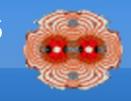


## Beam stability without eigenvalues (ongoing work)

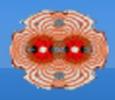


X. Buffat, N. Mounet, T. Pieloni, S. White

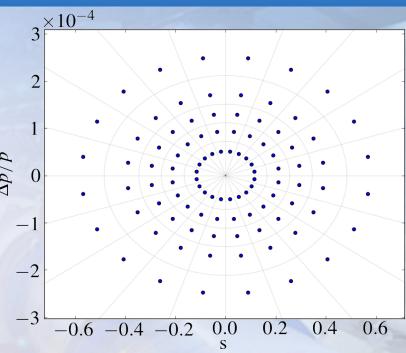
- Remainder on the circulant matrix model
- Multibunch beam breakup / coupled bunch instabilities
- Non-diagonalizable system and pseudospectrum
- First results for the LHC case



#### Circulant matrix model

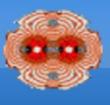


- Discretized distribution in longitudinal plane
- Build and diagonalize one turn map
- Especially usefull to study beambeam and impedance (S. White, et al, "Beam beam and Impedance", BB2013)
- Python module "BimBim"
  - Single/multi bunch Impedance based on wake tables
  - 4D/6D Beam-beam interactions
  - Perfect BbyB damper
  - Any filling scheme / IP configuration





## Multibunch beam breakup



#### 2 bunch model

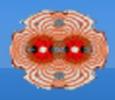
- For simplicity, let us consider 2 bunches with 1 slice,
   1 ring (i.e. Rigid bunches)
- One turn map for a single bunch : B
- One turn map for the two bunches, assuming that the distance from b1 to b2 << b2 to b1 (e.g. Train of two bunches in the LHC)

$$M = \begin{pmatrix} B & 0 \\ 0 & B \end{pmatrix} \cdot \begin{pmatrix} 1 & 0 \\ Z & 1 \end{pmatrix}$$

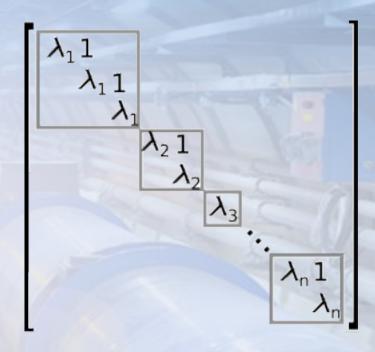
- This matrix is not diagonalizable!
  - ightarrow Eigenvectors do not provide a complete basis for  $\mathbb{C}^{\dim(M)}$
  - → The dynamic of the system is not fully described by the eigenvalues/eigenvectors



#### Jordan normal form

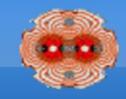


- One can always find a basis such that the matrix has the Jordan normal form →
- Vectors of this basis, which are not eigenvectors, are generalized eigenvectors
- Behavior under n<sup>th</sup> power of M:
  - Eigenvectors:  $\vec{v}_n = e^{2\pi i Qn} \vec{v}_0$
  - Generalized eigenvectors :  $\vec{v}_n = e^{2\pi i Qn} \vec{v}_0 + \dots$





#### Powers of M 2 bunch model



$$M_d = \begin{pmatrix} \lambda_1 & 1 & 0 & 0 \\ 0 & \lambda_1 & 0 & 0 \\ 0 & 0 & \lambda_2 & 1 \\ 0 & 0 & 0 & \lambda_2 \end{pmatrix}$$
 Consider a vector in the subspace associated to  $\lambda_1$  
$$\vec{V} = a_1 \vec{e}_1 + a_2 \vec{e}_2 \quad \text{with} \quad \vec{e}_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \vec{e}_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

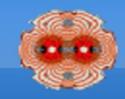
Consider a vector in the subspace associated to  $\lambda_i$ :

$$\vec{V} = a_1 \vec{e}_1 + a_2 \vec{e}_2$$
 with  $\vec{e}_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$   $\vec{e}_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ 

$$\vec{V}_{n} = \lambda_{i}^{n} a_{1} \vec{e}_{1} + \lambda_{i}^{n} a_{2} \vec{e}_{2} + \sum_{k=0}^{n-1} (\lambda_{i} a_{2})^{k} \vec{e}_{1}$$



