

TOWARD QUENCH-FREE PERFORMANCE IN SRF CAVITIES

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&
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Outline of the talk

- **What we actually know** about quench
- **Thin film** cavities (almost) never quench
- Thinking about the **impossible** is cost-free
- Need to go deeper into **Cryogenics**

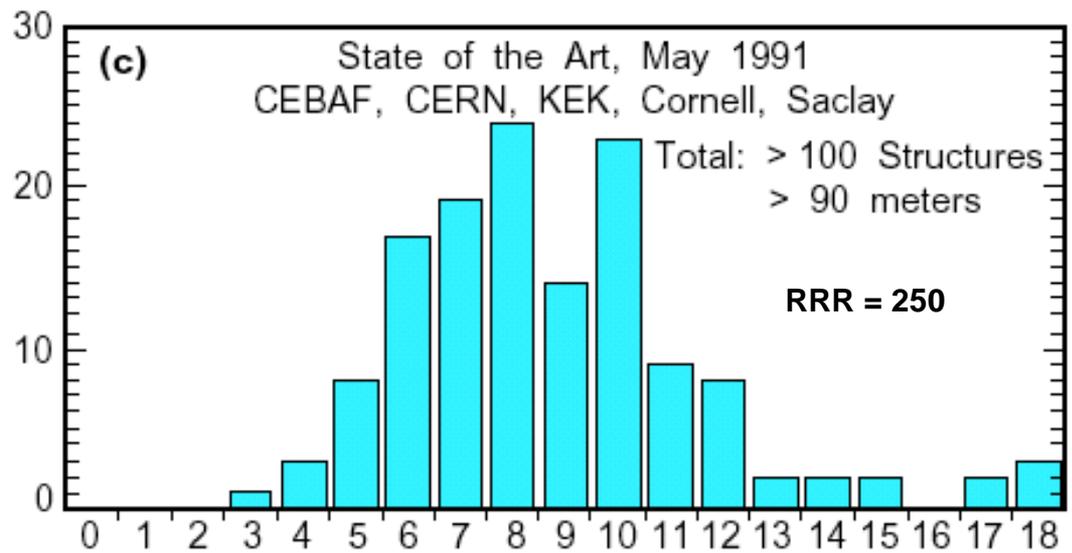
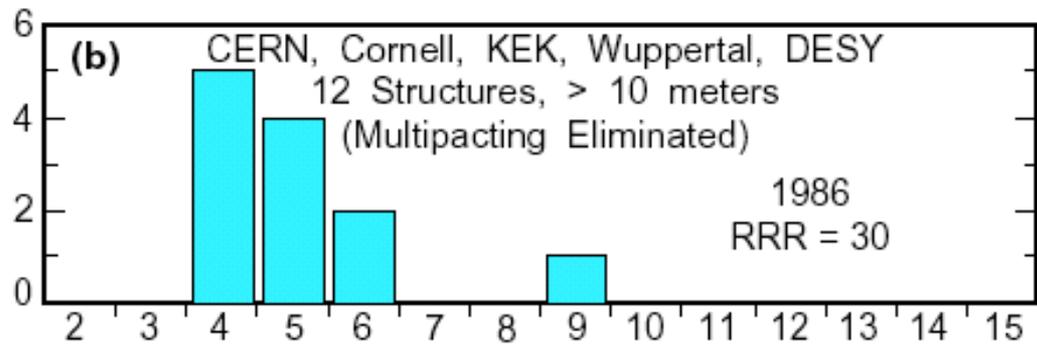
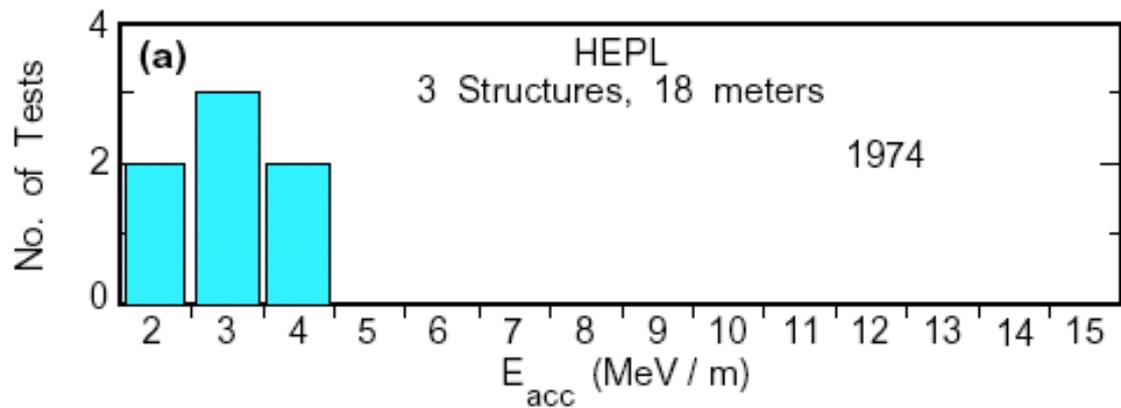
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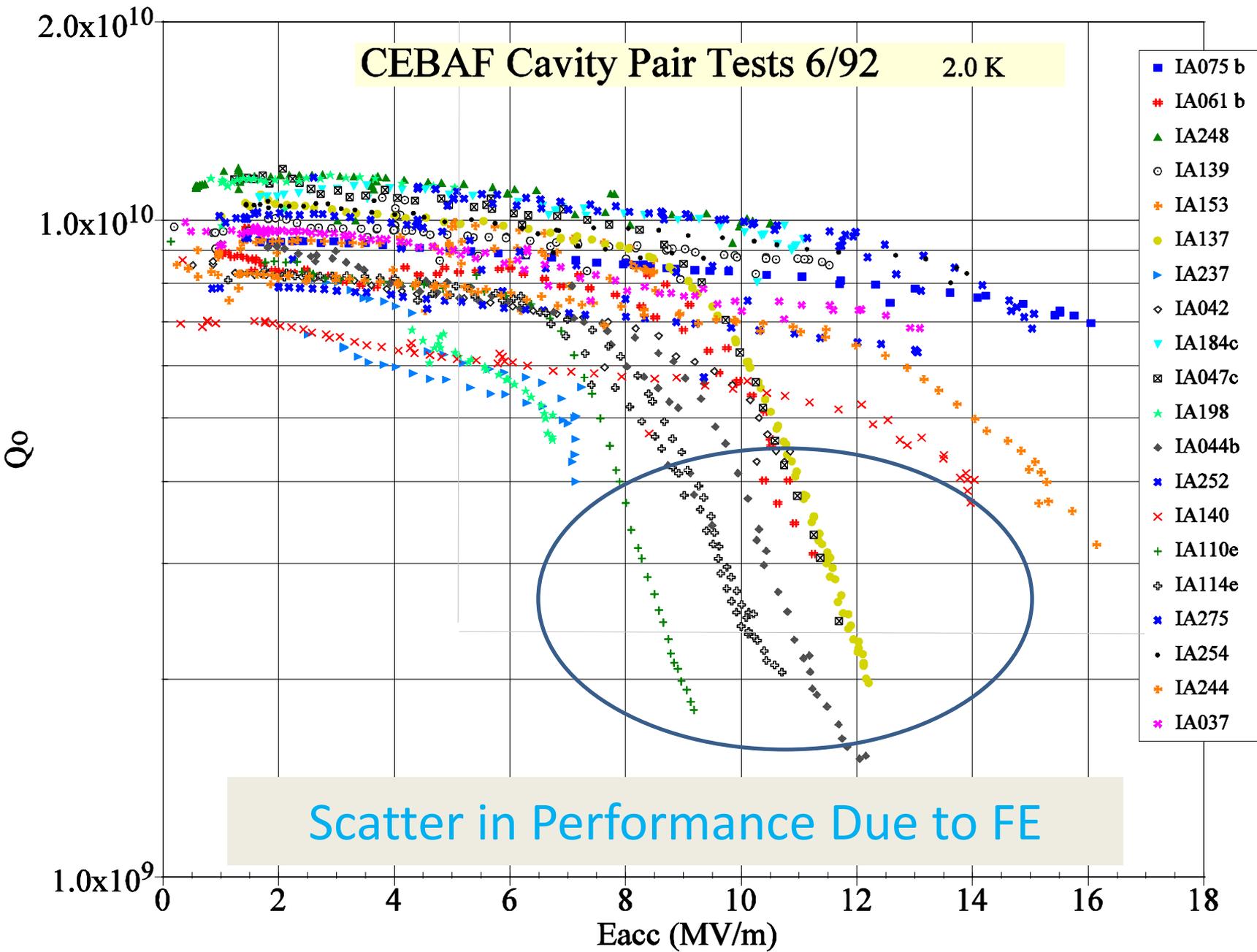
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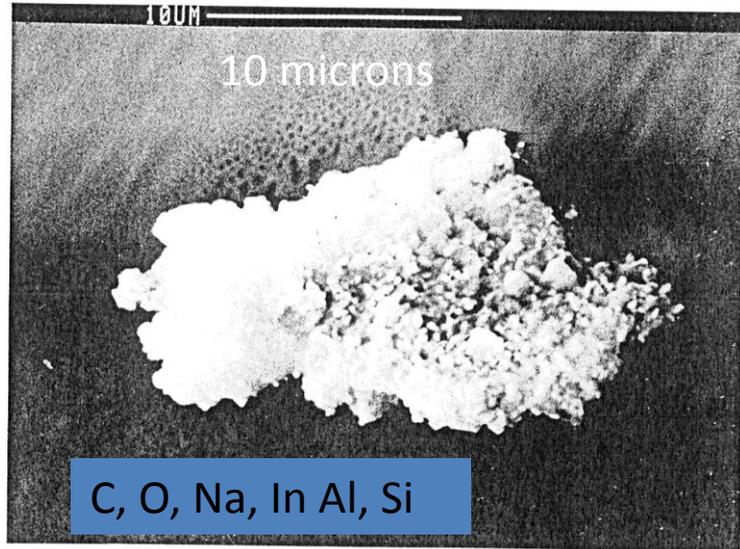
Multipacting

Thermal Breakdown

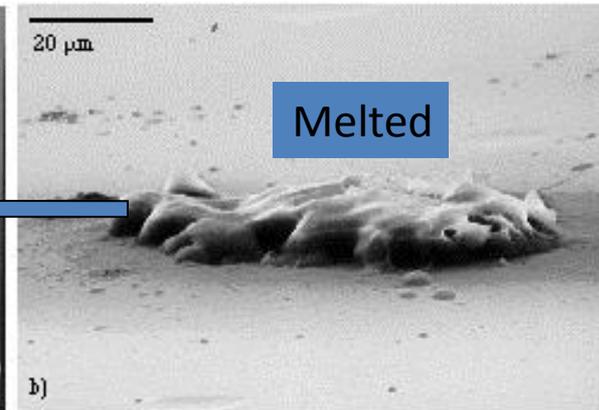
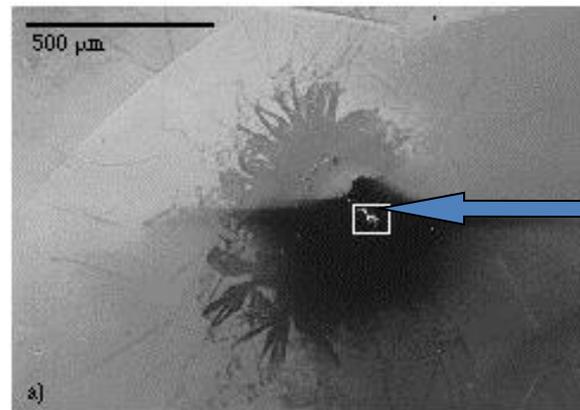
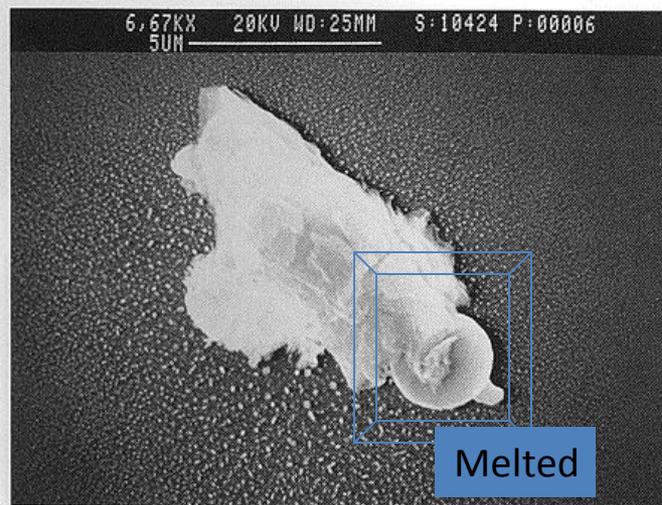
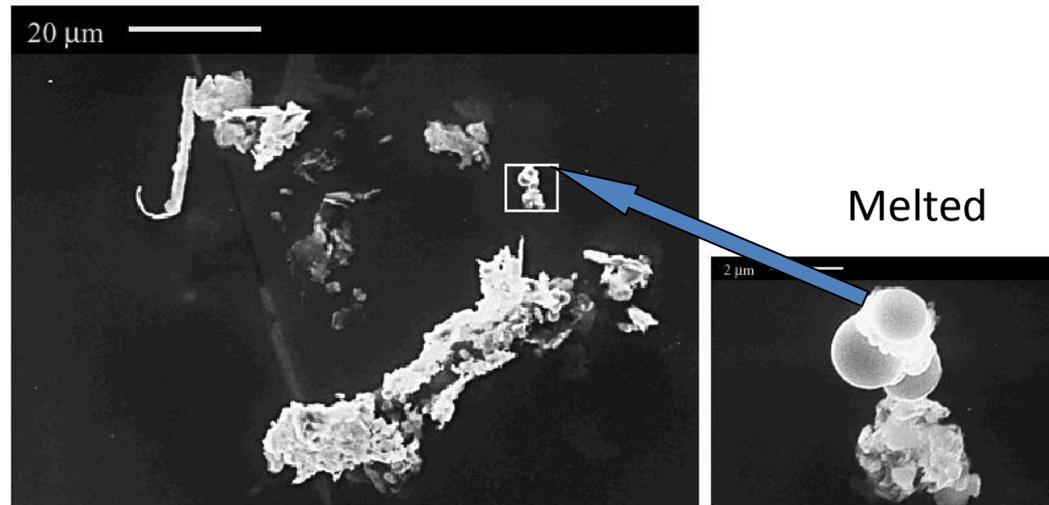
Field Emission







