



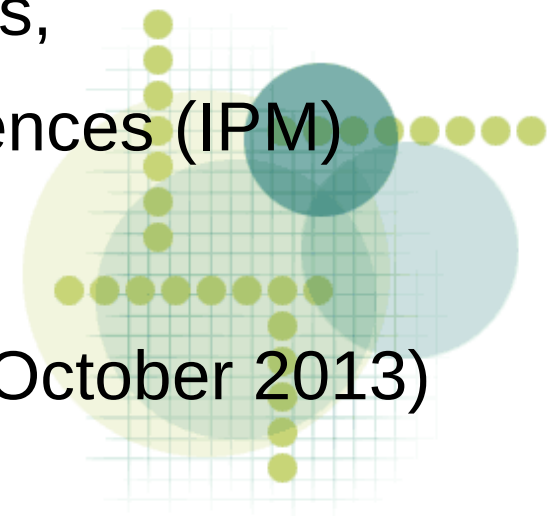
Hadronic Search for SUSY with MT_2 variable

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on behalf of the CMS Collaboration

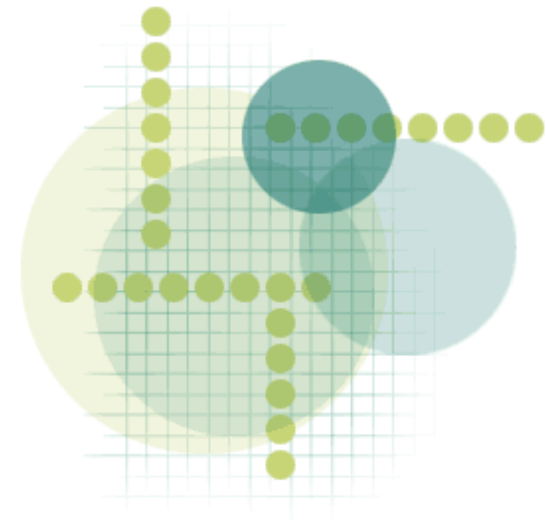
School of Particles & Accelerators,
Institute for Research in Fundamental Sciences (IPM)

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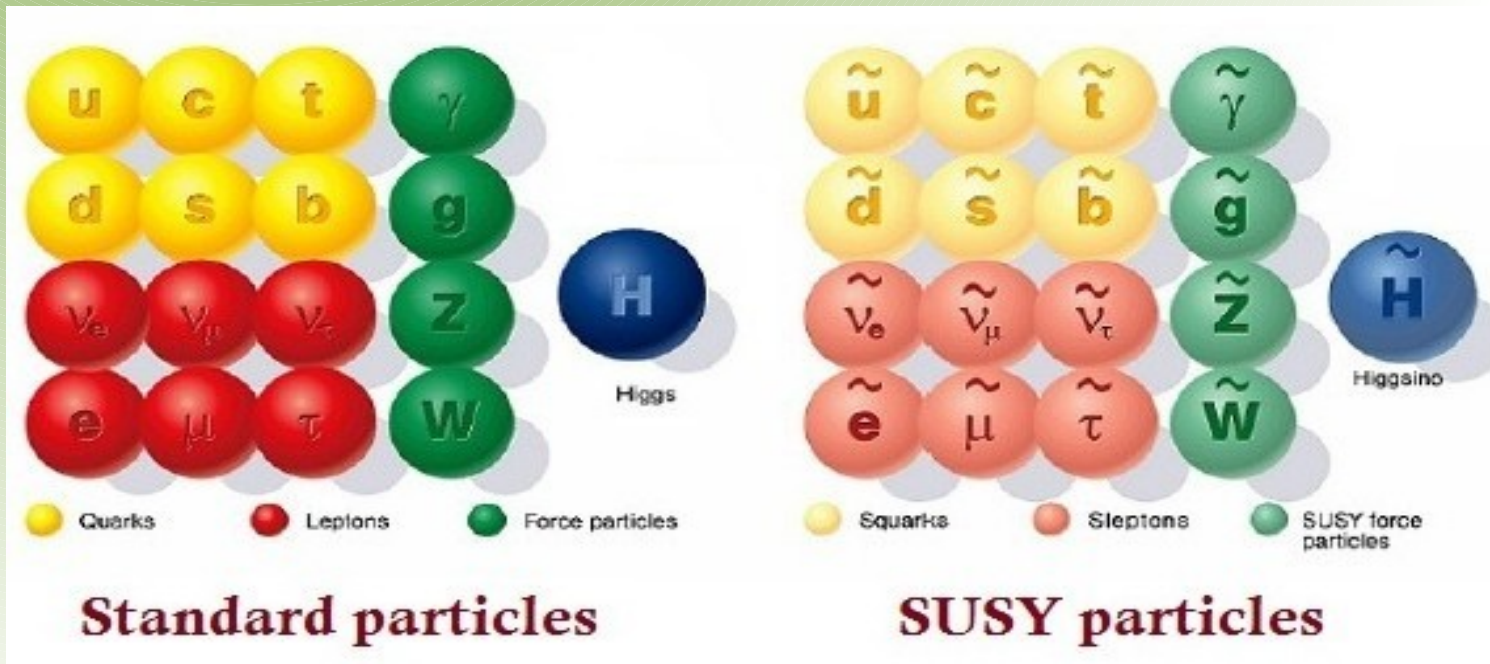
Outline

- Supersymmetry and Its Motivations
- Why Hadronic search? Why $MT2$?
- From MT to $MT2$
- Interpretation of $MT2$
- Search Strategy
- Backgrounds Estimation
- Results
- Conclusion

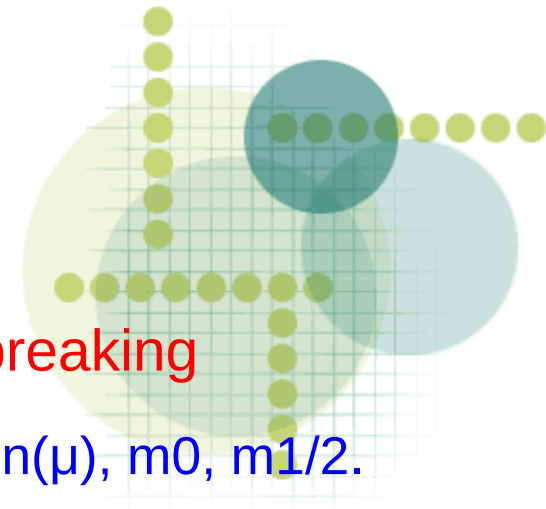


Supersymmetry and Its Motivations

- SuperSYmmetry (SUSY) is a symmetry between fermions \leftrightarrow bosons,
 $Q |\text{boson}\rangle = |\text{fermion}\rangle$, $Q |\text{fermion}\rangle = |\text{boson}\rangle$

