

Recent Results in Polarized Proton-Proton Elastic Scattering at STAR

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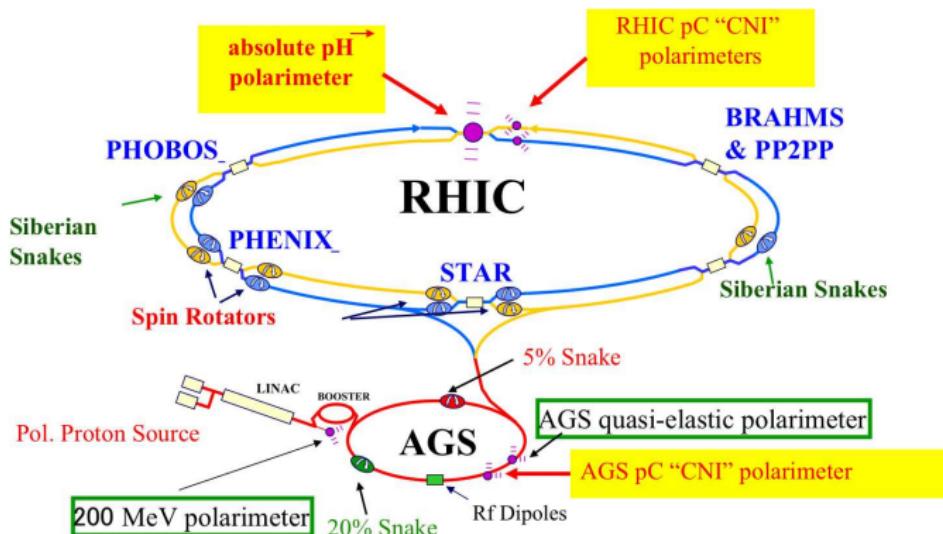
On behalf of STAR Collaboration

September 9, 2013

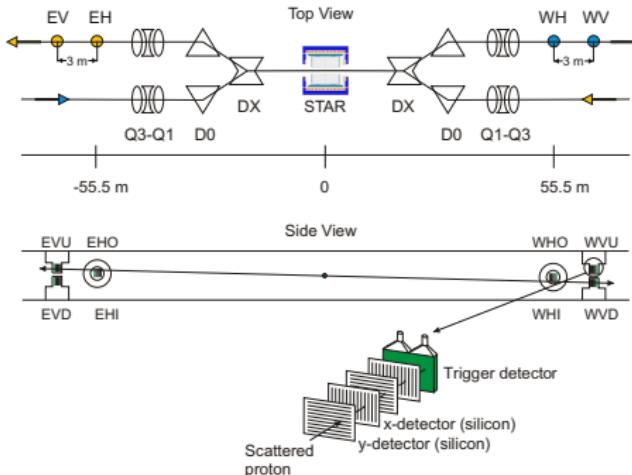
RHIC

AA: Au-Au, Cu-Cu, Cu-Au, d-Au, U-U up to $\sqrt{s_{NN}} = 200$ GeV
polarized proton-proton: up to $\sqrt{s} = 510$ GeV

this talk: $p + p \rightarrow p + p$; at $\sqrt{s} = 200$ GeV



Forward proton tagging at STAR



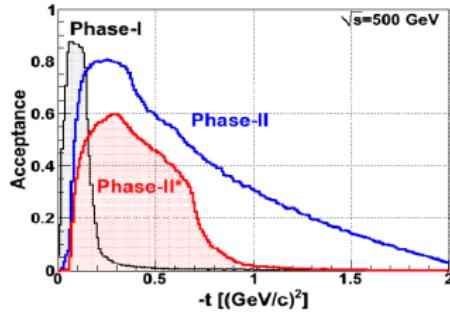
current setup(I)

- 8 Roman Pots (4 stations)
- with 4 silicon strip layers
- full angular coverage
- need special beam optics
- data at $\sqrt{s} = 200 \text{ GeV}$
- $0.003 < -t < 0.035 \text{ GeV}^2$

Upgrade to be installed 2014 (II*)

- 4 Roman Pots (2 stations)
- possible extension (II)
- 6 Roman Pots (3 stations)
- limited angular coverage
- no special runs needed
- larger statistics
- $0.1 < -t < 1.2 \text{ GeV}^2$