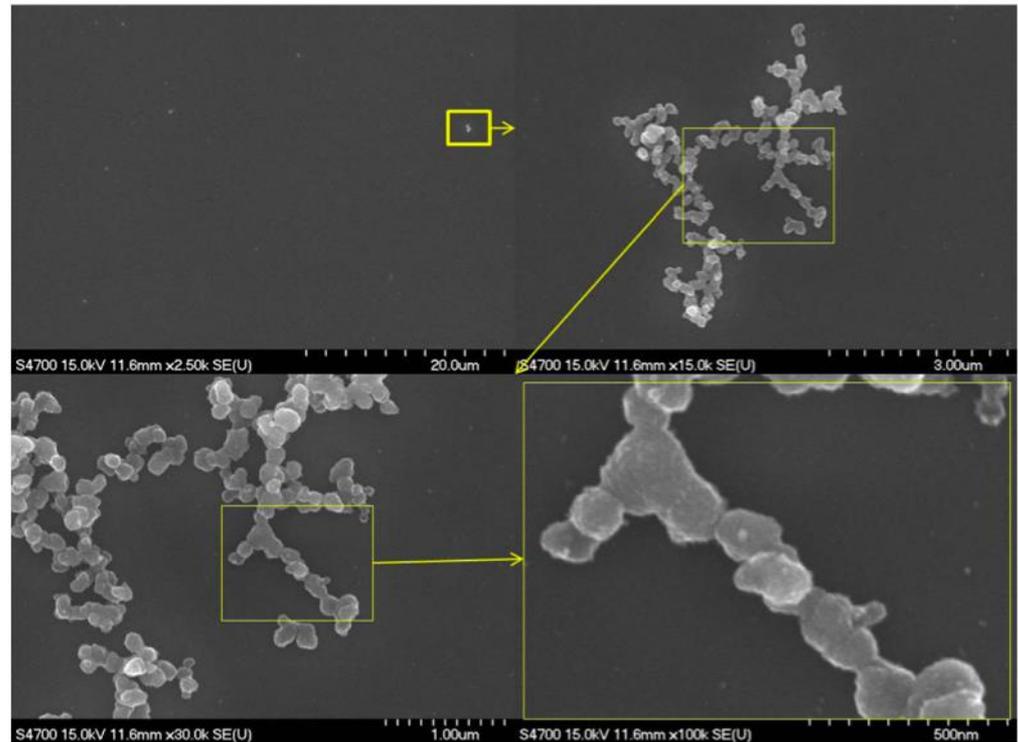
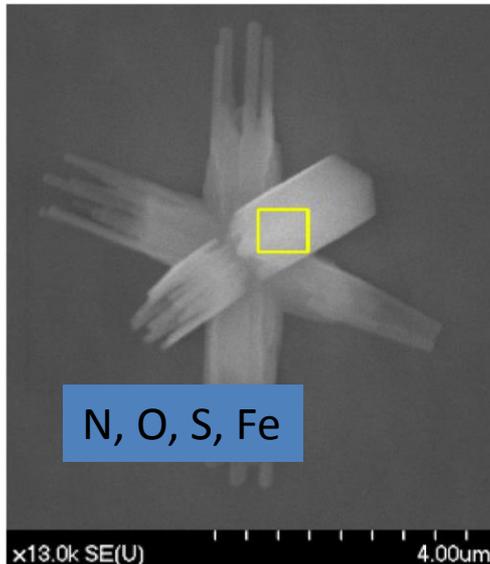
Stainless steel



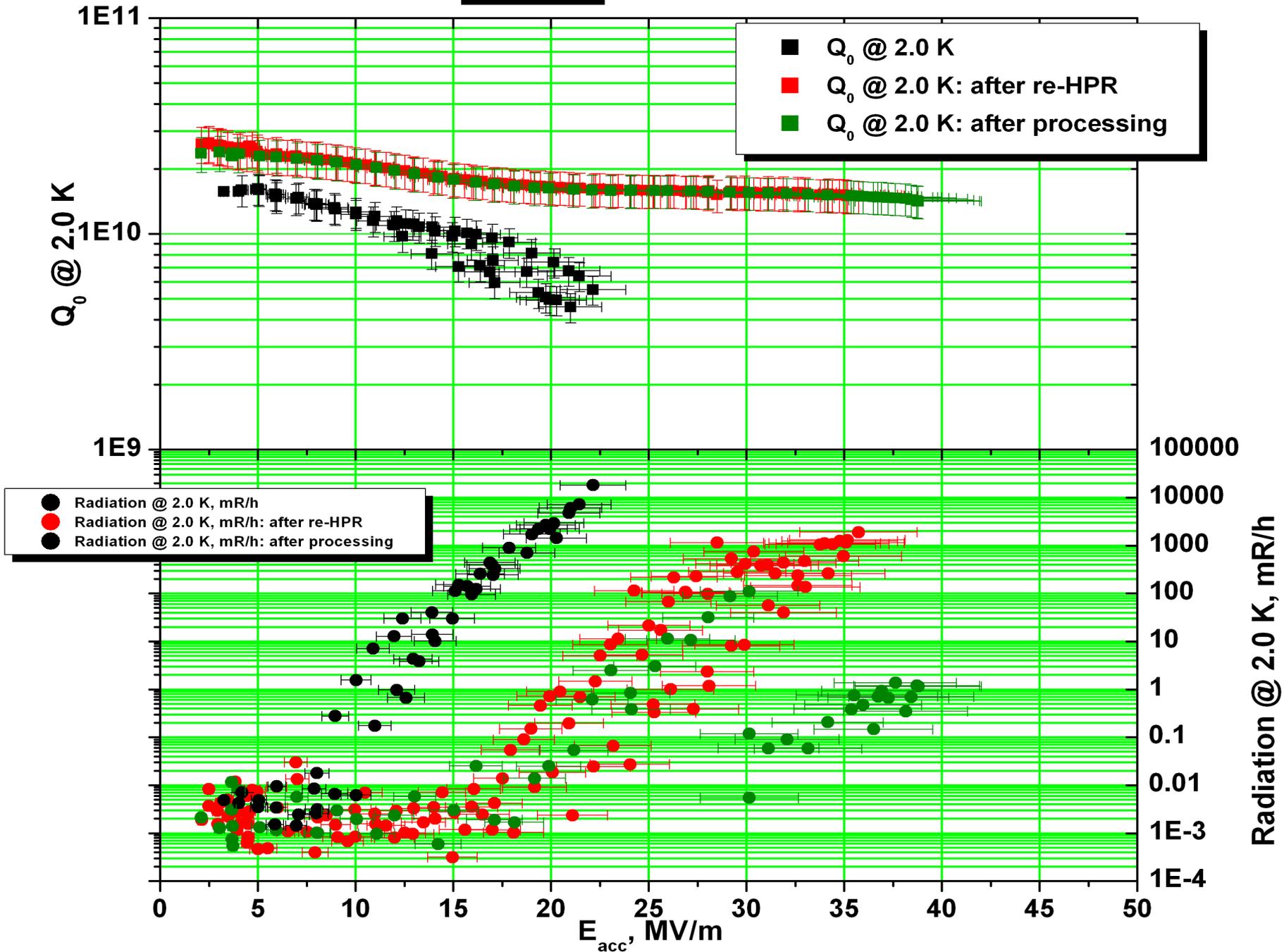
Emitters Found in Niobium SC Cavities

Field Emitters

- Contaminants from surface processing
 - Niobium oxide granules (electropolished surface)
 - Sulfur
 - And other elements