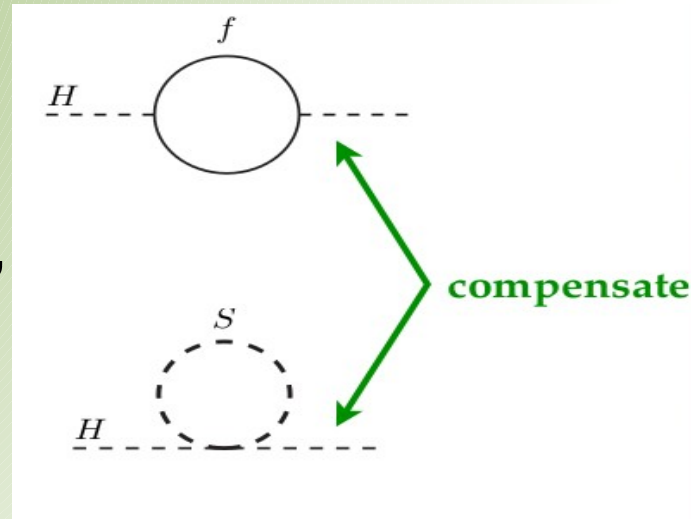


- SUSY: double number of particles (MSSM)
- Spin of SUSY particles differs by 1/2 a unit
- Must be broken, 105 free parameters due to SUSY breaking
- mSUGRA/cMSSM has 5 free parameters: A_0 , $\tan\beta$, $\text{sgn}(\mu)$, m_0 , $m_{1/2}$.

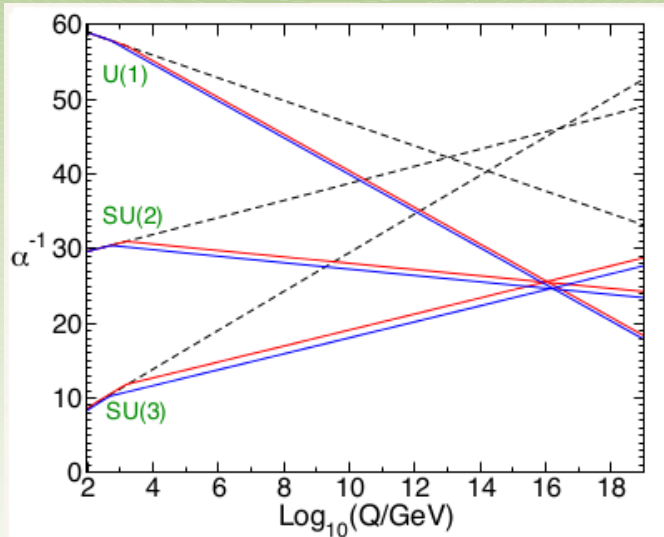


Supersymmetry and Its Motivations

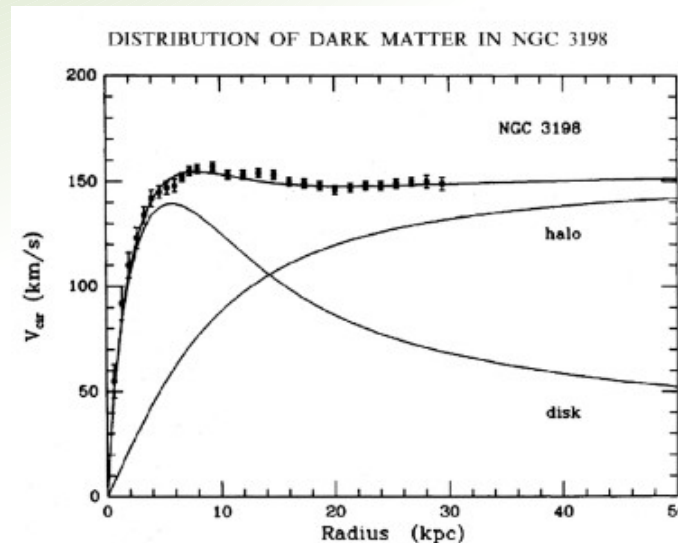
- **SUSY** is very interesting, as it:
 1. gives a
 “**Solution to the Hierarchy Problem**”



2. facilitates a
 “**Gauge Couplings Unification**”



3. provides a
 “**Dark Matter candidate**”



Astrophysical Journal,
 vol. 295, Aug. 15, 1985,
 p. 305-313.

Why Hadronic search? Why MT2?

If **SUSY**, **conserving R-parity**, comes with a Dark Matter candidate:

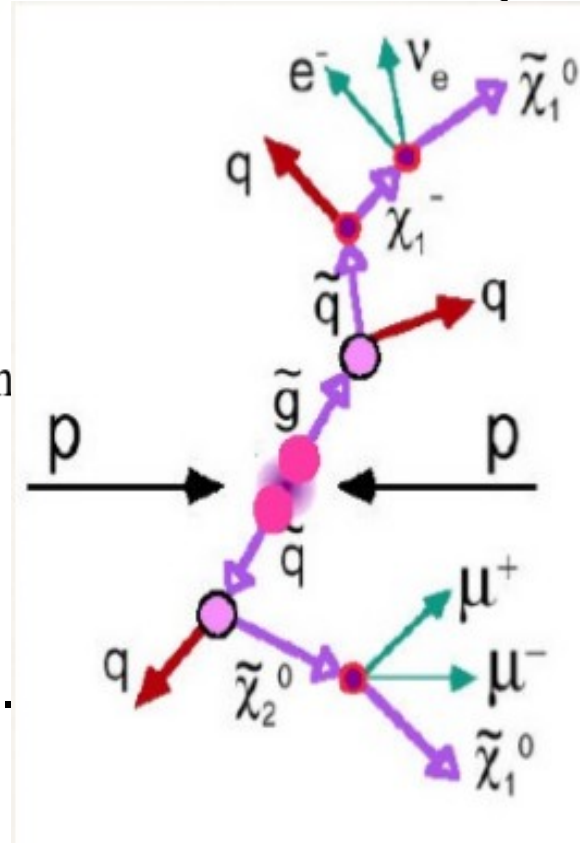
- **SUSY** particles must be **produced in pairs** and the **lightest sparticle (LSP)** as a **Dark matter candidate** is **stable**.
- **Colored SUSY** particles cascade down to **LSP** with **emission of jets** and sometimes leptons, therefore in search for SUSY, **MET** and **HT** are very useful variables:

$$H_T = \sum_{jets} |\vec{p}_T|$$

$$E_T^{miss} = \left| - \sum_{particles} \vec{p}_T \right|$$

← Hadronic activity
← Missing transverse momentum

- Classical **SUSY** searches are based on these two variables but it can be defined **a clever kinematical variable** reflecting these properties which is **"MT2"**.



