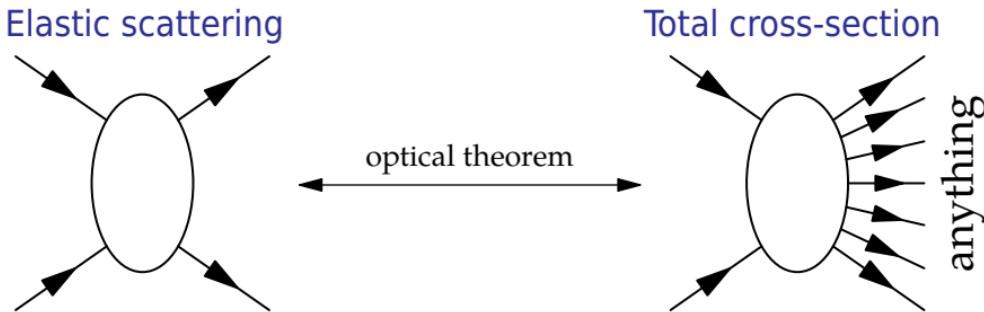


TOTEM Results on Elastic Scattering and Total Cross-Section

Jan Kašpar
on behalf of the TOTEM collaboration



EDS Blois 2013, Saariselkä, Finland
9 September, 2013



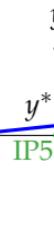
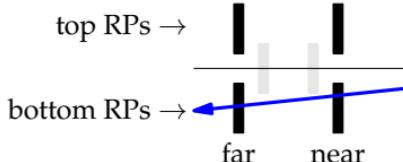
Outline

- 1) Introduction: *Roman Pot detectors, ...*
- 2) Elastic scattering: *analysis method and results*
- 3) Total cross-section: *analysis method and results*
- 4) Study of Coulomb-nuclear interference

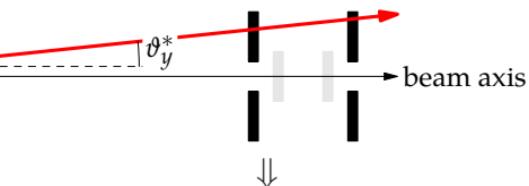
TOTEM and elastic scattering

- TOTEM shares the LHC interaction point (IP5) with CMS
- elastic scattering = 2 anti-collinear protons from the same vertex:

Left: RP station at -220 m



Right: RP station at +220 m



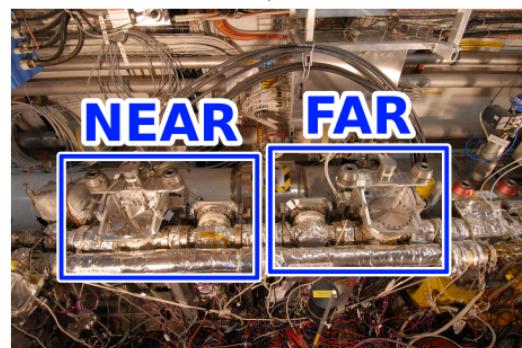
four-momentum transfer squared: t

scattering angle: $\vartheta^* \simeq \sqrt{t/p}$

azimuthal angle: φ^*

horizontal angle: $\vartheta_x^* = \vartheta^* \cos \varphi^*$

vertical angle: $\vartheta_y^* = \vartheta^* \sin \varphi^*$

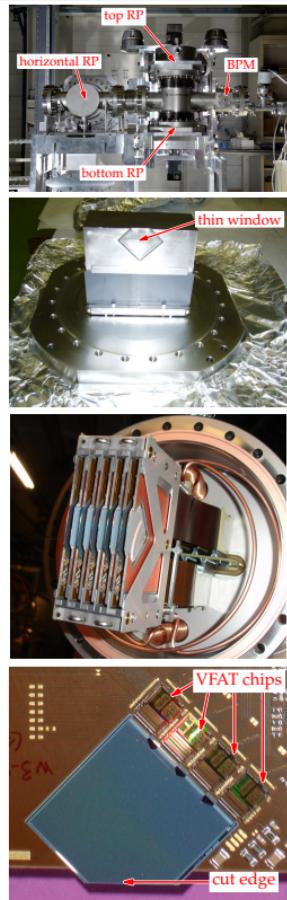


- 2 diagonals \Rightarrow control of systematics
 - left *bottom* - right *top*
 - left *top* - right *bottom*

- 2 units \Rightarrow improved:
 - event selection
 - kinematics reconstruction

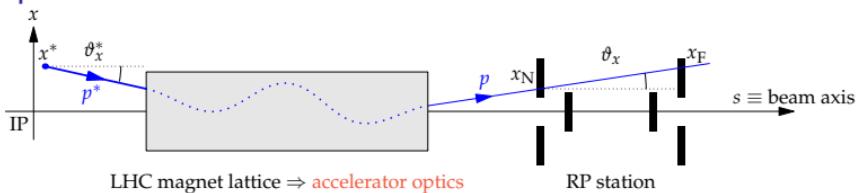
Roman Pot detectors

- each station: near and far units
- each unit: top, bottom and horizontal Roman Pots
- Roman Pot
 - movable beam-pipe insertion
 - retracted when beam unstable
 - close to beam for data taking
 - contains: 5×2 back-to-back mounted silicon sensors
- edge-less silicon sensors
 - insensitive edge (facing beam): $\approx 50 \mu\text{m}$
 - strips with pitch $66 \mu\text{m}$ oriented at 45° wrt. active edge
- VFAT: trigger-capable read-out chip



Proton measurement with RPs

- proton transport IP5 → RP detectors:

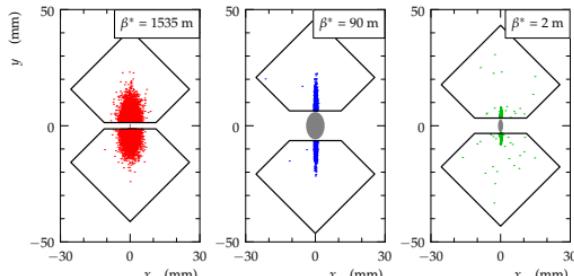


- optics

hit position at RP optical functions proton kinematics at IP

$$x(\text{RP}) = (\text{effective length } L_x) \cdot (\text{scattering angle } \vartheta_x^*) + (\text{magnification } v_x) \cdot (\text{vertex } x^*) + (\text{dispersion } D_x) \cdot (\text{rel. momentum loss } \xi \equiv \frac{\Delta p}{p})$$

