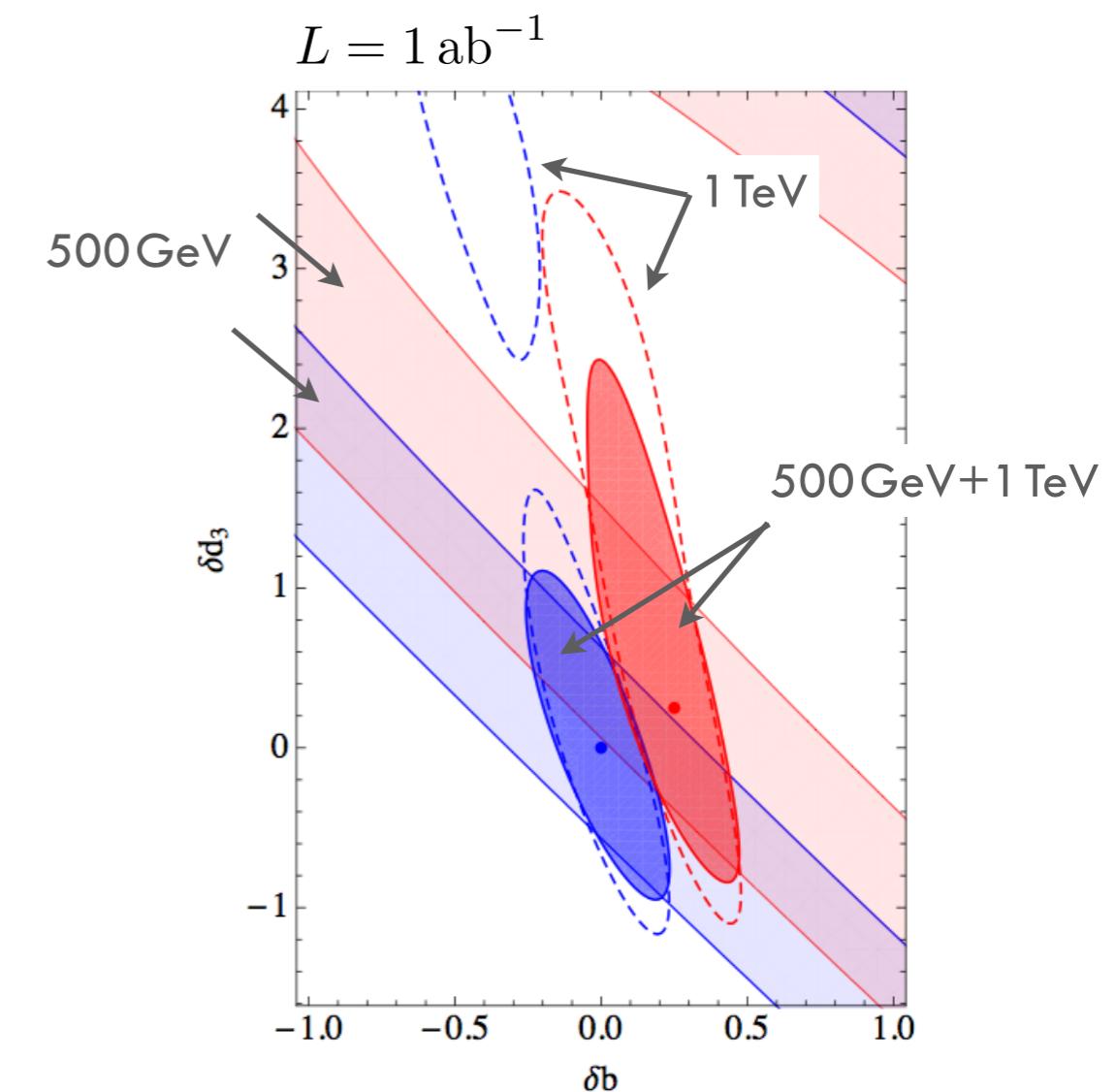
with  $\sqrt{s} = 500 \text{ GeV} - 1 \text{ TeV}$