From MT to MT2

In $W \rightarrow e\nu$ decay, **transverse mass MT** has an endpoint at the true W-mass:

$$m_W^2 = m_l^2 + m_\nu^2 + 2(E_T^l E_T^\nu \cosh \Delta\eta - \mathbf{p}_T^l \cdot \mathbf{p}_T^\nu) \geq$$

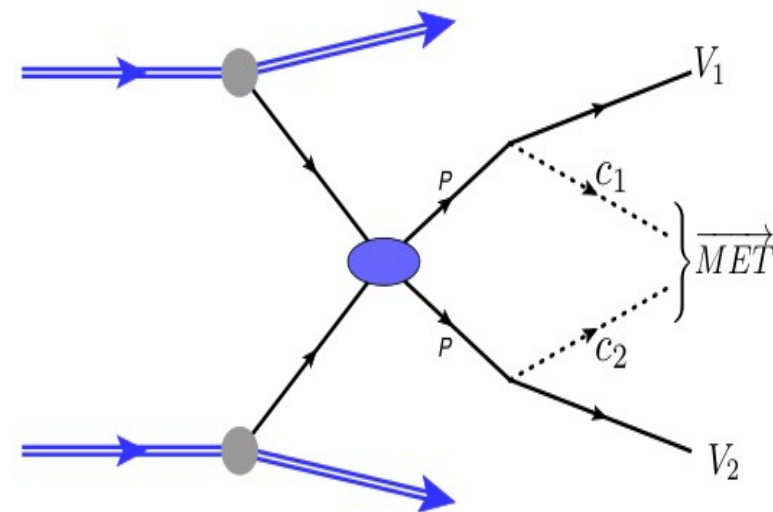
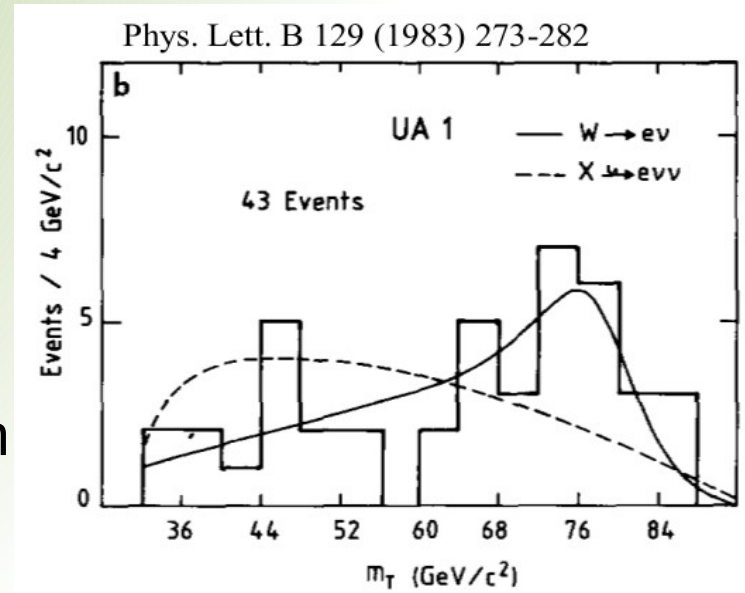
$$m_T^2 = m_l^2 + m_\nu^2 + 2(E_T^l E_T^\nu - \mathbf{p}_T^l \cdot \mathbf{p}_T^\nu)$$

At the LHC, assuming R-parity conservation, **SUSY** events give rise to two decay chains with an unobserved child (c1 and c2) at each end.

The **“stransverse” mass MT2**: extension of MT for the **SUSY** case of two unobserved particles:

$$M_{T2}(m_c) = \min_{\mathbf{p}_T^{c(1)} + \mathbf{p}_T^{c(2)} = \mathbf{p}_T^{miss}} \left[\max \left(m_T^{(1)}, m_T^{(2)} \right) \right]$$

If M_c were known, the endpoint of MT2 would correspond to the parent mass M_p .



Interpretation of MT2

- In case of **no initial state radiation (ISR)** and **zero masses**:

$$M_{T2}^2 = 2 p_T^{(1)} p_T^{(2)} (1 + \cos \phi_{1,2})$$

which $p_T(i)$ is the transverse momenta of the visible systems.