- example: elastic sample seen with 3 different optics:



⇒ *optics knowledge essential*
⇓
TOTEM can improve optics accuracy

- entirely data-driven
- two diagonals, several LHC fills \simeq different experiments \Rightarrow control of systematics

1. Alignment

- prior to data-taking: collimator-like beam-based alignment
- offline alignment: *relative* (analysis of track fit residuals) and *absolute wrt. beam* (symmetries of elastic scattering)

2. Kinematics reconstruction

- tracks in RPs \rightarrow kinematics at IP ($\xi = 0 \Rightarrow$ relatively easy)
- choice of formulae (using *Near* and *Far* RPs) \rightarrow minimisation of systematics, typically:

$$\theta_x^* = \frac{x^F - x^N}{L_x^F - L_x^N}, \quad \theta_y^* = \frac{1}{2} \left(\frac{y^N}{L_y^N} + \frac{y^F}{L_y^F} \right)$$

3. Elastic tagging

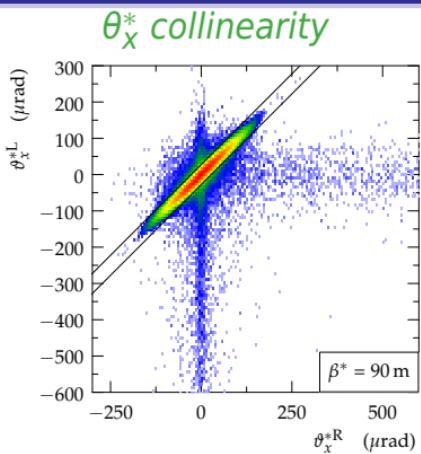
- angles left = angles right (tolerance set by beam divergence: higher β^* \Rightarrow more stringent cut)
- vertex left = vertex right
- protons $\xi \approx 0 \Rightarrow$ correlation hit position vs. track angle at RP

4. Background subtraction

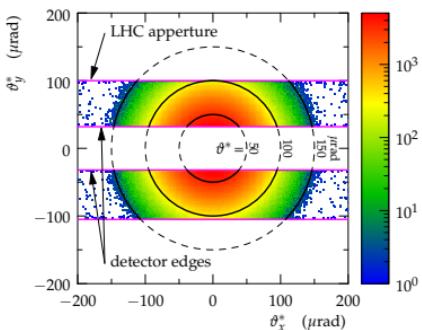
- typically needed only for low β^* optics
- interpolation of event distribution surrounding the signal (tagged) region

5. Acceptance corrections

- RP sensors have finite size \Rightarrow low $|\theta_y^*|$ cut
- LHC apertures \Rightarrow high $|\theta_y^*|$ cut
- azimuthal symmetry (verified) \Rightarrow geometrical correction (+ smearing around edges)



acceptance correction



6. Unfolding of resolution effects

- angular resolution (better for high β^*): left-right proton comparison
- Monte Carlo calculation \Rightarrow impact on t -distribution

7. Inefficiency corrections

- uncorrelated 1-RP inefficiencies: repeat tagging with 3 RPs only and check the signal in 4th RP
- near-far correlated RP inefficiencies (showers from near to far RP)
- “pile-up” = elastic event + another track in a RP (prob. from zero-bias stream)

8. Luminosity

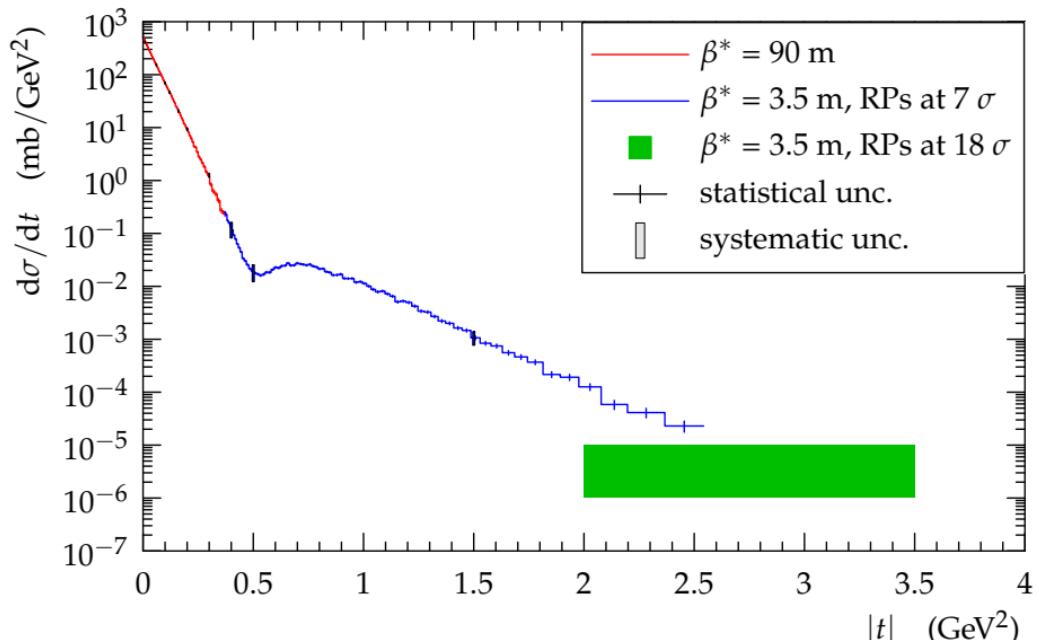
- from CMS (if available), uncertainty $\approx 4\%$
- from TOTEM (details later on)