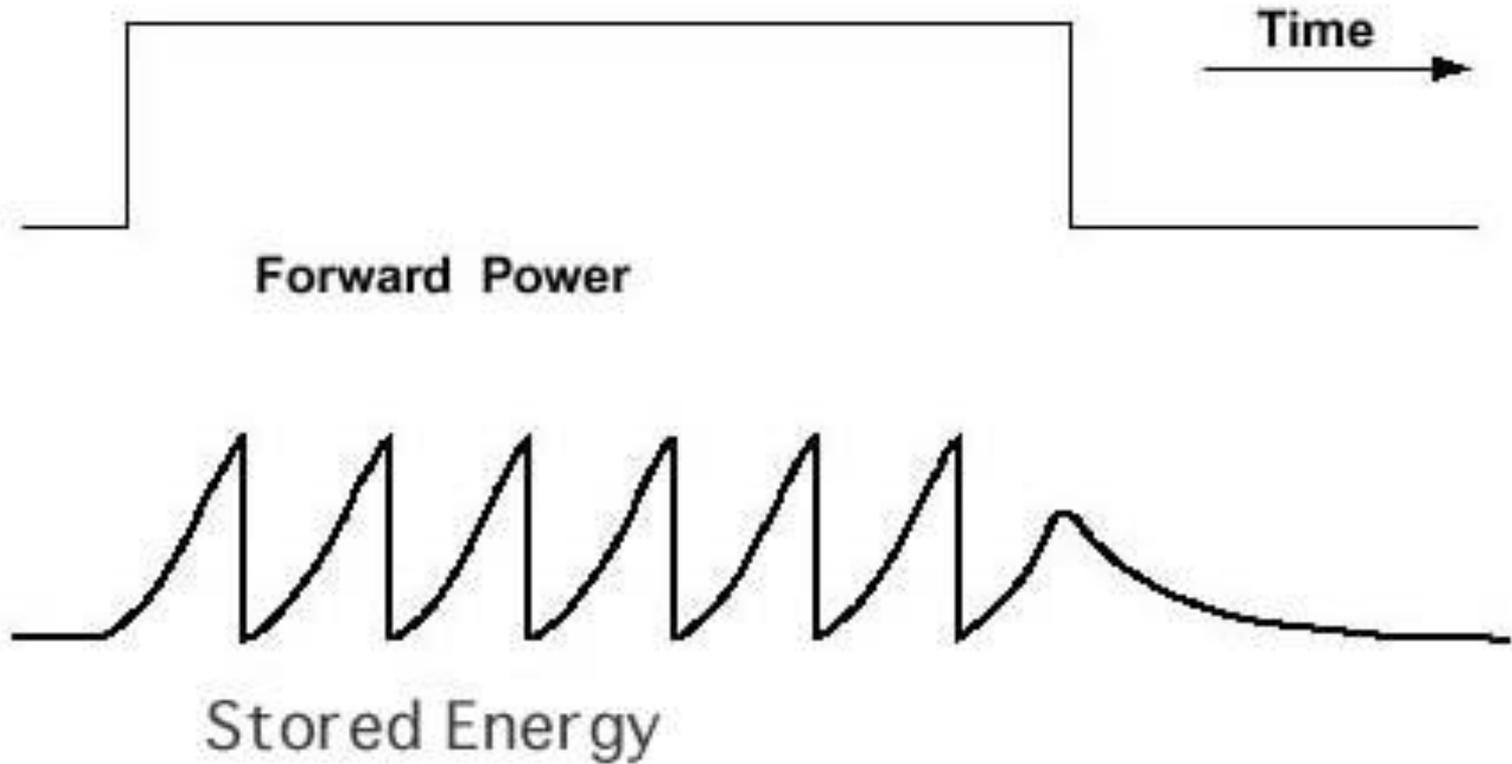


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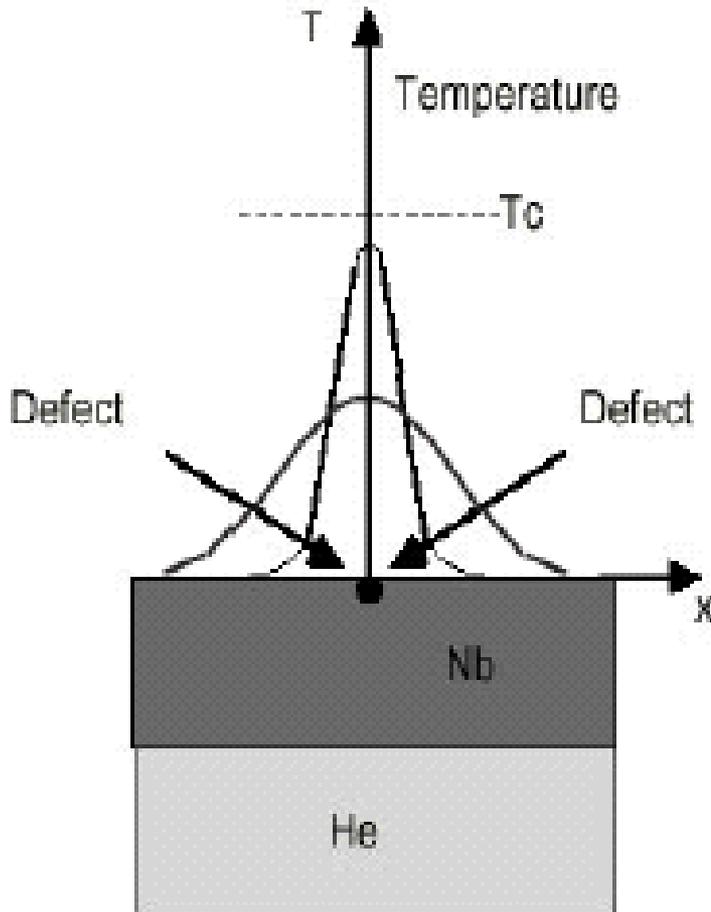
Thermal Breakdown

(Very strong heating observed over extended areas.)



Theory of thermal breakdown of NC Defect

A normal conducting spot triggers quench, when it heats the Nb above T_c .



Breakdown field given by (very approximately):

$$H_{tb} = \sqrt{\frac{4\kappa_T (T_c - T_b)}{r_d R_d}}$$

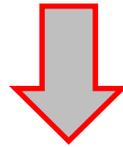
κ_T : Thermal conductivity of Nb

R_d : Defect surface resistance

T_c : Critical temperature of Nb

T_b : Bath temperature

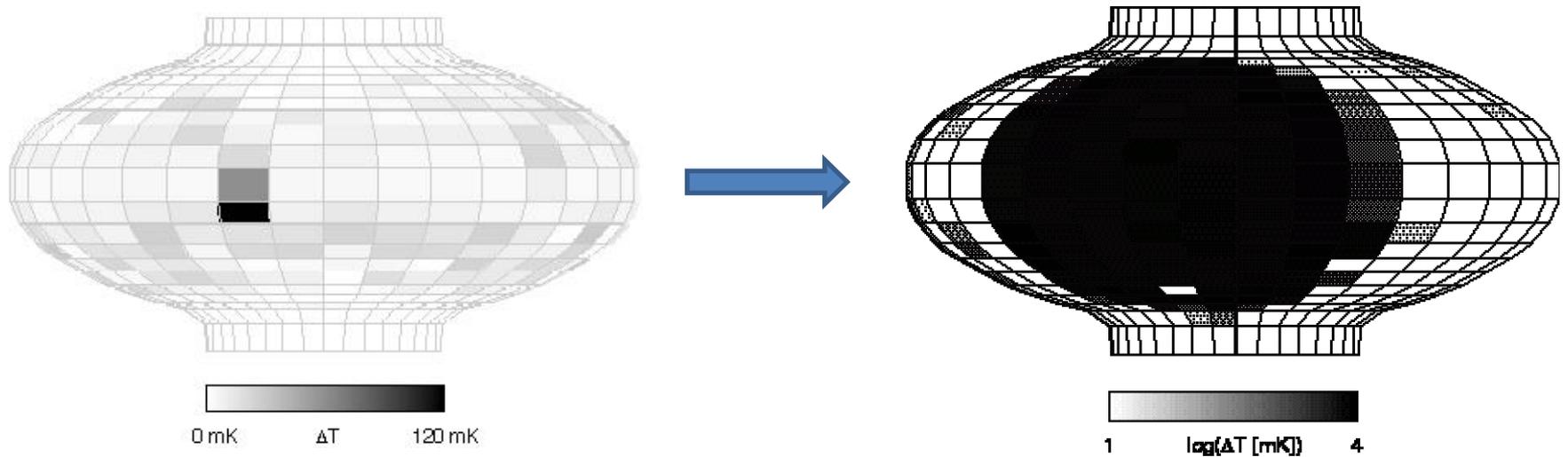
*Tolerate unavoidable defects,
but “neutralize” them
by thermally stabilizing them*



*Improve **Bulk Thermal Conductivity**
by raising Nb purity*

Thermal Breakdown

(Very strong heating observed over extended areas.)



Field cannot be increased!

STUDY OF THERMAL EFFECTS IN SRF CAVITIES

J.Lesrel, S.Bousson, T.Junquera, A.Caruelle, M.Fouaidy

I.P.N. (CNRS-IN2P3-Univ.Paris XI) ORSAY

The response time of both thermometer types are much longer than the time needed for the cavity temperature rise (Q_0 drop) during a quench so they are not suited for a reliable study of such transients.

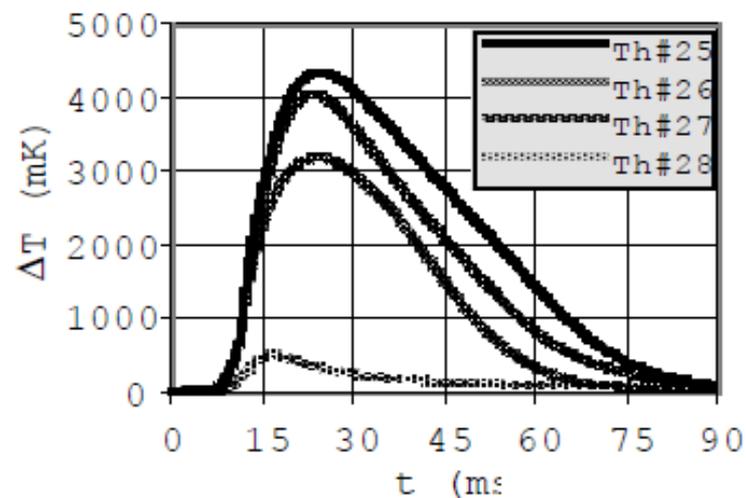
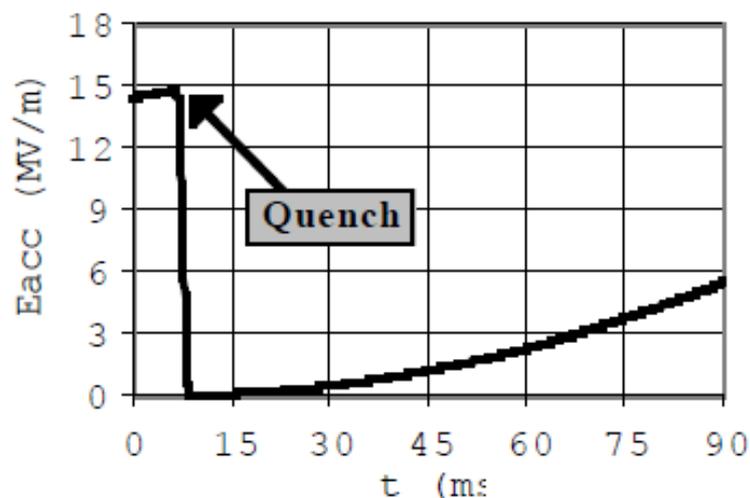


Fig. 7: Variation of E_{acc} during a quench in a 3 Ghz cavity.

Fig. 8: Heatings measured on 4 thermometers during the quench.

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Without FE, the dissipated power P_{diss} is related to the magnetic field H , the cavity surface S_{cav} and the surface resistance R_s by the relation:

$$P_{diss}(t) = \frac{1}{2} \cdot \iint_{S_{cav}} H^2(t) \cdot R_s \cdot dS$$

During a quench, a normal resistive surface S_N (having a normal surface resistance R_s^N) grows. Then P_{diss} has two contributions: power dissipation in the superconducting area ($= S_{cav} - S_N$) and in the normal conducting region S_N .

$$P_{diss}(t) = \frac{1}{2} \cdot \iint_{S_{cav} - S_N} H^2(t) \cdot R_s \cdot dS + \frac{1}{2} \cdot \iint_{S_N} H^2(t) \cdot R_s^N \cdot dS \quad (1)$$

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.... Finally by combining the previous equations, the normal conducting surface is deduced:

$$S_N(t) = \frac{2 \cdot \left(\frac{P_i \cdot \omega}{\sqrt{Q_{ext}}} - \frac{d\sqrt{U(t)}}{dt} \right)}{\omega \cdot \sqrt{U(t)}} \cdot \frac{1}{Q_{ext}} \cdot \frac{R_S \cdot S}{R_S^N}$$

$$\frac{1}{2} \cdot \alpha^2 \cdot \left(\frac{1}{l \sqrt{Q}} \sqrt{\frac{r}{Q}} \right)^2 \cdot R_S^N$$

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Assuming a circular shape for the expanding normal region $S_N = \pi \cdot R_{NC}^2$, we have an immediately measurement of the normal conducting surface radius $R_{NC}(t)$ (Fig. 12).

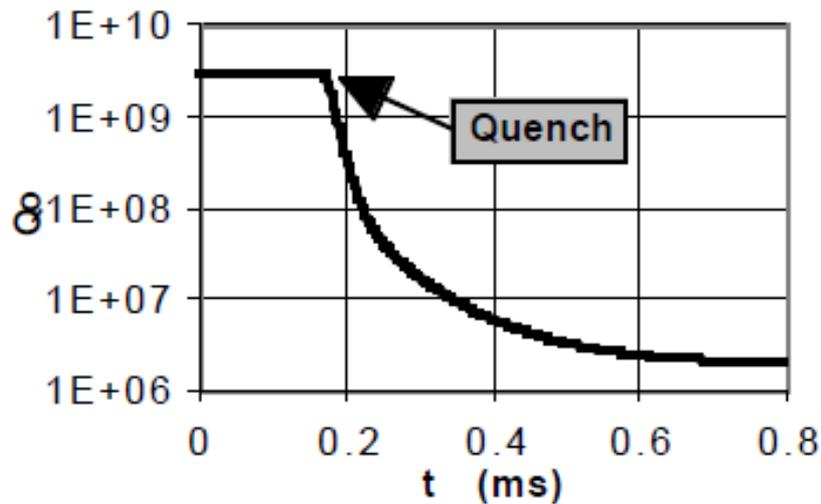


Fig. 11: Unloaded quality factor during the cavity quench.

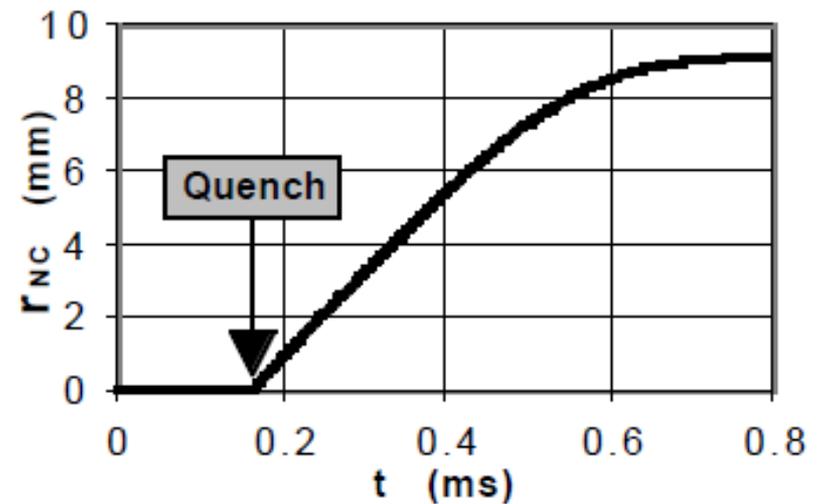


Fig. 12: Increase of the radius of the normal conducting area.