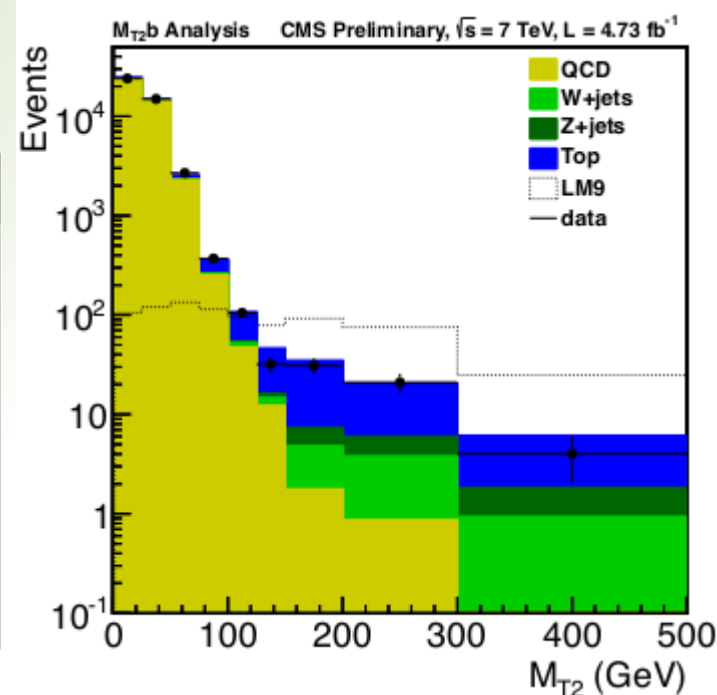
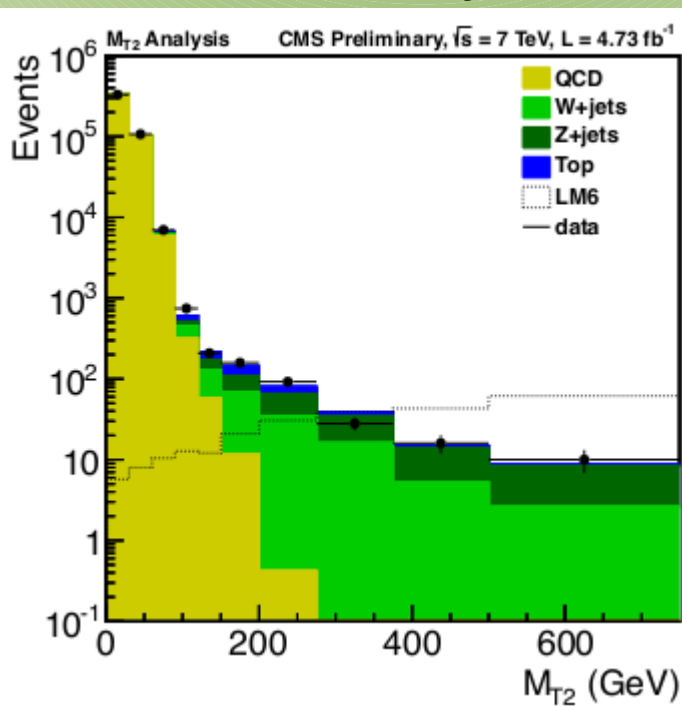
- MT2 = MET** for symmetric systems, $p_T(1)=p_T(2)$

$$(E_T^{\text{miss}})^2 = (p_T^{(1)} - p_T^{(2)})^2 + 2p_T^{(1)}p_T^{(2)} (1 + \cos \phi_{12})$$

- MT2 = 0** GeV for **back-to-back systems**, like QCD di-jets events.
- MT2 is similar to MET in signal region, but **more robust against jet energy mismeasurements than MET**. Therefore QCD multijet events accumulated at low MT2.
Events with real MET can have large values in MT2.
- Multijet events are divided into a 2 pseudo-jets topology using a **hemisphere algorithm**.

Search Strategy

- Search in fully hadronic final states with 2011 pp collision data at 7 TeV collected by the CMS detector, corresponding to 4.73 fb⁻¹.



Baseline selection

$$HT \geq 750 \text{ GeV}$$

- electron and muon veto
- filters against detector noise

MT2 analysis

- At least 3 jets
- $\min\Delta\phi(\text{jets}, \text{MET}) > 0.3$
- $MT_2 > 150 \text{ GeV}$

sensitive to high/medium squark and gluino masses

sensitive to light gluinos with heavy squarks

MT2b analysis

- At least 4 jets
- At least one jet b-tagged
- $\min\Delta\phi(\text{j1-j4}, \text{MET}) > 0.3$
- $MT_2 > 125 \text{ GeV}$

Backgrounds Estimation

- **QCD multijets** have no genuine MET → **small MT_2** . MC predicts that the signal regions are QCD free.
- Mismeasured jets can lead to larger MT_2 . These jets are aligned with MET. Use **correlation of $\min\Delta\phi(\text{jets}, \text{MET})$ and MT_2** to conservatively estimate the QCD contamination in the signal region from data.