9. Study of systematic uncertainties

- final $d\sigma/dt \Rightarrow$ input to Monte-Carlo simulation
- any analysis parameter: discrepancy simulation vs. reconstruction \Rightarrow study impact on t -distribution

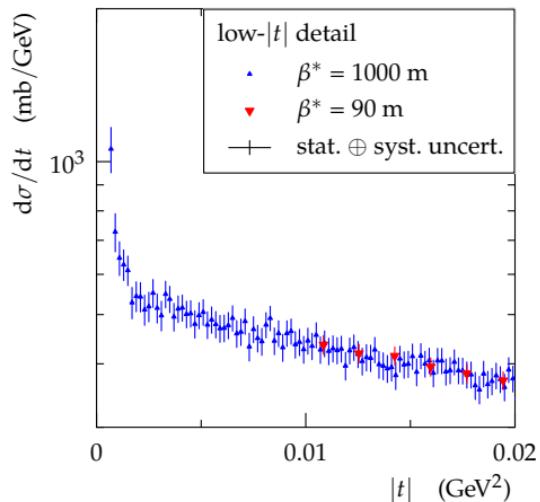
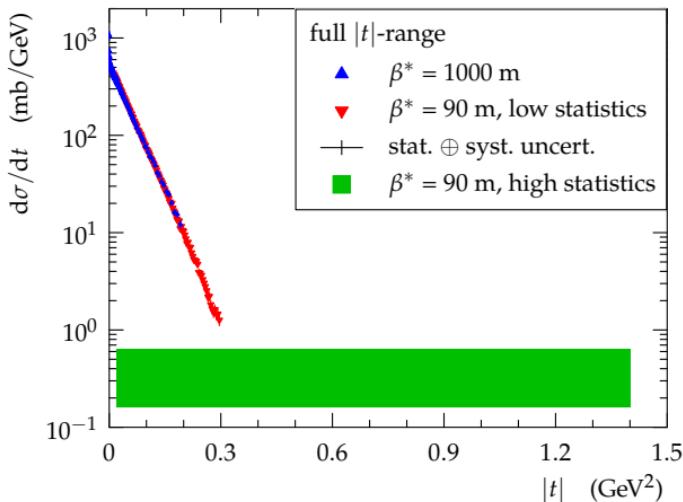
Elastic scattering results : $\sqrt{s} = 7 \text{ TeV}$

β^*	RP approach	$ t $ range	el. events	status
90 m	4.8 to 6.5 σ	0.005 to 0.4 GeV^2	1 M	[EPL 101 (2013) 2100]
3.5 m	7 σ	0.4 to 2.5 GeV^2	66 k	[EPL 95 (2011) 41001]
3.5 m	18 σ	\approx 2 to 3.5 GeV^2	10 k	anal. advanced



Elastic scattering results : $\sqrt{s} = 8 \text{ TeV}$

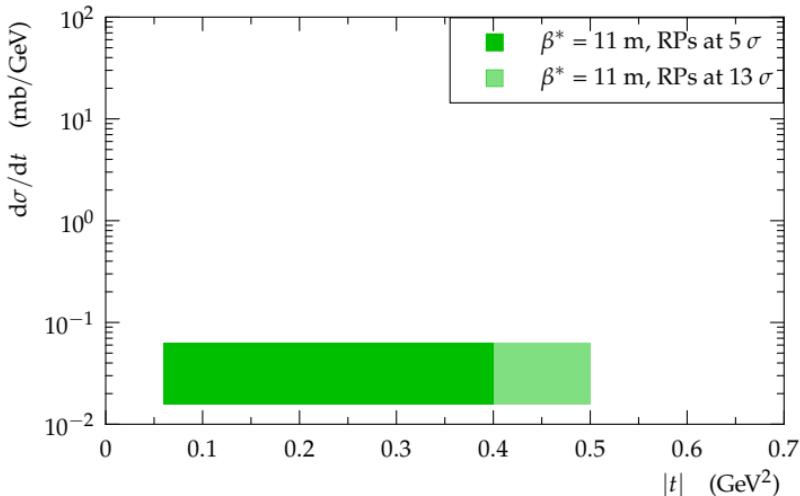
β^*	RP approach	$ t $ range	el. events	status
1000 m	3 or 10 σ	0.0006 to 0.2 GeV^2	352 k	publ. in prep.
90 m	6 to 9.5 σ	0.01 to 0.3 GeV^2	0.68 M	[PRL 111 (2013)]
90 m	9.5 σ	0.02 to 1.4 GeV^2	7.2 M	anal. advanced



- dip well visible in the combined $\beta^* = 90 \text{ m}$ data

Elastic scattering results : $\sqrt{s} = 2.76 \text{ TeV}$

β^*	RP approach	$ t $ range	el. events	status
11 m	5σ	0.06 to 0.4 GeV^2	45 k	anal. started
11 m	13σ	0.4 to 0.5 GeV^2	2 k	anal. started



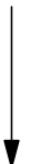
- $\beta^* = 11 \text{ m}$ optics tuning in progress ($\rightarrow t$ values preliminary)
- LHC aperture(s) at $\approx 14 \sigma$
- dip (expected at $|t| \approx 0.6 \text{ GeV}^2$) unlikely to be visible

3 complementary methods:

$$\rho \equiv \left. \frac{\Re \mathcal{A}_{\text{el}}}{\Im \mathcal{A}_{\text{el}}} \right|_{t=0}$$

elastic observables only:

$$\sigma_{\text{tot}}^2 = \frac{16\pi}{1 + \varrho^2} \frac{1}{\mathcal{L}} \left. \frac{dN_{\text{el}}}{dt} \right|_0$$



$$\sigma_{\text{tot}}$$

q -independent:

$$\sigma_{\text{tot}} = \frac{1}{\mathcal{L}} (N_{\text{el}} + N_{\text{inel}})$$

luminosity-independent:

$$\sigma_{\text{tot}} = \frac{16\pi}{1 + \varrho^2} \frac{dN_{\text{el}}/dt|_0}{N_{\text{el}} + N_{\text{inel}}}$$

