15th International Conference on Elastic \& Diffractive Scattering (15th "Blois Workshop")

FINLAND (SAARISELKÄ)
September, 09-13

Some problems of the determination of sigma_tot at LHC
O.V. Selyugin BLTPh, JINR

## Contents

Introduction

* Do we need take into account the electromagnetic interaction?
* Extraction of the real part of scattering amplitude
* Non-exponential form of the imaginary part
* Additional normalization
* The Donachi-Landshoff Pomeron
* The new method of the determination real part and normalization
* Conclusion


## Elastic scattering amplitude

## $p p \rightarrow p p$ $p \bar{p} \rightarrow p \bar{p}$

$$
\begin{aligned}
& \frac{\mathrm{d} \sigma}{\mathrm{dt}}=2 \pi\left|\Phi_{1}^{2}+\left|\Phi_{2}^{2}+\left|\Phi_{3}^{2}+\left|\Phi_{4}^{2}+4\right| \Phi_{5}^{2}\right]\right.\right. \\
& \Phi_{i}(s, t)=\Phi_{i}^{h}(s, t)+\Phi_{i}^{e}(t) e^{i \omega \varphi} \\
& \varphi(s, t)=\mp\left[\gamma+\ln (\mathrm{B}(\mathrm{~s}, \mathrm{t})|\mathrm{t}| / 2)+v_{1}+v_{2}\right]
\end{aligned}
$$

$\gamma=0,577 \ldots$.. (the Euler constant ) $\quad v_{1}$ and $v_{2}$ are small correction terms

$$
\begin{align*}
\frac{d N}{d t}= & \mathcal{L}\left[\frac{4 \pi \alpha^{2}}{|t|^{2}} G^{4}(t)-\frac{2 \alpha\left(\rho(s, t)+\phi_{C N}(s, t)\right) \sigma_{t o t} G^{2}(t) e^{-\frac{B(s, t)|t|}{2}}}{|t|}\right. \\
& \left.+\frac{\left.\sigma_{t o t}^{2}\left(1+\rho(s, t)^{2}\right) e^{-B(s, t)|t|}\right]}{16 \pi}\right]
\end{align*}
$$

Usual assumptions

$$
\operatorname{Im} F_{N}(s, t)=\sigma_{t o t} /(0.389 \square 4 \pi) e^{B t / 2}
$$

$\operatorname{Re} F_{N}(s, t) \square \sigma_{t o t} /(0.389 \square 4 \pi) e^{B t / 2} \quad B^{\mathrm{Re}}(s, t)=B^{\mathrm{Im}}(s, t)$

$$
\sigma_{t o t}\left(s_{j}\right)=f\left(s_{j}\right)_{\exp .}
$$

$\sigma_{\text {tot }}\left(s_{j}\right)-$ fix. from other experiment UA4/2

> TOTEM
or $\rho\left(s_{j}\right)-$ fix. from other experiments or - fix. from theory

Corrections: Form-factors; Coulomb-hadron phase
Problem: contribution of the spin-flip amplitude

$$
F_{1}^{e m}(t)=\alpha f_{1}^{2}(t) \frac{s-2 m^{2}}{t} ; \quad F_{3}^{e m}(t)=F_{1}^{e m}
$$

and for spin-flip amplitudes:

$$
\begin{aligned}
& F_{2}^{e m}(t)=\alpha \frac{f_{2}^{2}(t)}{4 m^{2}} s ; \quad F_{4}^{e m}(t)=-F_{2}^{e m}(t) \\
& F_{5}^{e m}(t)=\alpha \frac{s}{2 m \sqrt{|t|}} f_{1}(t) f_{2}(t)
\end{aligned}
$$

where the form factors are:

$$
\begin{array}{r}
f_{1}(t)=\frac{4 m_{p}^{2}-(1+k) t}{4 m_{p}^{2}-t} G_{d}(t) \\
f_{2}(t)=\frac{4 m_{p}^{2} k}{4 m_{p}^{2}-t} G_{d}(t)
\end{array}
$$