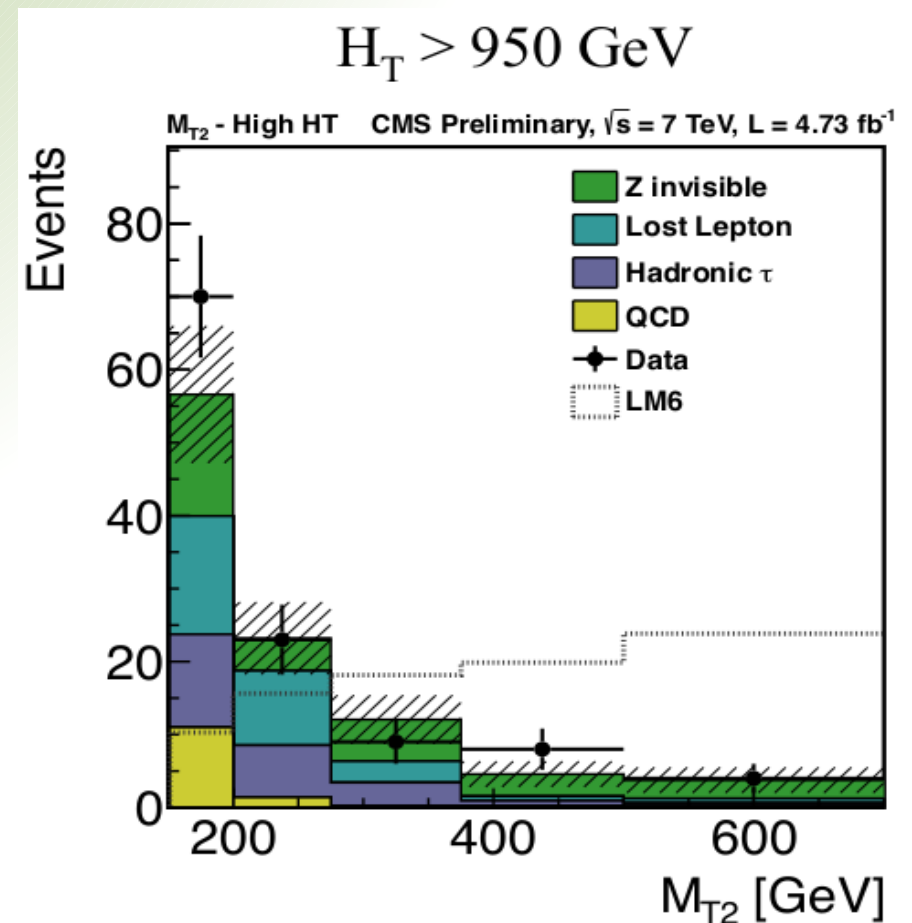
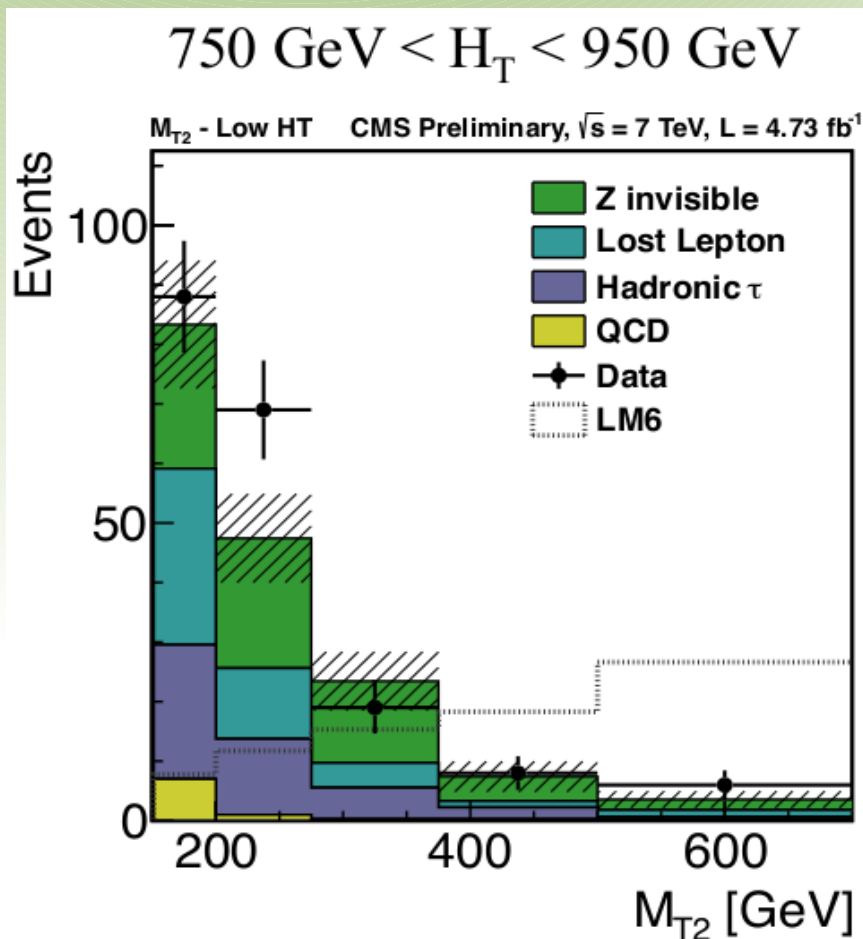
- **Z bosons decaying into neutrinos** are signal like with large MET.
- This background is predicted by using **$W(\rightarrow \mu\nu)+\text{jets}$ and $\text{photon} + \text{jets}$** events.
- In both cases the visible vector boson p_T is added to the MET to mimic $Z(\rightarrow \nu\nu)+\text{jets}$ events.

- **Leptonic $W+\text{jets}$ and $\text{Top}+\text{jets}$** events have real MET. Largely reduced due to lepton veto.
- Enter to signal region if the charged lepton is not reconstructed or out of acceptance (= is lost).
- Remaining (**lost lepton**) background estimated in EWK control region from the number of events with a found lepton in data and corrected for the probability to lose a lepton.

Robust estimation of the SM backgrounds contribution to all signal bins.

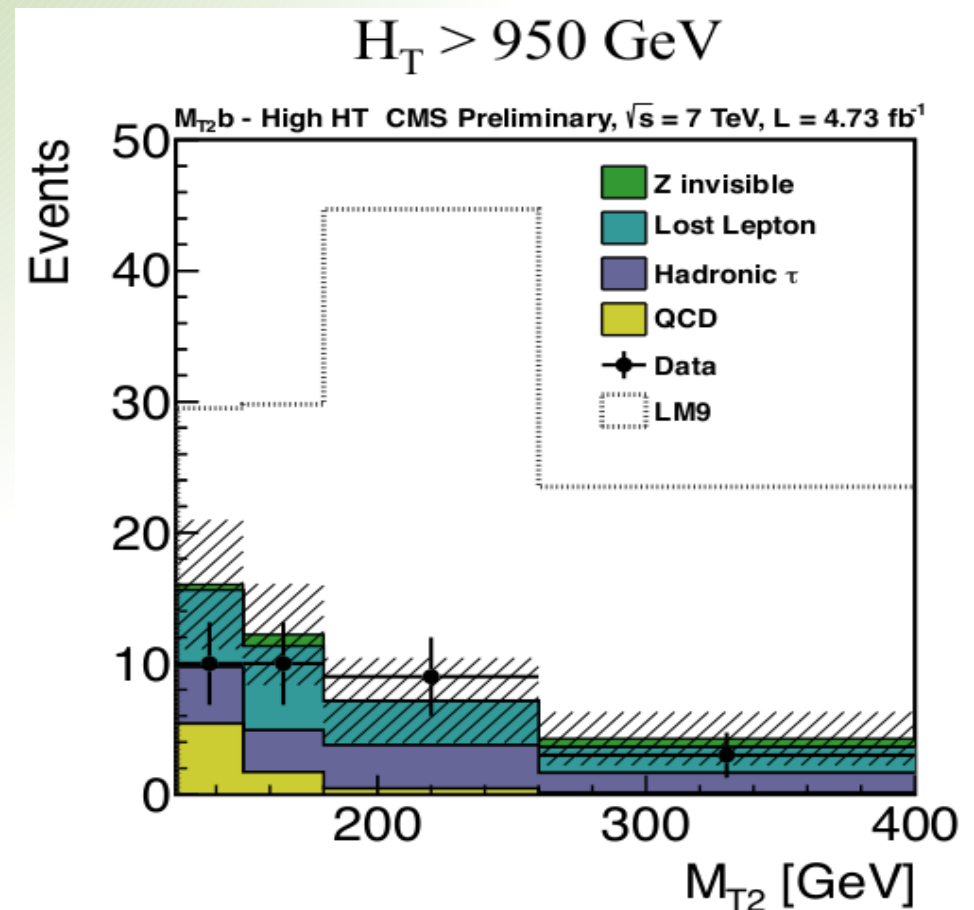
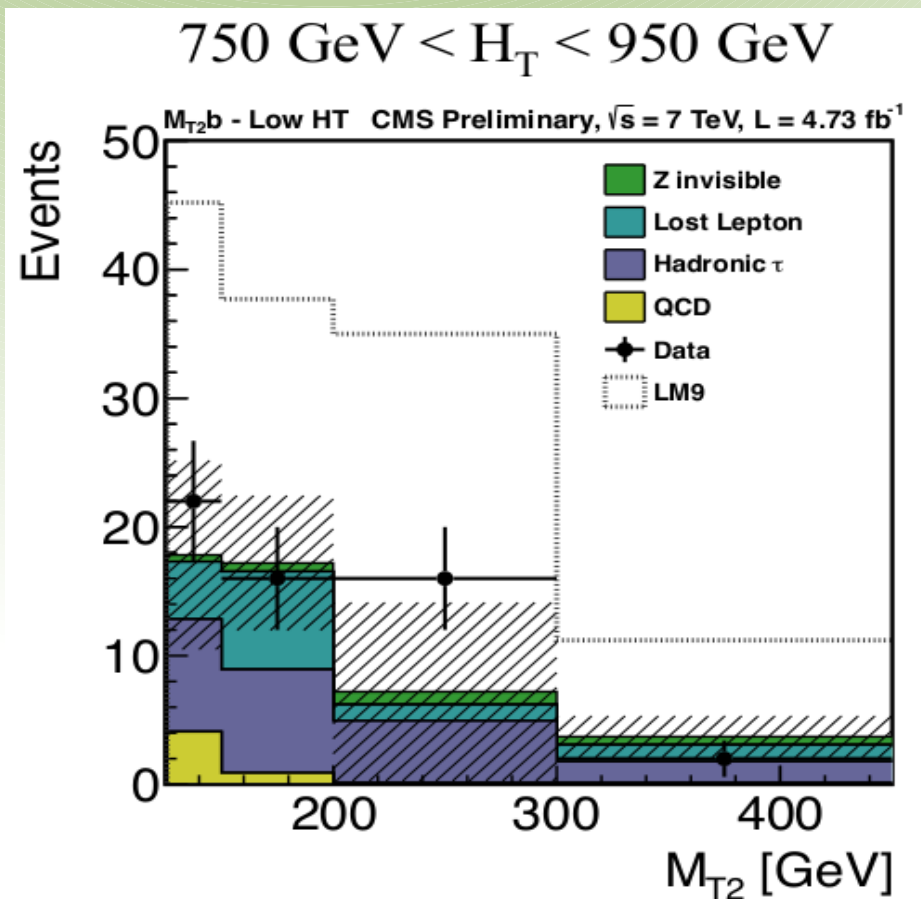
Backgrounds Estimation

