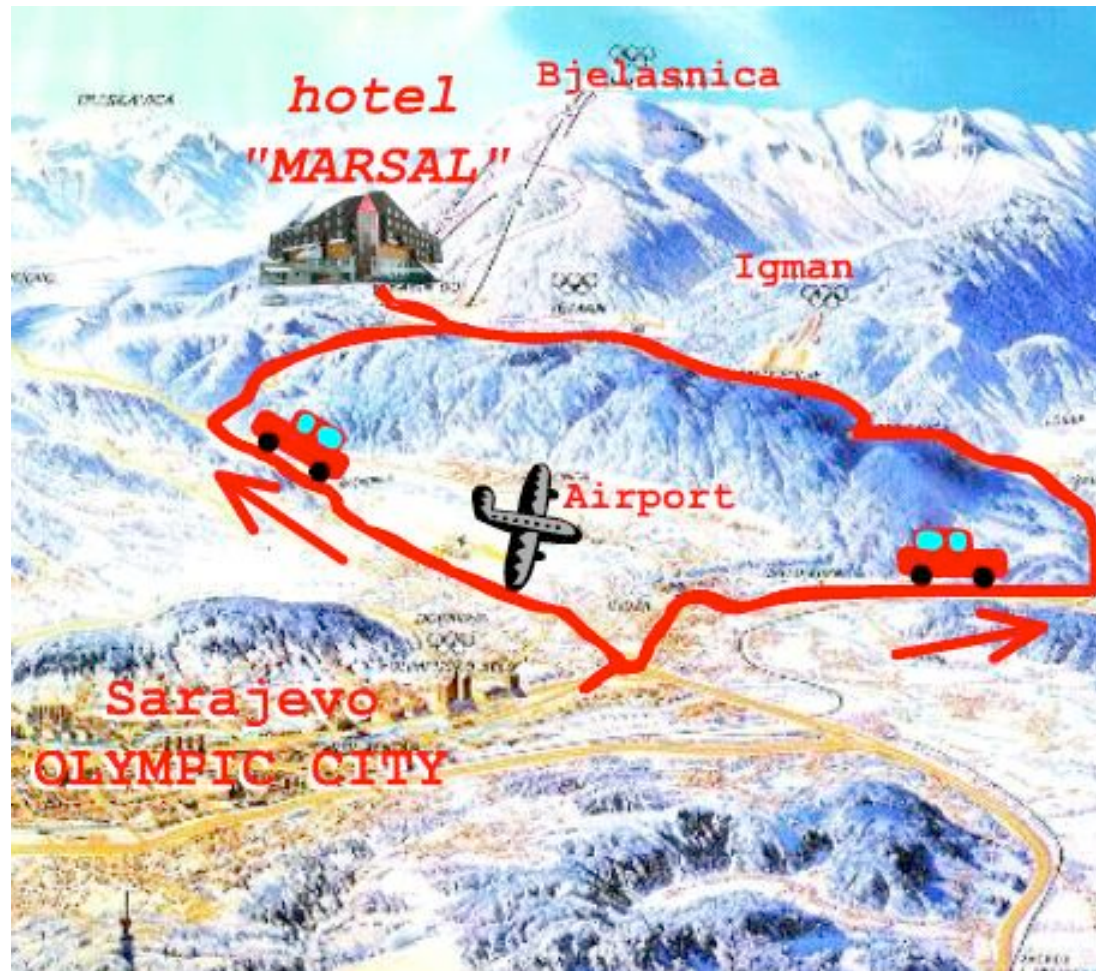


# Exotic Hadrons and Large $N_c$ QCD



# An Overview

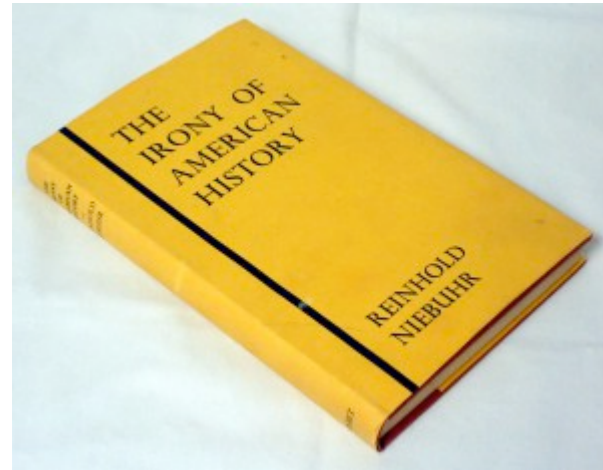


# An overview

- Introduction
  - QCD vs the Quark Model
    - Exotics
  - Why Large  $N_c$ ?
    - Glueballs, Hybrids
- The Conventional Wisdom
  - Tetraquarks can't exist at large  $N_c$
- Weinberg's critique
  - "Proof" that narrow tetraquarks don't exist is flawed
- QCD(AS)
  - Narrow tetraquarks must exist at large  $N_c$
- QCD(F) & Implications for the real world

# Introduction

- History is full of Irony



- The history of QCD is no exception: **the naïve quark model was an essential ingredient in the development of QCD, but given the existence of QCD it very hard to understand why the quark model works at all.**

# What does “quark” mean?

- It is a type of cheese which exists in the US but is really popular in Germany
- It is a nonsense word invented by James Joyce.
- It is an effective degree of freedom in the quark model.
- It is a fundamental degree of freedom in QCD.



**All these meanings are fundamentally different**



# In what sense is the quark model a “model”?



	mass →	charge →	spin →					
	≈2.3 MeV/c <sup>2</sup>	2/3	1/2	<b>u</b>	≈1.275 GeV/c <sup>2</sup>	2/3	1/2	<b>c</b>
				up				charm
	≈173.07 GeV/c <sup>2</sup>	2/3	1/2	<b>t</b>				<b>g</b>
				top				gluon
								<b>H</b>
								Higgs boson
<b>QUARKS</b>	≈4.8 MeV/c <sup>2</sup>	-1/3	1/2	<b>d</b>	≈95 MeV/c <sup>2</sup>	-1/3	1/2	<b>s</b>
				down				strange
	≈4.18 GeV/c <sup>2</sup>	-1/3	1/2	<b>b</b>				<b>γ</b>
				bottom				photon
	0.511 MeV/c <sup>2</sup>	-1	1/2	<b>e</b>	105.7 MeV/c <sup>2</sup>	-1	1/2	<b>μ</b>
				electron				muon
	1.777 GeV/c <sup>2</sup>	-1	1/2	<b>τ</b>	91.2 GeV/c <sup>2</sup>	0	1	<b>Z</b>
				tau				Z boson
<b>LEPTONS</b>	≈2.2 eV/c <sup>2</sup>	0	1/2	<b>ν<sub>e</sub></b>	≈0.17 MeV/c <sup>2</sup>	0	1/2	<b>ν<sub>μ</sub></b>
				electron neutrino				muon neutrino
	≈15.5 MeV/c <sup>2</sup>	0	1/2	<b>ν<sub>τ</sub></b>	80.4 GeV/c <sup>2</sup>	±1	1	<b>W</b>
				tau neutrino				W boson
								<b>GAUGE BOSONS</b>