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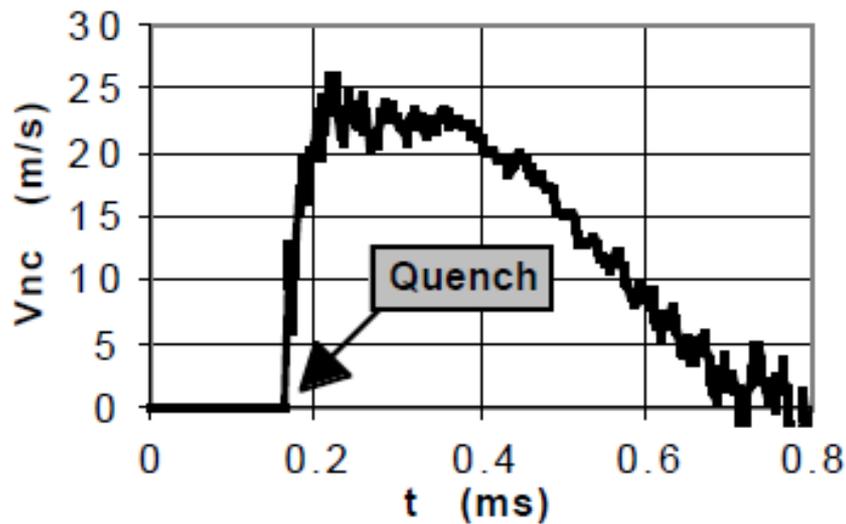


Fig. 13: Expansion velocity of the normal conducting area.

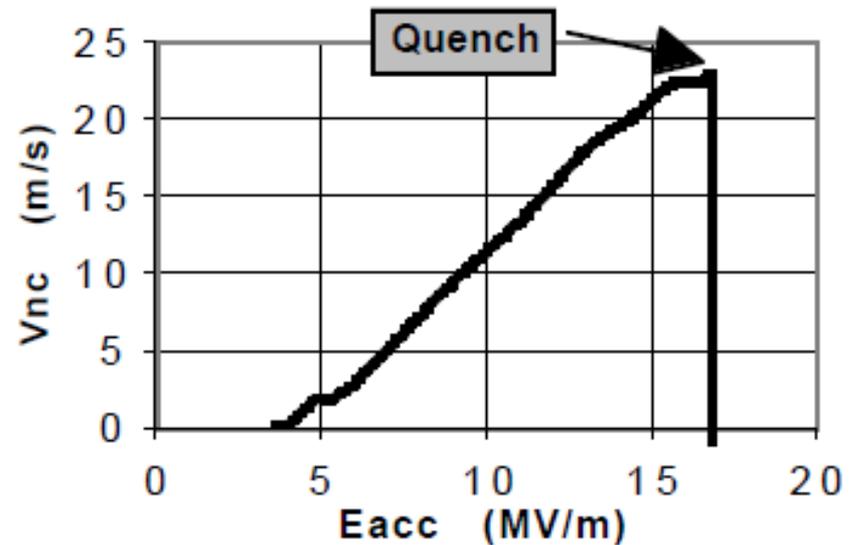


Fig. 14: Expansion velocity as a function of the accelerating field.

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One important result is that V_{NC} does not depend on the defect radius R_d in the studied range $50 \mu\text{m}$ - $1000 \mu\text{m}$.

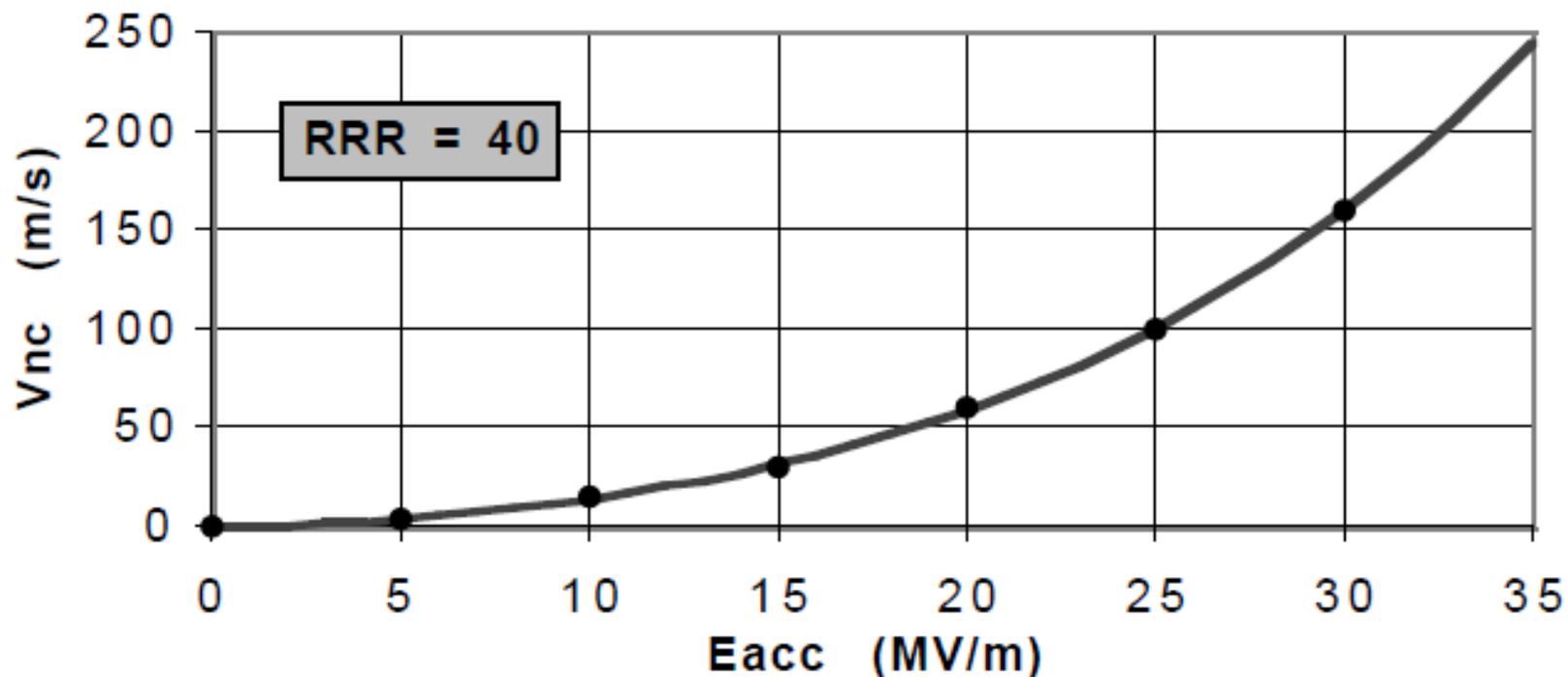


Fig. 15: Calculated expansion velocity as a function of E_{acc} .

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A REVIEW OF THE NIOBIUM ON COPPER SPUTTERING TECHNOLOGY FOR SUPERCONDUCTING RF CAVITIES

S. Calatroni*, CERN, 1211 Geneva 23, Switzerland

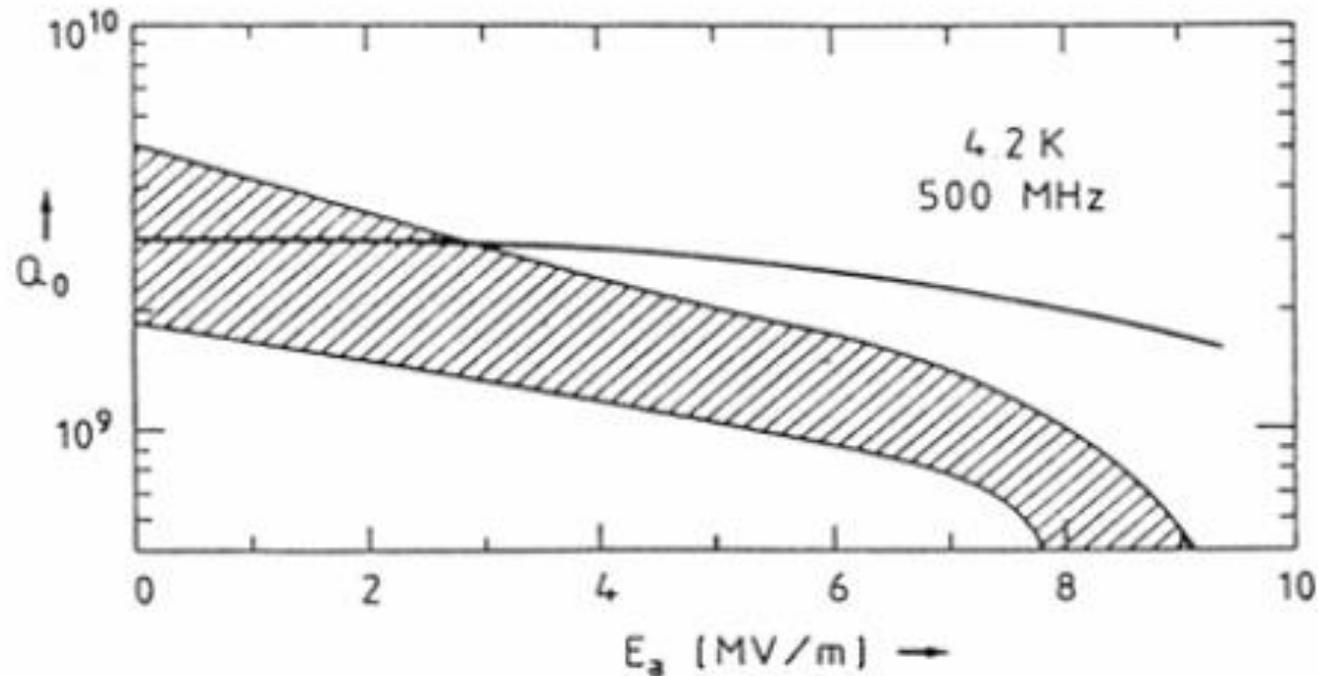


Figure 1: Typical $Q(E_a)$ curves of niobium sheet metal (line) and niobium sputter coated copper cavities (hatched) in the 80's.

