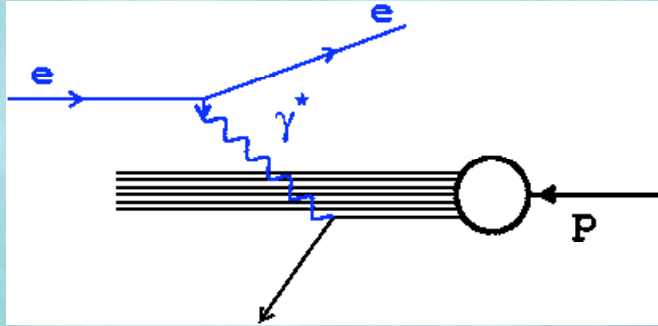


# Low-x Physics Results from HERA



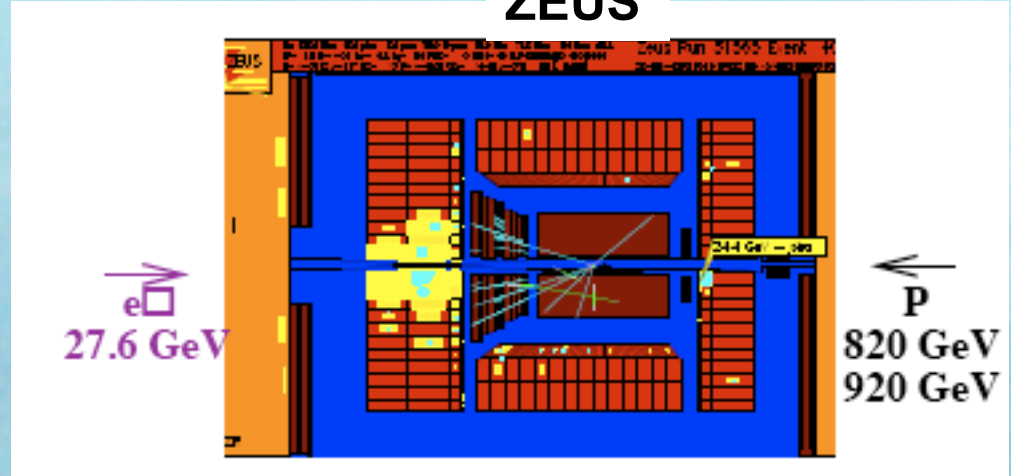
**Henri Kowalski**  
on behalf of H1 and ZEUS Collaborations  
EDS Blois 2013, Saariselka, Finland  
9th of September, 2013

# Inclusive Scattering



# DIS at HERA at low x

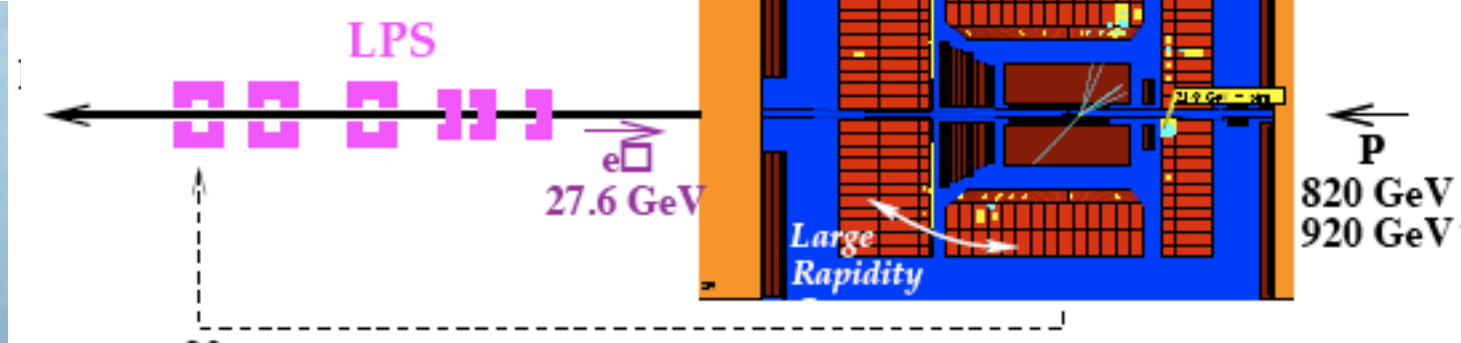
## ZEUS



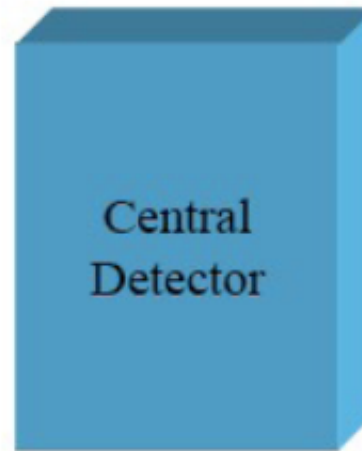
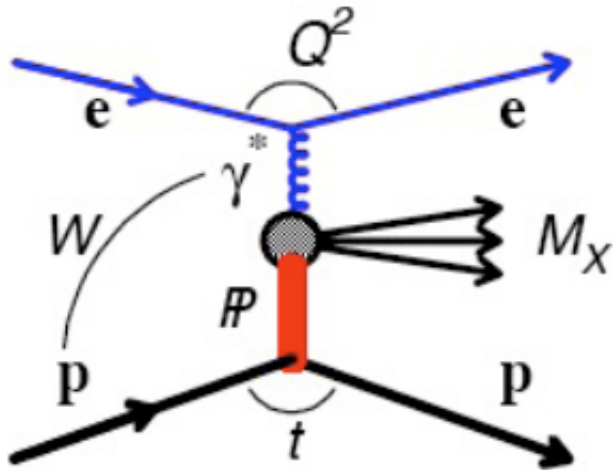
# Diffractive Scattering

expectation before HERA  
~ 0.01%

seen ~20% at  $Q^2 = 4 \text{ GeV}^2$   
~10% at  $Q^2 = 20 \text{ GeV}^2$



# DIS Reactions



**Rapidity Gaps**  
 $\Delta Y = \ln(W^2/M_X^2) \approx \Delta\eta$

**Forward protons**  
 with  $x_L = 1 - x_{IP} > 95\%$   
 $x_L \sim$  longitudinal fraction of proton momentum

## Inclusive variables:

$Q^2$  - virtuality of the incoming photon

$W$  - CMS energy of the incoming photon-proton system

$x$  -  $\approx Q^2 / W^2$

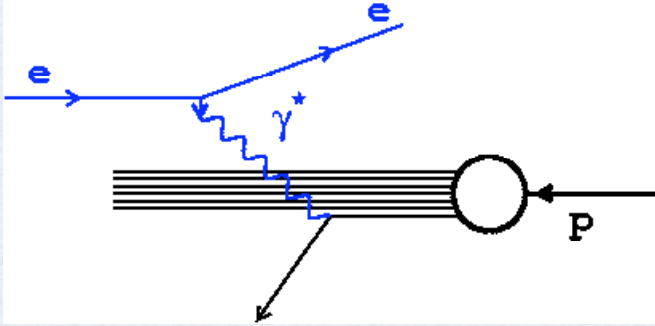
## Diffractive variables:

$M_X$  - invariant mass of all particles seen in the detector

$t$  - momentum transfer to the diffractively scattered proton

# Partons vs Dipoles at low- $x$

Infinite momentum frame: Partons



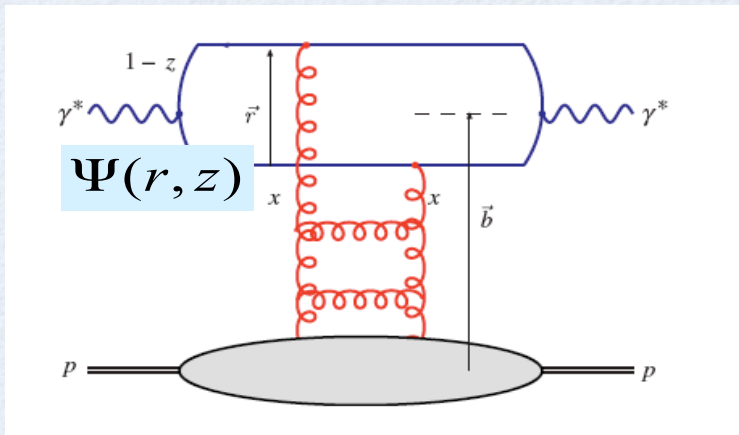
$F_2$  measures parton density at a scale  $Q^2$

$$F_2 = \sum_f e_f^2 xq(x, Q^2)$$

Proton rest frame: Dipoles - long living quark pair interacts with the gluons of the proton