- Note that from the modern perspective, the quark model is a more like a “model” in the sense of a model ship rather than the standard model. It captures *some* aspects of the real thing but misses others.
- It is not a complete theory like the standard model.
- It is definitely not beautiful like a supermodel.

- However, the simple quark model still strongly influences the language that we use to describe hadrons and remains a basic way most hadronic physicists think about states.
- Exotic hadrons are ones which do not fit into a quark model description and are important in that they help clarify what QCD has and the quark model does not.
  - There are two types
    - Quantum number exotics. States which by their quantum numbers **cannot** be made in the quark model. (eg. an isospin 2 meson)
    - Cryptoexotics. States which by their quantum numbers can be made in the simple quark model but which dynamically are dominated by components which are not of the quark model type.



- Why large  $N_c$ ?

- Many of the issues associated with the status of whether states are “really” quark model states contain intrinsic ambiguity. The hadronic states in nature are not typically bound states but resonances, while the naïve quark model describes bound states. Moreover, there can be mixing between different kinds of states (eg. a meson and a glueball).
- In many cases these ambiguities vanish in the formal large  $N_c$  limit of QCD ('t Hooft 1973).

For example

- glueballs must exist as narrow states only weakly coupled to meson in the large  $N_c$  limit (Witten 1979)
- Quantum number exotic hybrids (quark+glue states) must exist as narrow resonance (Cohen 1998)

- In many cases the large  $N_c$  limit provides a crude cartoon of the real world of  $N_c=3$  and this enables one to get insights into the real world.

For example:

- The OZI rule becomes exact at large  $N_c$
- Explains why baryons are heavier than mesons
- Explains why decays with the smallest number of mesons typically dominate decays

- In some cases the large  $N_c$  limit provides a semi-quantitative understanding in a particular for baryons

- An contracted  $SU(2N_f)$  spin-flavor symmetry emerges at large  $N_c$  and makes quantitative predictions with corrections of order  $1/N_c$  or  $1/N_c^2$  Gervais and Sakita 1983; Dashen and Manhar 1993

- Given the usefulness of the large  $N_c$  limit and  $1/N_c$  expansion one might hope that insights from large  $N_c$  to the real world for exotics.
- What is the real world situation?
  - The existence of glueballs remains controversial
  - There does seem to be strong evidence for least one hybrid state, the  $\pi_1(1400)$
  - While there are believed to be heavy tetraquark states (but these probably are special to heavy quark physics; the heavy quark limit and large  $N_c$  limit do not, in general, commute) no known exotic tetraquarks composed of light quarks exist. However, there is a long history of identifying scalars (eg.  $f_0(980)$  ) as crypto-tetraquarks (Jaffe 1977)

# The Conventional Wisdom

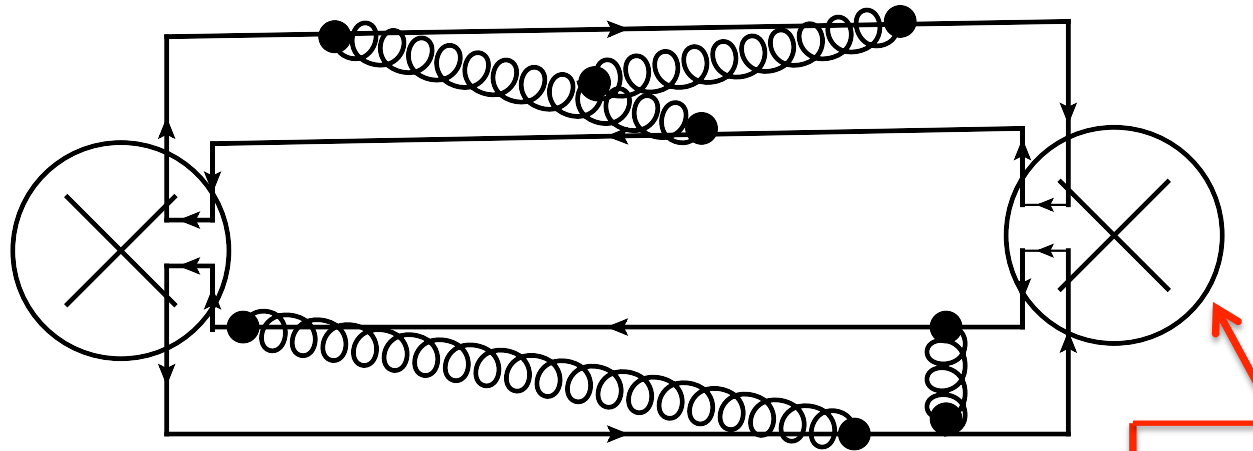
- **Tetraquarks do not exist at large  $N_c$ .** (Witten 1979; Coleman 1985)

## Basic argument:

The standard method to study hadrons at large  $N_c$  is via a study of the correlation functions for sources with the appropriate quantum numbers. It is easy to show that with a minimal tetraquark source of two bilinears at the same point, the leading order diagram ( $\mathcal{O}(N_c^2)$ ) is just a disconnected diagram which behaves like two non-interacting mesons. It does not act like a tetraquark.