The hadron spin non-flip amplitude was chosen in the form

$$
F^{h}(s, t)=(i+\rho) \frac{\sigma_{t o t}}{4 \pi 0.38938} e^{\left.B t / 2+C\left(\sqrt{4 \mu^{2}-t}-2 \mu\right)\right]}
$$

$$
0.005 \leq|t| \leq 0.31 G e V^{2} ; \quad N=86 .
$$

| i | $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $\sigma_{t o t}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| 1 | 86 | 287. | 0.14 fix | 20. | 0. fix | $98.87 \pm 0.1$ |
| 2 | 86 | 287 | 0.05 fix | 20. | 0. fix | $99.7 \pm 0.1$ |
| 3 | 86 | 287 | $0.146 \pm 0.3$ | 20. | 0. fix | $98.8 \pm 0.4$ |
| 4 | 86 | 220.5 | 0.14 fix | 21.7 | $-1.4 \pm 0.2$ | $97.9 \pm 0.2$ |
| 5 | 86 | 220. | $0.05 \pm 0.4$ | 21.8 | $-1.4 \pm 0.2$ | $98.76 \pm 4$. |

Table 1: The basic parameters of the model are determined by fitting experimental data without the electromagnetic contributions and with free $\sigma_{t o t}$.

$$
0.005 \leq|t| \leq 0.1 \mathrm{GeV}^{2} ; \quad N=47 .
$$

| i | $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $\sigma_{t o t}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| 1 | 47 | 64.96 | $0.176 \pm 0.2$ | 19.9 | 0.fix | $98.05 \pm 1.7$ |
| 2 | 47 | 64.96 | 0.15 fixed | 19.9 | 0. fix | $98.47 \pm 0.1$ |
| 3 | 47 | 64.96 | 0.14 fixed | 19.9 | 0. fix | $98.6 \pm 0.1$ |
| 4 | 47 | 64.96 | 0.1 fix | 19.9 | 0. fix | $99.1 \pm 0.1$ |
| 5 | 47 | 64.96 | 0.05 fix | 19.9 | 0. fix | $99.44 \pm 0.1$ |
| 6 | 47 | 64.96 | 0.0 fix | 19.9 | 0.fix | $99.57 \pm 0.1$ |
| 7 | 47 | 64.96 | -0.05 fix | 19.9 | 0. fix | $99.44 \pm 0.1$ |
| 8 | 47 | 61.09 | 0.14 fix | 18.5 | $1.05 \pm 0.54$ | $98.99 \pm 0.2$ |
| 9 | 47 | 61.09 | 0.1 fix | 18.5 | $1.06 \pm 0.54$ | $99.47 \pm 0.2$ |
| 10 | 47 | 61.09 | 0.0 fix | 18.5 | $1.07 \pm 0.54$ | $99.97 \pm 0.2$ |
| 11 | 47 | 60.08 | $-0.03 \pm 0.1$ | 18.4 | $1.07 \pm 0.54$ | $99.94 \pm 0.4$ |

Table 2: The basic parameters of the model are determined by fitting experimental data without the electromagnetic contributions and with free $\sigma_{t o t}$.

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $\sigma_{t o t}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| 86 | 281. | 0.14fixed | 20. | 0.fix | $99.4 \pm 0.1$ |
| 86 | 281. | 0.1 fix | 20. | 0.fix | $99.7 \pm 0.1$ |
| 86 | 288. | 0.fix | 20. | 0.fix | $99.8 \pm 0.1$ |
| 86 | 245. | 0.14 fix | 21.3 | $-1.03 \pm 0.2$ | $98.6 \pm 0.2$ |
| 86 | 215. | 0.0 fix | 21.8 | $-1.2 \pm 0.2$ | $98.7 \pm 0.2$ |
| 86 | 175. | $-0.41 \pm 0.1$ | 23.2 | $-2.77 \pm 0.2$ | $89.1 \pm 3$. |

Table 3: The basic parameters of the model are determined by fitting experimental data with free $\sigma_{t o t}$.

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $\sigma_{\text {tot }, ~}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| 47 | 87.1 | 0.2fixed | 20.1 | 0.fix | $98.5 \pm 0.1$ |
| 47 | 77.1 | 0.14fixed | 20. | 0.fix | $99.2 \pm 0.1$ |
| 47 | 71.6 | 0.1 fix | 20 | 0.fix | $99.5 \pm 0.1$ |
| 47 | 61.1 | $-0.069 \pm 0.05$ | 19.8 | 0. fix | $98.93 \pm 0.8$ |
| 47 | 61.2 | 0.1 fix | 17.7 | $1.66 \pm 0.54$ | $100.1 \pm 0.2$ |
| 47 | 60.6 | 0.0 fix | 18.8 | $0.82 \pm 0.54$ | $99.8 \pm 0.2$ |
| 47 | 60.6 | $0.01 \pm 0.1$ | 18.9 | $0.74 \pm 0.8$ | $99.7 \pm 0.8$ |

Table 4: The basic parameters of the model are determined by fitting experimental data with free $\sigma_{t o t}$.