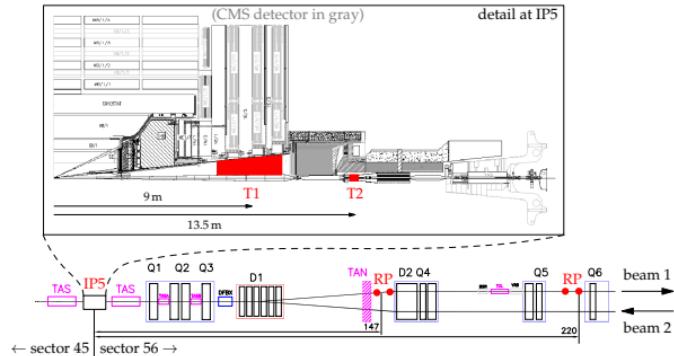
$$\mathcal{L} = \frac{1 + \varrho^2}{16\pi} \frac{(N_{\text{el}} + N_{\text{inel}})^2}{dN_{\text{el}}/dt|_0}$$

N_{el} from RPs

N_{inel} from T2

\mathcal{L} from CMS

ρ from COMPETE or TOTEM



Forward inelastic telescope T2

- detects charged particles at $5.3 < |\eta| < 6.5$
- $\approx 95\%$ of inelastic events seen (enough to detect 1 track!)

Inelastic cross-section analysis

1) *Raw rate*: event counting with T2

↓ experimental corrections: *trigger and reconstruction inefficiencies, beam-gas event suppression, pile-up consideration*

2) *Visible rate*: visible with T2 in perfect conditions

↓ recovery of events with no tracks in T2: *T1-only events, events with gap over T2, low-mass diffraction, cen. diff. without tracks in T1 and T2*

3) *Physics rate*: true rate of inelastic events

- only one major Monte-Carlo-based correction: *low-mass diffraction*
⇒ but can be constrained from data ($\sigma_{\text{tot}}^{\text{RP}} - \sigma_{\text{el}}^{\text{RP}} - \sigma_{\text{visible}}^{\text{T2}}$)

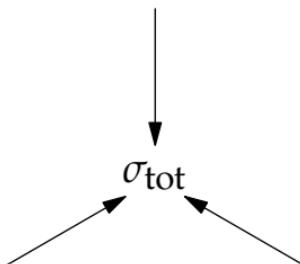
Total cross-section : $\sqrt{s} = 7$ and 8 TeV results

$\sqrt{s} = 7$ TeV

elastic observables only:

$$\sigma_{\text{tot}}^2 = \frac{16\pi}{1 + \varrho^2} \frac{1}{\mathcal{L}} \left. \frac{dN_{\text{el}}}{dt} \right|_0$$

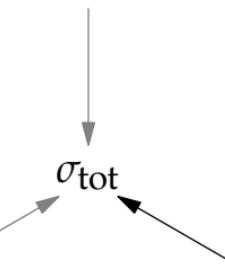
$$\sigma_{\text{tot}} = (98.6 \pm 2.3) \text{ mb}$$



$\sqrt{s} = 8$ TeV

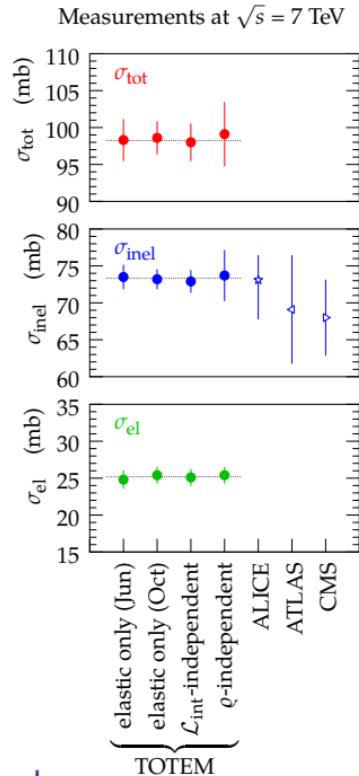
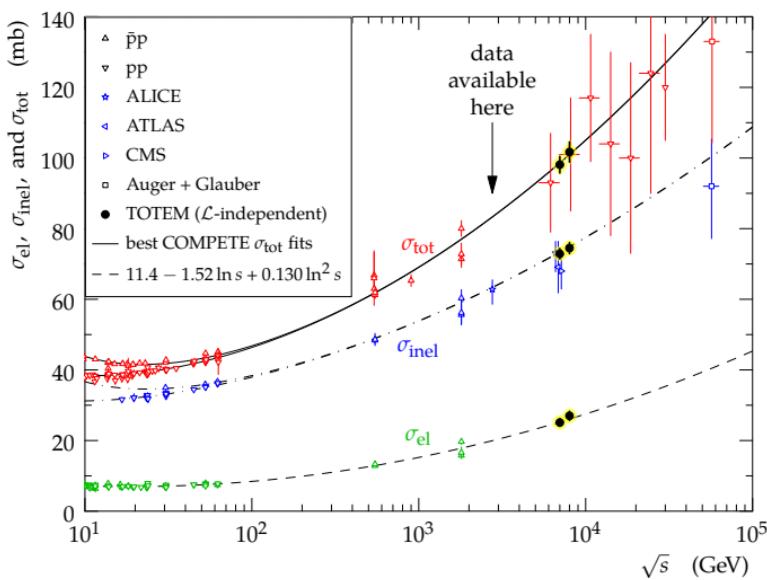
elastic observables only:

$$\sigma_{\text{tot}}^2 = \frac{16\pi}{1 + \varrho^2} \frac{1}{\mathcal{L}} \left. \frac{dN_{\text{el}}}{dt} \right|_0$$



- CMS luminosity unavailable
- \mathcal{L} from luminosity-independent method
⇒ normalisation of $d\sigma/dt$ both at $\beta^* = 90$ and 1000 m