# Disconnected graphs $\mathcal{O}(N_c^2)$

a typical diagram at  
quark/gluon level:  
dominated by loops  
with planar gluons  
inside



Source

$$J = \bar{q}^a(x)q_a(x)\bar{q}^a(x)q_b(x)$$

a, b are color indices



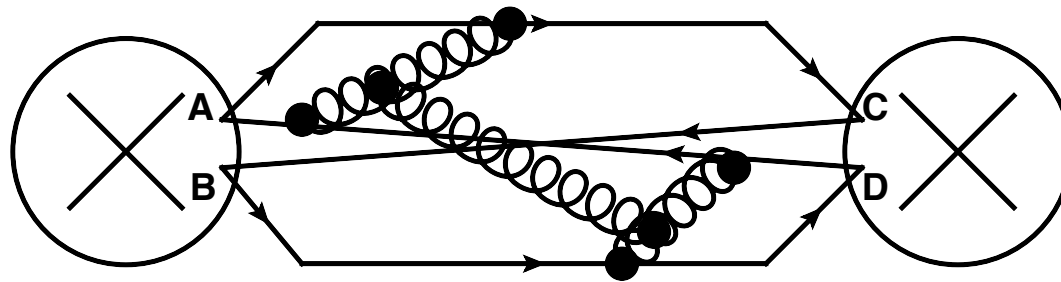
hadronic level  
two mesons



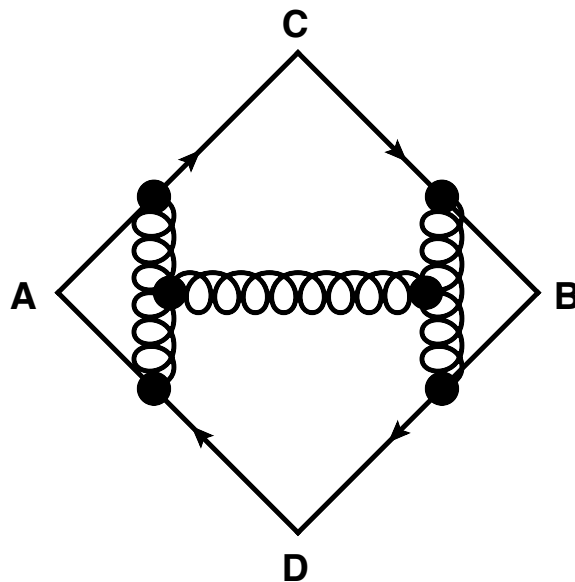
# Weinberg's Critique

- Recently this has been called into question. Weinberg pointed out in a PRL last year that the standard argument is not valid.
- The argument is wrong for a very simple reason: the fact that leading order correlator implies that that the tetraquark operator “makes two meson and nothing else” is irrelevant. One needs to look at the leading diagrams in which the four quarks all interact---i.e. the leading connected diagram---to see states which look like to interacting mesons or tetraquarks. *Whether or not these resonate into tetraquarks is separate question from whether the leading diagrams only make noninteracting mesons.*

# Connected graphs $\mathcal{O}(N_c)$



A typical diagram at quark/gluon level : dominated by a single loop with planar gluons inside. Written as a sensible looking space-time type diagram, it does not seem to be by a single loop with planar gluons inside.



But topologically it is, and the  $N_c$  counting only depends on the topology

- If tetraquarks do exist they can be found in the dynamics of these connected diagrams.
- Note that the logic by which tetraquarks must be absent since they do not appear in the  $\mathcal{O}(N_c^2)$  leading order contribution to the correlator must be wrong
  - The same argument could be applied to a 4-quark source with the nonexotic quantum number of two-pions combined to a vector-isovector. The  $\mathcal{O}(N_c^2)$  leading order contribution indeed just makes two non-interacting pions. However one cannot deduce from that a  $\rho$  meson does not exist. They do, and can be seen in the leading order connected contributions  $\mathcal{O}(N_c)$

- It is important to note, however, that Weinberg has **NOT** shown that tetraquarks do exist as narrow resonances at large  $N_c$ . Merely that the argument to disprove the existence of tetraquark is wrong.
- At this stage there is no known way to show that tetraquarks exist with the standard version of large  $N_c$  QCD. Indeed, I will argue at the end that they do not!!
- However there is a variant of large  $N_c$  QCD in which it is possible to show that quantum number exotics must exist and become narrow as  $N_c \rightarrow \infty$ .

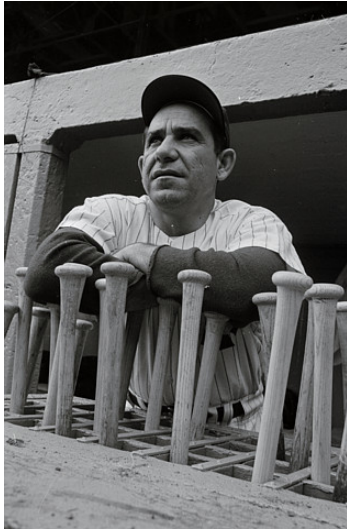
# QCD (AS)

- **The large  $N_c$  limit of QCD is not unique**
  - **For gluons there is a unique prescription  $SU(3) \rightarrow SU(N_c)$**
  - **However for quarks, we can choose different representations of the gauge group**
  - **Asymptotic freedom restricts the possibilities to the fundamental (F), adjoint (Adj), two index symmetric (S), two index anti-symmetric (AS).**
    - **Adj transforms like gluons (traceless fundamental color-anticolor); dimension  $N_c^2 - 1$ ; 8 for  $N_c = 3$  (unlike our world).**
    - **S transforms like two colors (eg fundamental quarks) with indices symmetrized; dimension  $\frac{1}{2}N_c(N_c + 1)$ ; 6 for  $N_c = 3$  (unlike our world).**
    - **AS transforms like two colors (eg fundamental quarks) with indices antisymmetrized; dimension  $\frac{1}{2}N_c(N_c - 1)$ ; 3 for  $N_c = 3$  (just like our world).**



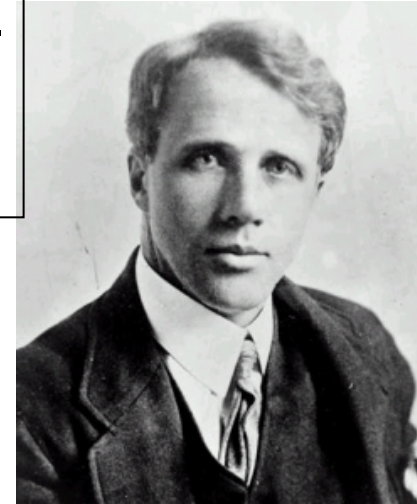
- Note that  $N_c=3$  quarks in the AS representation are indistinguishable from the (anti-)fundamental. (In essence antisymmetric  $r$   $b$  is the same as  $\bar{g}$  .)
- However quarks in the AS and F extrapolate to large  $N_c$  in **different ways**.
  - The large  $N_c$  limits are physically different
  - The  $1/N_c$  expansions are different.
  - A priori it is not obvious which expansion is better
  - It may well depend on the observable in question
- The idea of using QCD (AS) at large  $N_c$  is old
  - Corrigan & Ramond (1979)
  - Idea was revived in early 2000's by Armoni, Shifman and Veneziano who discovered a remarkable duality that emerges at large  $N_c$ .