Acceptance for $(pp \rightarrow ppX)$



Proton-Proton Elastic Scattering

Five helicity amplitudes describe proton-proton elastic scattering

- **non-flip**

$$\phi_1(s, t) \propto \langle ++ |M| ++ \rangle$$
$$\phi_3(s, t) \propto \langle +- |M| +- \rangle$$

- **double-flip**

$$\phi_2(s, t) \propto \langle ++ |M| -- \rangle$$
$$\phi_4(s, t) \propto \langle +- |M| -+ \rangle$$

- **single-flip**

$$\phi_5(s, t) \propto \langle ++ |M| +- \rangle$$

$$\bullet \phi_i = \phi_i^{em} + \phi_i^{had}$$

$$\bullet \phi_i^{had} = \phi_i^R + \phi_i^P$$

It is useful to use:

$$\bullet \phi_+ = \frac{1}{2}(\phi_1 + \phi_3)$$

$$\bullet \phi_- = \frac{1}{2}(\phi_1 - \phi_3)$$

Some observables (cross sections and spin asymmetries)

$$\sigma_{tot}(s) = \frac{2\pi}{s} \text{Im}(\phi_+)_{|t=0}$$

$$\frac{d\sigma(s)}{dt} = \frac{2\pi}{s^2} \left(|\phi_1|^2 + |\phi_2|^2 + |\phi_3|^2 + |\phi_4|^2 + 4|\phi_5|^2 \right)$$

$$A_N(s, t) \frac{d\sigma}{dt} = \frac{-4\pi}{s^2} \text{Im} (\phi_5^*(\phi_1 + \phi_2 + \phi_3 - \phi_4))$$

$$A_{NN}(s, t) \frac{d\sigma}{dt} = \frac{4\pi}{s^2} \left(2|\phi_5|^2 + \text{Re}(\phi_1^*\phi_2 - \phi_3^*\phi_4) \right)$$

$$A_{SS}(s, t) \frac{d\sigma}{dt} = \frac{4\pi}{s^2} \text{Re} (\phi_1\phi_2^* + \phi_3\phi_4^*)$$

STAR measurements

- published A_N
- preliminary A_{NN} and A_{SS}
- ongoing $d\sigma/dt$ and σ_{tot}

Spin assymmetries

For transversely polarized beams cross section can be expressed as:

$$\sigma = \sigma_0 \left(1 + A_N (\vec{P}_Y + \vec{P}_B) \cdot \vec{n} + A_{NN} (\vec{P}_Y \cdot \vec{n})(\vec{P}_B \cdot \vec{n}) + A_{SS} (\vec{P}_Y \cdot \vec{s})(\vec{P}_B \cdot \vec{s}) \right)$$

- \vec{n} - vector normal to scattering plane
- \vec{s} - vector in scattering plane normal to initial momentum
- \vec{P}_Y, \vec{P}_B - beam polarization vectors

For $\uparrow\uparrow$ combination of beam polarization:

$$2\pi \frac{d^2\sigma^{\uparrow\uparrow}}{dt d\phi} = \frac{d\sigma}{dt} \left(1 + A_N (P_Y^\uparrow + P_B^\uparrow) \cos(\phi) + P_Y^\uparrow P_B^\uparrow [A_{NN} \cos^2(\phi) + A_{SS} \sin^2(\phi)] \right)$$

For $|P_Y^\downarrow| = |P_Y^\uparrow| = P_Y$ and $|P_B^\downarrow| = |P_B^\uparrow| = P_B$ some combinations of measurements allow extraction of spin assymmetries from uncorrected for inefficiencies event numbers (N)

$$\frac{\sigma^{\uparrow\uparrow} - \sigma^{\downarrow\downarrow}}{\sigma^{\uparrow\uparrow} + \sigma^{\downarrow\downarrow}} = \frac{A_N (P_Y + P_B) \cos(\phi)}{1 + \delta(\phi)} = \frac{N^{\uparrow\uparrow}/L^{\uparrow\uparrow} - N^{\downarrow\downarrow}/L^{\downarrow\downarrow}}{N^{\uparrow\uparrow}/L^{\uparrow\uparrow} + N^{\downarrow\downarrow}/L^{\downarrow\downarrow}}$$

$$\frac{\sigma^{\uparrow\downarrow} - \sigma^{\downarrow\uparrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\downarrow\uparrow}} = \frac{A_N (P_B - P_Y) \cos(\phi)}{1 - \delta(\phi)} = \frac{N^{\uparrow\downarrow}/L^{\uparrow\downarrow} - N^{\downarrow\uparrow}/L^{\downarrow\uparrow}}{N^{\uparrow\downarrow}/L^{\uparrow\downarrow} + N^{\downarrow\uparrow}/L^{\downarrow\uparrow}}$$