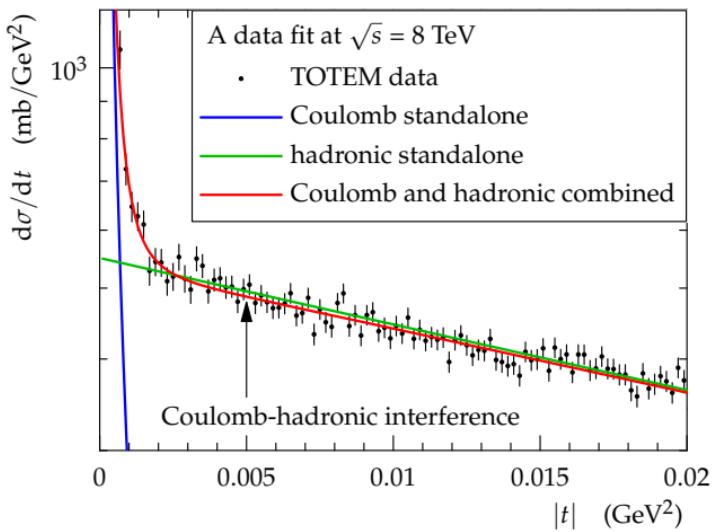
Total cross-section : Results in context



- analysis at $\sqrt{s} = 2.76$ TeV: all three methods planned
 - elastic analysis: ongoing
 - inelastic analysis: almost finished

Coulomb-nuclear interference at 8 TeV

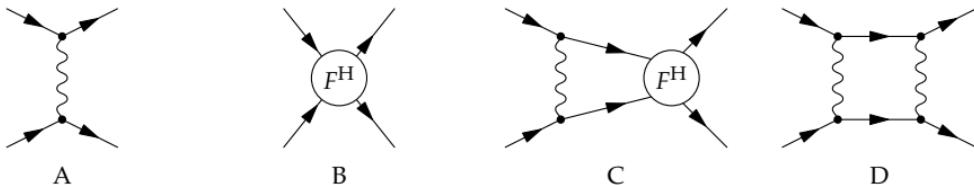
- $\beta^* = 1000$ m : $|t|$ as low as $6 \cdot 10^{-4}$ GeV $^2 \Rightarrow$ *observed Coulomb-nuclear interference* (between Coulomb/electromagnetic and nuclear/strong interactions):



- interesting aspects
 - interference \Rightarrow *determination of phase* of nuclear amplitude
 - separation of Coulomb/nuclear effects \Rightarrow *methodically better determination of σ_{tot}*

$$\sigma_{\text{tot}}^{(\text{nuclear})} \propto \Im A_{\text{el}}^{\text{nuclear}}(t = 0)$$

$$\frac{d\sigma}{dt} \propto |\mathcal{A}^{C+N}|^2, \quad \mathcal{A}^{C+N} = \text{interference formula}(\mathcal{A}^C, \mathcal{A}^N)$$



- *Coulomb amplitude \mathcal{A}^C* : well known (QED, form-factors measured)
- *Nuclear amplitude \mathcal{A}^N*
 - *modulus*: constrained by TOTEM data ⇒ parametrised:

$$\exp(b_1 t + b_2 t^2 + \dots) \quad N_b = \text{number of } b_i \text{ parameters} = 1 \text{ to } 3$$
 - *phase*: weak guidance from data ⇒ test a range of theoretical alternatives
- *interference formula*
 - simplified West-Yennie (SWY) [Phys. Rev. 172 (1968) 1413-1422]
 - traditional but
 - only compatible with constant phase and purely exponential modulus
 - Kundrák-Lokajíček (KL) [Z. Phys. C63 (1994) 619-630]
 - no \mathcal{A}^N limitations

general approach: exploration

- constant phase – the simplest choice

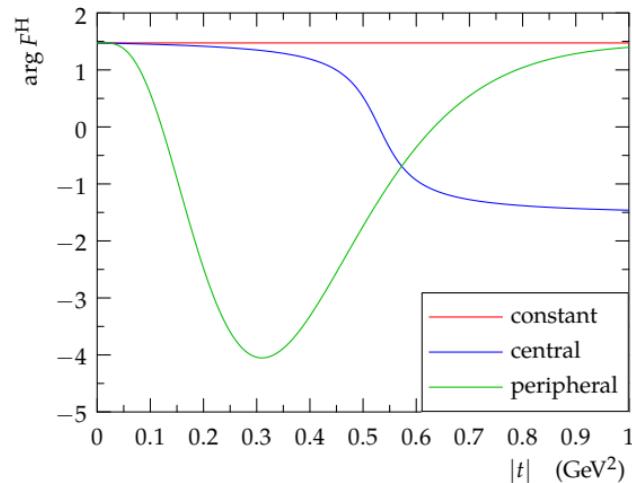
$$\arg \mathcal{A}^N = p_0$$

- central phase – similar shape as in many phenomenological models

$$\arg \mathcal{A}^N = \frac{\pi}{2} - \text{atan} \frac{\rho_0}{1 - \frac{t}{t_d}}, \quad \rho_0 = \frac{1}{\tan p_0}$$

$$t_d \approx -0.53 \text{ GeV}^2$$

- peripheral phase [Z. Phys. C63 (1994) 619-630] – expected order in impact parameter space: elastic collisions more peripheral than inelastic $\langle b^2 \rangle^{\text{el}} > \langle b^2 \rangle^{\text{inel}}$