# Two Roads to Large Nc QCD



Quarks in  
Fundamental

Quarks in 2-  
index anti-  
symmetric



“When you come to a  
fork in the road, take it.”

---Yogi Berra,  
American baseball  
player, coach and part-  
time philosopher

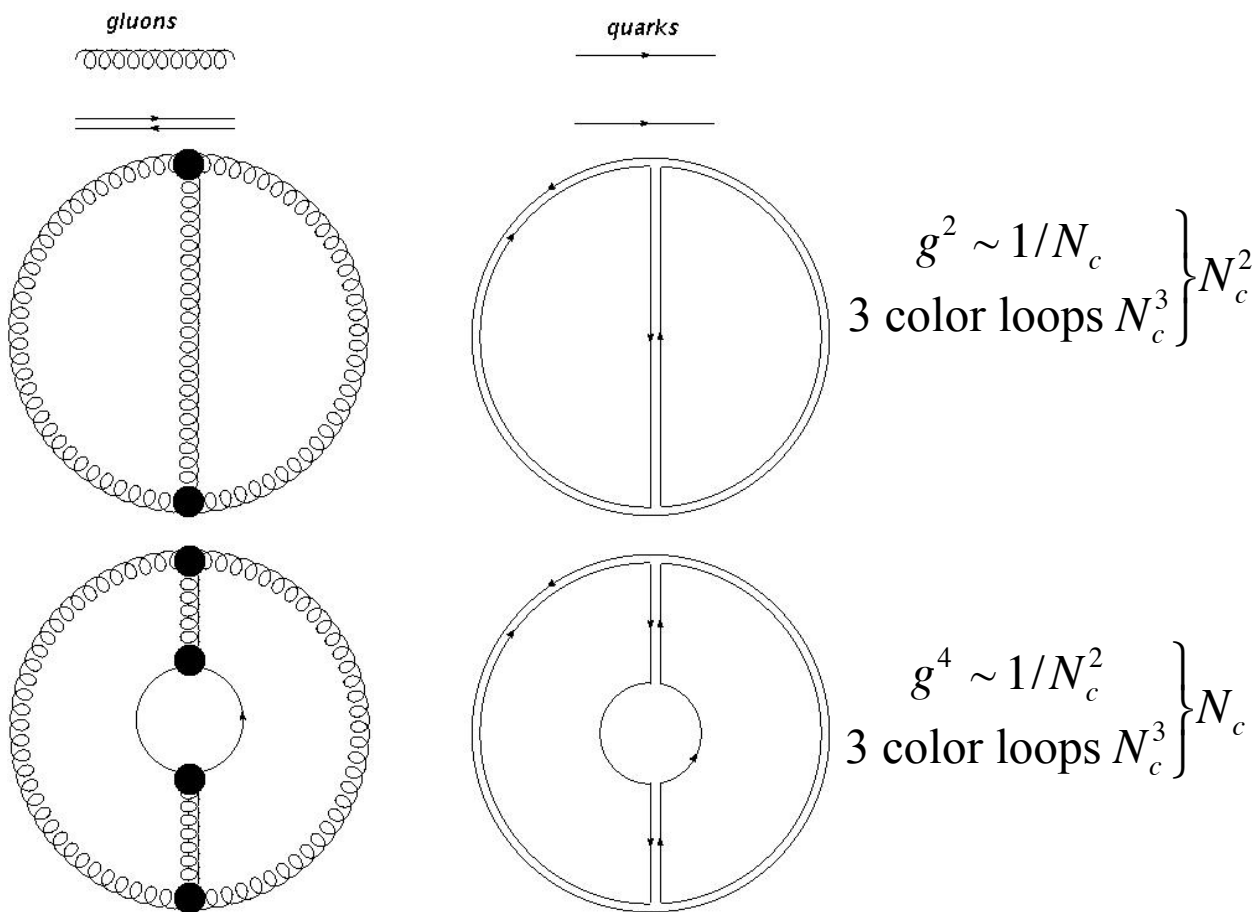
“Two roads diverged in a  
wood, and I—  
I took the one less traveled  
by And that has made all the  
difference.”

---Robert Frost,  
American poet

Large Nc QCD

Principal difference between QCD(AS) and QCD(F) at large  $N_c$  is in the role of quarks loops