*dipole life time  $\approx 1/(m_p x)$*

*$= 10 - 1000 \text{ fm at } x = 10^{-2} - 10^{-4}$*



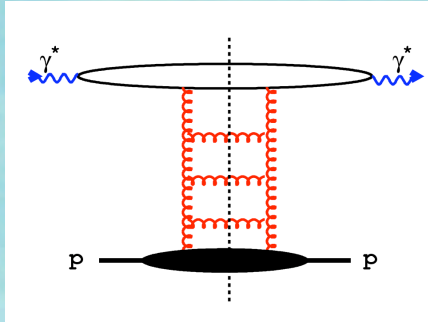
$$\sigma_{tot}^{\gamma^* p} = \int \Psi^* \sigma_{qq} \Psi ; \quad F_2 = \frac{Q^2}{4\pi^2 \alpha_{em}} \sigma_{tot}^{\gamma^* p}$$

for small dipoles, at low- $x$ , dipole picture is equivalent to the QCD parton picture

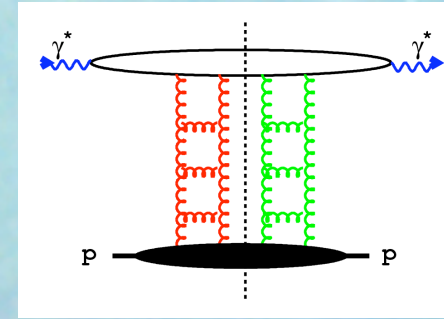
$$\sigma_{qq} \sim r^2 xg(x, Q^2)$$

# Low-x phenomena in DIS give access to the properties of the gluon density

- rise of F2 with decreasing x (this talk)
- diffractive reactions ( A. Valkarova talk)



← Optical Theorem →



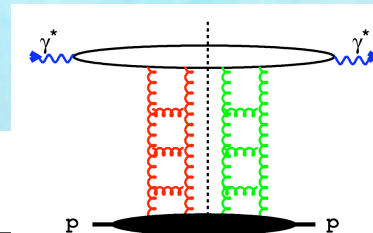
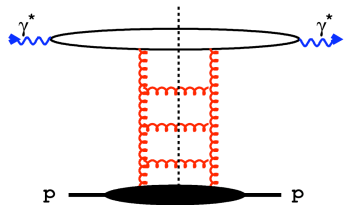
$$\sigma_{tot}^{\gamma^* p} = \int \Psi^* \sigma_{qq} \Psi$$

$$\frac{d\sigma^{\gamma^* p \rightarrow V p}}{dt} \sim \left| \int d^2b \Psi_V^* \Psi e^{-i\vec{b} \cdot \vec{\Delta}} \frac{d\sigma_{q\bar{q}}}{d^2b} \right|^2$$

$$\frac{d\sigma_{q\bar{q}}}{d^2b} \sim r^2 \alpha_s x g(x, \mu^2) T(b) \quad \text{for small dipole size}$$

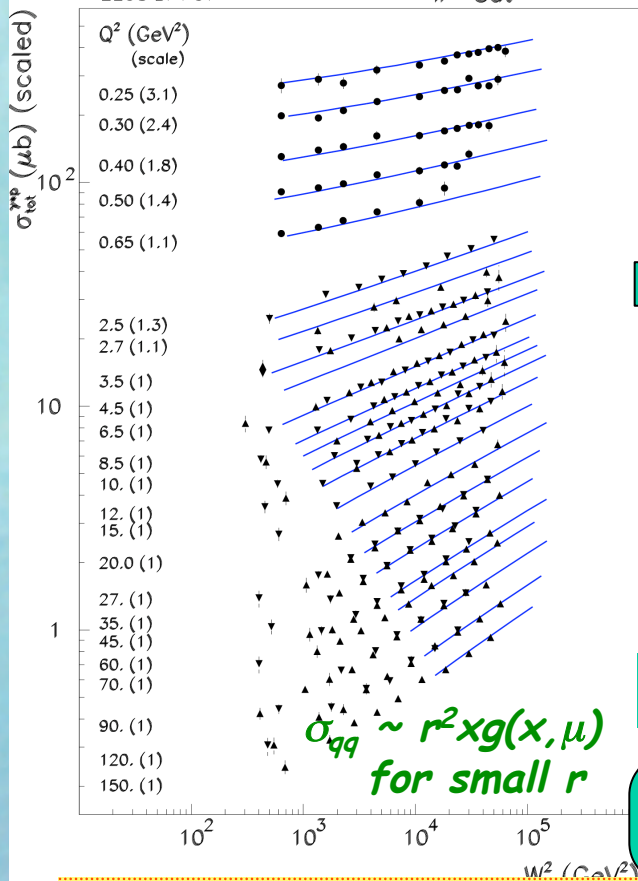
The same, universal, gluon density describes the properties of many reactions:  $F_2$ ,  $F_L$ , inclusive diffraction, exclusive J/Psi, Phi and Rho production, DVCS, diffractive jets

# Diffraction as a shadow of DIS



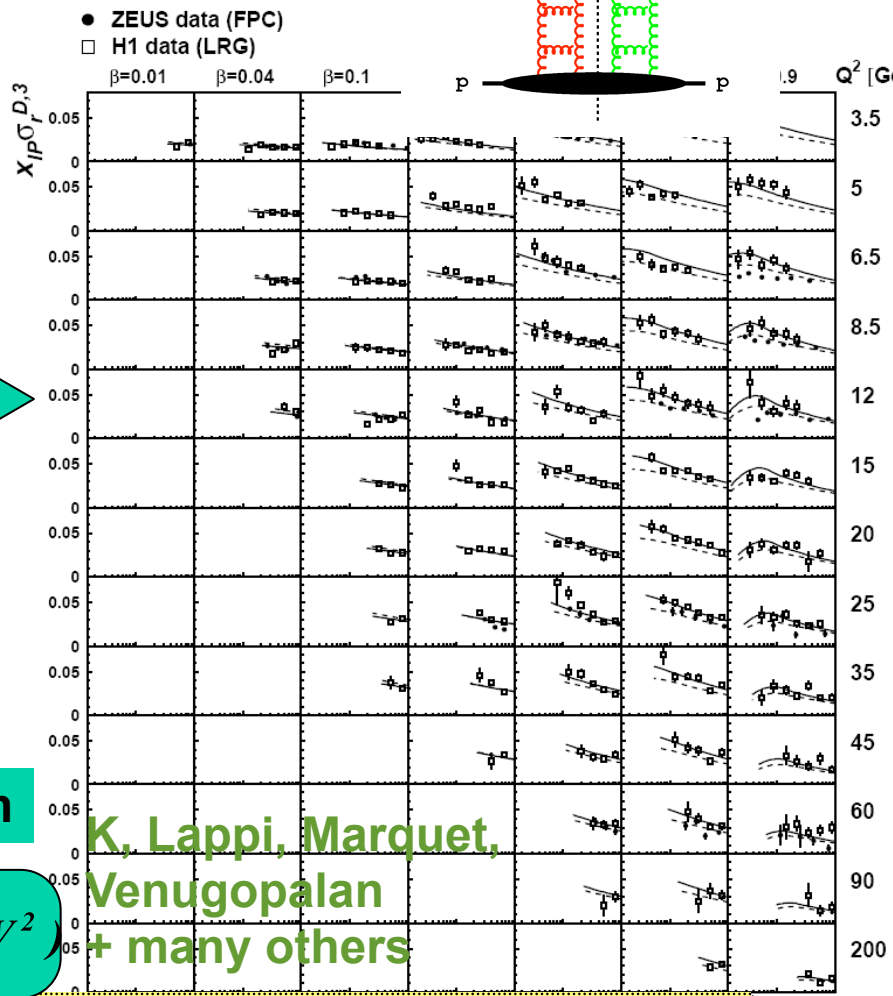
- ▼ H1 96-ε
- ▲ ZEUS 96
- ZEUS BPT 97

— IP-Sat



Optical Theorem

$$\sigma_{tot}^{\gamma^* p} = \frac{1}{W^2} \text{Im} A_{el}(W^2)$$



K, Lappi, Marquet,  
Venugopalan  
+ many others

$$\sigma_{tot}^{\gamma^* p} = \int d^2 \vec{r} \int_0^1 dz \Psi^* \sigma_{q\bar{q}}(x, r^2) \Psi$$

$$\frac{d\sigma_{diff}^{\gamma^* p}}{dt} \Big|_{t=0} = \frac{1}{16\pi} \int d^2 \vec{r} \int_0^1 dz \Psi^* \sigma_{q\bar{q}}^2(x, r^2) \Psi$$

# Determination of Gluon Density in pdf's

GD is determined from the increase of  $F_2$  with  $x$  and  $Q^2$  in low- $x$  region