through double Higgs-strahlung

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038



Cut on  $m_{hh}$  useful at  $\sqrt{s} = 1 \text{ TeV}$

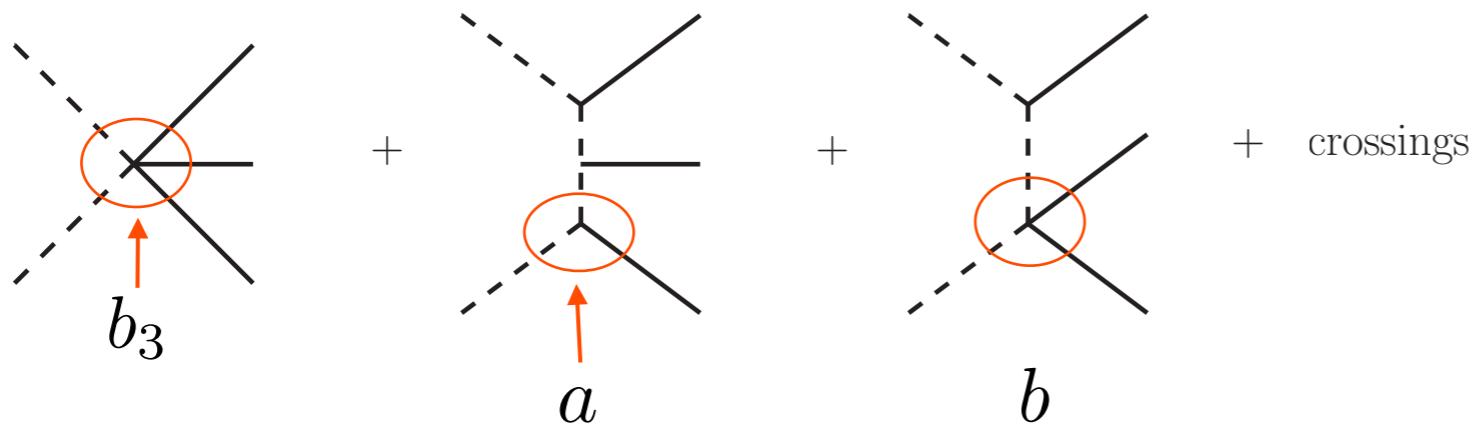


$$a^2(BR(b\bar{b})/BR(b\bar{b})_{SM}) = 1$$

## Further test of PNGB vs SILH (more difficult):

$$WW \rightarrow hh$$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038

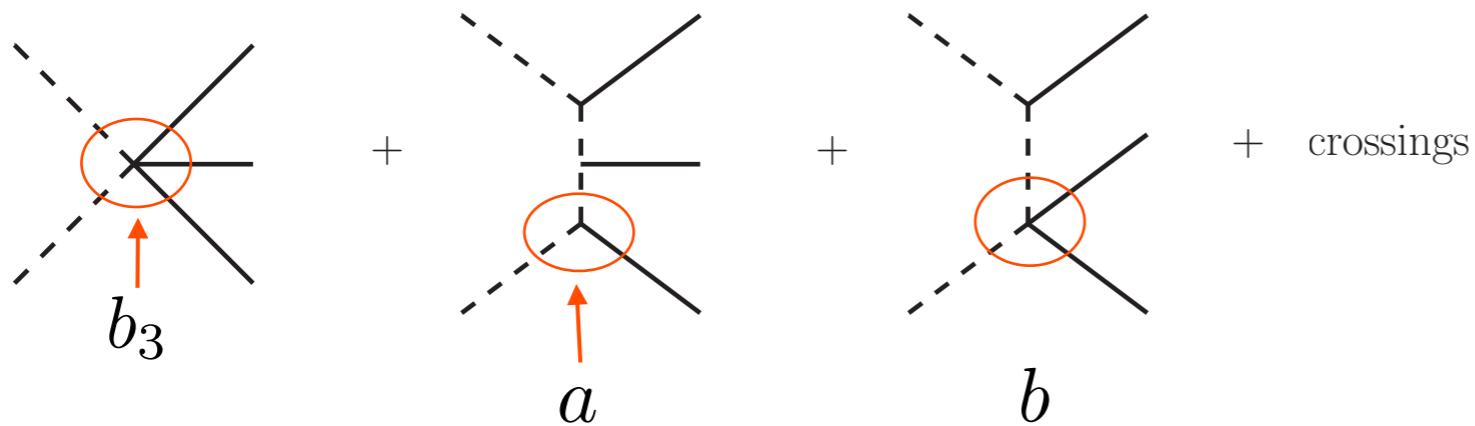


$$\mathcal{A}(\chi\chi \rightarrow hh) = \frac{i\hat{s}}{v^3} (4ab - 4a^3 - 3b_3) = 2i(c'_H - 2c_H) \frac{\hat{s}}{v^3} \left( \frac{v^4}{f^4} \right) + \dots$$