Easy to see this using 't Hooft color flow diagrams

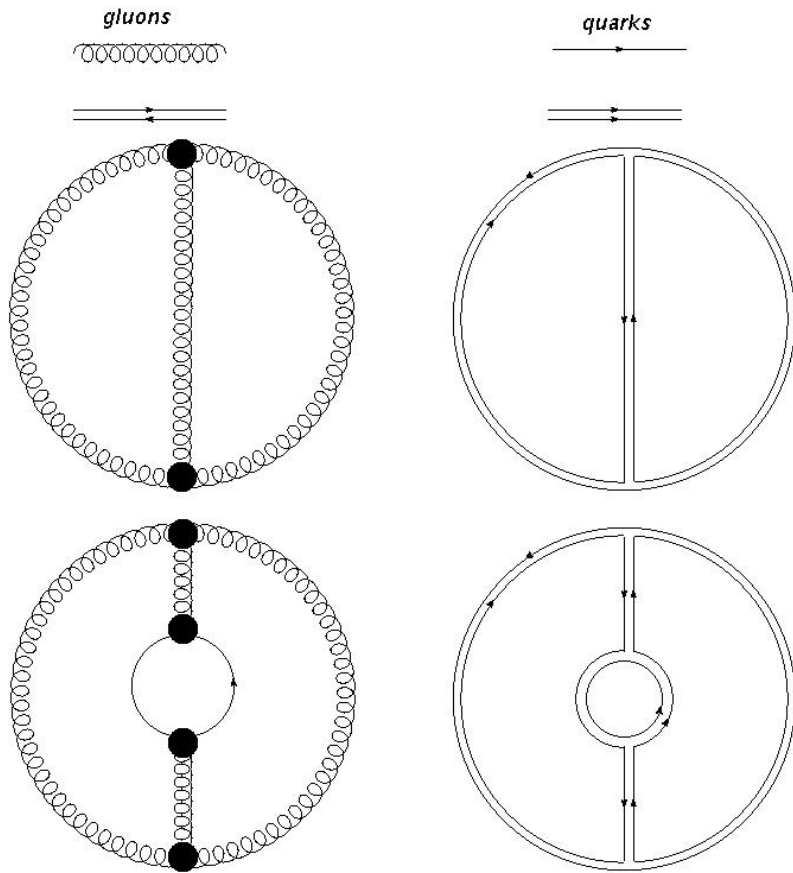


Recall the 't Hooft large  $N_c$  limit keeps  $g^2 N_c$  fixed  
 So  $g^2 \sim 1/N_c$

QCD(F)

Insertion of a planar quark loops yields a  $1/N_c$  suppression.

Leading order graphs are made of planar gluons



$$g^2 \sim 1/N_c \left. \begin{array}{l} 3 \text{ color loops } N_c^3 \end{array} \right\} N_c^2$$

$$g^4 \sim 1/N_c^2 \left. \begin{array}{l} 4 \text{ color loops } N_c^4 \end{array} \right\} N_c^2$$

## QCD(AS)

Insertion of a planar quark loops does not lead to a  $1/N_c$  suppression.

Leading order graphs are made of planar gluons and quarks

Principal phenomenological difference between the two is the inclusion of quark loop effects at leading order in QCD(AS). Whether this is a bug or a feature depends upon the observable. In baryon spectroscopy based on emergent symmetry, both QCD(AS) and QCD(F) appear to have predictive power (Cherman, Cohen & Lebed 2009, 2012).

- QCD(AS) naturally includes quark loops. Thus one might expect that in non-quantum number exotic channels tetraquarks will mix with ordinary mesons at leading order.
  - This can be shown to be correct.
- More interestingly, in quantum number exotic channels, QCD(AS) **MUST** have narrow tetraquarks at large  $N_c$  (i.e. narrow states which have at least 2 quarks and 2 antiquarks)Cohen&Lebed 2014.



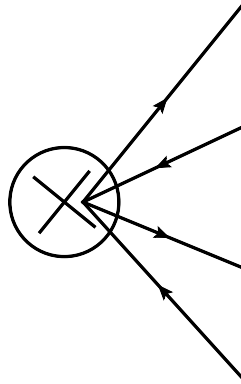
Key ingredient: there are single color trace tetraquark sources in QCD(AS). That is the source cannot be broken up into two separate color singlets (except for  $N_c^{-2}$  contributions). **This cannot be done in QCD(F)**

$$J(x) = \sum_{\substack{A,B \\ a,b,c,d}} C_{AB} \bar{q}^{ab}(x) \Gamma_A q_{bc}(x) \bar{q}^{cd}(x) \Gamma_B q_{da}(x)$$

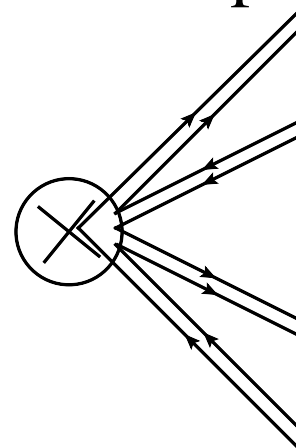
$\Gamma_A, \Gamma_B$  are matrices in Dirac-flavor space.

$a, b, c, d$  are fundamental color indices

choice of  $C_{AB}$  fixes quantum #s; for simplicity chose an exotic

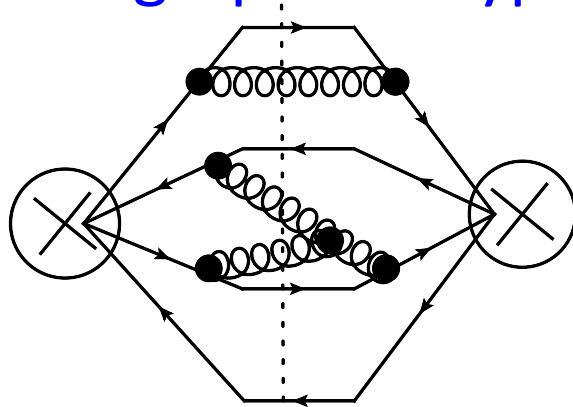


Source as a Feynman diagram

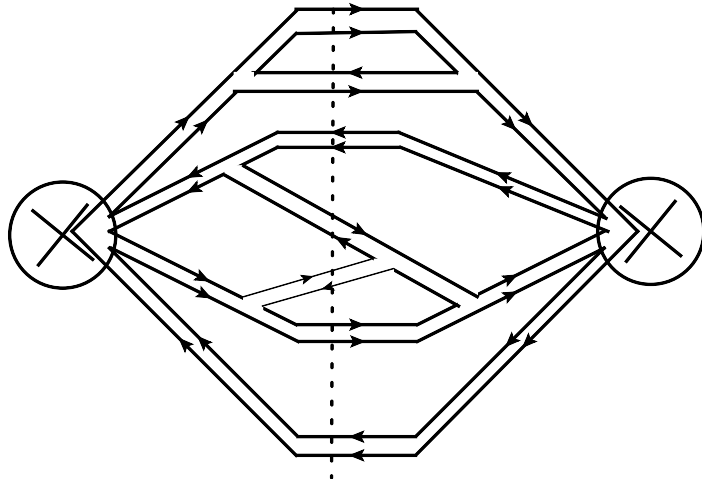


Source as a color-flow diagram

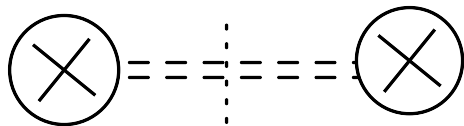
Look at the JJ correlation function. It is dominated by planar graphs. A typical diagram scales as  $N_c^4$



