

Non-perturbative features of the three-gluon vertex in Landau gauge

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The three-gluon vertex in Landau gauge, arXiv: 1402.1365

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Outline of talk

- 1 Introduction: Strong QCD.
- 2 Greens functions and Dyson-Schwinger equations (DSEs).
- 3 Greens functions of QCD.
- 4 Three-gluon vertex: calculation setup and results.
- 5 Summary and outlook.

Strong QCD and non-perturbative methods

- QCD coupling α_S large at low energies \rightarrow Strong QCD.
- Non-perturbative phenomena: $D_\chi SB$, confinement, ...
- Studies in this regime require non-perturbative tools.
 - 1 Lattice QCD.
 - 2 Functional renormalisation group (FRG).
 - 3 Effective field theories (ChPT...)
 - 4 Dyson-Schwinger/Bethe-Salpeter/Faddeev equations.
- These tools are used for hadron phenomenology.
 - 1 Masses.
 - 2 Decay constants.
 - 3 Form factors.

Greens functions

- Basic objects in QFTs are Greens functions.
- n -point Greens function encodes interactions of n 'particles'.
- Continuum version of lattice QCD correlation functions.
- Example: two-point Greens function (propagator) given by

$$\begin{aligned} G(x_1, x_2) &= \frac{1}{N} \int \mathcal{D}\phi e^{-i(S + \int_x J \cdot \phi)} \phi_1(x_1) \phi_2(x_2) \\ &= \langle \phi_1(x_1) \phi_2(x_2) \rangle \end{aligned} \quad (1)$$

$S =$ action, $J =$ 'sources'

Dyson-Schwinger equations

- Theory invariant under field translations $\phi(x) \rightarrow \phi(x) + \epsilon(x)$
 \Rightarrow Dyson-Schwinger equations (DSEs).

$$\left\langle \frac{\delta S[\phi]}{\delta \phi(x)} \right\rangle_J = J(x) \quad (2)$$

- DSEs connect various n -point Greens functions.
- Infinite tower of coupled integral equations.
- Exact, continuous formulation of a quantum field theory. ✓
- Perturbation theory can be extracted from DSEs. ✓
- Practical calculations necessitate truncations. ✗

Quark propagator DSE

$$\text{Quark Propagator}^{-1} = \text{Free Quark Propagator}^{-1} + \text{Quark Loop Diagram}$$

- Exact equation, with full gluon propagator and quark-gluon vertex.
- Satisfy own DSEs, with higher-point Greens functions $\rightarrow \infty$ tower.

QCD action

$$S = \text{quark propagator} + \text{quark-gluon vertex} + \text{gluon propagator} + \text{three-gluon vertex} + \text{four-gluon vertex}$$

The diagram shows the expansion of the QCD action S into five terms, each represented by a Feynman diagram:

- quark propagator:** A horizontal solid line connecting two red circular vertices. A coefficient of -1 is written above the line.
- quark-gluon vertex:** A red circular vertex at the top connected by a wavy line to a horizontal solid line connecting two other red circular vertices.
- gluon propagator:** A horizontal wavy line connecting two red circular vertices. A coefficient of -1 is written above the line.
- three-gluon vertex:** A central vertex connected by wavy lines to three red circular vertices.
- four-gluon vertex:** A central vertex connected by wavy lines to four red circular vertices.

QCD action (not gauge-fixed, no ghosts)

Quark propagator and quark-gluon vertex

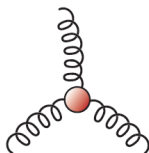
- Quark propagator
→ central object in bound state equations (Bethe-Salpeter, Fadeev).
- Quark-gluon vertex: connects matter and gluon fields.
 - 1 Essential ingredient in quark propagator DSE.
 - 2 Structure of bound state equations, Ward-Takahashi identity.

The diagram shows the Dyson-Schwinger equation (DSE) for the quark-gluon vertex. On the left is the full vertex, represented by a black circle with a gluon line (curly) entering from the top and two quark lines (arrows) exiting from the bottom. This is set equal to a sum of terms on the right. The first term is the tree-level vertex (black circle). The second term is a loop diagram with a dashed line (ghost) and a quark line. The third term is a loop diagram with a gluon line and a quark line, with a coefficient of $+\frac{1}{2}$. The fourth term is a loop diagram with a gluon line and a ghost line. The fifth term is a loop diagram with a gluon line and a quark line. The sixth term is a loop diagram with a gluon line and a ghost line, with a coefficient of $+\frac{1}{6}$.