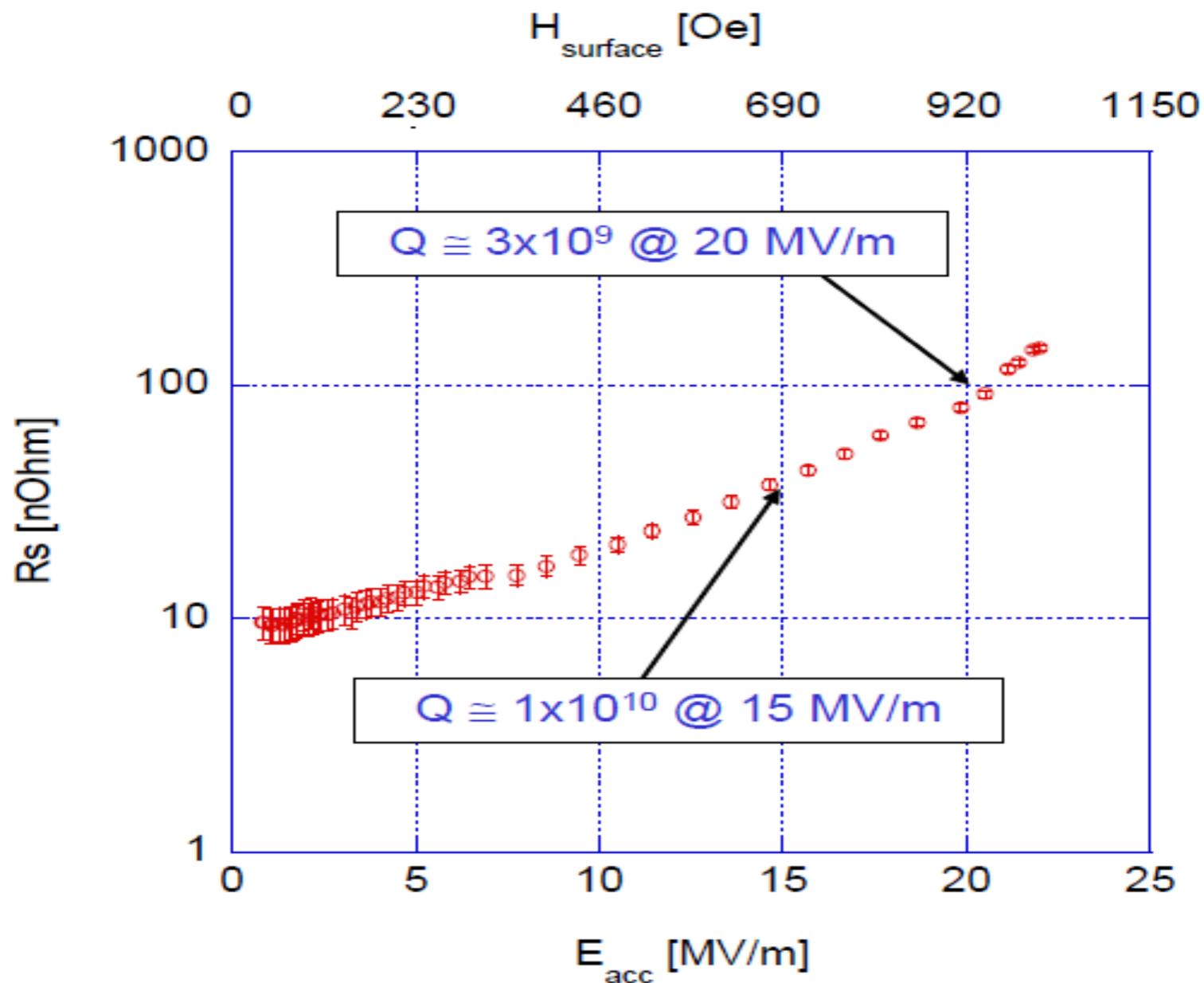
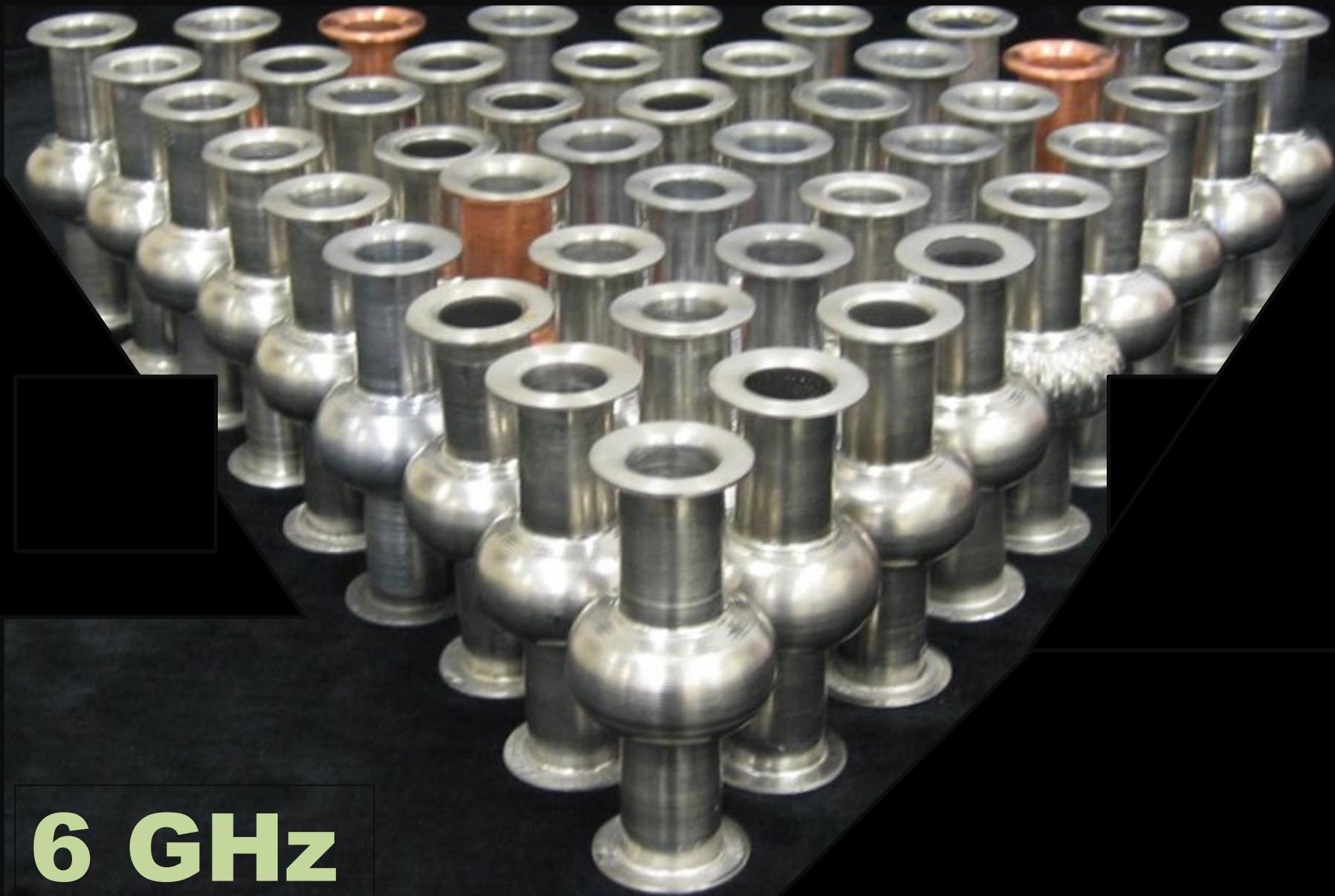


Figure 4: State-of-the art performance of Nb/Cu cavities at 1.7 K and 1500 MHz

on my personal opinion

*It is not written on the Bible that
«films are worse than the bulk!»*



6 GHz

In a small, but active, Research Group

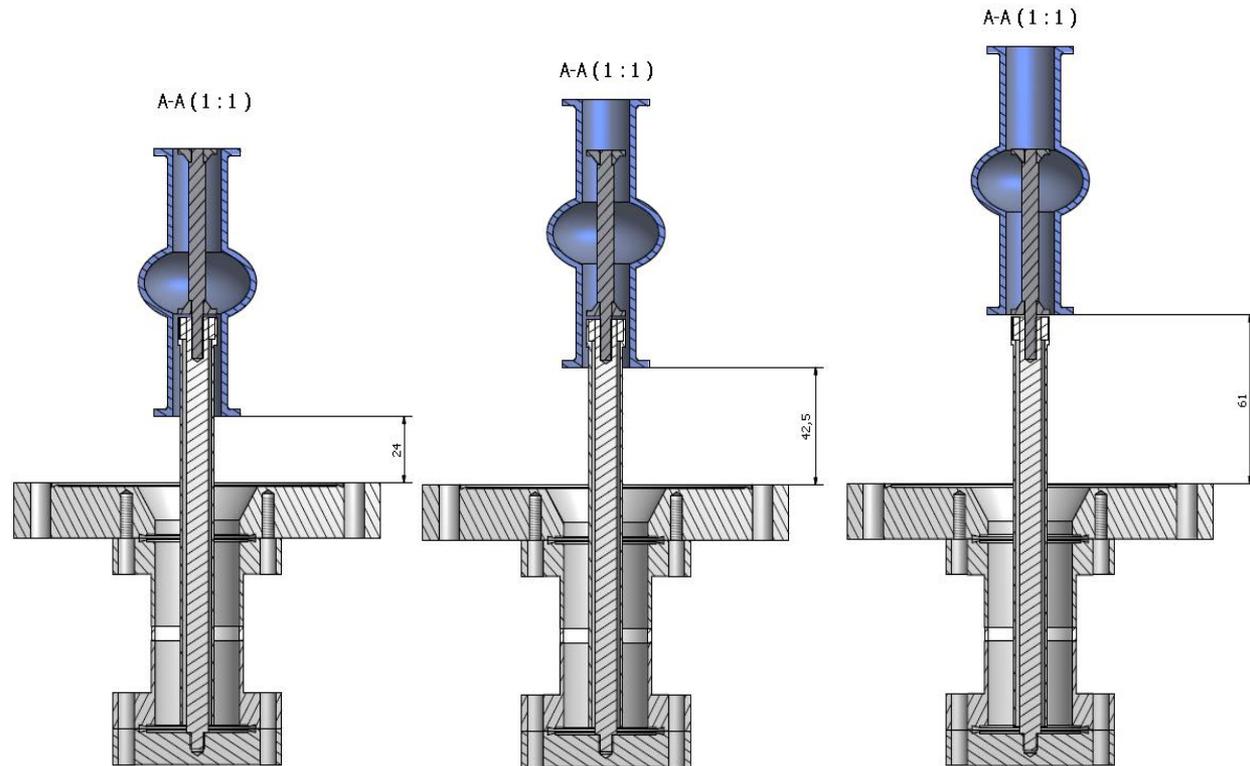
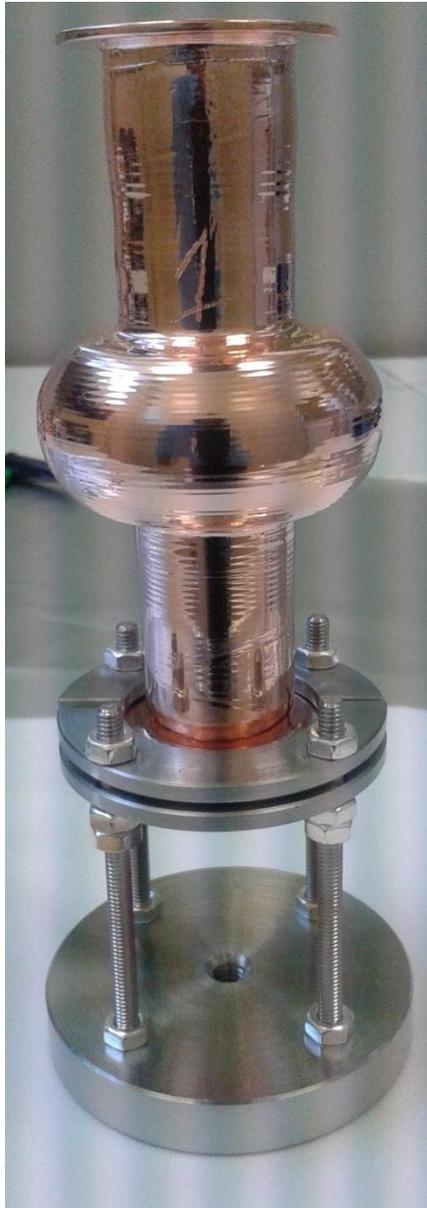
6 GHz Cavities

The Ideal Tool For Self-motivation !

Common sentences you can daily hear in our lab:

- «Bye, I go to spin some other **60 cavities!**»
- «This week we have **tested 12 cavities**»
- «We **just did** the EP, I go **now** to anneal the cavity, **tonight** we do the rf test!»
- «This year we tested **148 cavities!**»

2nd Nb Sputtering - Cu1 6GHz





Niobium-based sputtered thin films for Corrosion Protection of proton-irradiated liquid water targets for [18F] production

H. Skliarova, O. Azzolini, O. Dousset, R.R. Johnson, V. Palmieri

(Submitted on 29 Aug 2013)

Chemically inert Coatings on Havar entrance foils of the targets for [18F] production via proton irradiation of enriched water at pressurized conditions are needed to decrease the amount of ionic contaminants released from Havar. In order to find the most effective protective coatings, the Nb-based coating microstructure and barrier properties have been correlated with deposition parameters as: substrate temperature, applied bias, deposition rate and sputtering gas pressure. Aluminated quartz used as a substrate allowed us to verify the protection efficiency of the desirable coatings as diffusion barriers. Two modeling corrosion tests based on the extreme susceptibility of aluminum to liquid gallium and acid corrosion were applied. Pure Niobium coatings have been found less effective barriers than Niobium-titanium coatings. But Niobium oxide films, according to the corrosion tests performed, showed superior barrier properties. Therefore Multi-layered Niobium-Niobium oxide films have been suggested, since they combine the high thermal conductivity of Niobium with the good barrier properties of Niobium oxide.