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $\sigma_{\text {tot }}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| 40 | 78.8 | 0.2 fix | 20.1 | 0.fix | $98.6 \pm 0.12$ |
| 40 | 70.4 | 0.14 fix | 20. | 0. fix | 99.29 |
| 40 | 65.8 | 0.1 fix | 20 | 0. fix | $99.55 \pm 0.12$ |
| 40 | 56.6 | $-0.076 \pm 0.06$ | 19.8 | 0. fix | $98.83 \pm 0.12$ |
| 40 | 54.7 | 0.1 fix | 16.3 | $2.63 \pm 0.8$ | $100.3 \pm 0.27$ |
| 40 | 54.9 | 0.0 fix | 17.8 | $1.47 \pm 0.8$ | $99.96 \pm 0.26$ |
| 40 | 54.6 | $0.06 \pm 0.01$ | 16.9 | $2.17 \pm 1.4$ | $100.3 \pm 0.3$ |

Table 5: The basic parameters of the model are determined by fitting experimental data with free $\sigma_{t o t}$.


Figure 1: Size of $\sigma_{\text {tot }}$ a)(left) over $\rho$ and b) (right) over $C$ (hard line - without electromagnetic interaction and dashed line with electromagnetic interaction).

## O.S. - J. Nucl.Phys. (Yad.Phys.) v. 55 (1992)

$$
\begin{align*}
\operatorname{Re} F^{h}(t) & =-\operatorname{Re} F_{c}(t)  \tag{9}\\
& +\left[\left[\left.\frac{d \sigma}{d t}\right|_{\text {exp. }}-k \pi *\left(\operatorname{Im} F_{c}+\operatorname{Im} F_{h}\right)^{2}\right] /(k \pi)\right]^{1 / 2}
\end{align*}
$$

let us take the imaginary part of the hadron scattering amplitude in the simple exponential form with the parameters obtained by the TOTEM Collaboration

$$
\begin{equation*}
\operatorname{Im} F^{h}(t)=\sigma_{t o t} /(4 k \pi) e^{B t / 2}, \tag{10}
\end{equation*}
$$



Figure 2: Real part of the hadronic amplitude calculated by (eq.9)(triangles and squared without and with $F_{c}$; solid and empty represent real and imaginary parts of (eq.9) see text). (long dashed line - the calculations by (eq.11)).

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $\sigma_{t o t}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| 47 | 134.5 | 0.14fixed | 19.8 | 0.- fixed | 98.4 |
| 47 | 174.7 | 0.1 fix | 19.7 | 0.fix | 98.4 |
| 47 | 88.1 | $0.203 \pm 0.01$ | 20.1 | 0.fixed | 98.4 |
| 47 | 105.3 | 0.14 fixed | 22.9 | $-2.3 \pm 0.3$ fixed | 98.4 |
| 47 | 61.4 | $-0.105 \pm 0.02$ | 20. | $-0.14 \pm 0.4$ | 98.4 |

Table 6: The basic parameters of the model are determined by fitting experimental data with fixed $\sigma_{t o t}$.

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $\sigma_{\text {tot }}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| 40 | 88.17 | $0.203 \pm 0.007$ | 20.1 | $0 .-$ fixed | 98.4 |
| 40 | 88.2 | 0.2 fixed | 20. | 0. - fixed | 98.4 |
| 40 | 125.6 | 0.14fixed | 19.7 | $0 .-$ fixed | 98.4 |
| 40 | 157.7 | 0.1 fix | 19.6 | 0. fix | 98.4 |
| 40 | 102.4 | 0.14 fix | 22.4 | $-1.75 \pm 0.4$ | 98.4 |
| 40 | 56.9 | $-0.106 \pm 0.01$ | 19.8 | 0. fix | 98.4 |
| 40 | 56.8 | $-0.11 \pm 0.02$ | 19.7 | $0.1 \pm 0.7$ | 98.4 |

Table 7: The basic parameters of the model are determined by fitting experimental data with fixed $\sigma_{t o t}$.