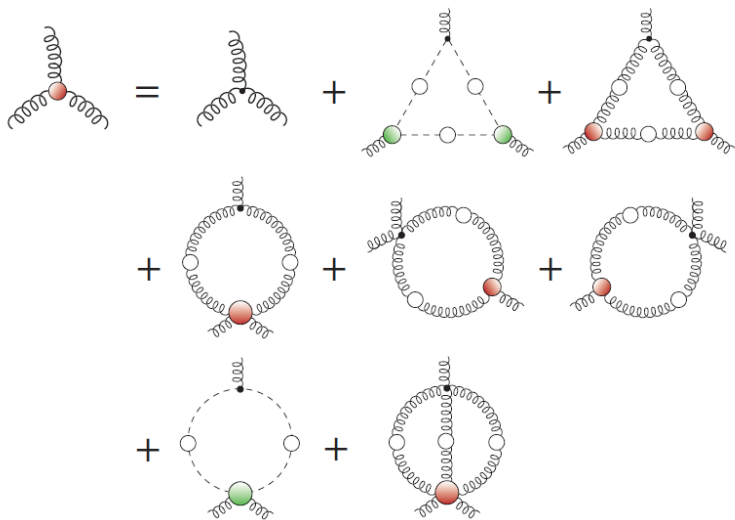
Full DSE for quark-gluon vertex.

Three-gluon vertex Γ_{3g}

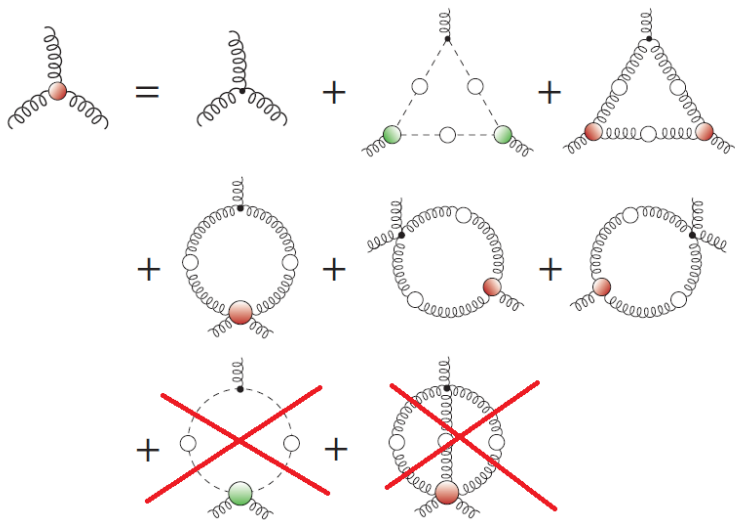
- Γ_{3g} and Γ_{4g} epitomise self-interaction of gluon field.
- Importance of three-gluon vertex:
 - 1 Driving term in quark-gluon vertex DSE.
 - 2 Enters bound state equations through quark-gluon vertex.
 - 3 Irreducible three-body force in Faddeev equation.
 - 4 Closing truncated system of DSEs in ghost/gluon sector.
- Γ_{3g} is the current focus of our study.



Three-gluon vertex DSE



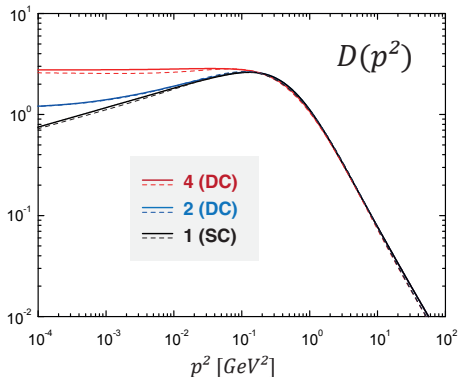
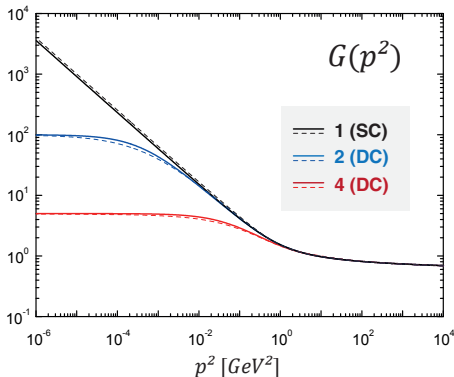
Truncated vertex DSE



Inputs for calculation

- 1 Ghost/gluon propagators \rightarrow use fit functions.
- 2 Ghost-gluon vertex \rightarrow taken to be bare.
A. Cucchieri, A. Maas, T. Mendes, PRD **77**, 2008
- 3 Three-gluon vertex \rightarrow self-consistently back-coupled.