As far as Nb sputtering is concerned,

The best films for SRF

**..... are the worst as diffusion barriers
against corrosion protection**

Nb films are porous

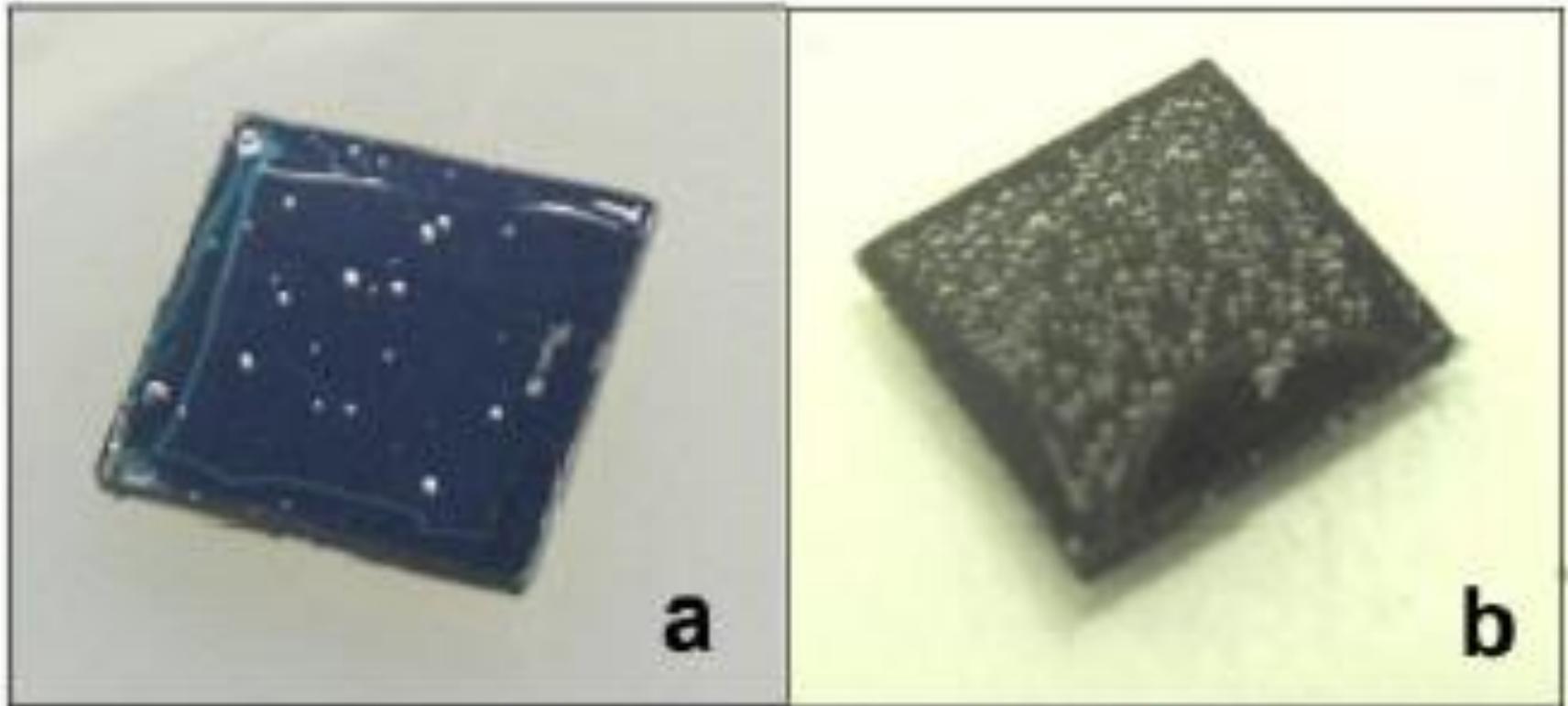


Figure 2 Acid solution resistance: a) high, b) low

Nb sputtered film resistance to Liquid metal Embrittlement



Figure 3 Liquid gallium test: corrosion,

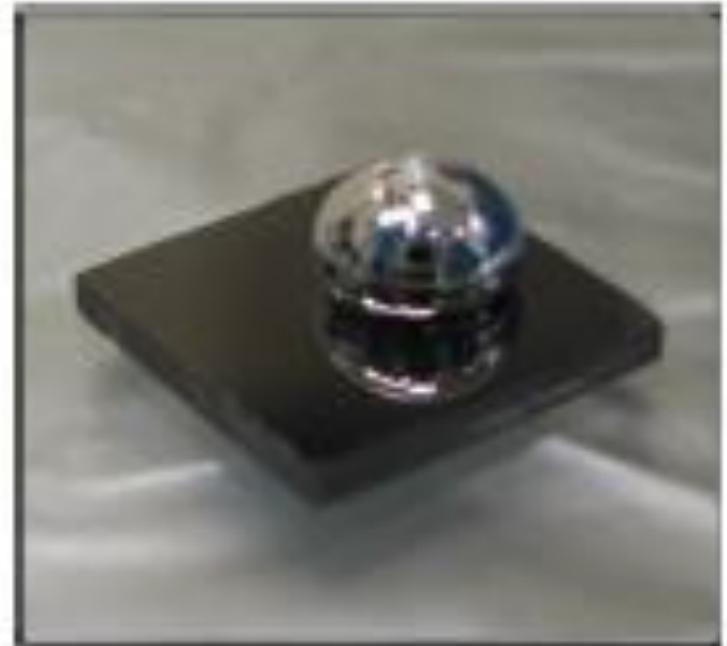


Figure 4 Liquid gallium test: resistance

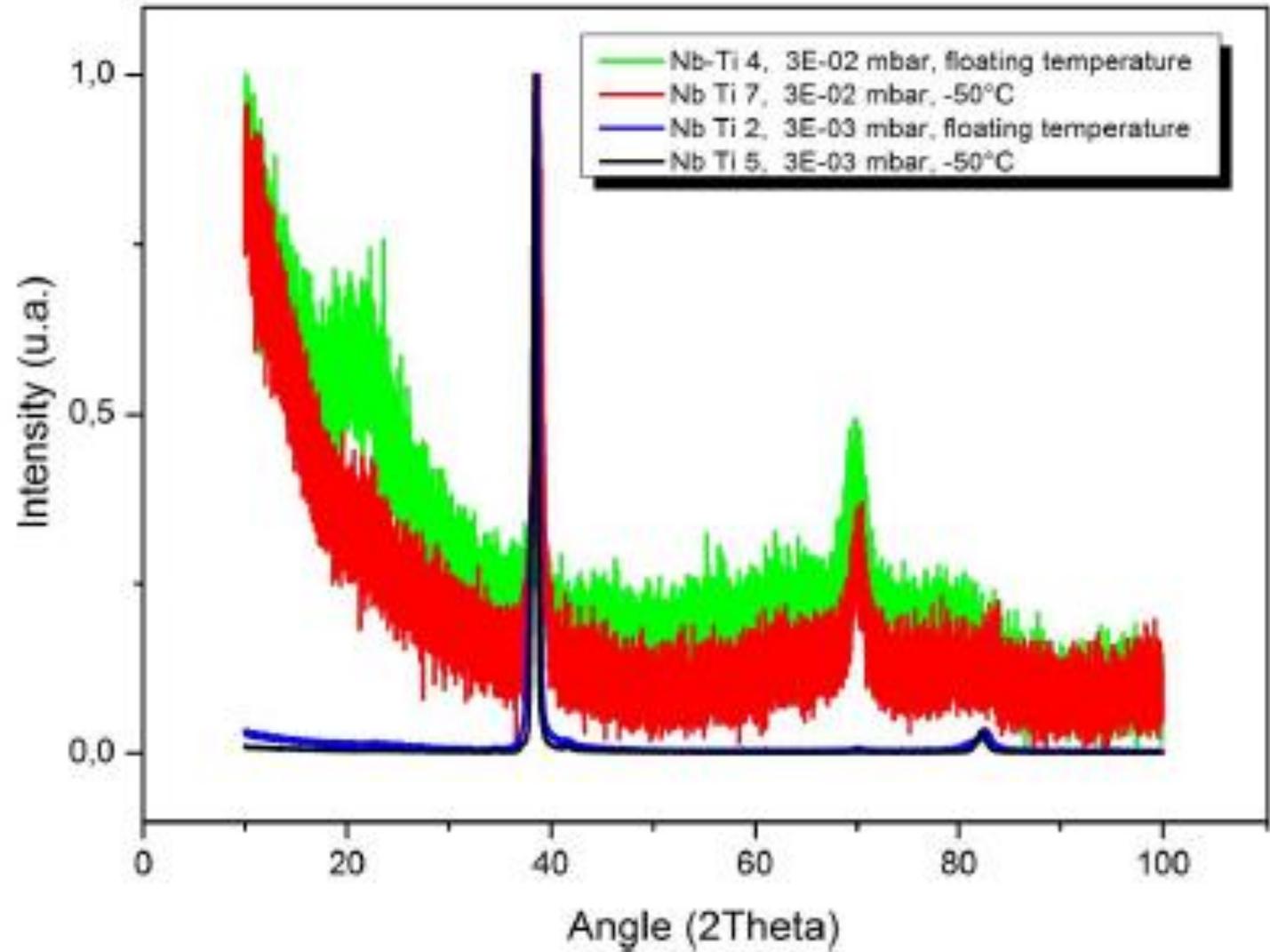
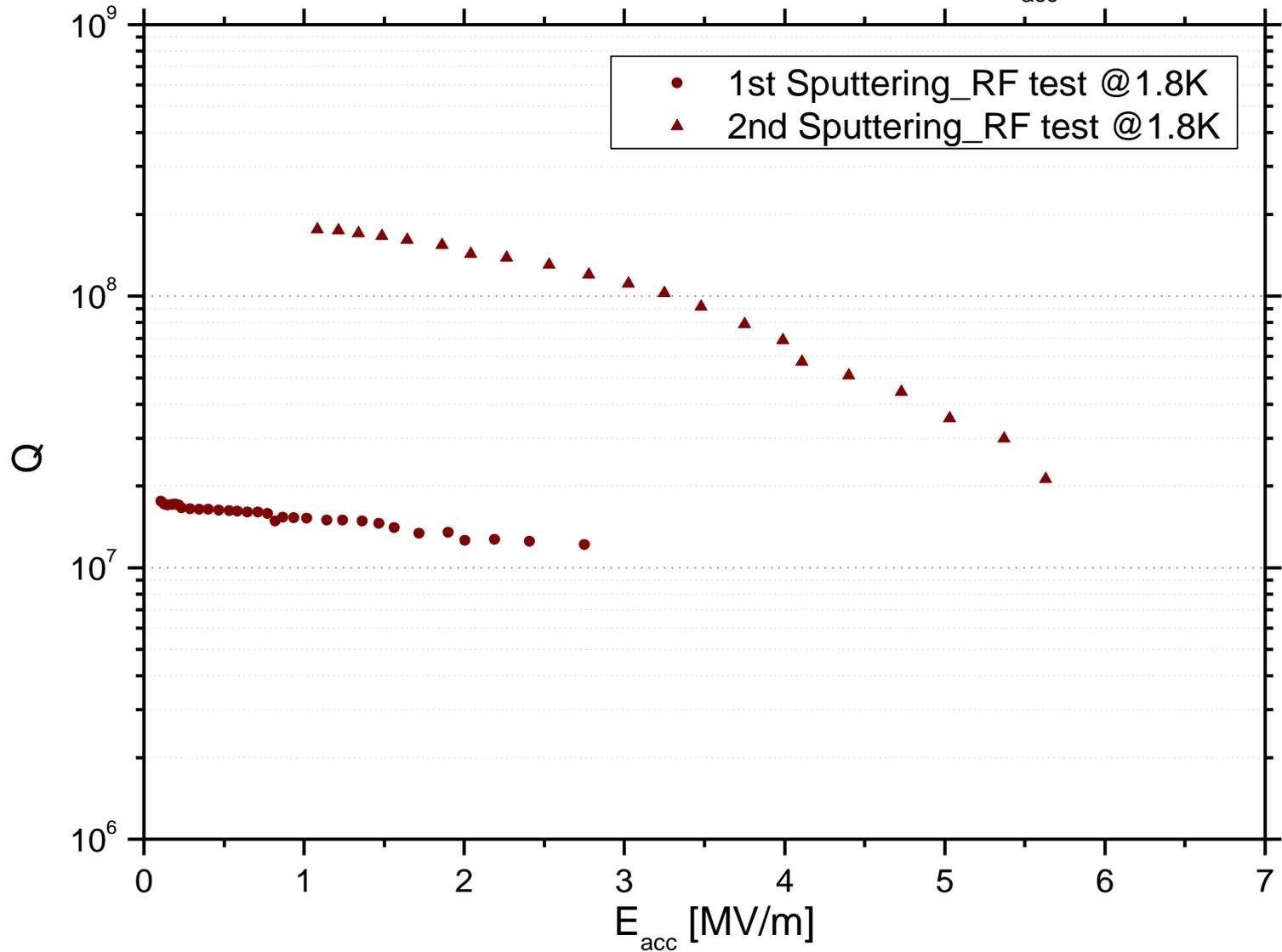