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $n$ | $\sigma_{\text {tot }, m b}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| 47 | 77.1 | 0.14 fixed | 20. | 0. - fixed | $1.017 \pm 0.002$ | 98.4 |
| 47 | 71.6 | 0.1 fix | 20. | 0. fix | $1.022 \pm 0.002$ | 98.4 |
| 47 | 63.1 | 0.14 fix | 17.2 | $2.1 \pm 0.5$ | $1.034 \pm 0.012$ | 98.4 |
| 47 | 61.1 | $-0.071 \pm 0.05$ | 19.8 | 0. fix | $1.01 \pm 0.013$ | 98.4 |
| 47 | 60.6 | $-0.01 \pm 0.09$ | 18.9 | $0.72 \pm 0.9$ | $1.02 \pm 0.02$ | 98.4 |

Table 8: The basic parameters of the model are determined by fitting experimental data with fixed $\sigma_{t o t}$ and with additional normalization coefficient $n$.

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $n$ | $\sigma_{t o t}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| 40 | 70.4 | 0.14fixed | 20. | 0.-fixed | $1.018 \pm 0.0025$ | 98.4 |
| 40 | 65.8 | 0.1fix | 20. | 0.fix | $1.024 \pm 0.0025$ | 98.4 |
| 40 | 55.1 | 0.14 fix | 15.7 | $3.1 \pm 0.77$ | $1.037 \pm 0.005$ | 98.4 |
| 40 | 56.6 | $-0.077 \pm 0.06$ | 19.8 | 0. fix | $1.009 \pm 0.015$ | 98.4 |
| 40 | 54.6 | $0.06 \pm 0.08$ | 16.9 | $2.18 \pm 0.4$ | $1.044 \pm 0.02$ | 98.4 |

Table 9: The basic parameters of the model are determined by fitting experimental data with fixed $\sigma_{t o t}$ and with additional normalization coefficient $n$.

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $n$ | $\sigma_{\text {tot }}, m b$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| 47 | 77.84 | 0.14fixed | 20.0 | $0 .-$ fix | 1.05 | $96.8 \pm 0.1$ |
| 47 | 71.65 | 0.1 ffix | 20. | 0. fix | 1.05 | $97.1 \pm 0.1$ |
| 47 | 66.3 | 0.05 fix | 20 | 0. fix | 1.05 | $97.2 \pm 0.1$ |
| 47 | 62.8 | 0. fix | 19.4 | 0. fix | 1.05 | $97.1 \pm 0.1$ |
| 47 | 63.1 | 0.14 fixed | 17.2 | $2.1 \pm 0.5$ | 1.05 | $97.56 \pm 0.2$ |
| 47 | 61.9 | 0.1 fix | 17.7 | $1.87 \pm 0.5$ | 1.05 | $97.7 \pm 0.2$ |
| 47 | 61.0 | 0.05 fix | 18.2 | $1.24 \pm 0.5$ | 1.05 | $97.7 \pm 0.2$ |
| 47 | 60.6 | 0. fix | 18.8 | $0.8 \pm 0.5$ | 1.05 | $97.4 \pm 0.2$ |
| 47 | 60.8 | -0.05 fix | 19.3 | $0.4 \pm 0.5$ | 1.05 | $96.9 \pm 0.3$ |
| 47 | 61.1 | $-0.064 \pm 0.05$ | 19.8 | $0 . f i x$ | 1.05 | $96.57 \pm 0.58$ |
| 47 | 60.6 | $-0.011 \pm 0.09$ | 18.9 | $0.7 \pm 0.9$ | 1.05 | $97.3 \pm 0.9$ |

Table 10: The basic parameters of the model are determined by fitting experimental data.


Figure 3: Size of $\sigma_{\text {tot }}$ over $n$ in two variants a) (hard line) - with free slope $C$ and b) (dashed line) with $C=0$.

| $F_{C}$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $n$ | $\sigma_{t o t}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| - | 64.96 | 0.15 fix | 19.9 | 0. fix | 1. | $98.47 \pm 0.1$ |
| - | 61.1 | $0.002 \pm 0.2$ | 18.4 | $1.07 \pm 0.5$ | 1. | $99.9 \pm 1.4$ |
| + | 60.6 | $-0.01 \pm 0.09$ | 18.9 | $0.74 \pm 0.8$ | 1. | $99.7 \pm 0.8$ |
| + | 60.6 | $-0.01 \pm 0.09$ | 18.9 | $0.72 \pm 0.9$ | $1.02 \pm 0.004$ | 98.4 fix |
| + | 60.6 | $-0.01 \pm 0.09$ | 18.9 | $0.7 \pm 0.9$ | 1.05 fix | $97.3 \pm 0.9$ |

Table 12: The basic parameters of the model are determined by fitting experimental data.