Ghost dressings and gluon 'propagators'



Solid lines: DSE solutions, data by C. Fischer
Dashed lines: Fit functions

Four-gluon vertex

- DSE study of Γ_{4g} , without full tensor/color basis.
C. Kellermann, C. Fischer, PRD **78**, 2008
- We model Γ_{4g} , dressed tree-level structure.

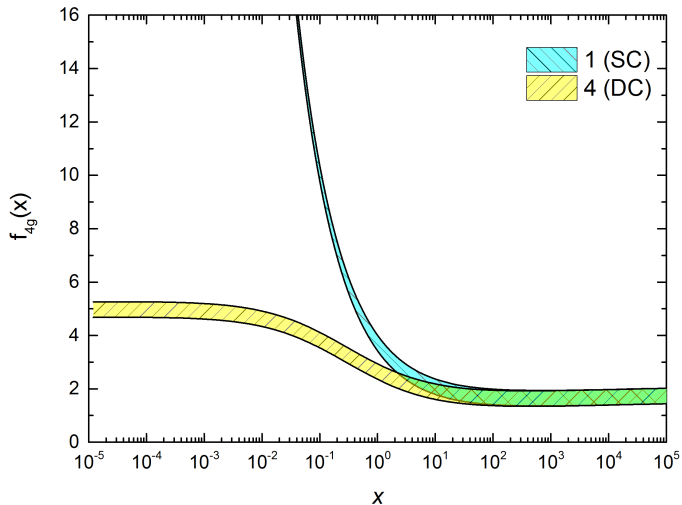
$$\Gamma^{\mu\nu\rho\sigma}(p_1, p_2, p_3, p_4) = f_{4g}(x) \cdot \Gamma_{(0)}^{\mu\nu\rho\sigma} \quad (3)$$

$$x \propto (p_1^2 + p_2^2 + p_3^2 + p_4^2)$$

- $f_{4g}(x)$ is 'correct' in IR and UV.
- Study dependence of results on Γ_{4g} model:

$$f_{4g}(x) \rightarrow f_{4g}(x) + (0 \dots 0.6) \quad (4)$$

Γ_{4g} dressing



Stability of solution

- For f_{4g} too weak, solution is unstable (gluon triangle diverges).
- Different iteration methods do not help.
- Obtained stable solutions by strenghtening f_{4g} .
- Through f_{4g} we account for:
 - 1 Neglected 2-loop terms in vertex DSE.
 - 2 Insufficient knowledge of 'true' four-gluon vertex.

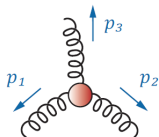
Basis for three-gluon vertex

- 3 Lorentz indices, 2 momenta \Rightarrow 14 tensor structures.
- In Landau gauge 4 tensors describe Γ_{3g} .
- Use orthonormal basis internally.

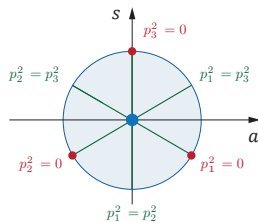
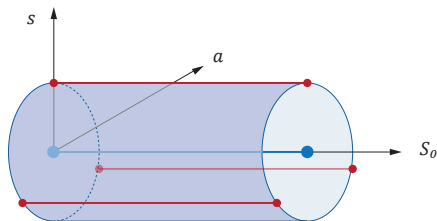
$$\tau_{\mu\nu\rho}^i \cdot \tau_{\mu\nu\rho}^j = \delta^{ij}, \quad i, j = 1 \dots 4 \quad (5)$$

- Rotate onto Bose-symmetric BC-like basis for results.
J. S. Ball, T. W. Chiu, PRD **22**, 2550 (1980)
- DSE calculation of Γ_{3g} with 1 tensor structure:
A. Blum, M. Huber, M. Mitter, L. von Smekal, arXiv:1401.0713