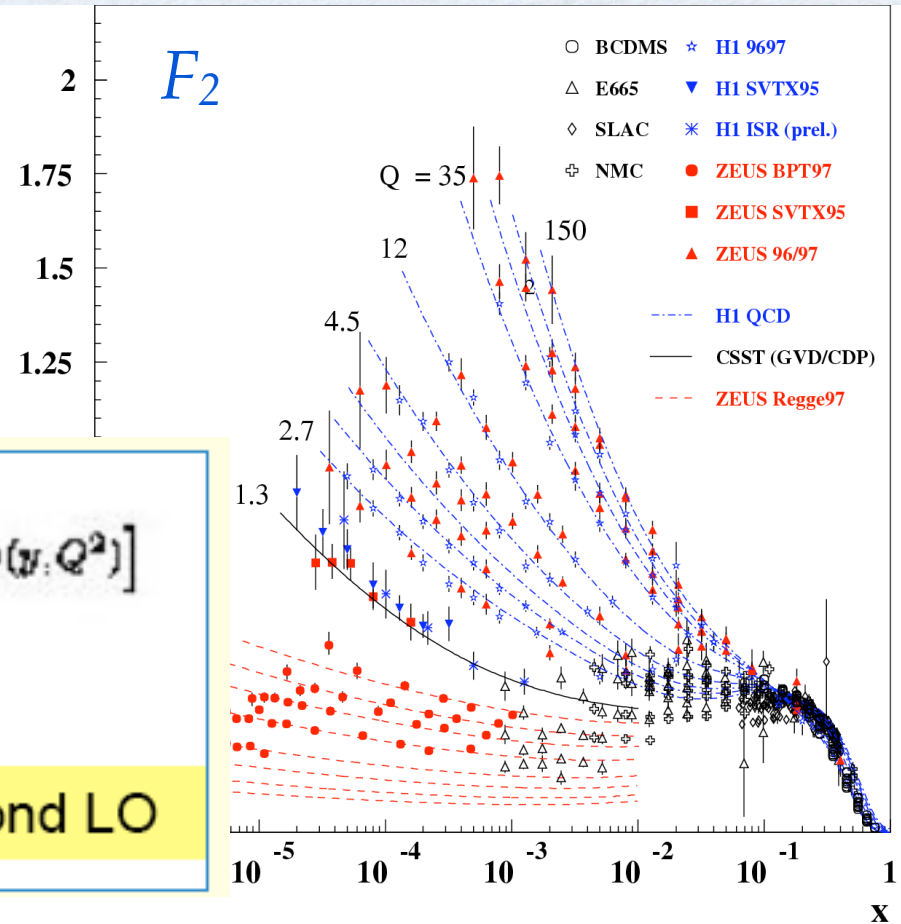
Determine pdf's densities from the  $\chi^2$  fit to the data

$$\frac{F_2(x, Q^2)}{x} = \int_0^1 \frac{dy}{y} \left[ \sum_i C_2(z, \alpha_s) q_i(x, Q^2) + C_g(z, \alpha_s) g(y, Q^2) \right]$$

$$C_2(z, \alpha_s) = e_s^2 [\delta(1-z) + \alpha_s f_2(z)]$$

$$C_g(z, \alpha_s) = \alpha_s f_g(z)$$

beyond LO



Who? ABM, MSTW, CT(EQ), HERAPDF, (G)JR, NNPDF

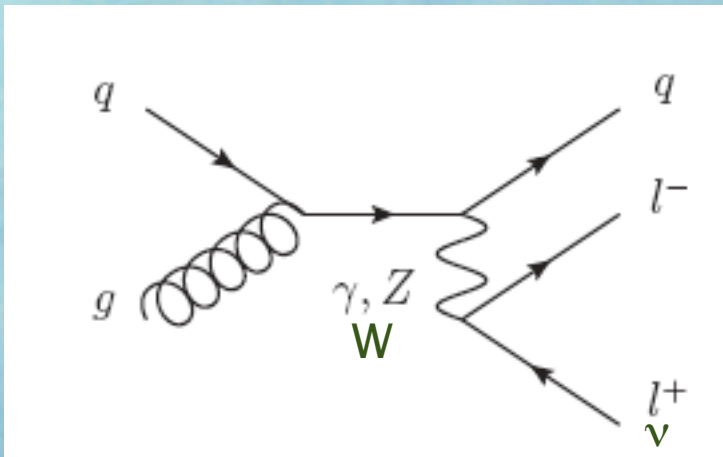
How? Start from parametrized form of  $g(q)(x, Q_0^2)$  at  $Q_0^2$  1-7  $GeV^2$   
use N(N)LO DGLAP, MSbar factorisation, Heavy quark scheme

# DATA

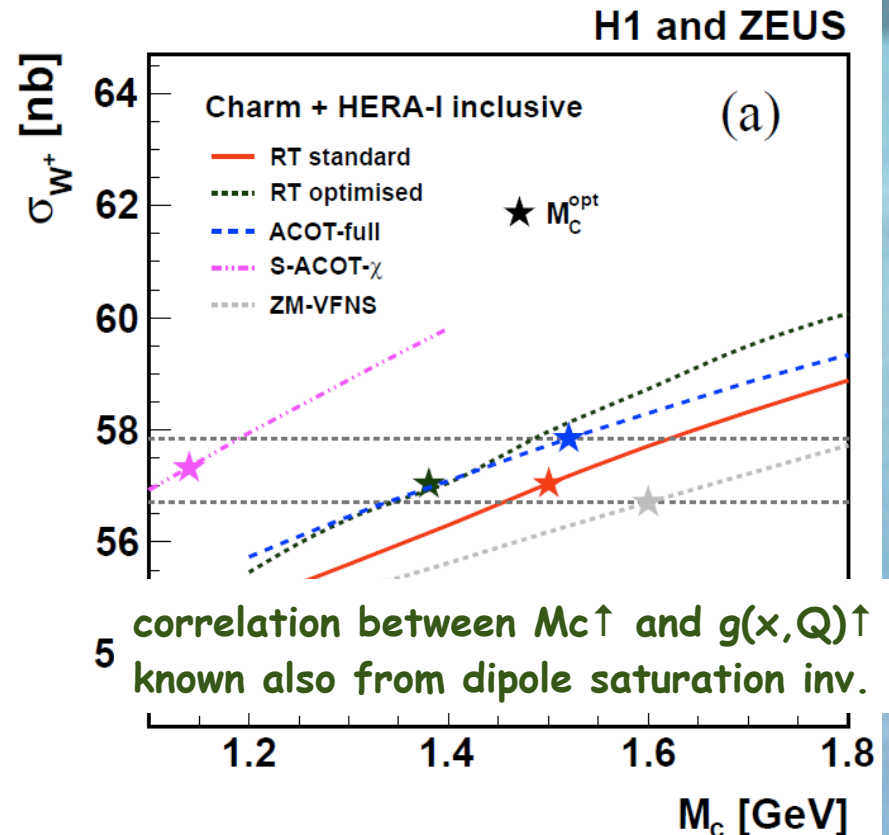
DIS (HERA and fixed target..), Drell-Yan processes (fixed target)  
 High  $E_T$  jets (Tevatron), W,Z rapidity (Tevatron)  
 $\nu N$  dimuon (CCFR, NuteV)....

HERAPDF 1.0 uses combined H1 and ZEUS HERA I xs data  
 HERAPDF 1.5 uses, in addition, combined HERA charm data ..

Excellent test reaction  
 W, Z production at LHC



W, Z production at LHC is  
 a low-x effect





## $\chi^2$ function

→ nuisance parameters:  $\chi^2 = \sum_i \frac{(D_i - T_i^*)^2}{(\delta_i^{unc})^2}$

$$T_i^* = T_i + \sum_j \xi_j \delta_i^{cor,j}$$

↑ Nuisance parameter
 ← Correlated error

→ covariance matrix:  $\chi^2 = \sum_{i,j} (D_i - T_i) Cov_{i,j}^{-1} (D_j - T_j)$

→ mixed

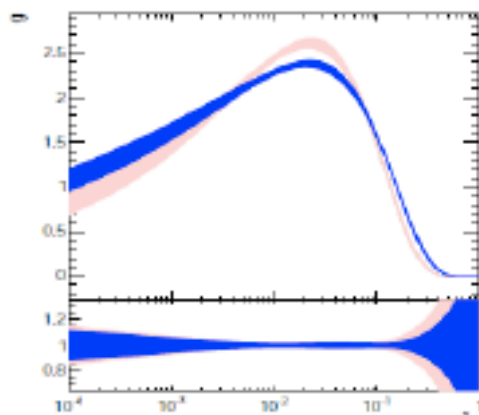
Various types of uncertainty treatment for experimental data:

→ Hessian, Monte Carlo, Offset

### Hessian

Error inflation by a tolerance parameter (nuisance) to accommodate inconsistencies between data sets