$$\frac{(\sigma^{\uparrow\uparrow} + \sigma^{\downarrow\downarrow}) - (\sigma^{\uparrow\downarrow} + \sigma^{\downarrow\uparrow})}{(\sigma^{\uparrow\uparrow} + \sigma^{\downarrow\downarrow}) + (\sigma^{\uparrow\downarrow} + \sigma^{\downarrow\uparrow})} = \overbrace{P_Y P_B [A_{NN} \cos^2(\phi) + A_{SS} \sin^2(\phi)]}^{\delta(\phi)} = \frac{(N^{\uparrow\uparrow}/L^{\uparrow\uparrow} + N^{\downarrow\downarrow}/L^{\downarrow\downarrow}) - (N^{\uparrow\downarrow}/L^{\uparrow\downarrow} + N^{\downarrow\uparrow}/L^{\downarrow\uparrow})}{(N^{\uparrow\uparrow}/L^{\uparrow\uparrow} + N^{\downarrow\downarrow}/L^{\downarrow\downarrow}) + (N^{\uparrow\downarrow}/L^{\uparrow\downarrow} + N^{\downarrow\uparrow}/L^{\downarrow\uparrow})}$$

Square root formula for single spin asymmetry

Symmetry properties: $\sigma^{\uparrow\uparrow}(\phi) = \sigma^{\downarrow\downarrow}(\pi - \phi)$ and $\sigma^{\uparrow\downarrow}(\phi) = \sigma^{\downarrow\uparrow}(\pi - \phi)$ allow extraction of A_N using non-normalized event numbers only

$$\frac{\sigma^{\uparrow\uparrow}(\phi) - \sigma^{\downarrow\downarrow}(\phi)}{\sigma^{\uparrow\uparrow}(\phi) + \sigma^{\downarrow\downarrow}(\phi)} = \frac{\sqrt{\sigma^{\uparrow\uparrow}(\phi)\sigma^{\downarrow\downarrow}(\pi - \phi)} - \sqrt{\sigma^{\downarrow\downarrow}(\phi)\sigma^{\uparrow\uparrow}(\pi - \phi)}}{\sqrt{\sigma^{\uparrow\uparrow}(\phi)\sigma^{\downarrow\downarrow}(\pi - \phi)} + \sqrt{\sigma^{\downarrow\downarrow}(\phi)\sigma^{\uparrow\uparrow}(\pi - \phi)}}$$

$$\frac{A_N(P_Y + P_B) \cos(\phi)}{1 + \delta(\phi)} = \frac{\sqrt{N^{\uparrow\uparrow}(\phi)N^{\downarrow\downarrow}(\pi - \phi)} - \sqrt{N^{\downarrow\downarrow}(\phi)N^{\uparrow\uparrow}(\pi - \phi)}}{\sqrt{N^{\uparrow\uparrow}(\phi)N^{\downarrow\downarrow}(\pi - \phi)} + \sqrt{N^{\downarrow\downarrow}(\phi)N^{\uparrow\uparrow}(\pi - \phi)}} = \epsilon_N(\phi)$$

Similary

$$\frac{A_N(P_B - P_Y) \cos(\phi)}{1 - \delta(\phi)} = \frac{\sqrt{N^{\uparrow\downarrow}(\phi)N^{\downarrow\uparrow}(\pi - \phi)} - \sqrt{N^{\downarrow\uparrow}(\phi)N^{\uparrow\downarrow}(\pi - \phi)}}{\sqrt{N^{\uparrow\downarrow}(\phi)N^{\downarrow\uparrow}(\pi - \phi)} + \sqrt{N^{\downarrow\uparrow}(\phi)N^{\uparrow\downarrow}(\pi - \phi)}} = \epsilon'_N(\phi)$$