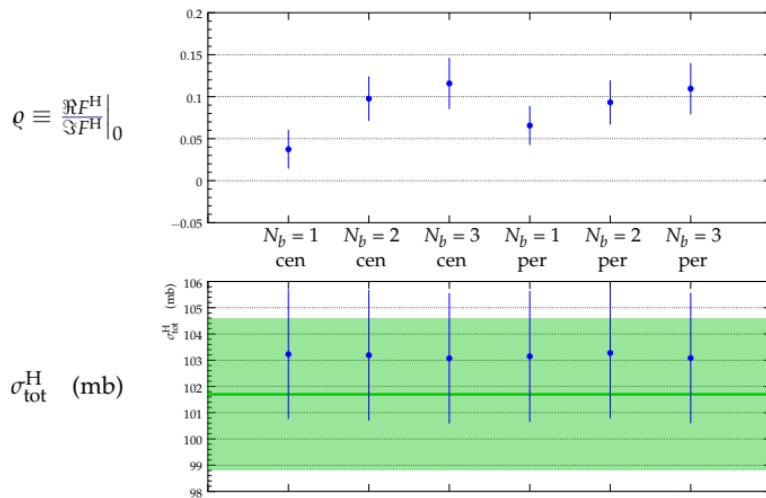


$$\arg \mathcal{A}^N = p_0 + A \exp \left[\kappa \left(\ln \frac{t}{t_m} - \frac{t}{t_m} + 1 \right) \right]$$

$$A \approx 5.53, \quad \kappa \approx 4.01, \quad t_m \approx -0.310 \text{ GeV}^2$$

Coulomb-nuclear interference : Data fits

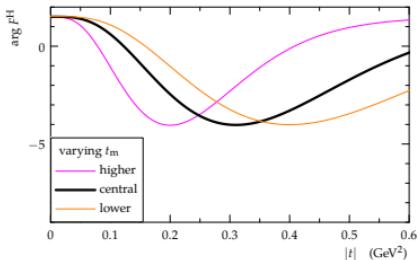
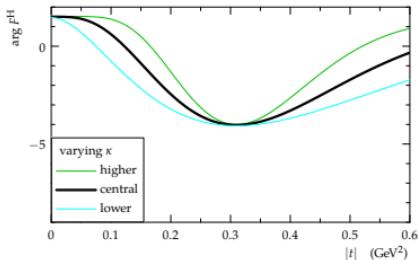
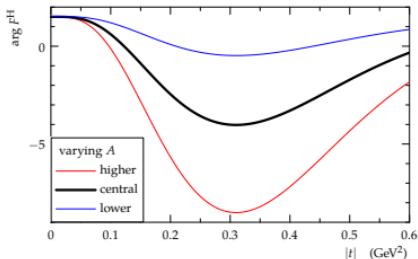
- data fits → for every parameter: value and uncertainty
 - full $|t|$ -range: $6 \cdot 10^{-4}$ to 0.2 GeV^2
 - generalised χ^2 (full covariance matrix)
 - typical $\chi^2/\text{"ndf"} \approx 1$
 - nuclear phase: only p_0 (value at $t = 0$) free parameter
- fits with constant and central phase: undistinguishable
- fits with $N_b = 1$ and KL or SWY interference formula: undistinguishable



- indications that $N_b = 1$ insufficient to describe data
- σ_{tot} : very stable results
- green line and band:
previous $\beta^* = 90 \text{ m}$ results
[PRL 111 (2013) 012001]

Coulomb-nuclear interference : Variation of peripheral phase

- data fits: only p_0 is free
- probe influence of phase shape: vary peak amplitude, width and position:

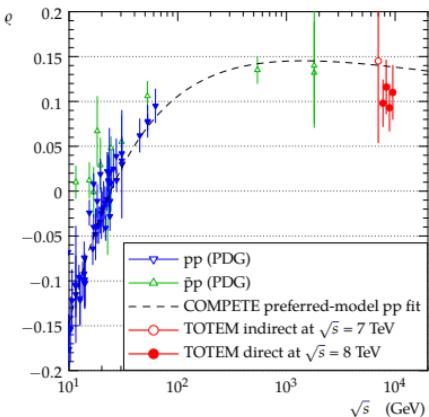


⇓

range of fit results

Coulomb-nuclear interference : Results

$\rho \rightarrow$



COMPETE: preferred model
and band from all models

TOTEM: red = fit
uncertainty, cyan = band
from varying peripheral
phase

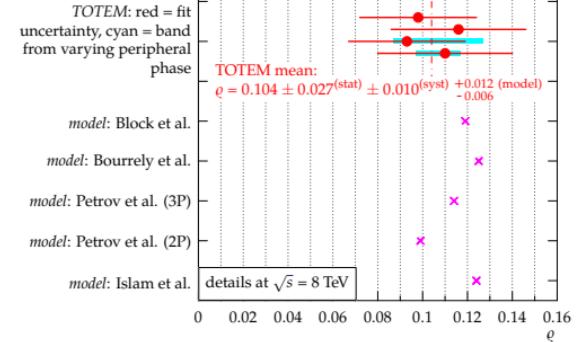
model: Block et al.

model: Bourrely et al.

model: Petrov et al. (3P)

model: Petrov et al. (2P)

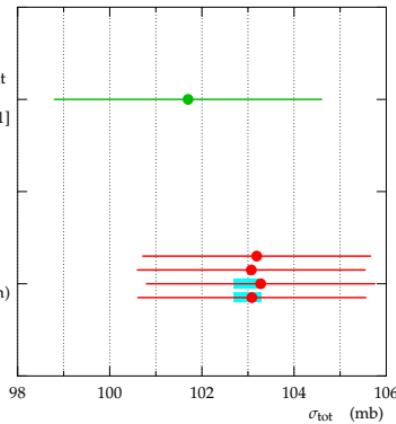
model: Islam et al.
details at $\sqrt{s} = 8$ TeV



$\sigma_{\text{tot}} \rightarrow$

previous measurement
($\beta^* = 90$ m)
[PRL 111 (2013) 012001]

this analysis
($\beta^* = 1000$ m)



Elastic differential cross-section

- **7 TeV**
 - $\beta^* = 90$ m and medium- $|t|$ at $\beta^* = 3.5$ m: published
 - high- $|t|$ at $\beta^* = 3.5$ m: advanced analysis
- **8 TeV**
 - $\beta^* = 1000$ m: publication ongoing
 - $\beta^* = 90$ m: advanced analysis
- **2.76 TeV**
 - $\beta^* = 11$ m: analysis ongoing

Total cross-section

- **7 TeV**
 - $\beta^* = 90$ m: published
- **8 TeV**
 - $\beta^* = 90$ m: published
 - $\beta^* = 1000$ m: publication ongoing (+ separation Coulomb/nuclear effects)
- **2.76 TeV**
 - $\beta^* = 11$ m: elastic analysis started, inelastic ready