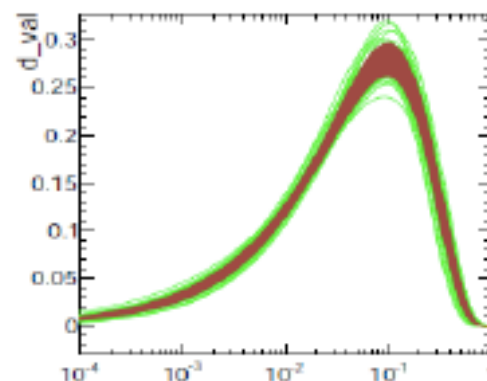
[Phys.Rev. D65 \(2001\) 014013](#), [[hep-ph/0101032](#)]



### Monte Carlo

MC replica method shifting data cross section points randomly within their uncertainties

[Phys.Rev. D58 \(1998\) 094023](#), [[hep-ph/9803393](#)]



## Various forms of parametrisation ansatz

→ HERAPDF, CTEQ style, Chebyshev, bi-log normal

[JHEP 1001:109 \(2010\)](#) [arXiv:1302.5246](#) [Phys. Let. B 695 \(2011\) 238](#)

## Bayesian Reweighting technique

[Nucl.Phys. B855, 608 \(2012\)](#), [[arXiv:1108.1758](#)],  
[JHEP 1208, 052 \(2012\)](#), [[arXiv:1205.4024](#)]

→ a method to study data sensitivity on PDFs without fitting the data

## Heavy Quark treatment in pdf's

Fixed Flavour Number Scheme - FFNS (exact calculation at fixed order)  
 Variable Flavour - VFNS (approx. eval.:  $M_c=0$ , resums large logs)

$x_{\min}$	$x_{\max}$	$Q_{\min}^2$	$Q_{\max}^2$	$\chi_{\text{tot}}^2(\text{FFN} - \text{VFN})$	$N_{\text{dat}}^{\text{tot}}$	$\chi_{\text{hera}}^2(\text{FFN} - \text{VFN})$	$N_{\text{dat}}^{\text{hera}}$
$10^{-6}$	1.0	3.0	$10^6$	28.26	2936	37.88	592
$10^{-6}$	1.0	3.0	$10^6$	68.88	1055	39.73	405
$10^{-6}$	1.0	3.0	$10^6$	28.54	422	10.65	202
$10^{-6}$	1.0	$10^2$	$10^6$	38.80	620	46.67	412
$10^{-6}$	0.1	10	$10^6$	49.67	583	32.43	350
$10^{-6}$	0.1	$10^2$	$10^6$	45.92	321	47.26	227
$10^{-6}$	0.1	10	$10^3$	31.17	510	13.52	298
$10^{-6}$	0.1	$10^2$	$10^3$	27.21	248	28.11	175

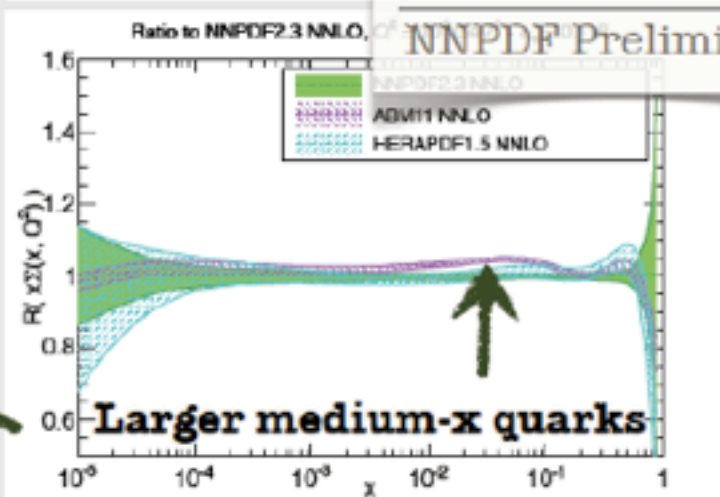
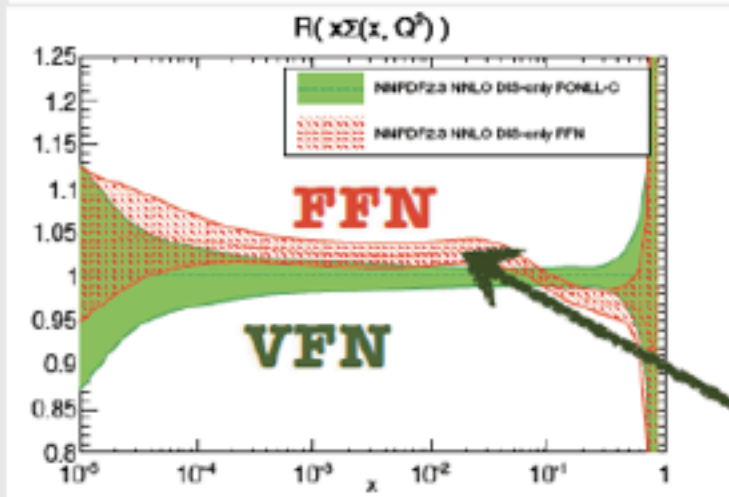
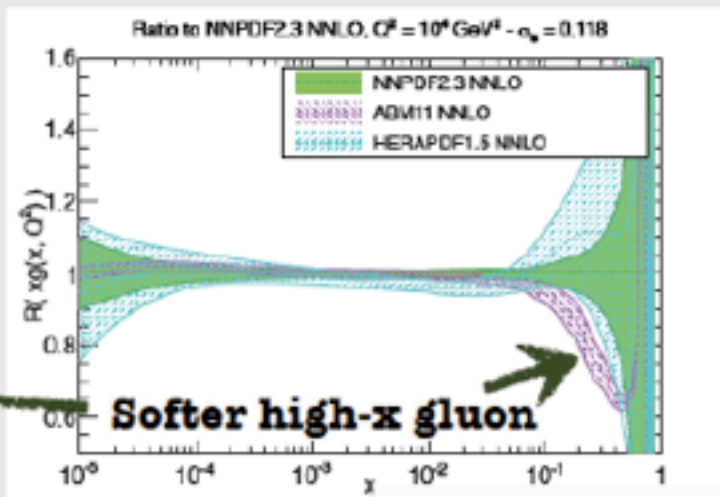
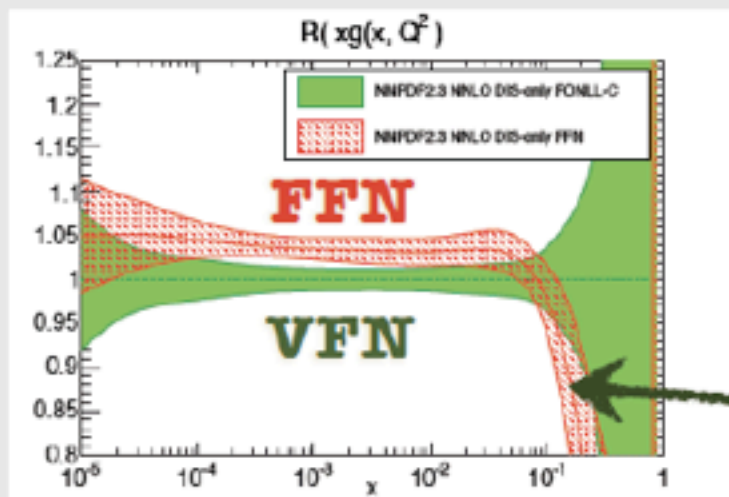
**VFNS provides consistently better fits than FFNS**  
 => importance of resummation

# Variable vs Fixed Flavor Number Schemes

☞ The impact of FFN vs GM-VFN

slide from a recent talk of *M. Cooper-Sarkar*

☞ Similar trend observed as between NNPDF2.3 and ABM11: softer large- $x$  gluon, harder medium- $x$  quarks