In our measurement

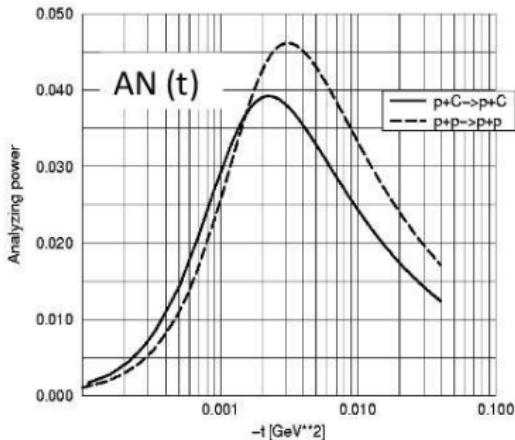
- $\delta(\phi) \ll 1$
- $P_B = 0.604 \pm 0.026 \quad P_Y = 0.618 \pm 0.028$
- $\epsilon_N \approx A_N(P_Y + P_B) \cos(\phi)$
- $\epsilon'_N \approx 0$ (consistency check)

A_N and hadronic single-flip contribution

- Double-flip spin contributions are small

$$\Rightarrow A_N \frac{d\sigma}{dt} \approx \frac{-8\pi}{s^2} \text{Im} \left(\phi_5^{em*} \phi_+^{had} + \phi_5^{had*} \phi_+^{em} \right)$$

- ϕ_+^{had} constrained by $\sigma_{tot}, \rho, \dots$
- em amplitudes fully calculable in QED.
- presize prediction of A_N for $\phi_5^{had} = 0$



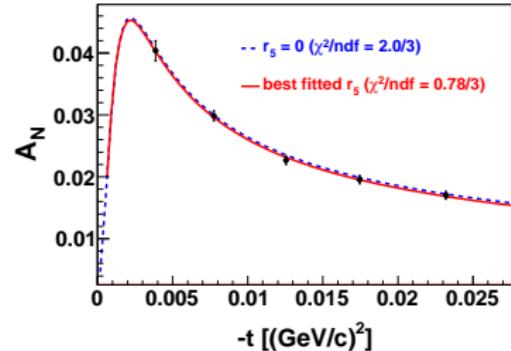
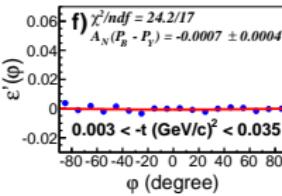
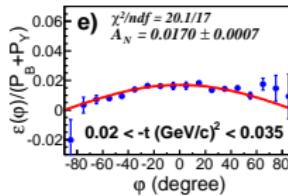
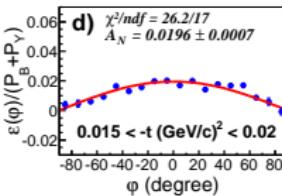
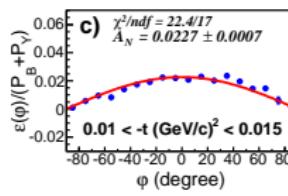
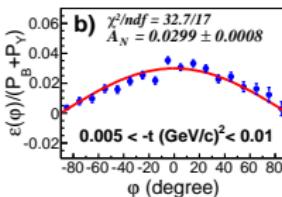
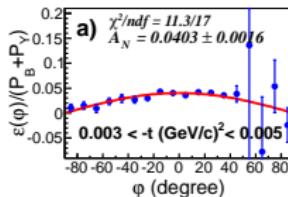
- Any difference from the above indicates contribution from hadronic spin-flip process caused by Regeon or Pomeron exchange.
- Usually hadronic spin-flip amplitude is parameterised using non-flip amplitude

$$\phi_5^{had}(s, t) = r_5(s) \frac{\sqrt{-t}}{m} \text{Im} \phi_+^{had}$$

- $\text{Im}r_5$ and $\text{Re}r_5$ can be obtained from the t -dependence of A_N
- Measurement of r_5 constrains models predicting hadronic spin-flip contribution.