#### Powers of M 2 bunch model



$$M_d = \begin{pmatrix} \lambda_1 & 1 & 0 & 0 \\ 0 & \lambda_1 & 0 & 0 \\ 0 & 0 & \lambda_2 & 1 \\ 0 & 0 & 0 & \lambda_2 \end{pmatrix}$$
 Consider a vector in the subspace associated to  $\lambda_1$  
$$\vec{V} = a_1 \vec{e}_1 + a_2 \vec{e}_2 \quad \text{with} \quad \vec{e}_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \vec{e}_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Consider a vector in the subspace associated to  $\lambda$ :

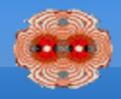
$$\vec{V} = a_1 \vec{e_1} + a_2 \vec{e_2}$$
 with  $\vec{e_1} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$   $\vec{e_2} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ 

$$\vec{V}_{n} = \lambda_{i}^{n} a_{1} \vec{e}_{1} + \lambda_{i}^{n} a_{2} \vec{e}_{2} + \sum_{k=0}^{n-1} (\lambda_{i} a_{2})^{k} \vec{e}_{1}$$

- Linear growth which depends on the initial condition
- The behavior of the system under a small perturbation is no longer independant of the perturbation
- More complicated behavior expected for higher number of bunches

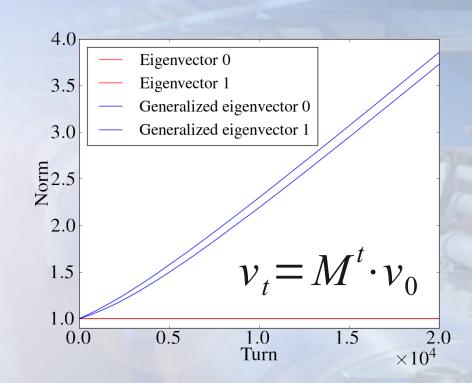


## Multibunch beam breakup



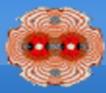
#### 2 bunch model

- Matrix model confirms this behavior
- However, mathematically, and physically, the matrix can be rendered diagonalizable
  - Multiturn wake
  - Equidistant bunches
  - Beam-beam
- Does the physcis change?

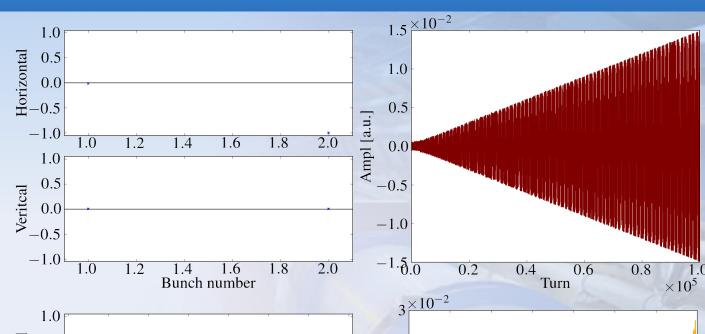


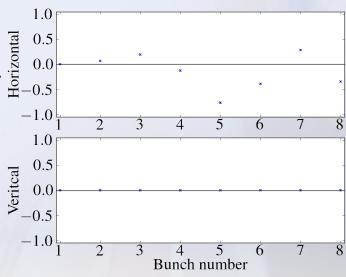


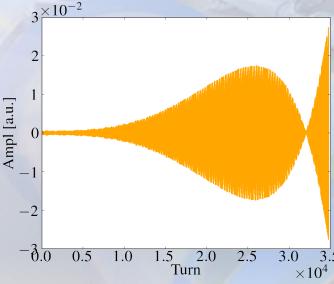
#### Coupled bunch instabilities



- Simulation with COMBI
- LHC impedance model
  - Dipolar wake
  - Quadrupolar wake
  - Multiturn wake
- Long term behavior is well described by eigenvalue approach, but short term is dominated by transient effects

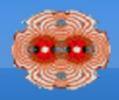








### Pseudo spectrum



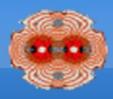
This phenomenon can be described by the pseudo spectrum

$$Spectrum(M) = \{\lambda \in \mathbb{C} | \exists \vec{v} : (M - \lambda I) \cdot \vec{v} = 0\}$$