NNPDF Preliminary

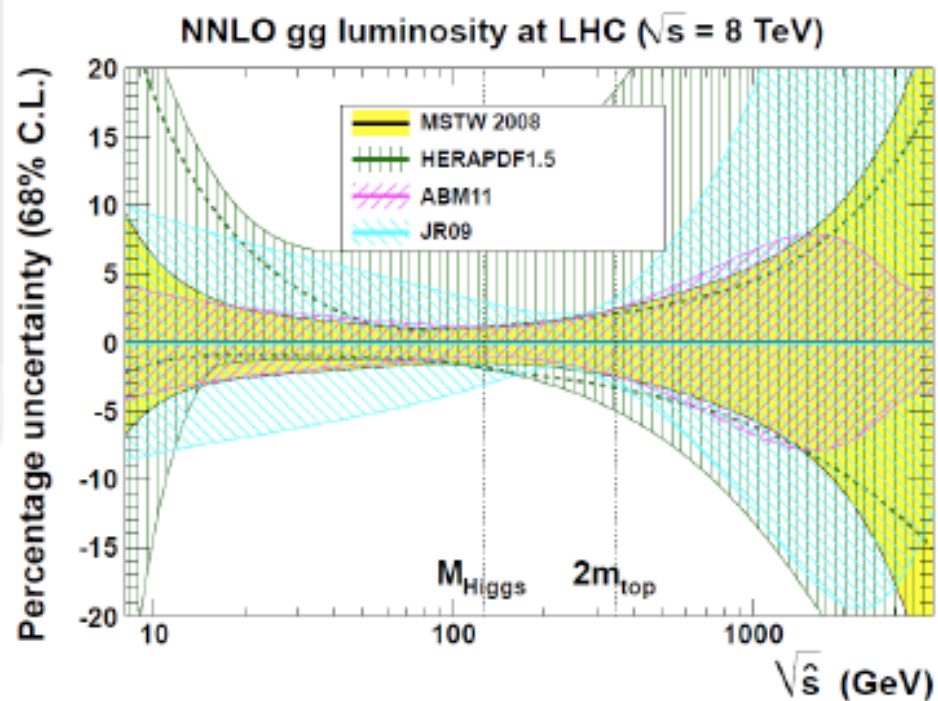
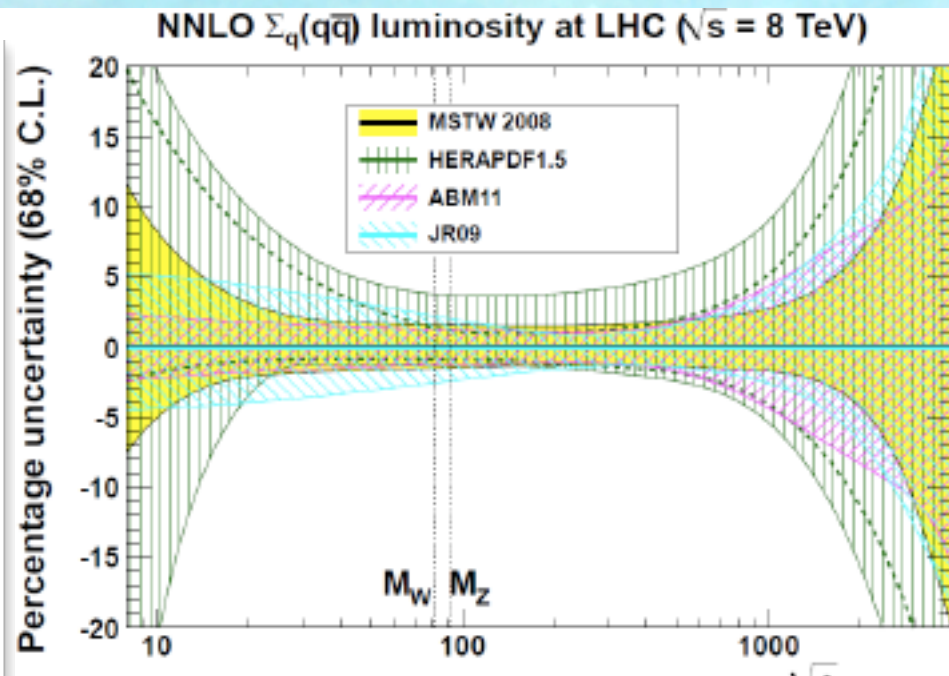
These fits are done with the same value of  $\alpha_s$  - so the PDF shape change does not come just from difference of  $\alpha_s(M_Z)$

## Sources of uncertainty

example: HERAPDF

- st. and sys. errors of data
- variation of  $Q_0^2$   
range - 1.5 to 2.5  $\text{GeV}^2$
- variation of the  $Q_{\text{cut}}^2$   
range - 2.5 to 5  $\text{GeV}^2$

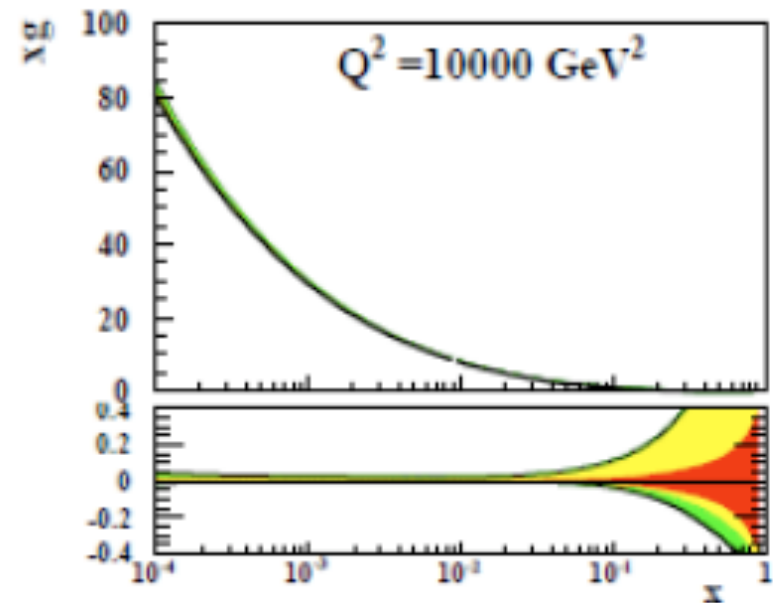
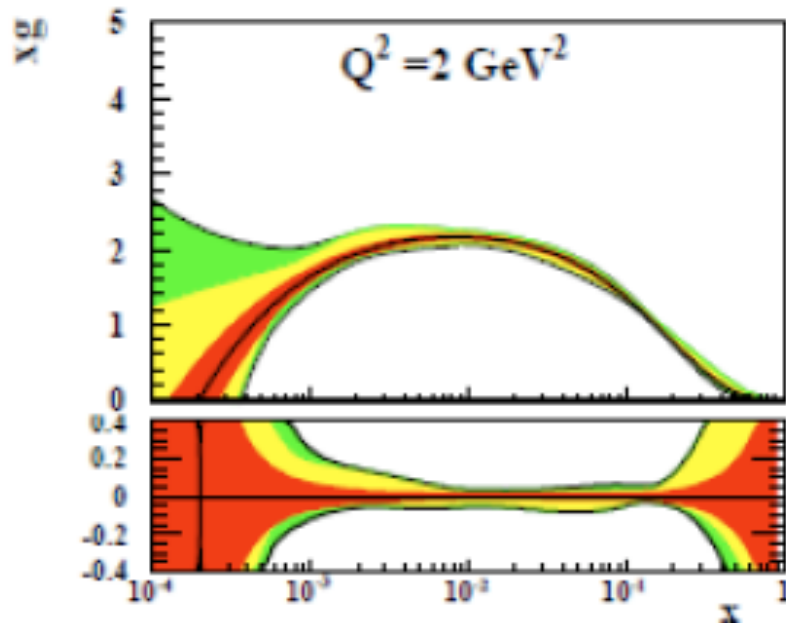
(relatively large uncertainties of HERAPDF are due presumably to the use of HERA data only)



from a recent talk of M. Cooper-Sarkar  
(DESY QCD Workshop, 2nd of Sep. 2013)