Figure 10 Influence of temperature and pressure on Nb-Ti alloy film deposition

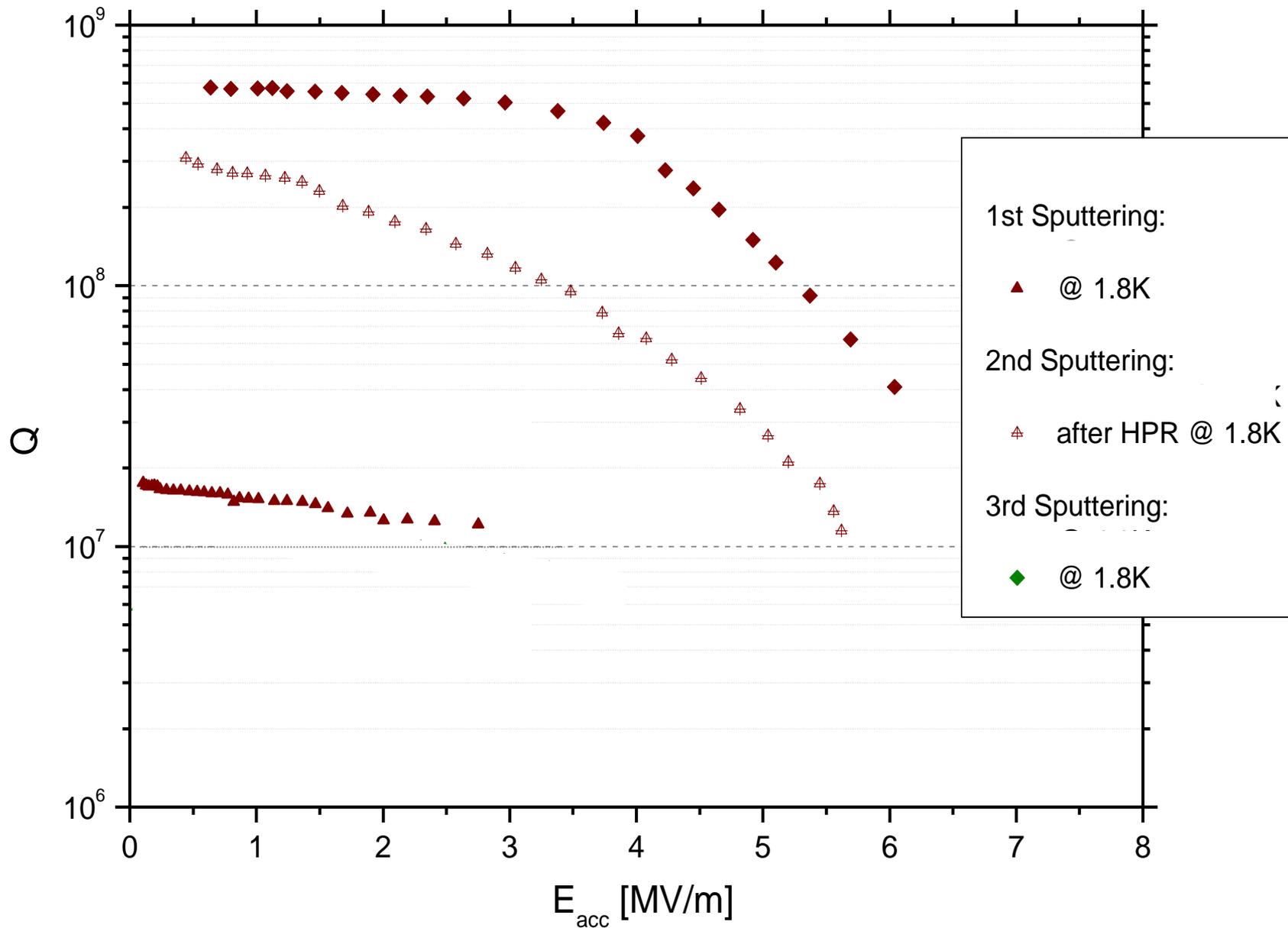
	Nb	Nb-Ti	Nb₂O₅
Low Ar pressure	Crystalline	Crystalline	Amorphous
High Ar Pressure	Crystalline	Amorphous-like	Amorphous
Negative bias	Crystalline	Crystalline	Amorphous
- 50 – 0 °C	Amorphous-like	Depends on Ar pressure	Amorphous

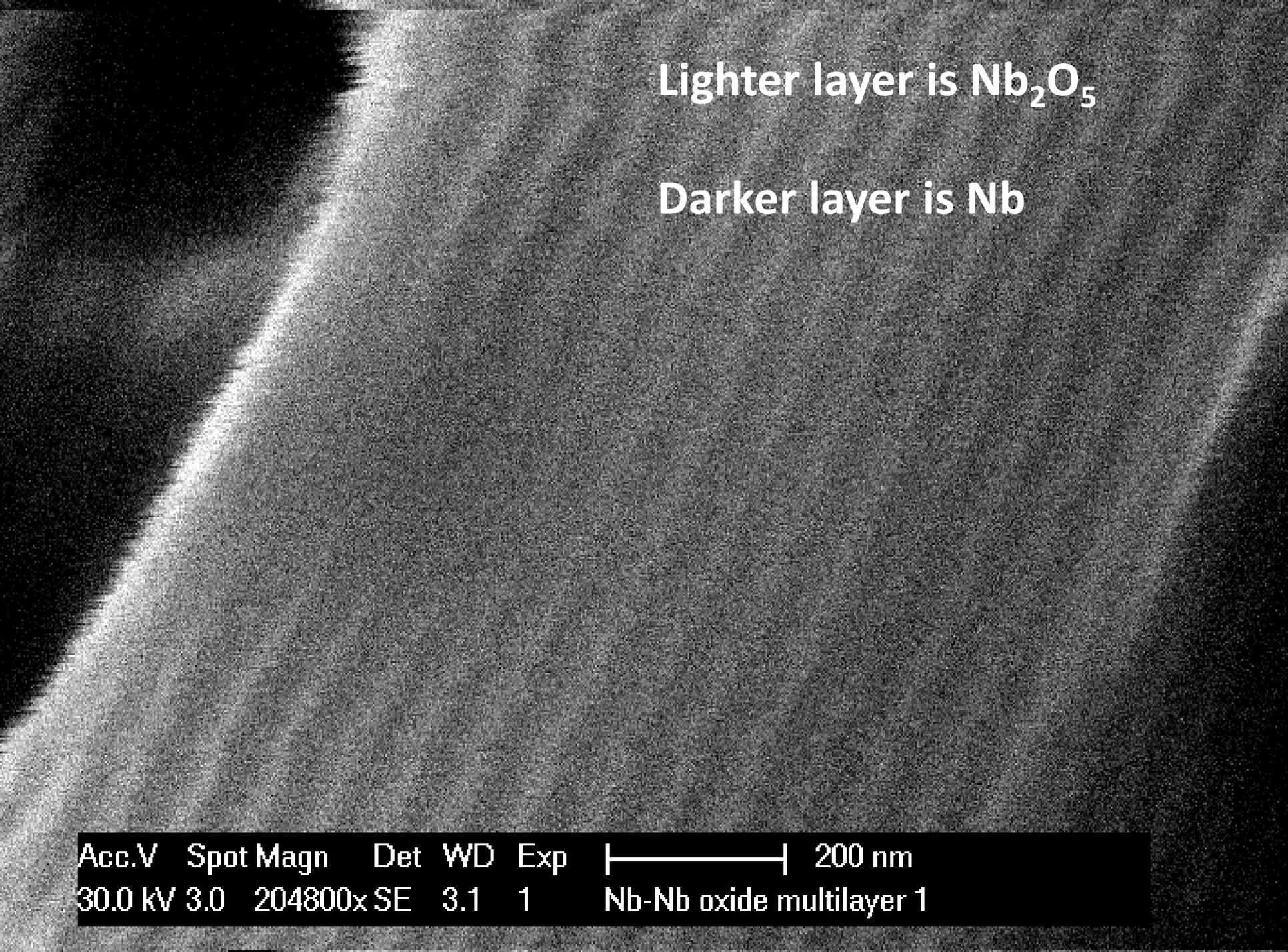
RF test Cu1

Comparison 1st and 2nd Sputtering_Q vs E_{acc} @1.8K



Comparison between 1st, 2nd and 3rd Nb/Cu Sputtering



Scanning electron micrograph (SEM) showing a cross-section of a Nb-Nb oxide multilayer. The image displays a dark, textured region on the left and a lighter, smoother region on the right, separated by a distinct boundary. The lighter region is identified as Nb2O5 and the darker region as Nb.

Lighter layer is Nb_2O_5

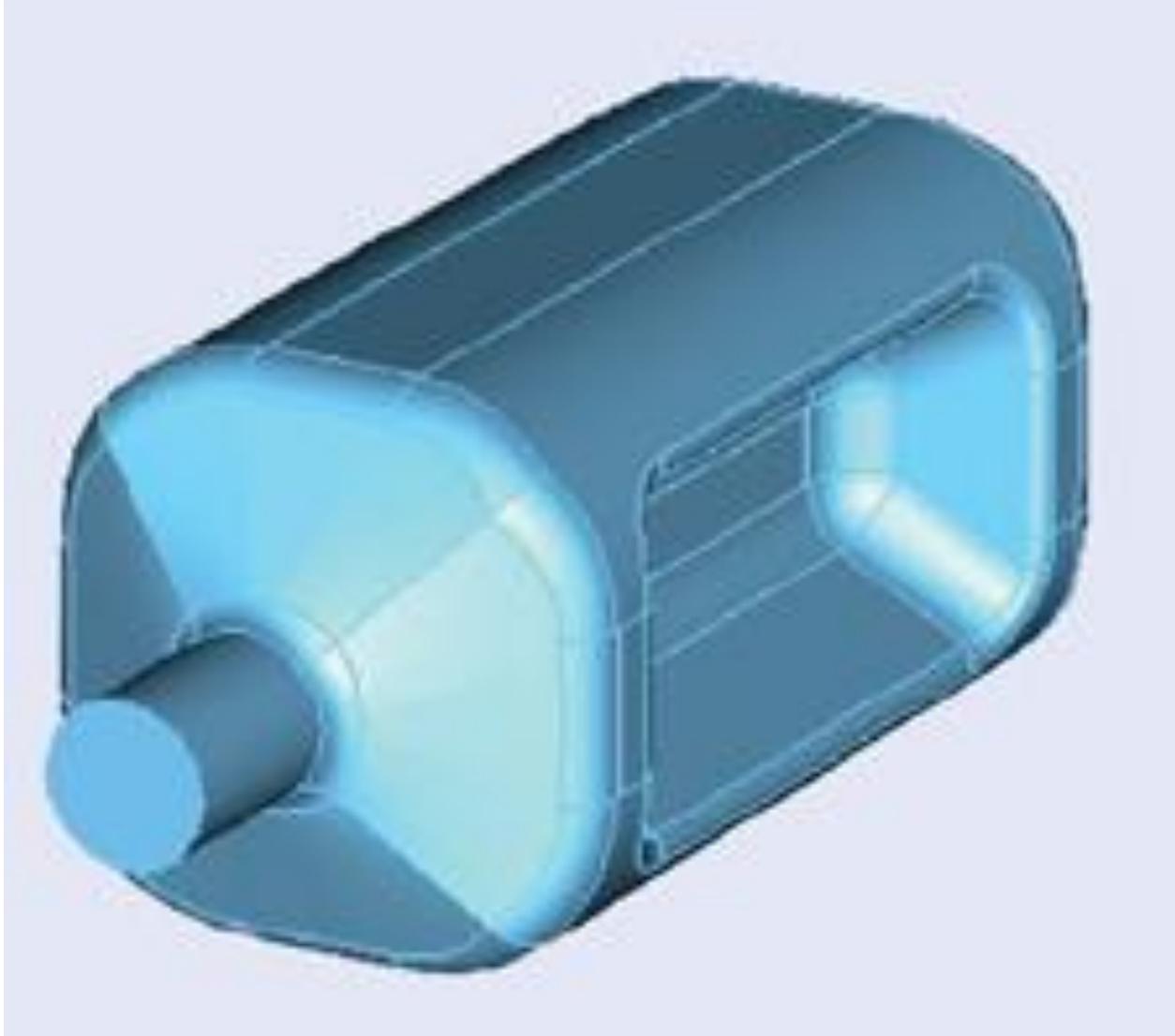
Darker layer is Nb

Acc.V	Spot	Magn	Det	WD	Exp	—————	200 nm
30.0 kV	3.0	204800x	SE	3.1	1	Nb-Nb oxide multilayer 1	

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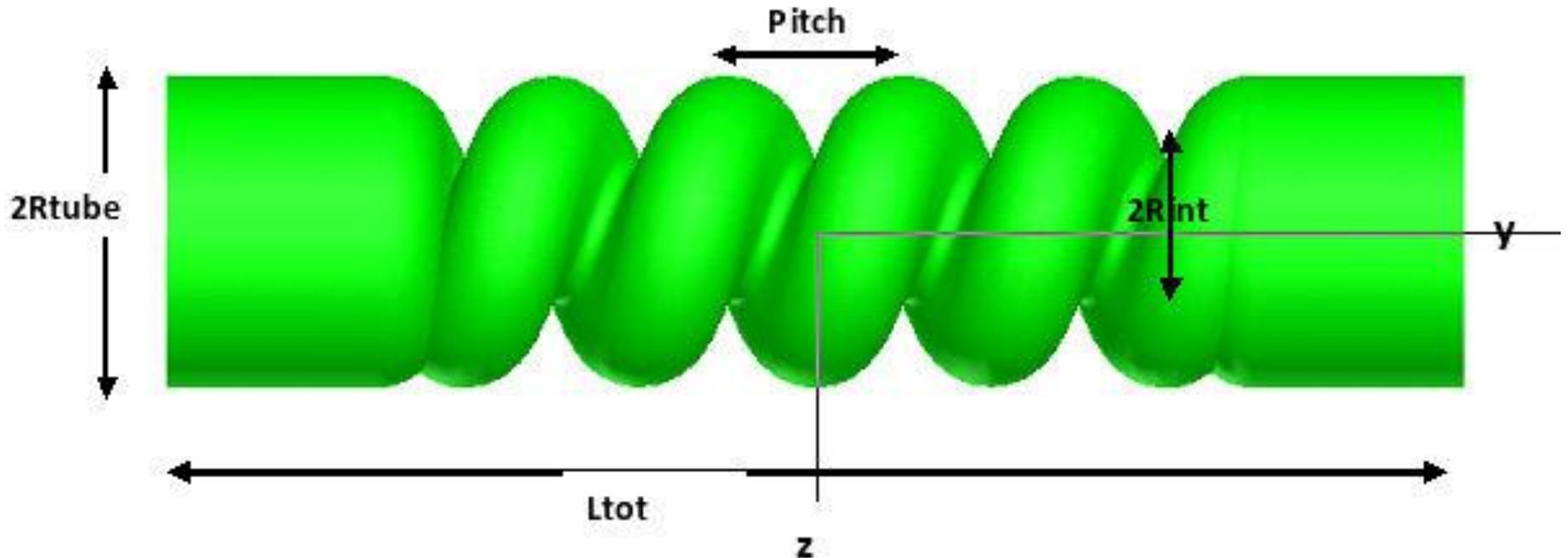
Take the most simple half wave cavity