## Further test of PNGB vs SILH (more difficult):

$$WW \rightarrow hh$$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038



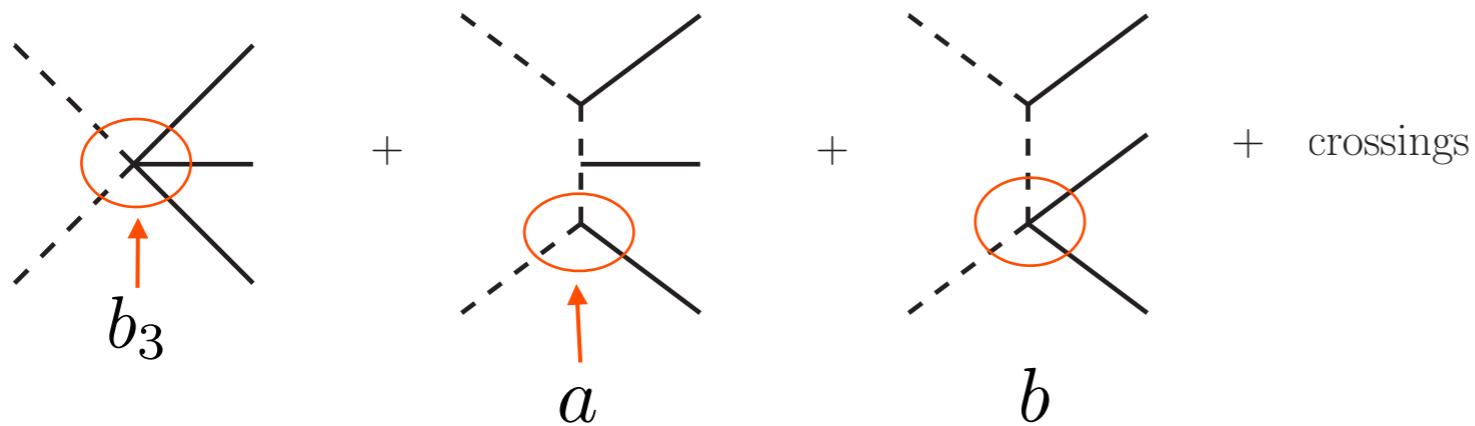
$$\mathcal{A}(\chi\chi \rightarrow hh) = \frac{i\hat{s}}{v^3} (4ab - 4a^3 - 3b_3) = 2i(c'_H - 2c_H) \frac{\hat{s}}{v^3} \left( \frac{v^4}{f^4} \right) + \dots$$

Test dim-8  
operators



## Further test of PNGB vs SILH (more difficult): $WW \rightarrow hh$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038



$$\mathcal{A}(\chi\chi \rightarrow hh) = \frac{i\hat{s}}{v^3} (4ab - 4a^3 - 3b_3) = 2i(c'_H - 2c_H) \frac{\hat{s}}{v^3} \left( \frac{v^4}{f^4} \right) + \dots$$



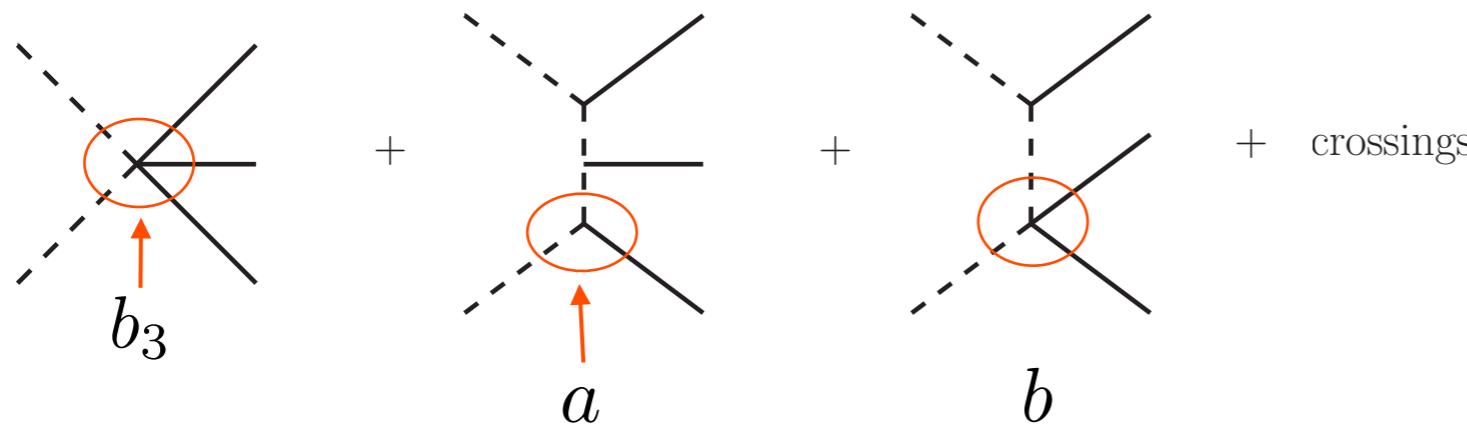
Test dim-8 operators

vanishes for a PNGB (with symmetric coset) due to  $Z_2$  parity  $\pi \rightarrow -\pi$

## Further test of PNGB vs SILH (more difficult):

$$WW \rightarrow hh$$

RC, Grojean, Pappadopulo, Rattazzi, Thamm arXiv:1309.7038



$$\mathcal{A}(\chi\chi \rightarrow hh) = \frac{i\hat{s}}{v^3} (4ab - 4a^3 - 3b_3) = 2i(c'_H - 2c_H) \frac{\hat{s}}{v^3} \left( \frac{v^4}{f^4} \right) + \dots$$



Test dim-8 operators

vanishes for a PNGB (with symmetric coset) due to  $Z_2$  parity  $\pi \rightarrow -\pi$

$\sigma$ [ab]	$\xi$						
	0	0.05	0.1	0.2	0.3	0.5	0.99
PNGB	0.32	0.46	0.71	1.47	2.41	4.13	0.30
SILH	0.32	0.71	0.87	7.56	42.89	407.9	7808

For  $\xi \gtrsim 0.3$  detectable for a SILH (PNGB disproved)