## Study of uncertainties

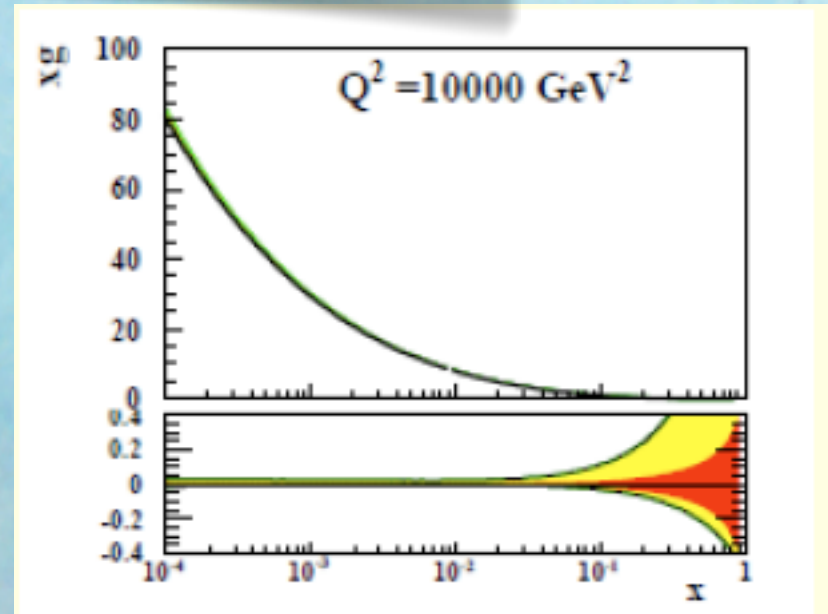
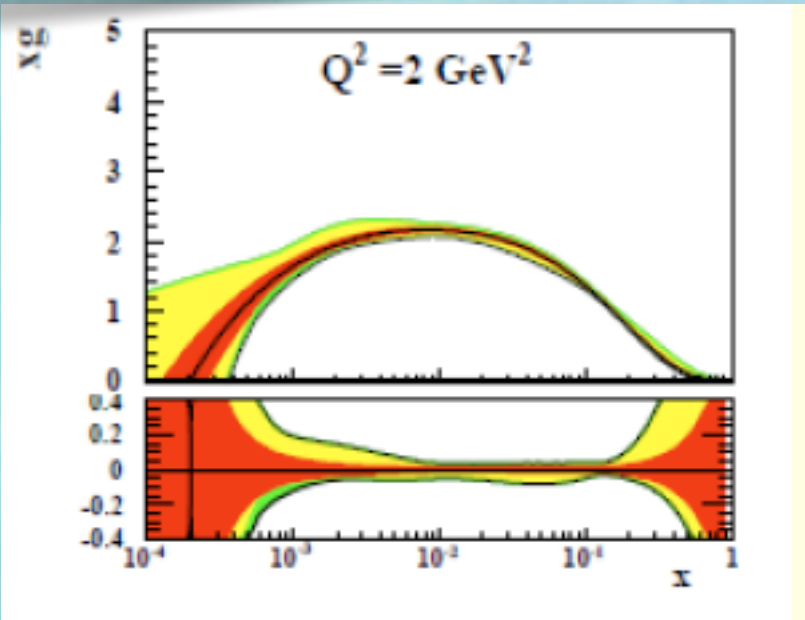
- st. and sys. errors of data (red)
- variation of  $Q_0^2$ , range - 1.9 to 2.5  $\text{GeV}^2$  (green)



► behaviour of gluon density at large  $x$  and/or large  $Q^2$ 's is strongly correlated with its behaviour at small  $x$  and small  $Q^2$ 's

## Study of uncertainties

- variation of the  $Q^2_{\text{cut}}$ , range - 2.5 to 5  $\text{GeV}^2$  (yellow) (most of the effect is coming from the change from 3.5 to 5  $\text{GeV}^2$ )

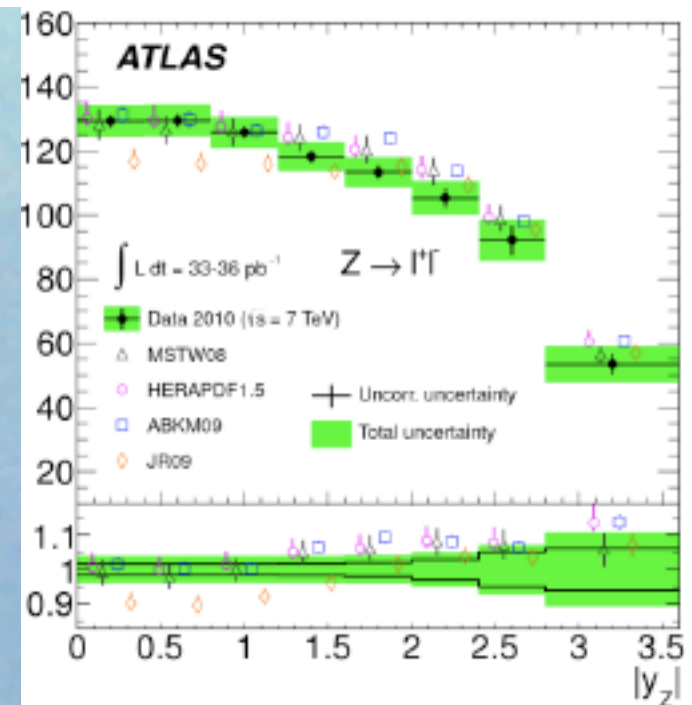
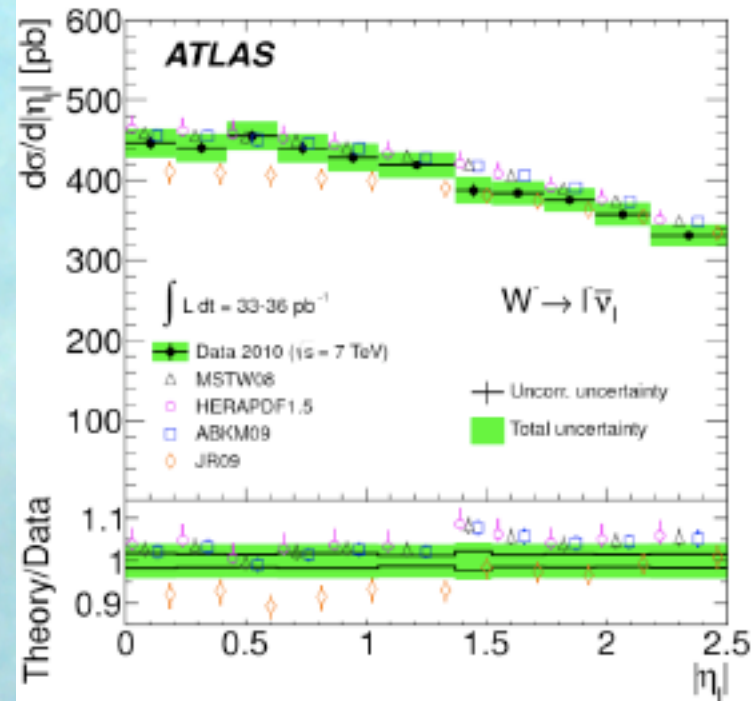
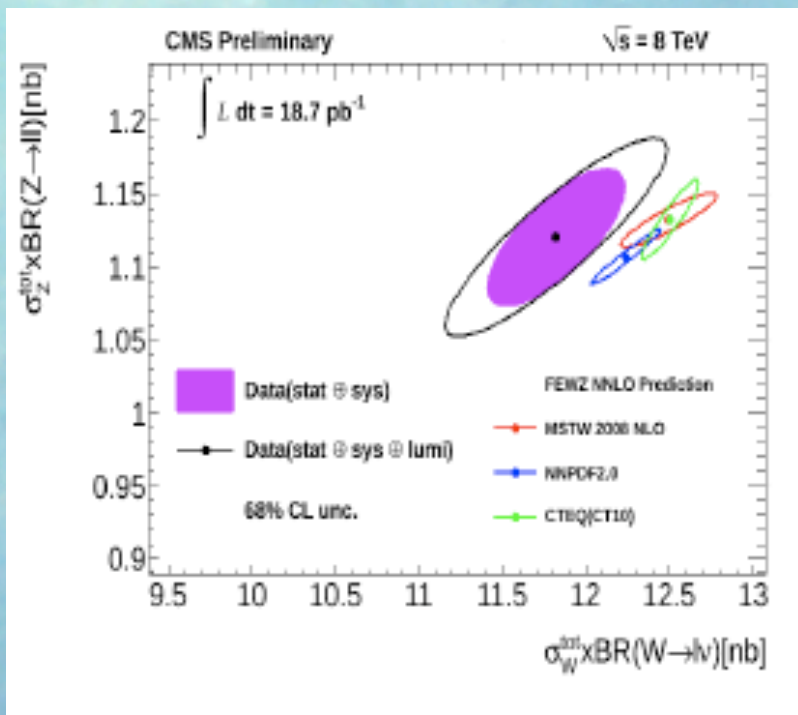


► behaviour of gluon density in the saturation region could be correlated with its large  $Q^2$  behaviour

more investigations are in progress, wait for HERAPDF2.0 which will also use HERA II inclusive combined data