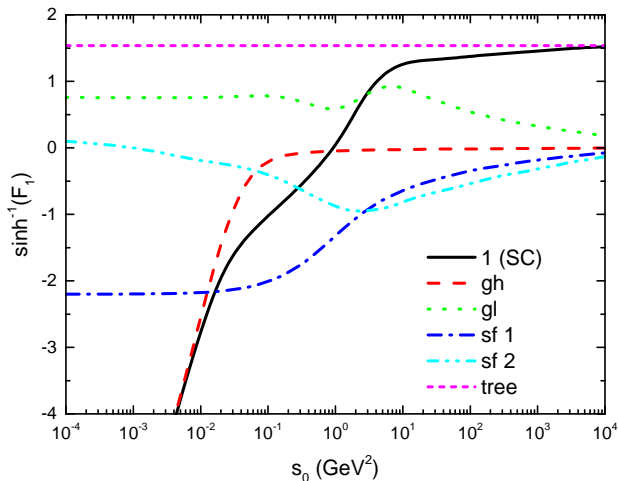
Vertex kinematics



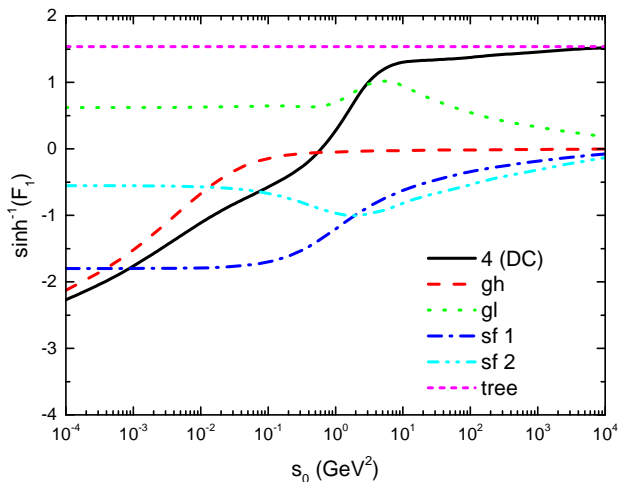
$$S_0 = \frac{1}{6}(p_1^2 + p_2^2 + p_3^2), \quad s = \sqrt{3} \frac{p_2^2 - p_1^2}{p_1^2 + p_2^2 + p_3^2}, \quad a = \frac{p_1^2 + p_2^2 - 2p_3^2}{p_1^2 + p_2^2 + p_3^2} \quad (6)$$



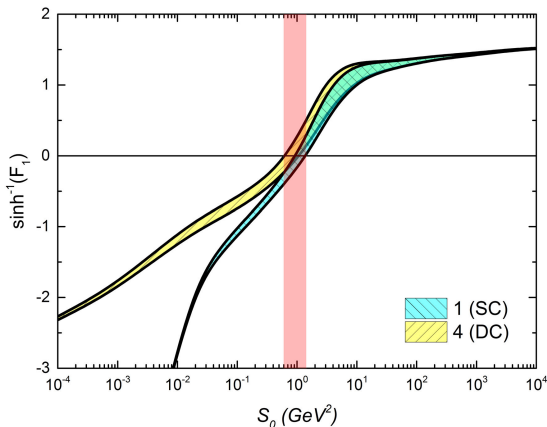
Contributions to $\Gamma_{(0)}^{\mu\nu\rho}$ dressing, scaling



Contributions to $\Gamma_{(0)}^{\mu\nu\rho}$ dressing, decoupling



Zero crossing of $\Gamma_{(0)}^{\mu\nu\rho}$ dressing



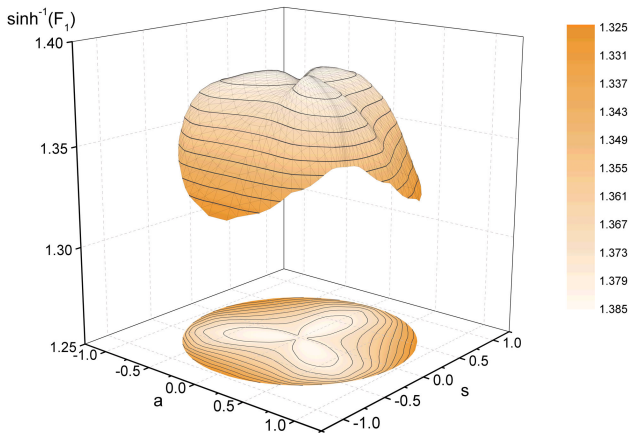
Zero crossing seen in 3D lattice and 4D continuum studies.