- The backgrounds estimation for the MT2 analysis are summarized here.
- Shaded region is uncertainty on the background estimation.
- The observed data are shown.
- A possible SUSY signal is overlaid to show sensitivity of search region.



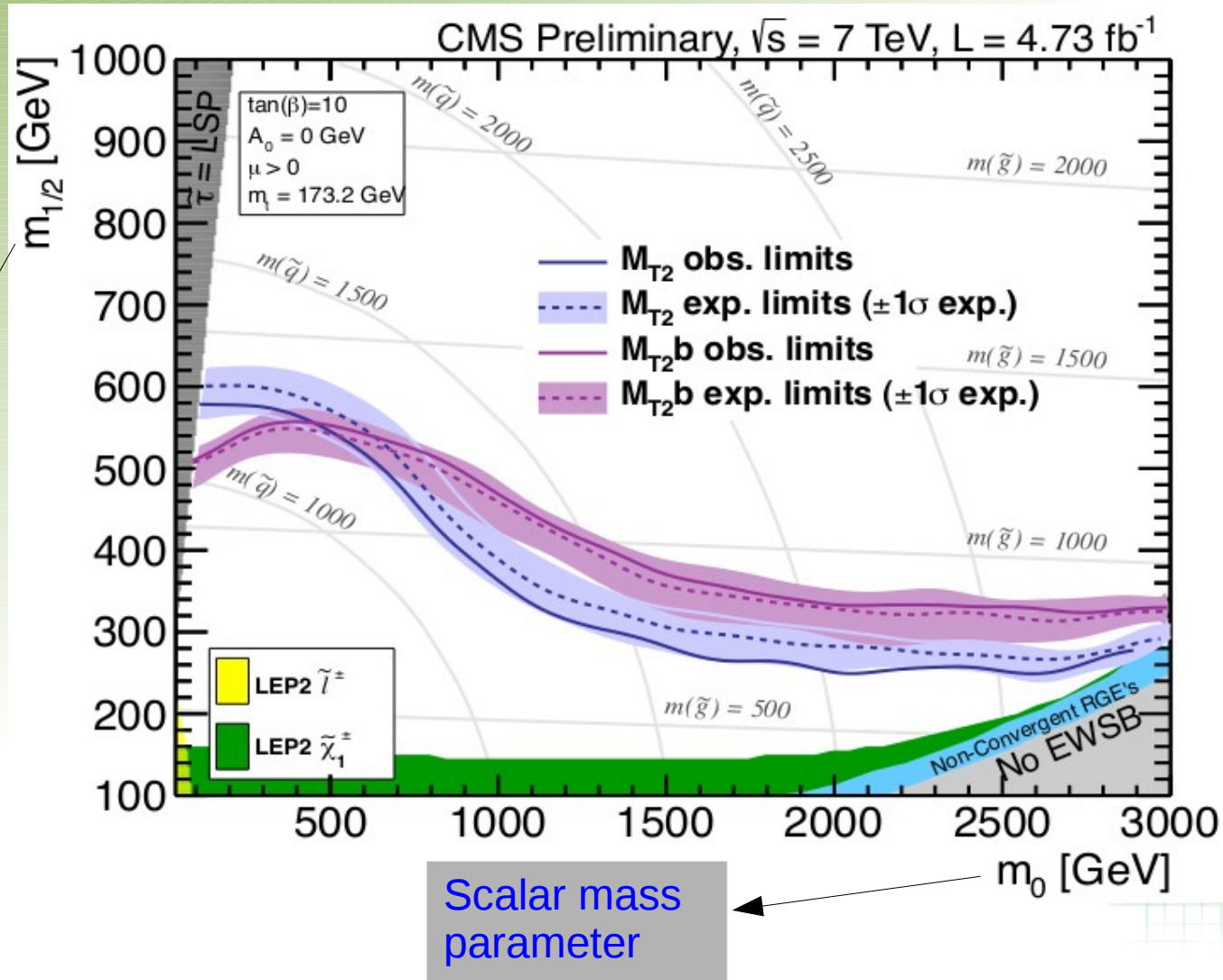
Backgrounds Estimation

- The backgrounds estimation for the MT2b analysis are summarized here.
- Shaded region is uncertainty on the background estimation.
- The observed data are shown.
- A possible SUSY signal is overlaid to show sensitivity of search region.