# Comparison with the W and Z production at LHC

combined muon and electron data

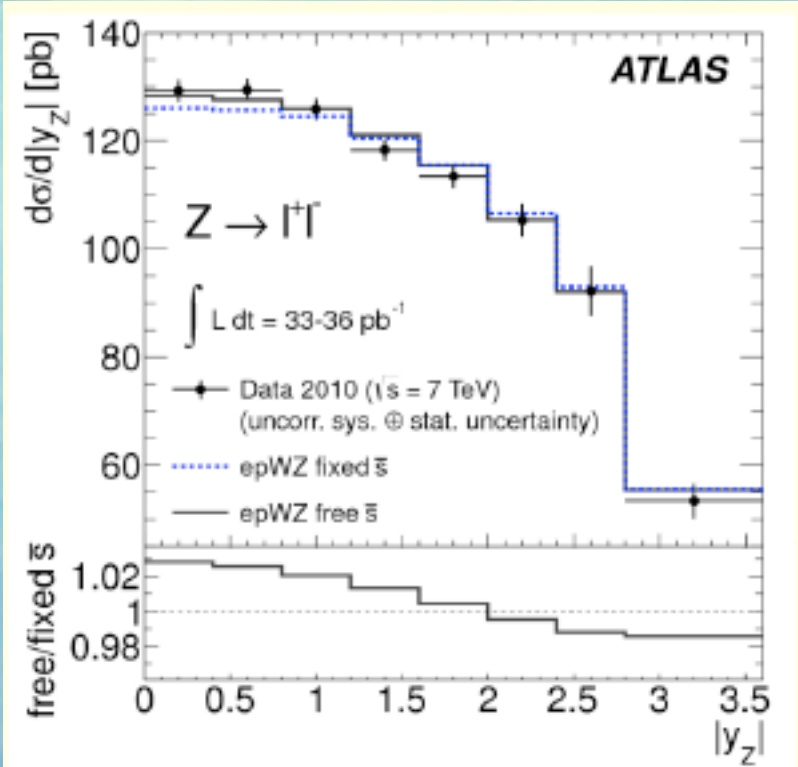


# Strangeness suppression at low-x?

dimuon production in  $\nu N$  DIS suggested that  $s/d \sim 0.5$

this was accepted by e.g.: MSTW08, NNPDF2.3, but contested by CTEQ (also difficult to understand in the dipole picture)

this is a 4% effect at LHC, could we see it?

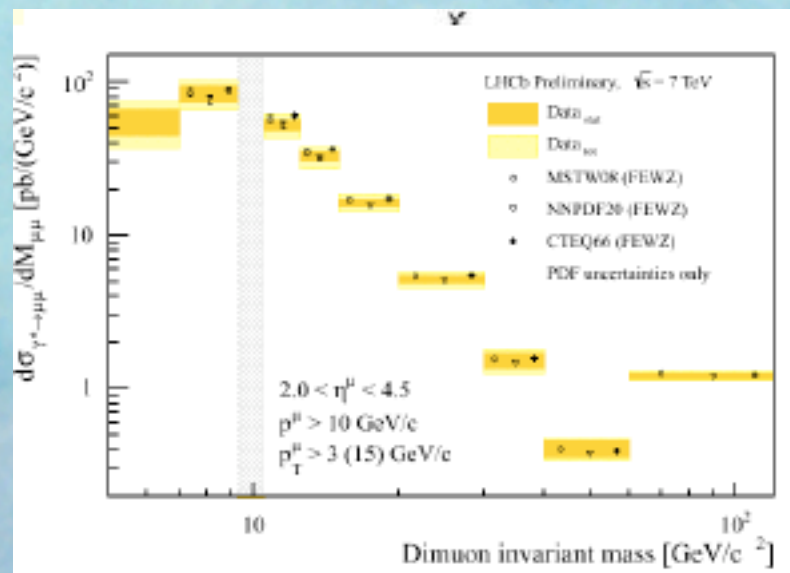
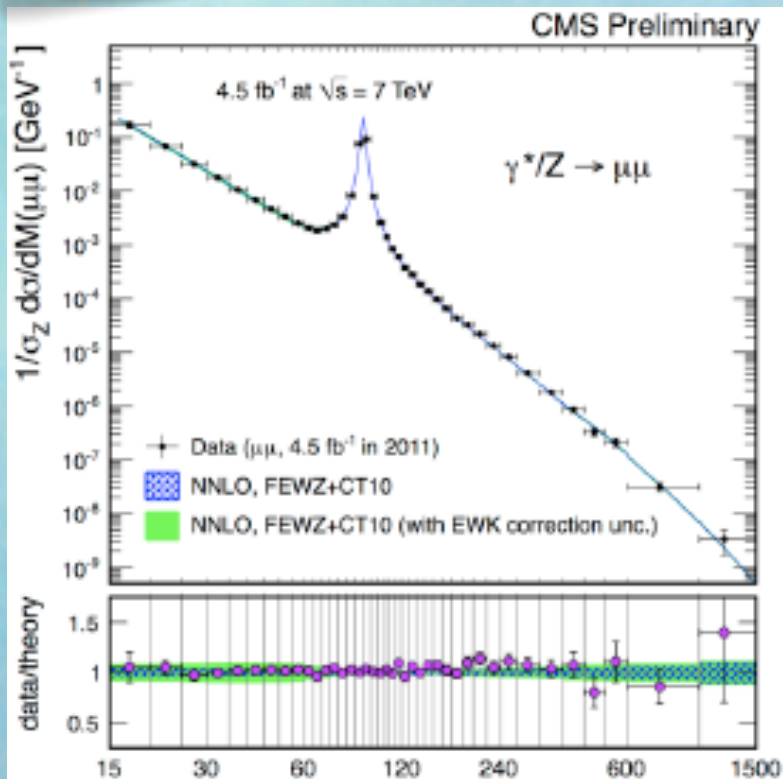
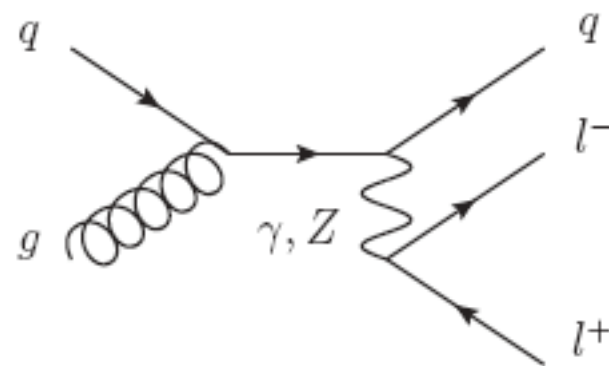


fit within HERAFitter:  
 $s/d = 1.0 \pm 0.25$



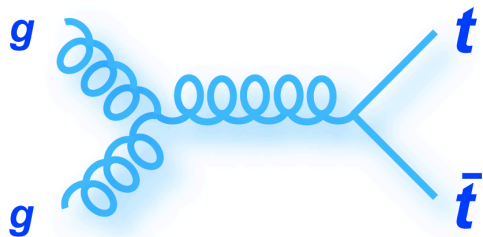
# Drell-Yan at LHC

(a potentially very interesting reaction for low- $x$  physics)



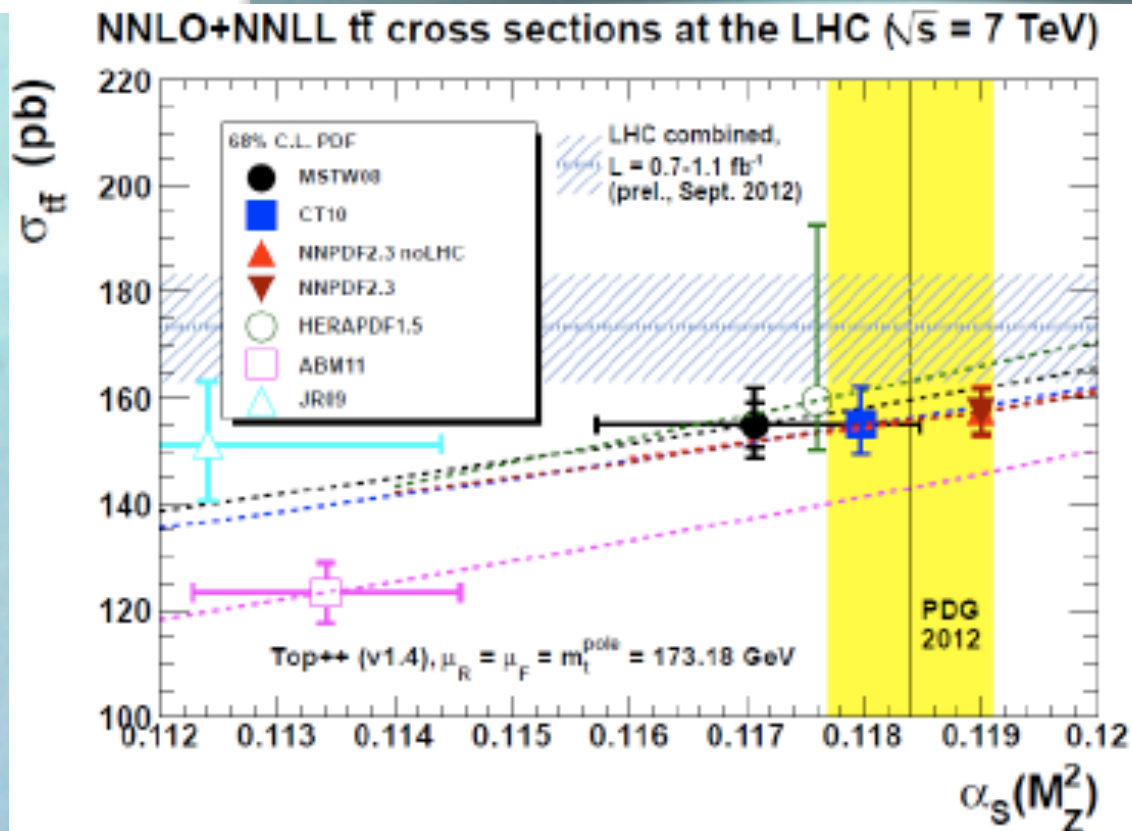
good agreement at low mass but error on data is still large

very good agreement between theory and data at high mass



## top production at LHC

test of gluon density,  $\alpha_s(M_Z)$  dependence and top mass



top production is also very sensitive to top mass, which in turn determines the running of the Higgs potential to high scales, which in turn determines the stability of our universe