Feynman diagram

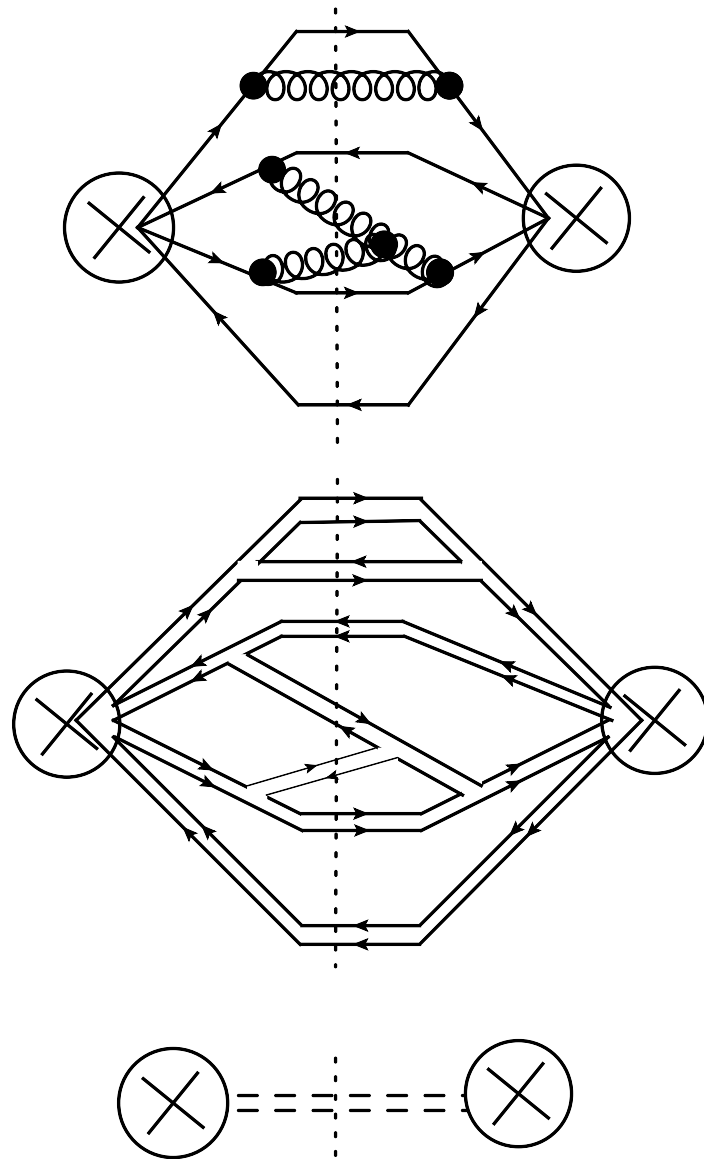


Color-flow diagram; 7 color loops  $\sim N_c^7$ ; 6 factors of  $g \sim N_c^{-3}$ ; overall scaling  $\sim N_c^4$



Hadronic level diagram: propagation of a single tetraquark

The reason this corresponds to a single tetra hadrons can be understood in terms of a cut of the diagram.



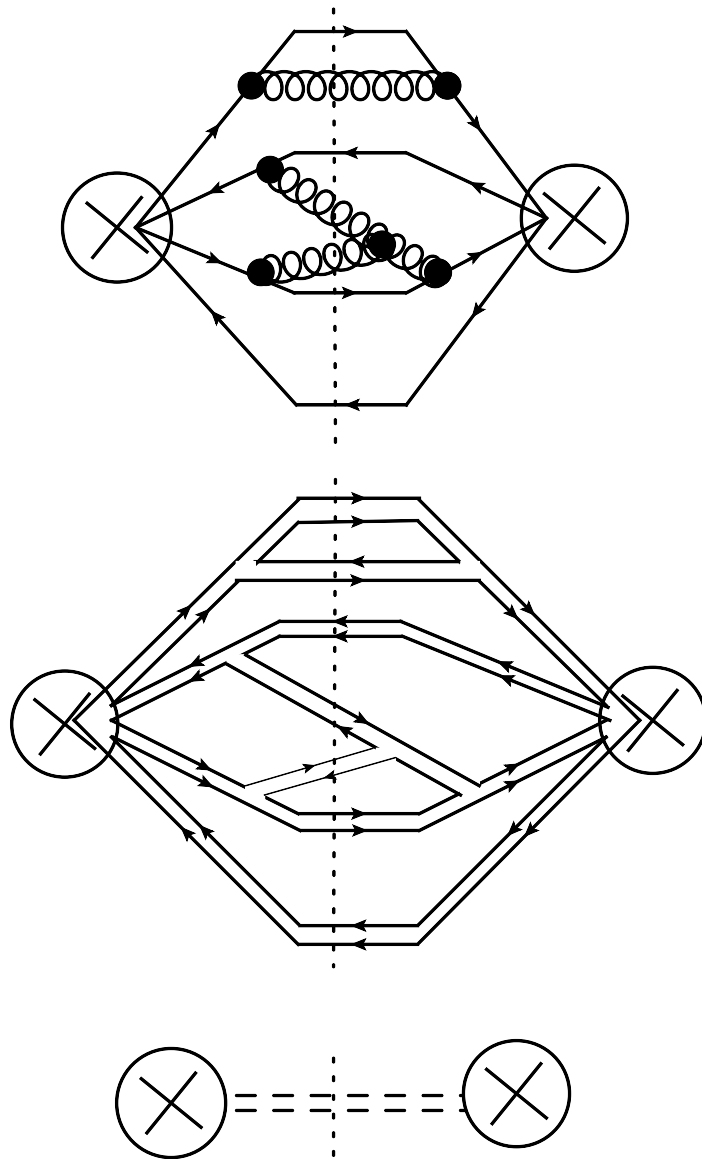
Short dashed line indicates a cut which reveals the intermediate state structure of the diagram.