$$
F^{h}(s, t)=i h \hat{s}^{\Delta} f_{1}(t)^{2} \frac{\sigma_{t o t}}{4 \pi 0.38938} e^{\alpha t+\alpha_{2}\left(\sqrt{4 \mu^{2}-t}-2 \mu\right) L n(\hat{s})}
$$

$h \operatorname{Re} \widehat{s}=1$
with the electromagnetic form factor $f_{1}(t)(7)$ and $\hat{s}=s e^{-i * \pi / 2} ; \mu$ is the pion mass.

|  | $\Delta=0.1$ | $\rho(\sqrt{s}=7 \mathrm{TeV}, t=0)=0.156$ |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\alpha_{1}^{\prime}$ | $\alpha_{2}^{\prime}$ | $\sigma_{\text {tot }}, m b$ |
|  |  |  |  |  |
| 47 | 65.2 | 0.325 | 0. fix | $99.7 \pm 0.15$ |
| 47 | 65.1 | 0.328 | $-0.002 \pm 0.015$ | $99.6 \pm 0.2$ |
| 40 | 58.1 | 0.324 | $0 . f i x$ | $99.6 \pm 0.05$ |
| 40 | 55.8 | 0.276 | $0.037 \pm 0.02$ | $99.9 \pm 0.26$ |
| 47 | 204 | 0.314 | $0 . f i x$ | 98.4 fix |
| 47 | 95.3 | 0.427 | $-0.075 \pm 0.008$ | 98.4 fix |
| 40 | 154 | 0.312 | $0 . f i x$ | 98.4 fix |
| 40 | 90.1 | 0.437 | $-0.08 \pm 0.01$ | 98.4 fix |

Table 13: The basic parameters of the Regge amplitude are determined by fitting experimental data with free $\sigma_{\text {tot }}$.

| $\Delta=0.08$ |  |  |  | $\rho(\sqrt{s}=7 \mathrm{TeV}, t=0)=0.128$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\alpha_{1}^{\prime}$ | $\alpha_{2}^{\prime}$ | $\sigma_{\text {tot }}, m b$ |  |  |
|  |  |  |  |  |  |  |
| 47 | 64.4 | 0.325 | 0. fix | $100.0 \pm 0.1$ |  |  |
| 47 | 64.3 | 0.338 | $0.008 \pm 0.003$ | $99.9 \pm 0.1$ |  |  |
| 40 | 56.4 | 0.323 | $0 . f$ ix | $99.9 \pm 0.12$ |  |  |
| 40 | 55.3 | 0.299 | $0.005 \pm 0.004$ | $100.2 \pm 0.3$ |  |  |
| 47 | 285 | 0.310 | 0. fix | 98.4 fix |  |  |
| 47 | 106.7 | 0.455 | $-0.096 \pm 0.008$ | 98.4 fix |  |  |
| 40 | 211 | 0.307 | 0. fix | 98.4 fix |  |  |
| 40 | 99.6 | 0.473 | $-0.108 \pm 0.011$ | 98.4 fix |  |  |

Table 14: The basic parameters of the model are determined by fitting experimental data with free $\sigma_{t o t}$.
D.S. - New methods for calculating parameters of the diffraction scattering amplitude, "VI Intern. Conf. On Diffraction...", Blois, France,(1995).
D.S. "Additional ways to determination of structure of high energy elastic scattering amplitude" arxiv.org:[hep-ph/0104295]
P. Gauron, B. Nicolescu, O.S. "A New Method for the Determination of the Real Part of the Hadron Elastic Scattering Amplitude at Small Angles and High Energies"

Phys.Lett. B629 (2005) 83-92

$$
\Delta_{R}(t)=\left[\operatorname{Re} F^{h}(t)+\operatorname{Re} F^{\tilde{N}}(t)\right]^{2}=\left[\left.\frac{d \sigma}{d t}\right|_{\text {exp. }}-k \pi\left(\operatorname{Im} F^{h}(t)+\operatorname{Im} F^{C}(t)\right] /(k \pi)\right.
$$


O.S. - Talk on X-th Blois Workshop (Xelsinki-2003)-hep-ph: hep-ph/0306256

$$
\mathrm{pp}-4 \mathrm{TeV}
$$




- TOTEM parameters
-     - $\quad=96.4 \mathrm{mb} \quad \mathrm{B}=20.3 \mathrm{GeV}-2 ; \mathrm{C}=-0.05 ; \mathrm{n}=1.08$
—— TOTEM parameters; $\square=0.141$
$---\quad$-tot=96.4 mb; B=19.9 GeV-2; $\square=0.1 \quad-----$
$\square=$
0
;
-tot $=96.4 \mathrm{mb} ; \quad B=19.9 \mathrm{GeV}-2 ; \quad \square=-0.05$;