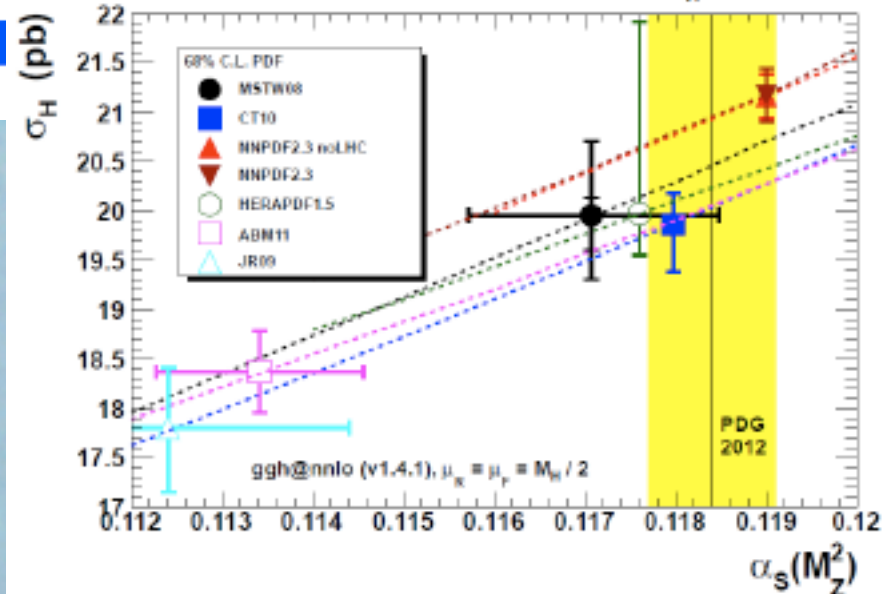
# Higgs production

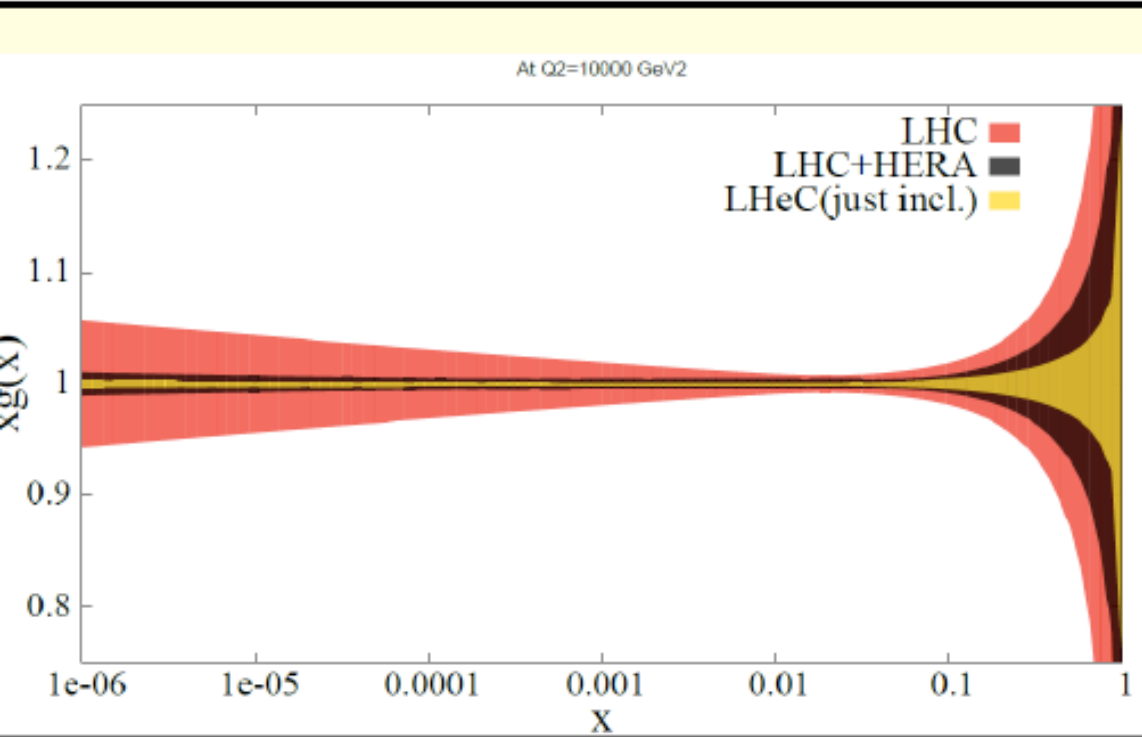
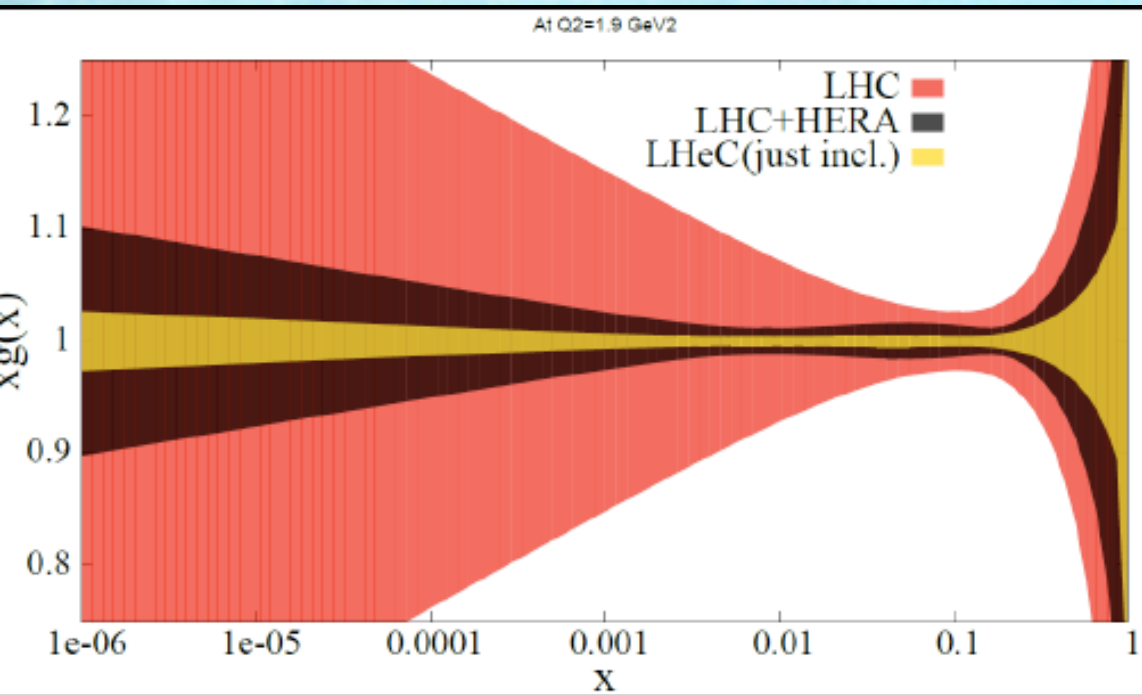
Higgs xsection is strongly dependent on the gluon density and  $\alpha_s(M_Z)$ , as the top xsection but at smaller x

		$\sigma$ (8 TeV)	uncertainty	
NNLL QCD +NLO EW	gg $\rightarrow$ H	19.5 pb	14.7%	
	VBF	1.56 pb	2.9%	
NNLO QCD +NLO EW	WH	0.70 pb	3.9%	
	ZH	0.39 pb	5.1%	
NLO QCD	ttH	0.13 pb	14.4%	

■ scale  
■ PDF+ $\alpha_s$

NNLO gg $\rightarrow$ H at the LHC ( $\sqrt{s} = 8$  TeV) for  $M_H = 126$  GeV





**Instead of  
conclusions**

**HERA data combined  
with LHC data will  
teach us a lot about  
gluon density**