The cut shown here corresponds to a state of the form

$$\bar{q}^{ab} q_{bc} A_d^c A_e^d \bar{q}^{ef} A_f^g q_{ga}$$

This is a single color-trace object. It can not be divided into two separate color singlets (except by a  $1/N_c^2$  contribution)

This is generic: all cuts yield single-color trace objects

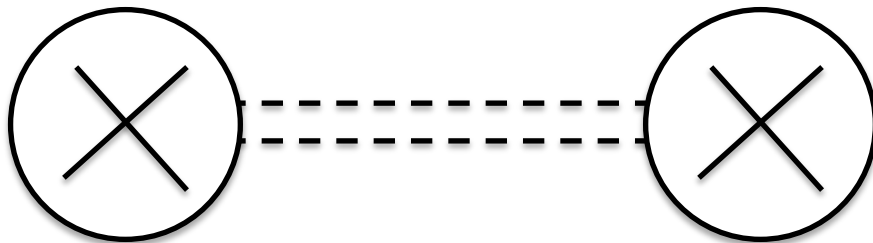


**If one includes confinement, this implies that the state must be a single hadron at leading order. It cannot break up into two color singlet hadrons since all intermediate states consist of a single indivisible color singlet.**

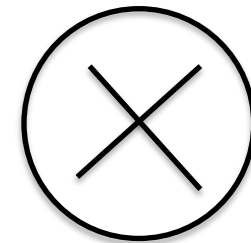
**It must be narrow as components with more than one hadron are suppressed in the  $1/N_c$  expansion.**

- This can be seen to be true self-consistently
    - One can use standard kind of large  $N_c$  analysis for correlators with appropriate changes to account for QCD(AS) to deduce that a generic multi-hadron vertex scales as  $N_c^{2-n}$  where  $n$  is the number of hadrons (mesons, glueballs, hybrids & tetraquarks).
    - Thus, tetraquark width  $\sim N_c^{-2}$
- Thus as advertised the tetraquark is narrow

An example: the tetraquark-2 meson vertex.  
Proceed by studying the scaling of the appropriate correlator.



Recall at quark-gluon level we deduced that the tetraquark two point function scaled as  $N_c^4$

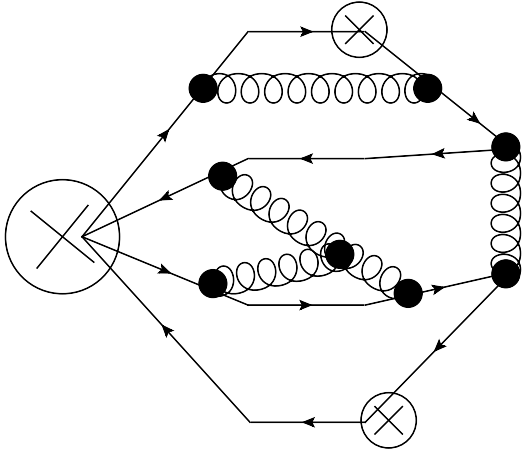


Thus the source scales as  $N_c^2$



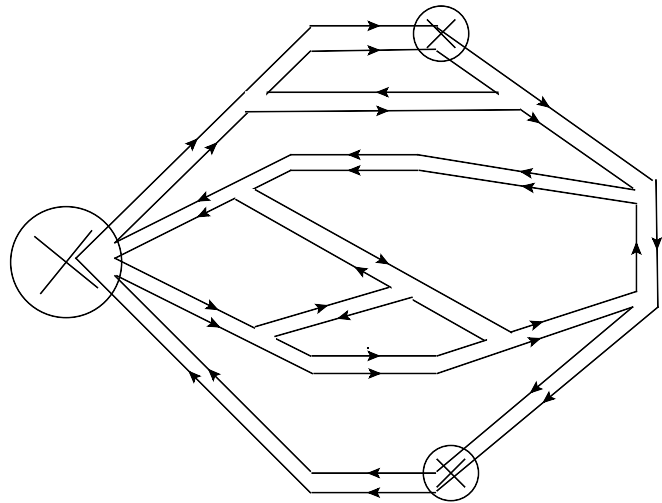
The meson source scales as  $N_c^1$  by analogous reasoning

## A typical graph

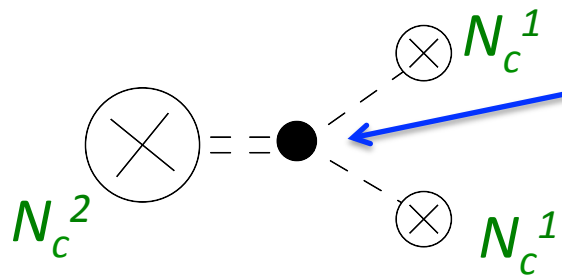


Feynman  
diagram: 8  
factors of  $g$

Overall  $N_c$   
scaling  $N_c^3 =$   
 $(N_c^{-4})(N_c^7)$



Color flow  
diagram: 7  
color loops

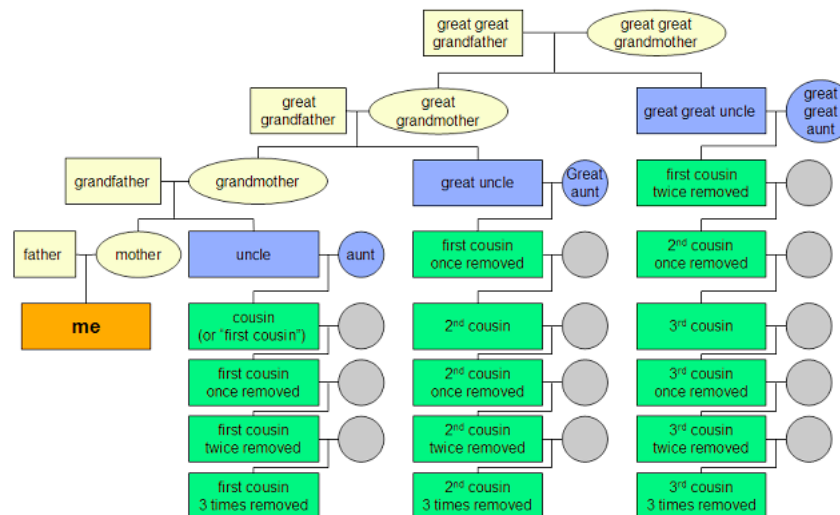


Hadronic  
level diagram

To match scaling at QCD  
level the tetraquark-2  
meson vertex must scale  
as  $N_c^{-1}$  yielding a width of  
 $N_c^{-2}$ .

# Implications For the Real World

- Minimally, this result shows that nothing in the structure of gauge theories such as QCD excludes exotic tetraquark states made from light quarks as narrow resonances---a cousin of QCD has been shown to have them.
- The question is “how close a cousin”?