## Summary

* The analysis of the new experimental data (TOTEM) shows there are some additional specific moments which are to be taken into account to determine the size of the total cross sections
I. we can not neglect the electromagnetic interaction;

2. the deviation of the form of the scattering amplitude at small $t$ can be taken into account by the part of the slope proportional $q$;
3. the errors of the luminocity can be taken into account by an additional normalization coefficient;
4. it is need to check out the obtained, during the fitting procedure, the real part by using the t -dependence of extracted

$$
\text { R }=(\text { ReF_h+ReF_c })^{\wedge} 2 .
$$

5. Our analysis of the TOTEM data at 7 TeV shows that the size of $\square$ 㶾 $\mathrm{t} \square$ (D) , very likely, near zero or -0.05 .

It is most likely that there is or some problems with normalization or there is some problems with the form of scattering amplitude
(for example, the oscillation term)

* The best way to decrease the impact of the different assumption consist in the determination of the sizes of and $\square$ simultaneously in one experiment.


## END

## THANKS <br> FOR YOUR ATTENTION

| $N$ | $\sum_{i=1}^{N} \chi_{i}^{2}$ | $\rho$ | $B$ | $C$ | $n$ | $\sigma_{\text {tot }}, m b$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| 47 | 77.35 | 0.14fixed | 20.03 | $0 .-$ fix | 1.08 | $95.49 \pm 0.1$ |
| 47 | 71.6 | 0.1 fix | 20. | 0. fix | 1.08 | $95.76 \pm 0.1$ |
| 47 | 62.75 | 0. fix | 19.4 | 0. fix | 1.08 | $95.74 \pm 0.1$ |
| 47 | 60.89 | 0.14 fixed | 19.0 | $0.45 \pm 0.11$ | 1.08 | $97.69 \pm 0.1$ |
| 47 | 60.56 | 0.1 fix | 19.15 | $0.37 \pm 0.11$ | 1.08 | $97.57 \pm 0.1$ |
| 47 | 60.49 | 0. fix | 19.52 | $0.17 \pm 0.11$ | 1.08 | $96.54 \pm 0.1$ |
| 47 | 60.9 | -0.05 fix | 19.7 | $0.07 \pm 0.11$ | 1.08 | $95.71 \pm 0.1$ |
| 47 | 61.14 | $-0.064 \pm 0.05$ | 19.8 | 0. fix | 1.08 | $97.18 \pm 0.2$ |
| 47 | 60.6 | $-0.007 \pm 0.09$ | 18.9 | $0.73 \pm 0.9$ | 1.08 | $96.0 \pm 0.8$ |

Table 11: The basic parameters of the model are determined by fitting experimental data.

$$
\begin{aligned}
\frac{d N}{d t}= & \mathcal{L}\left[\frac{4 \pi \alpha^{2}}{|t|^{2}} G^{4}(t)-\frac{2 \alpha\left(\rho(s, t)+\phi_{C N}(s, t)\right) \sigma_{t o t} G^{2}(t) e^{-\frac{B(s, t)|t|}{2}}}{|t|}\right. \\
& \left.+\frac{\left.\sigma_{t o t}^{2}\left(1+\rho(s, t)^{2}\right) e^{-B(s, t)|t|}\right]}{16 \pi}\right]
\end{aligned}
$$

$I m F^{h}(t)=\sigma_{t o t} /(4 k \pi) e^{B t / 2}$,
$\operatorname{Re} F^{h}(t)=\rho \sigma_{t o t} /(4 k \pi) e^{B t / 2}$,

## Standard definitions

$$
\begin{aligned}
& G_{E p}(0)=1 ; G_{E n}(0)=0 ; \quad G_{M}(0)=\left(G_{E}(0)+k\right)=\mu ; \\
& \mu_{p}=(1+1.79) \frac{e}{2 M} ; \quad k_{p}=1.79 ; \\
& F_{1}^{D}(t)=\frac{4 M_{p}^{2}-t \mu_{p}}{4 M_{p}^{2}-t} G_{D}(t) ; \quad F_{2}^{P}(t)=\frac{1}{1-t / 4 M_{p}^{2}} G_{D}(t) ; \\
& \quad G_{D}(t)=\frac{\Lambda^{2}}{(\Lambda-t)^{2}} ; \quad \Lambda=0.71 G e V^{2} ;
\end{aligned}
$$