Results

- The results are interpreted in a full SUSY model constrained to five parameters (mSUGRA/cMSSM).
- In the plane below three of those parameters are fixed: $A_0 = 0$, $\tan\beta = 10$, $\mu > 0$

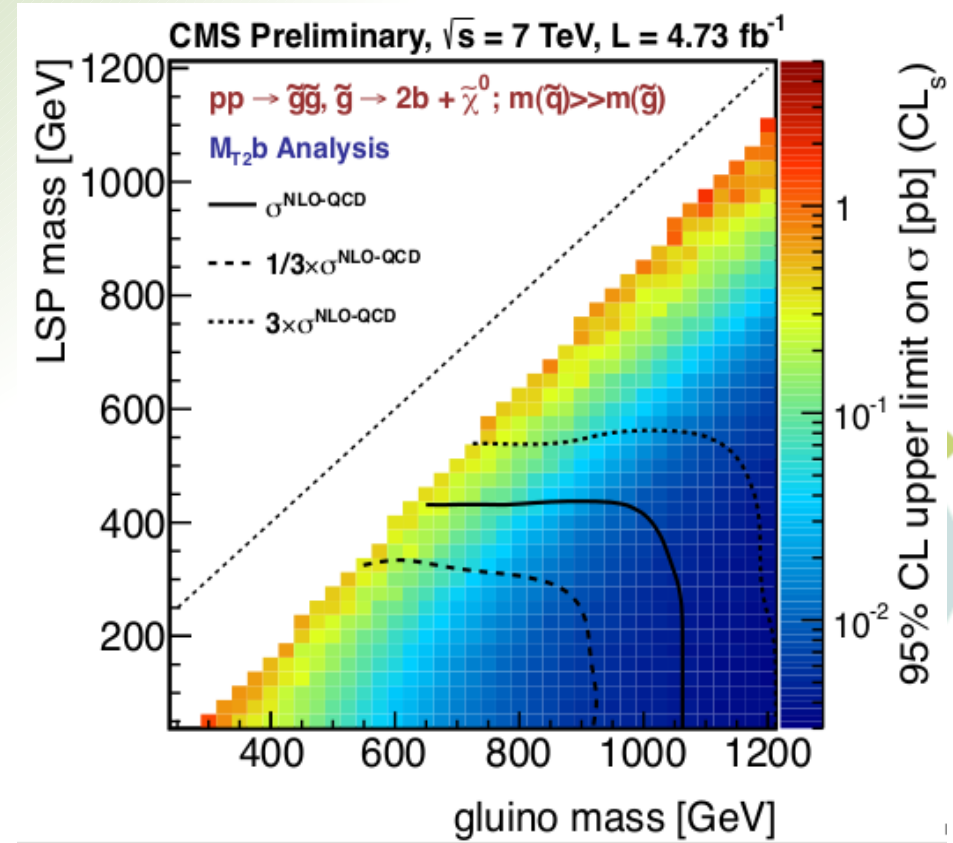
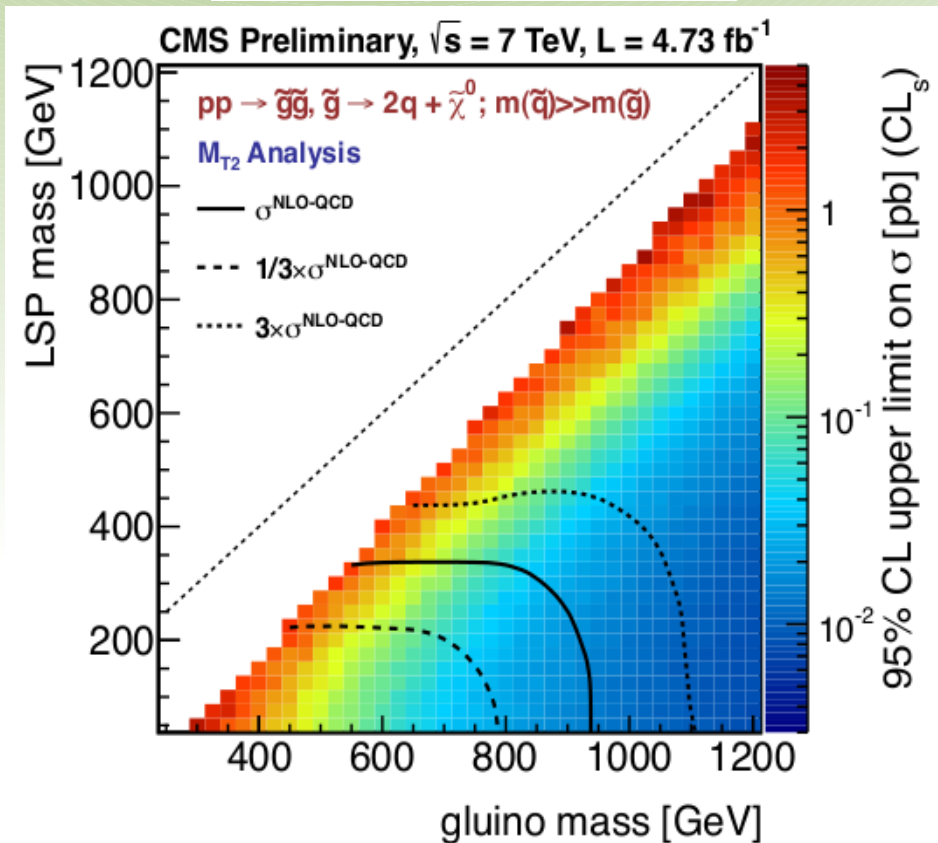
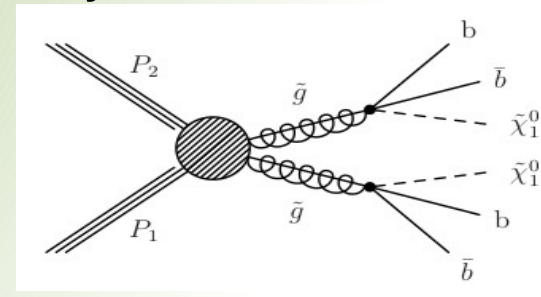
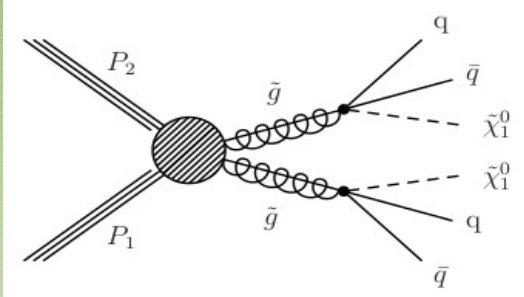