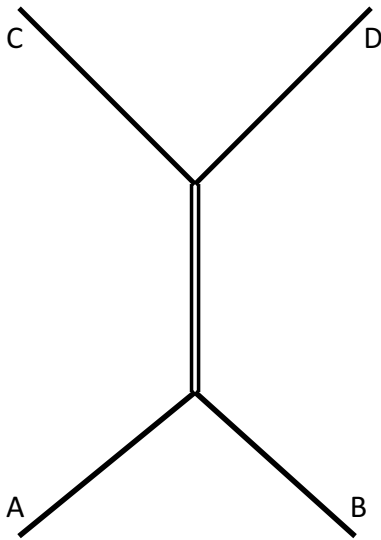




- One hint on this would be the behavior of QCD(F) (the more standard type of large  $N_c$  QCD)
- There is a somewhat subtle argument that despite Weinberg's critique of Coleman and Witten tetraquarks in fact, do not exist at large  $N_c$  in QCD(F) cohen & Lebed in preparation
  - Recall that Weinberg merely showed that the Witten/Coleman argument against tetraquarks is flawed, not that tetraquarks themselves exist.
  - The argument is too involved to be given in detail here.

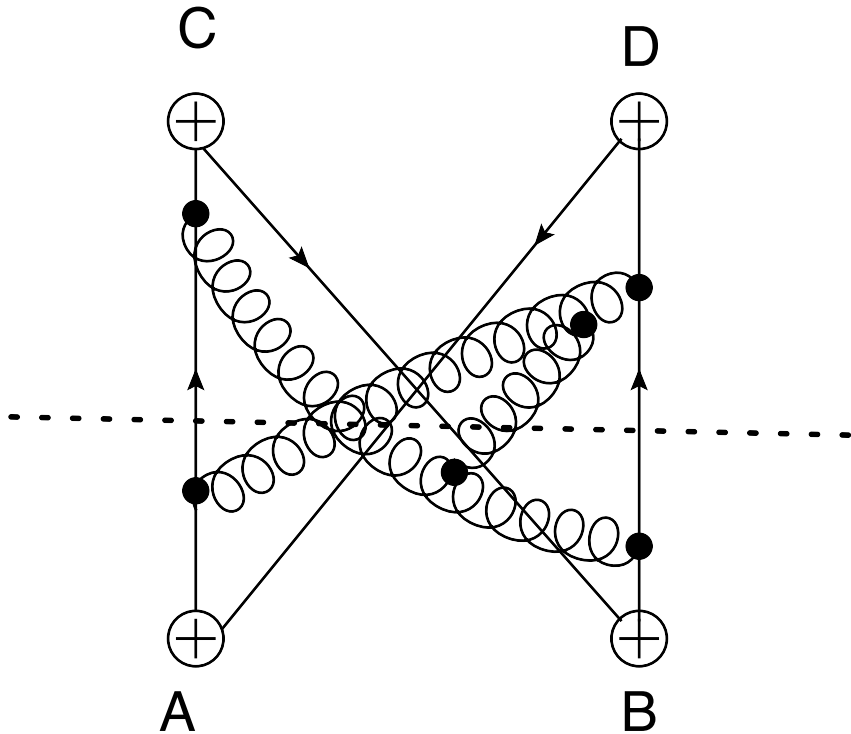
## A brief sketch of the argument

If exotic tetraquarks exist they will couple to ordinary meson with a coupling strength  $\sim N_c^{-1/2}$ . Thus it must appear as a singularity in the s-channel of scattering for incident mesons.



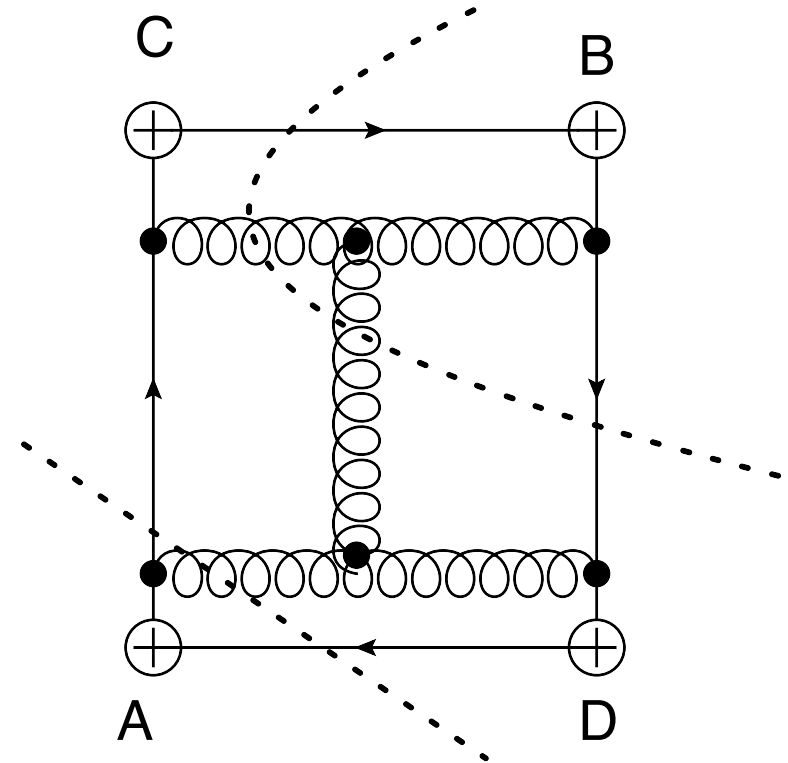
There is a topological argument that the connected 4-point function at leading order **every diagram** in an exotic channel only has singularities in the s-channel associated with the asymptotic mesons (either initial or final) in the sense that the cut has two color singlets carry the initial four momenta of each; thus there are no singularities associated with intermediate object.

## A typical contribution



Space-time diagram

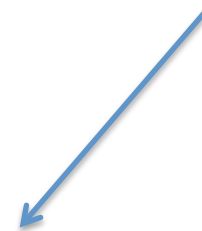
This behavior is generic and occurs in all diagrams associated with exotic channels



Topologically equivalent planar graph. Note that the cut has broken in two and each part carries the 4-momenta inserted at A or B

- Given this argument, whether tetraquarks exist in the real world depends on whether the real world is closer (in this aspect) to QCD(F) at large  $N_c$  or to QCD(AS) at large  $N_c$ . This is a dynamical question.

The bottom line



---

Whether narrow tetraquarks exist in the real world is a matter of dynamical detail which generic large  $N_c$  arguments cannot answer. Large  $N_c$  arguments show they cannot be excluded as incompatible with QCDlike theories