Results on A_N and r_5

A_N from fits to ϵ'_N in 5 t -bins



$$\text{Re}r_5 = 0.0017 \pm 0.0017 \text{(stat.)} \pm 0.0061 \text{(syst.)}$$

$$\text{Im}r_5 = 0.007 \pm 0.030 \text{(stat.)} \pm 0.049 \text{(syst.)}$$

systematic uncertainty dominated by polarization error

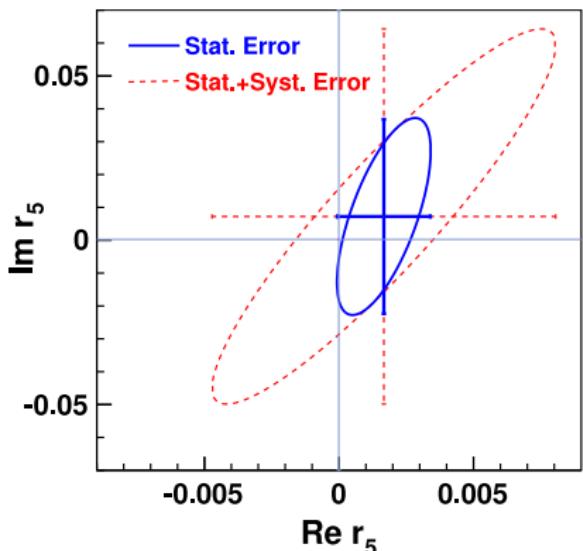
From fit to ϵ'_N

$$A_N(P_B - P_Y) = 0.0007 \pm 0.0004$$

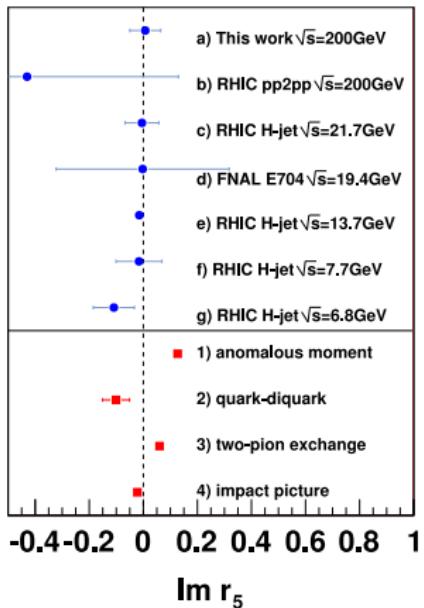
Result on r_5

$$\text{Re} r_5 = 0.0017 \pm 0.0017(\text{stat.}) \pm 0.0061(\text{syst.})$$

$$\text{Im} r_5 = 0.007 \pm 0.030(\text{stat.}) \pm 0.049(\text{syst.})$$



spin-flip contribution at high energy consistent with zero



Precise measurement of r_5 at high energy

A_{NN} , A_{SS} and double-flip contributions and Odderon

- Spin-flip contribution is small \Rightarrow

$$A_{NN}(s, t) \frac{d\sigma}{dt} \approx \frac{4\pi}{s^2} \operatorname{Re}(\phi_1^* \phi_2 - \phi_3^* \phi_4)$$

$$A_{SS}(s, t) \frac{d\sigma}{dt} = \frac{4\pi}{s^2} \operatorname{Re}(\phi_1 \phi_2^* + \phi_3 \phi_4^*)$$

- or

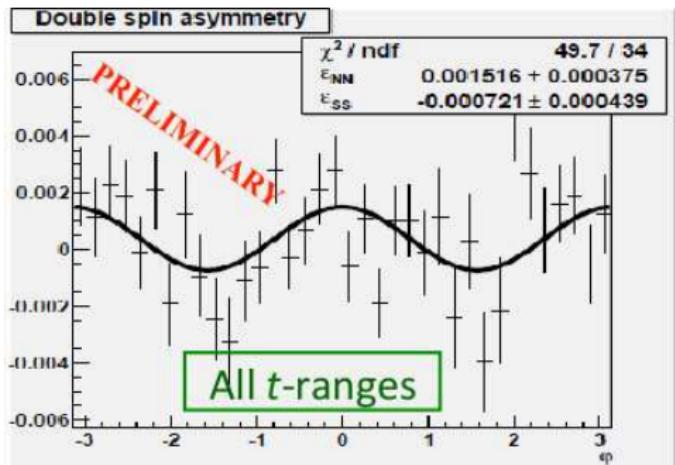
$$\frac{A_{NN} + A_{SS}}{2} \frac{d\sigma}{dt} \approx \frac{4\pi}{s^2} \operatorname{Re}(\phi_1 \phi_2^*)$$

$$\frac{A_{NN} - A_{SS}}{2} \frac{d\sigma}{dt} \approx -\frac{4\pi}{s^2} \operatorname{Re}(\phi_3 \phi_4^*)$$

- A_{NN} and A_{SS} sensitive to the Odderon contribution
for example arXiv:hep-ph/0604153v1 (2006)