A. Blum, M. Huber, M. Mitter, L. von Smekal, arXiv:1401.0713

A. Aguilar, D. Binosi, D. Ibanez, J. Papavassiliou, arXiv:1312.1212v1

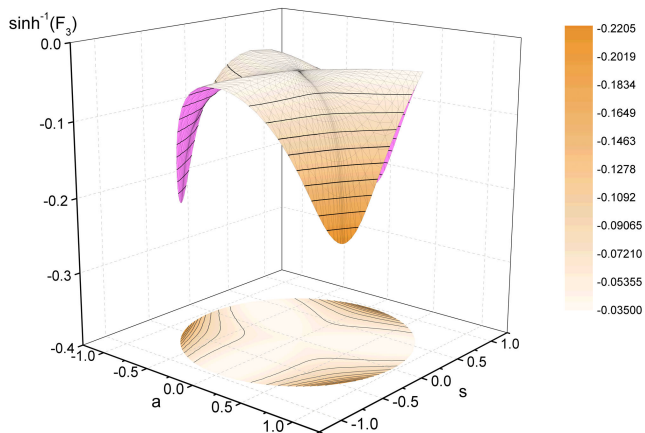
Model and truncation dependent!

Angular dependence of F_1 dressing



Angular dependence of F_1 dressing function at $S_0 = 102 \text{ GeV}^2$.

Angular dependence of F_3 dressing



Angular dependence of F_3 dressing function at $S_0 = 102 \text{ GeV}^2$.

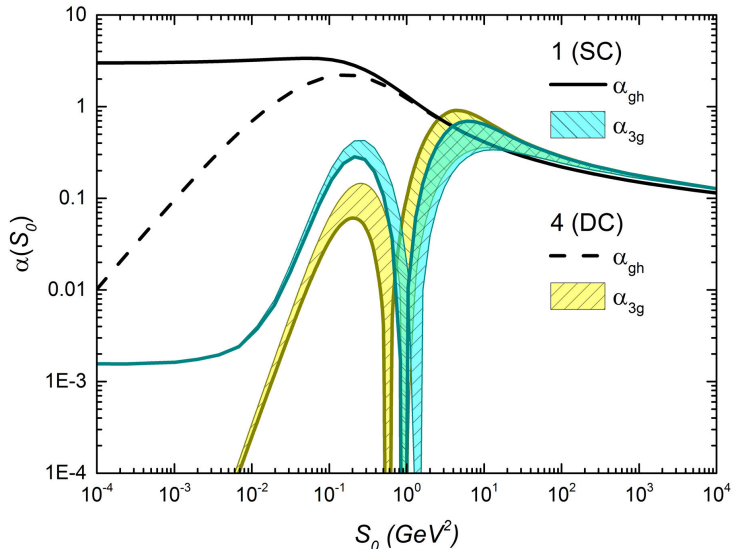
Running couplings

- RG invariant running couplings for different vertices:

$$\alpha_{gh} = \frac{g^2}{4\pi} Z \cdot G^2$$
$$\alpha_{3g}^{(n)} = \frac{g^2}{4\pi} Z \cdot G^2 \left[\frac{Z \cdot F_1}{G} \right]^n \quad (7)$$

- We calculate $\alpha_{3g}^{(n)}$ for $n = 2$.
- In scaling scenario, IR fixed points of running couplings.
R. Alkofer, C. Fischer, F. Llanes-Estrada, PLB **611**, (2005)
- Our result for fixed point: $\alpha_{3g}^{(2)}(0) = 0.0016$.

Results for α_{gh} and α_{3g}



Summary

- 1 First DSE study of Γ_{3g} with full tensor structure.
- 2 Expected all-scales divergences in the IR.
- 3 Soft collinear divergences in subleading components.
- 4 Zero crossing in F_1 at hadronic scales $\Lambda \approx 1$ GeV:
 - Mild $4g$ vertex model dependence.
 - Possible significant dependence on truncation.

Outlook: Applications for Γ_{3g}

- 1 Impact of zero crossing on hadronic observables.
- 2 Meson spectroscopy beyond rainbow ladder:
 - Excited states.
 - Exotic mesons, hybrids and glueballs.
- 3 Irreducible three-body force in baryons:
 - Two- vs. three-quark dominance in excited states.
 - Baryonic hybrids.
- 4 Investigations of (near) conformal window of QCD-like theories.
see also talk by Markus Hopfer

THANK YOU FOR YOUR ATTENTION!