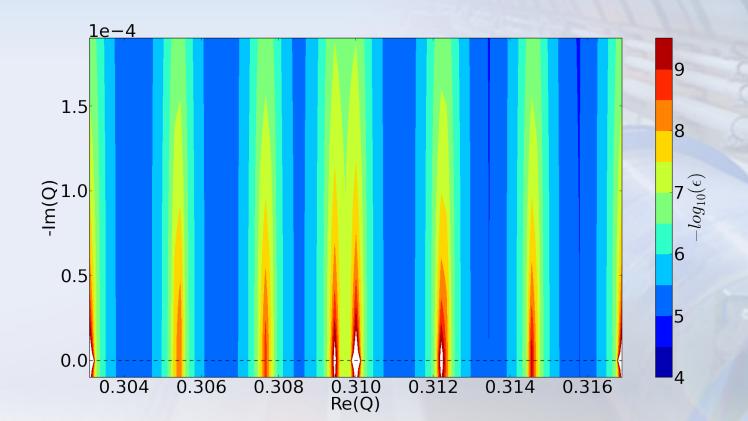
$$Pseudo spectrum(M, \epsilon) = \{\lambda \in \mathbb{C} | \exists \vec{v} : ||(M - \lambda I) \cdot \vec{v}|| < \epsilon\}$$

• For a given point of the complex plane z, the corresponding  $\varepsilon$  is given by the smallest singular value of (M-zI)

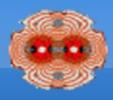




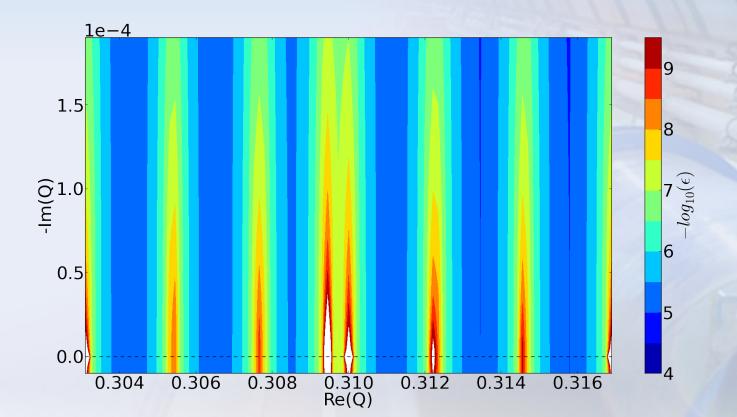
- 1 bunch
- 0.0 chromaticity, no damper







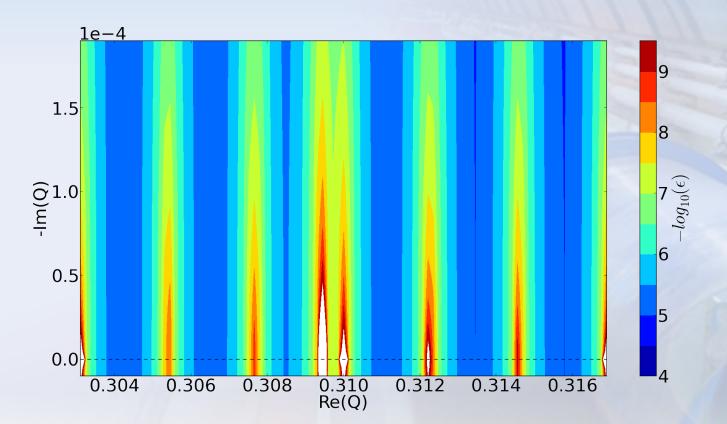
- 4 bunches
- 0.0 chromaticity, no damper



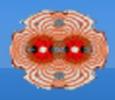




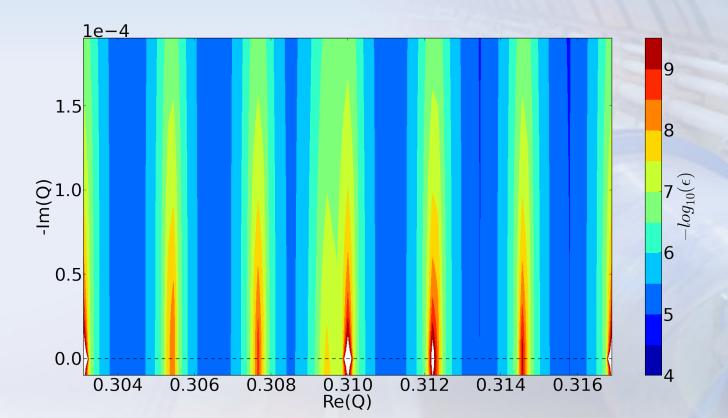
- 8 bunches
- 0.0 chromaticity, no damper