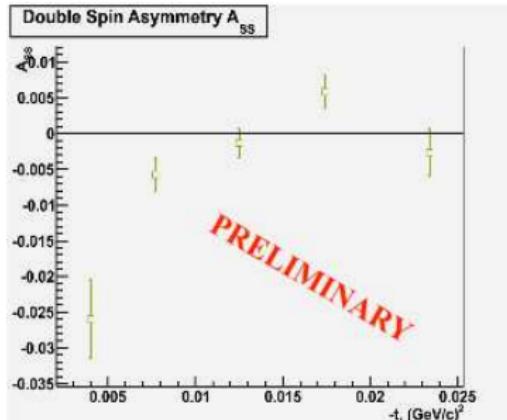
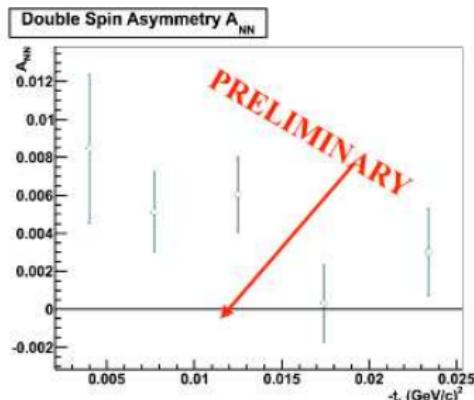
Preliminary results on A_{NN}, A_{SS}

$$\begin{aligned}\delta(\phi) &= \frac{(N^{\uparrow\uparrow}/L^{\uparrow\uparrow} + N^{\downarrow\downarrow}/L^{\downarrow\downarrow}) - (N^{\uparrow\downarrow}/L^{\uparrow\downarrow} + N^{\downarrow\uparrow}/L^{\downarrow\uparrow})}{(N^{\uparrow\uparrow}/L^{\uparrow\uparrow} + N^{\downarrow\downarrow}/L^{\downarrow\downarrow}) + (N^{\uparrow\downarrow}/L^{\uparrow\downarrow} + N^{\downarrow\uparrow}/L^{\downarrow\uparrow})} \\ &= P_Y P_B [A_{NN} \cos^2(\phi) + A_{SS} \sin^2(\phi)] \\ &= P_Y P_B \left[\frac{A_{NN} + A_{SS}}{2} + \frac{A_{NN} - A_{SS}}{2} \cos(2\phi) \right]\end{aligned}$$



$$P_B P_Y = 0.372 \pm 0.023$$

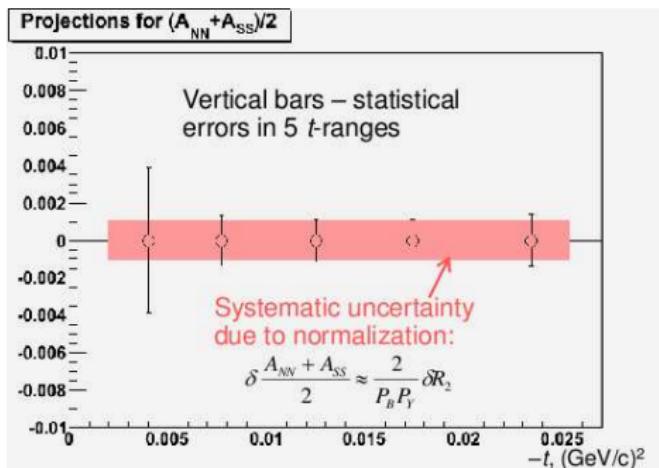
Preliminary results on A_{NN}, A_{SS}



- Both A_{NN} and A_{SS} are very small
- Need much better analysis of the normalization (uncertainties of standard method is of the order of effect)

Projection for $\frac{A_{NN}+A_{SS}}{2}$

Further normalization studies use independent ratios instead of luminosity and we have considered using STAR detectors such as beam-beam counters (BBC), vertex position detector(VPD), zero degree calorimeter(ZDC) for external normalization .



Achieved uncertainty on $\frac{A_{NN}+A_{SS}}{2}$ of the order of 8×10^{-4}

Summary and outlook

- We have collected 20 million good elastic events in polarized pp scattering at $\sqrt{s} = 200$ GeV, highest \sqrt{s} to date in $-t$ range $0.003 - 0.035$ GeV^2
- We published high precision measurements of A_N indicating spin-flip contribution at high energy consistent with zero
- Preliminary results on A_{NN} , A_{SS} have been obtained. Measurement indicates that transverse double spin asymmetries are small but non-zero. Results planned to be submitted for publication in near future.
- Program of elastic scattering measurements will continue. Planning a Phase II which allows to collect more data at $\sqrt{s} = 200(500)$ GeV