Coulomb-nuclear interference studies

- **8 TeV**
 - $\beta^* = 1000$ m: publication ongoing

Backup

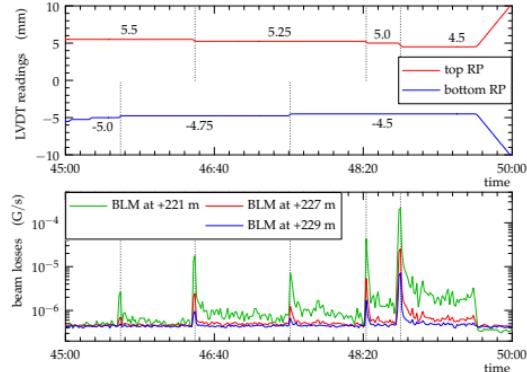
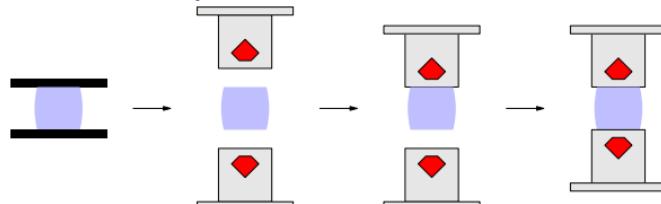


- RPs = movable insertions \Rightarrow each run at different positions
- required angular precision $\mu\text{rad} \Rightarrow \mu\text{m}$ alignment precision needed
 \Downarrow
- two types of alignment needed
 - alignment of mechanical RP edges \rightarrow for machine protection
 - alignment of RP sensors \rightarrow for physics
- need alignment *wrt. the beam*
 \Downarrow

3-step alignment procedure:

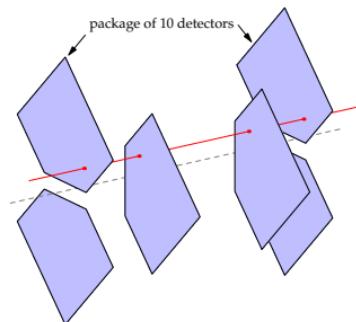
- 1) *Collimation alignment*: RP alignment wrt. beam, rough sensor alignment

- procedure prior to data taking
- standard procedure for LHC collimators

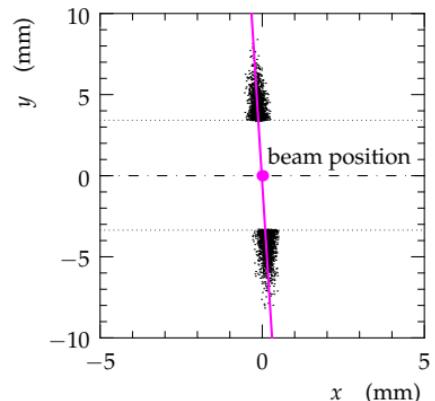
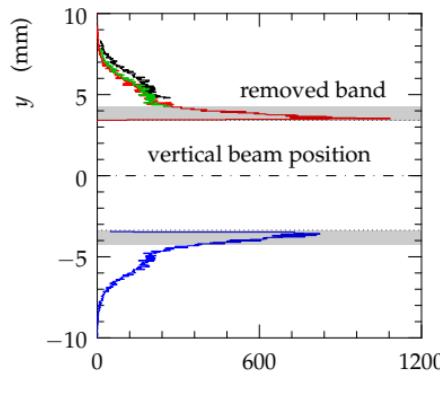
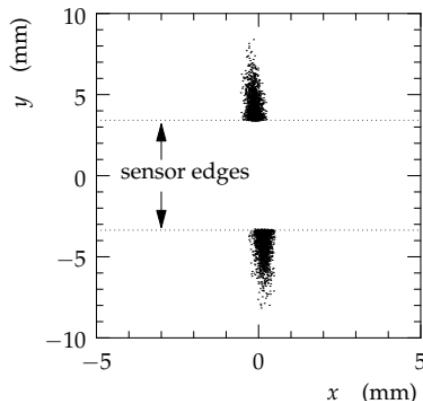


2) *Track-based alignment*: relative alignment among sensors

- RP station: no magnetic field → straight tracks
- misalignments → residuals
- residual analysis → alignment corrections
- overlap between horizontal and vertical RPs → relative alignment among all sensors
- singular/weak modes: e.g. overall shift/rotation
⇒ need further alignment step

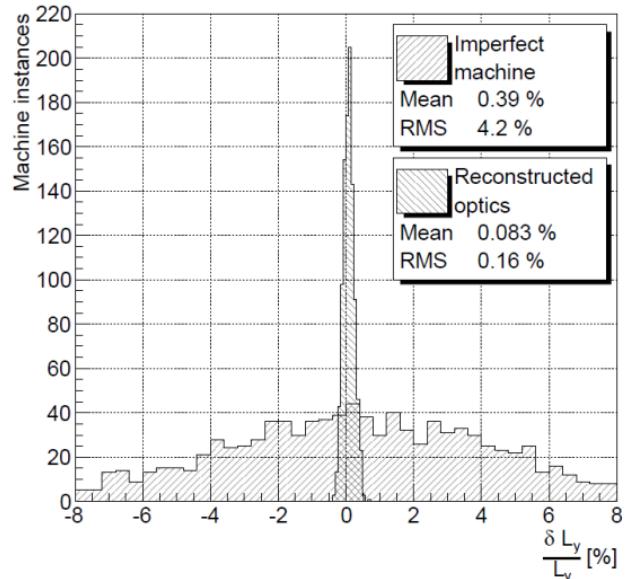
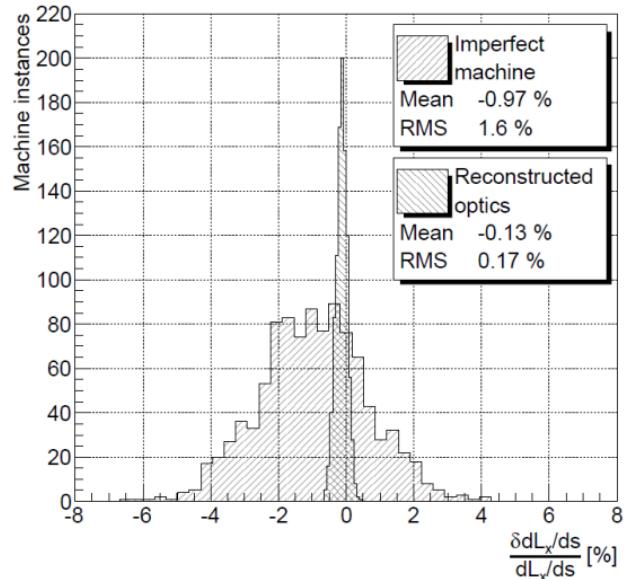