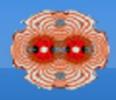




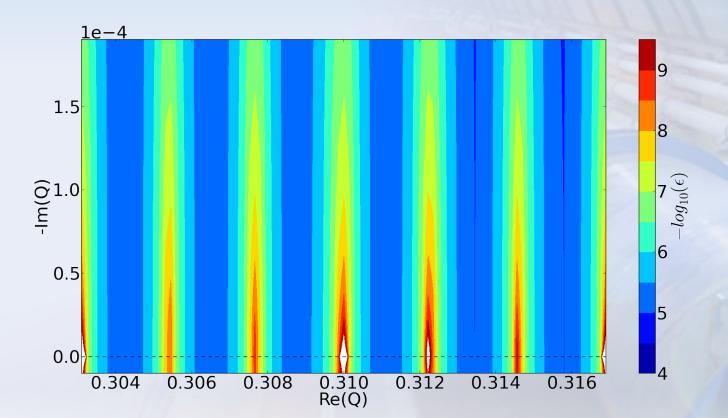
- 4 bunches
- 0.0 chromaticity, 1000 turn damper





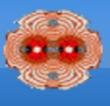


- 4 bunches
- 0.0 chromaticity, 100 turn damper

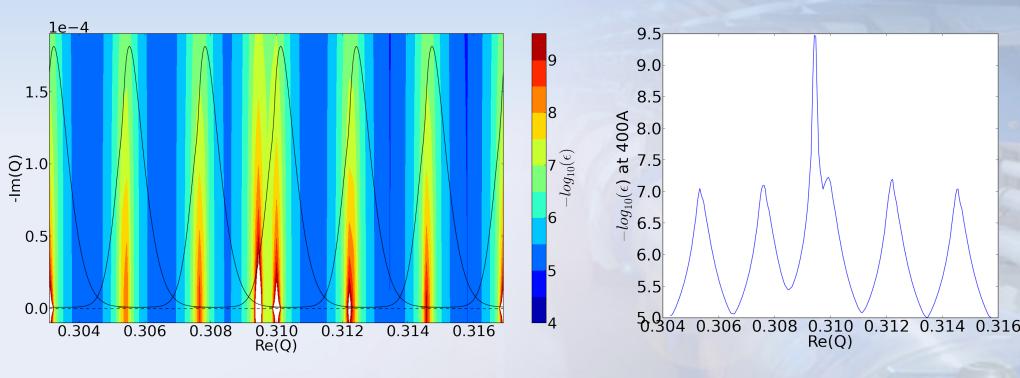




# A naive trial to include transverse non-linearities



Small tune shift: each side band can be treated separately



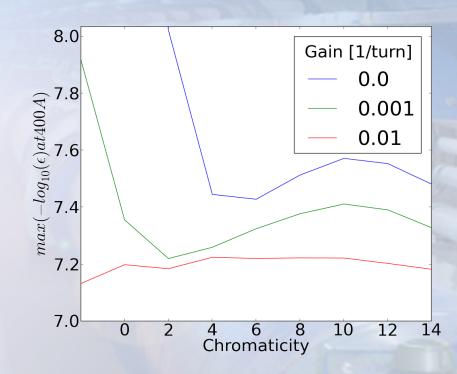
 If all eigenvalues are in the stable area, one can characterize the stability of the beam by the maximum of -log(ε) on the stability diagram



#### Chromaticity / damper



- Results in accordance with expectations
- Not valid for most cases of interest, where mode coupling is not negligible





#### Conclusion



- Coupled bunch instabilities cannot be fully treated using the standard eigenvalue approach
  - Transient growth may be expected even in systems with only decaying eigenmodes
  - Behavior depends on initial condition / external excitation
- The pseudo spectrum provides information on the behavior of such non-normal system
- Real life application less obvious than eigenvalues
- Including the effect of transverse non-linearities is also not trivial