Gaugino mass parameter

Scalar mass parameter

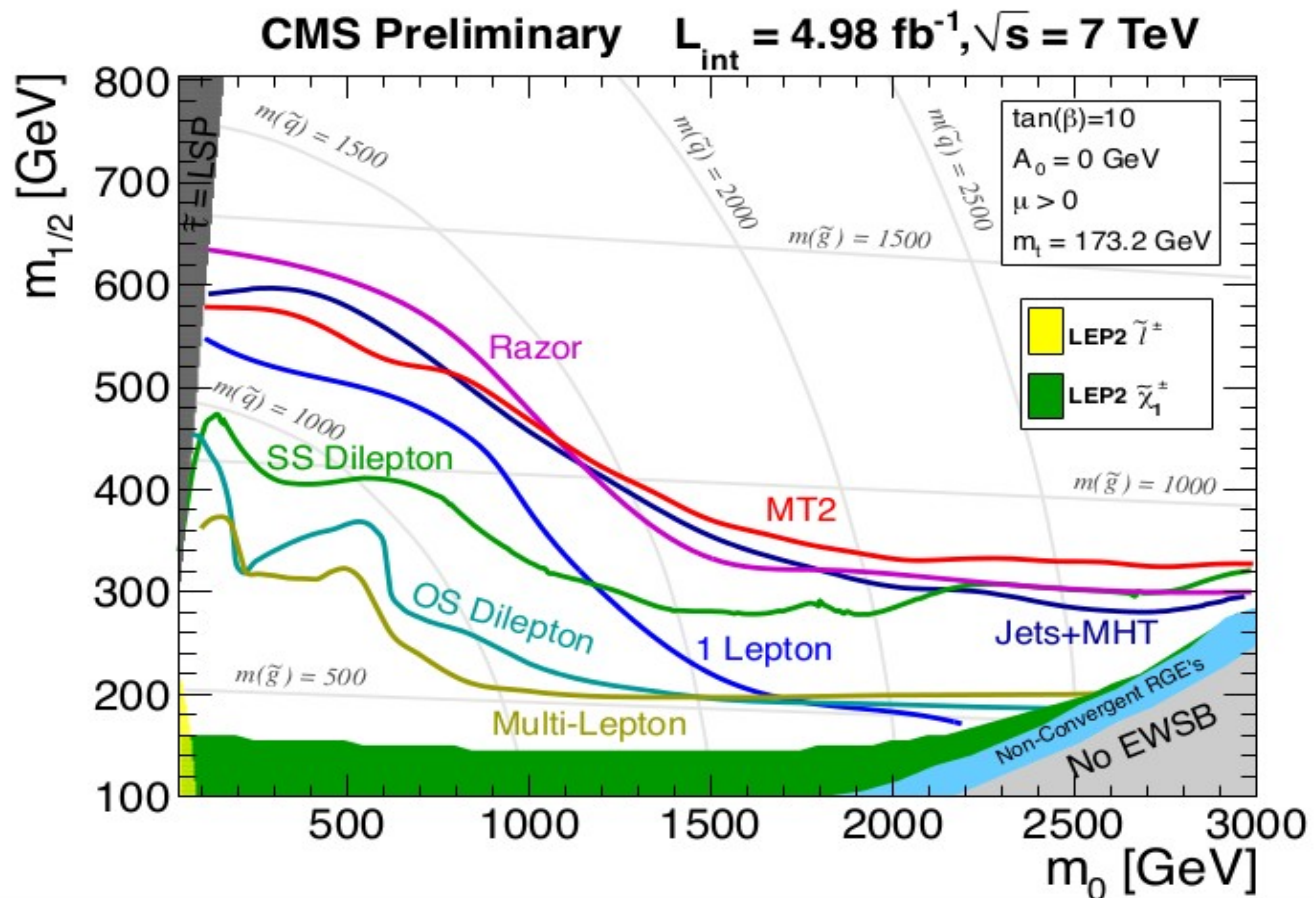
Results

- The analyses are also interpreted in simplified models.
- Models are reduced to one SUSY decay chain only.



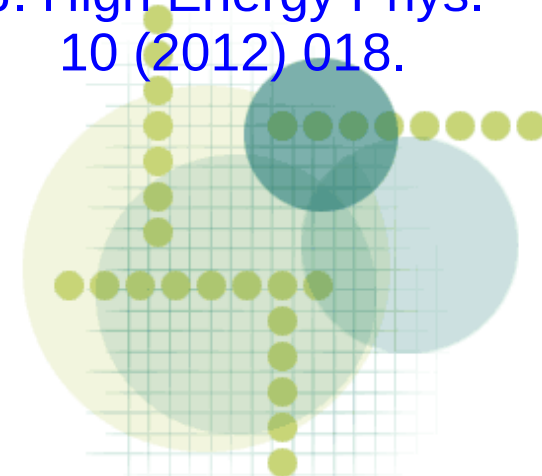
Conclusion

- A search for supersymmetry in fully hadronic final states with 2011 pp collision data collected by the CMS detector has been performed.
- **No excess over the SM predicted background** has been found.
- Limits in various signal model spaces have been set. Absolute mass limits in the mSUGRA/cMSSM scenario for $A_0 = 0$, $\tan\beta = 10$, $\mu > 0$ are found to be $m(\text{squark}) > 1160 \text{ GeV}$ and $m(\text{gluino}) > 860 \text{ GeV}$, and $m(\text{squark}) = m(\text{gluino}) > 1200 \text{ GeV}$ assuming equal squark and gluino masses.



For more details:

CMS-PAS SUS-12-002;
arXiv:1207.1798;
Published in :
J. High Energy Phys.
10 (2012) 018.



Thank you.



"One day, all of these will be supersymmetric phenomenology papers."