3) *Alignment with elastic scattering*: sensor alignment wrt. beam



- optics imperfection sources
 - power-converter error: $\Delta I/I \approx 10^{-4}$
 - magnet transfer function: $\Delta B/B \approx 10^{-3}$
 - magnet rotation (< 1 mrad) and displacements (< 0.5 mm)
 - magnet harmonics ($\Delta B/B \approx 10^{-4}$)
 - beam momentum offset: $\Delta p/p \approx 10^{-3}$
 - beam crossing-angle uncertainty
- optics determination
 - direct measurement – difficult
 - indirect from TOTEM observables
- TOTEM optics determination – variation of magnet/beam parameters (within tolerances) to match TOTEM observables:
 - L_y^L/L_y^R
 - $\frac{dL_y}{ds}/L_y$
 - $s(L_x = 0)$
 - xy coupling (tilts in xy plane)
 - ...

example for $\beta^* = 3.5 \text{ m optics}$



- optics uncertainty reduced:

x projection: from 1.6% to 0.17%

y projection: from 4.2% to 0.16%

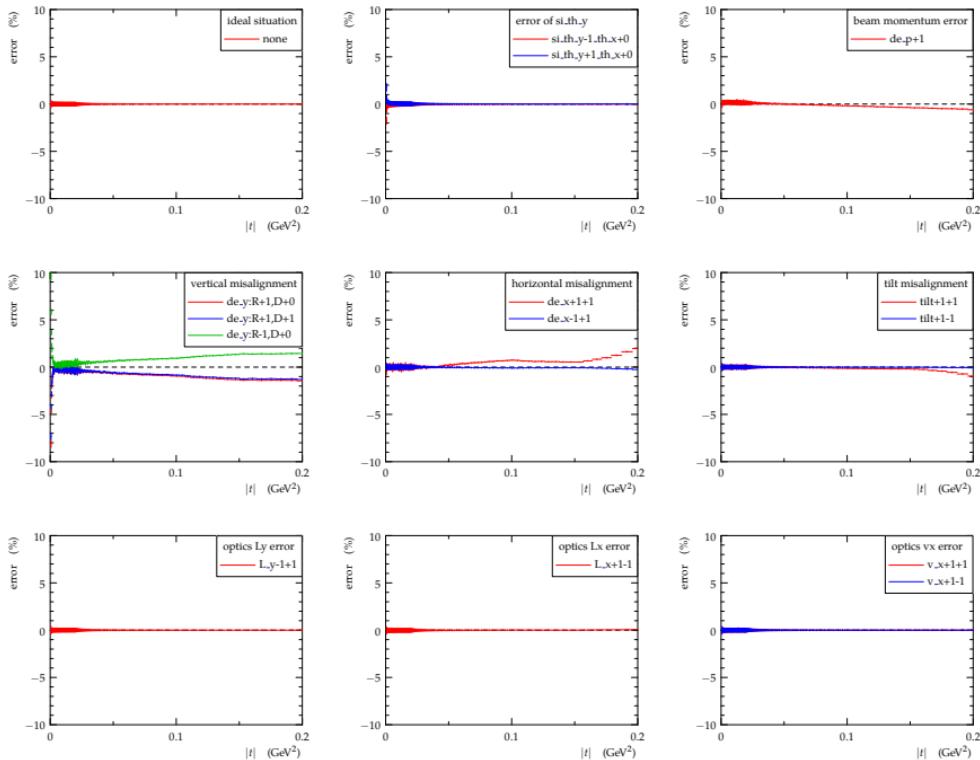
[H. Niewiadomski, Roman Pots for beam diagnostic, OMCM, CERN, 20-23.06.2011]

[H. Niewiadomski, F. Nemes, LHC Optics Determination with Proton Tracks, IPAC'12, Louisiana, USA, 20-25.05.2012]

Backup : Systematic studies

- Monte-Carlo: offset in 1 parameter \Rightarrow impact on t -distribution

example for $\beta^* = 1000 \text{ m optics}$



- simplified West-Yennie formula (SWY)
 - *limitation*: derived for *constant slope B* (1 b_i parameter only) and *constant hadronic phase*
 - acts as simple interference phase (i.e. ϕ is real-valued)

$$F^{C+H} = F^C e^{i\alpha\phi} + F^H, \quad \phi = -\left(\frac{B|t|}{2} + \gamma\right)$$

- Kundrát-Lokajíček formula (KL)
 - any slope B , any hadronic phase
 - more complicated effect (ψ complex in general)

$$F^{C+H} = F^C + F^H e^{i\alpha\psi}$$

$$\begin{aligned} \psi(t) &= \mp \int_{t_{\min}}^0 dt' \ln \frac{t'}{t} \frac{d}{dt'} \mathcal{F}^2(t') \pm \int_{t_{\min}}^0 dt' \left(\frac{F^H(t')}{F^H(t)} - 1 \right) \frac{I(t, t')}{2\pi} \\ I(t, t') &= \int_0^{2\pi} d\phi \frac{\mathcal{F}^2(t'')}{t''}, \quad t'' = t + t' + 2\sqrt{tt'} \cos\phi \end{aligned}$$