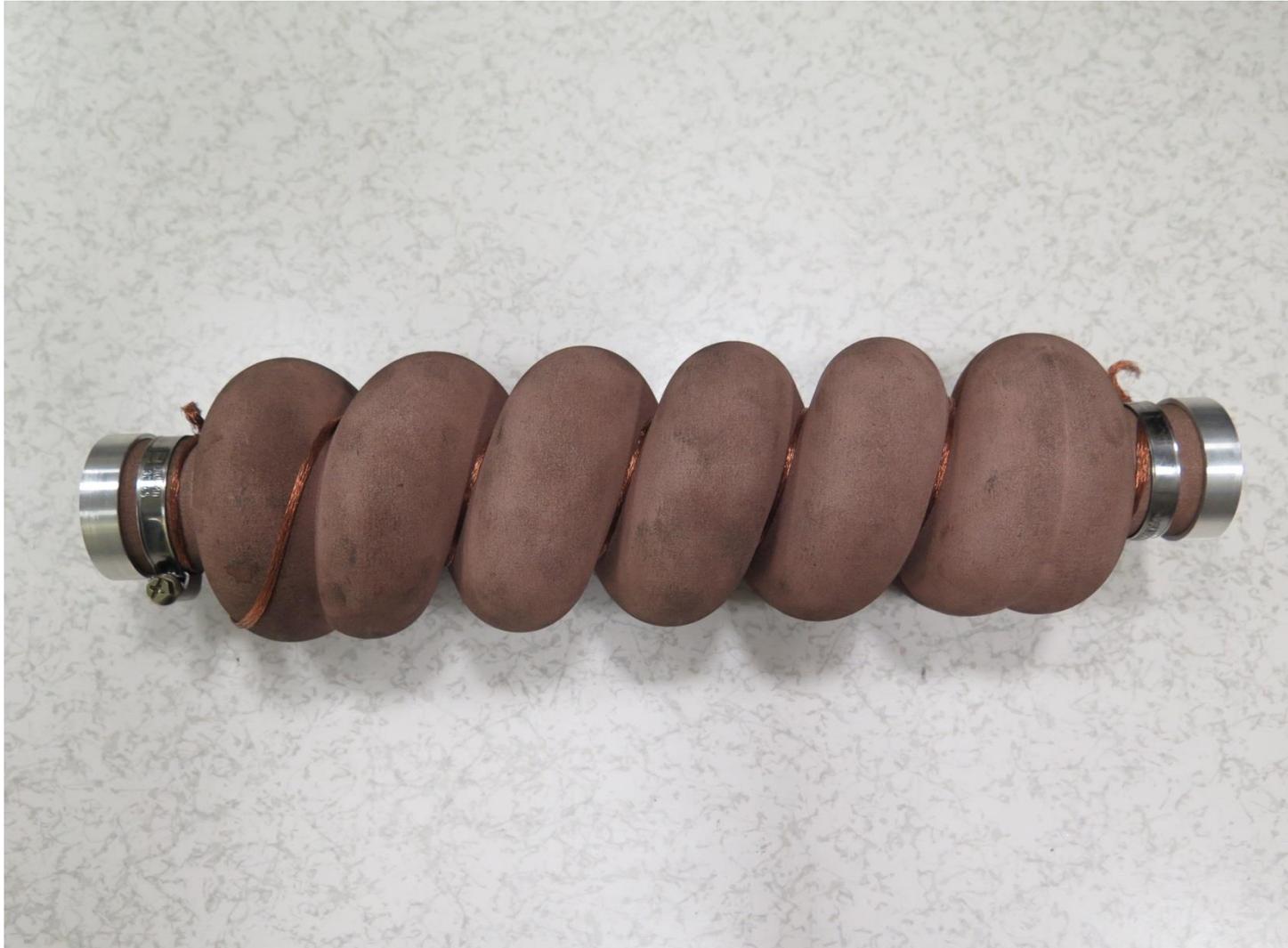
Crazy Idea n.º 1

The BCS cavity

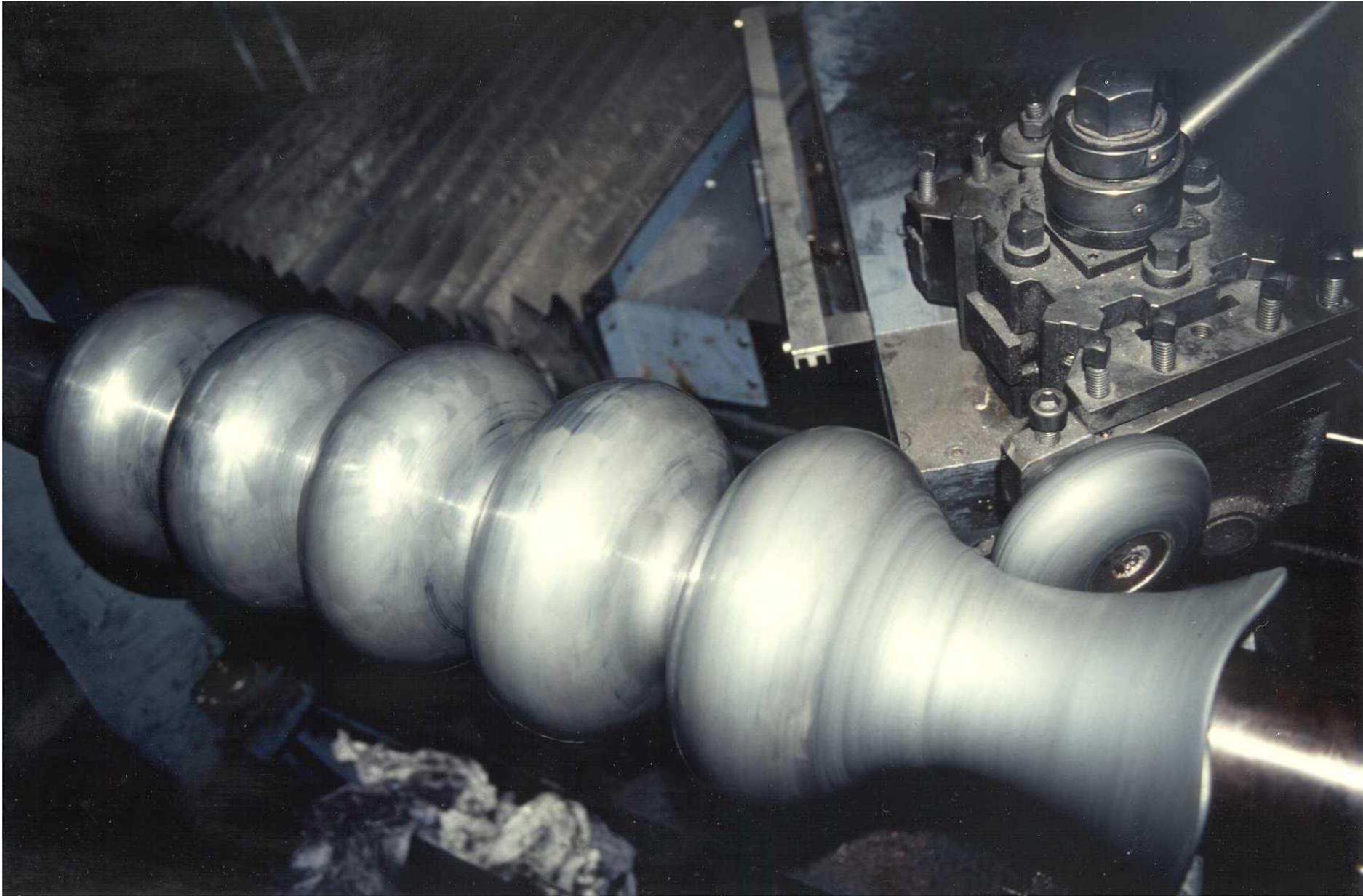
(Breaking the Cylindrical Symmetry)



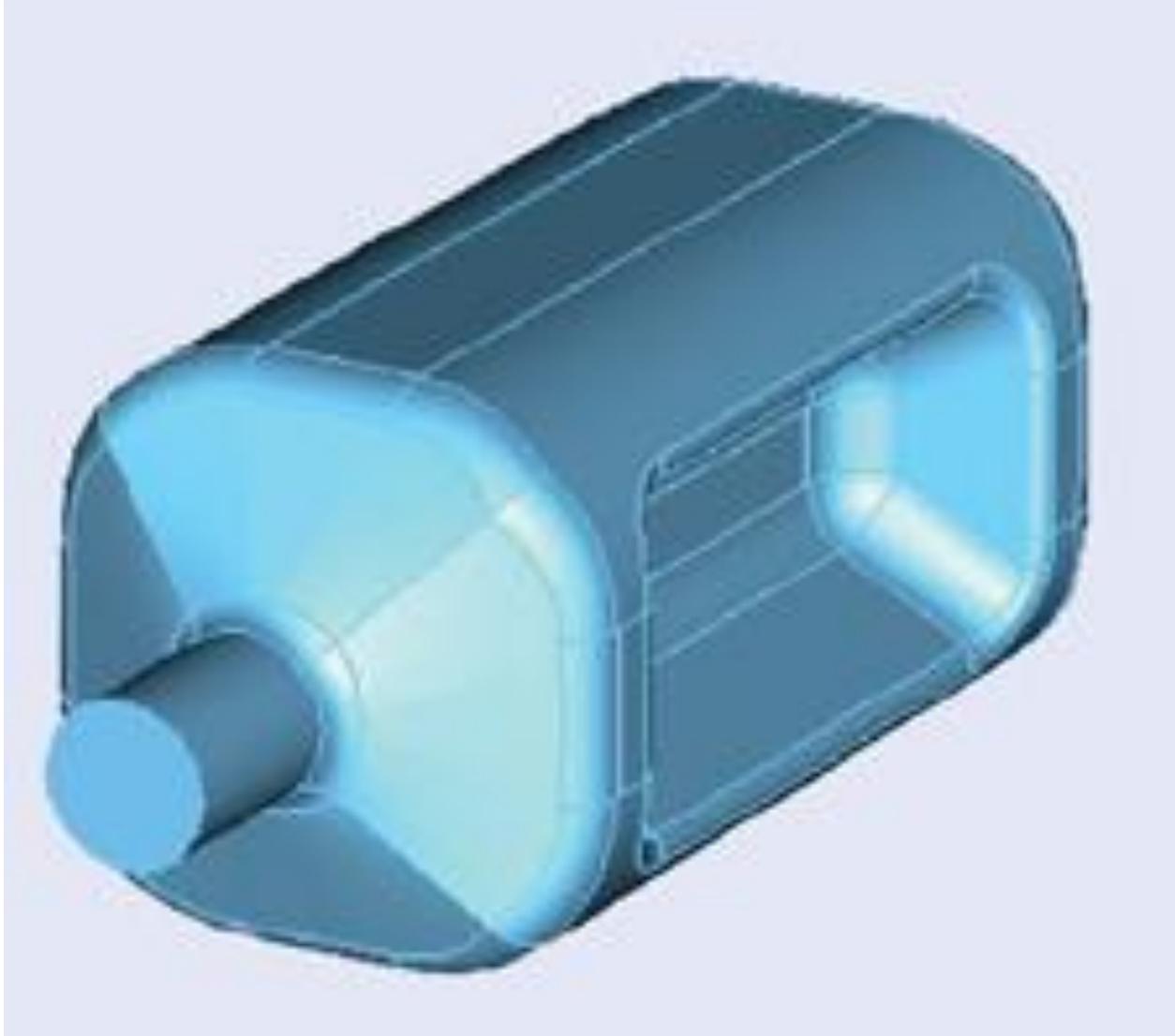
3D printing



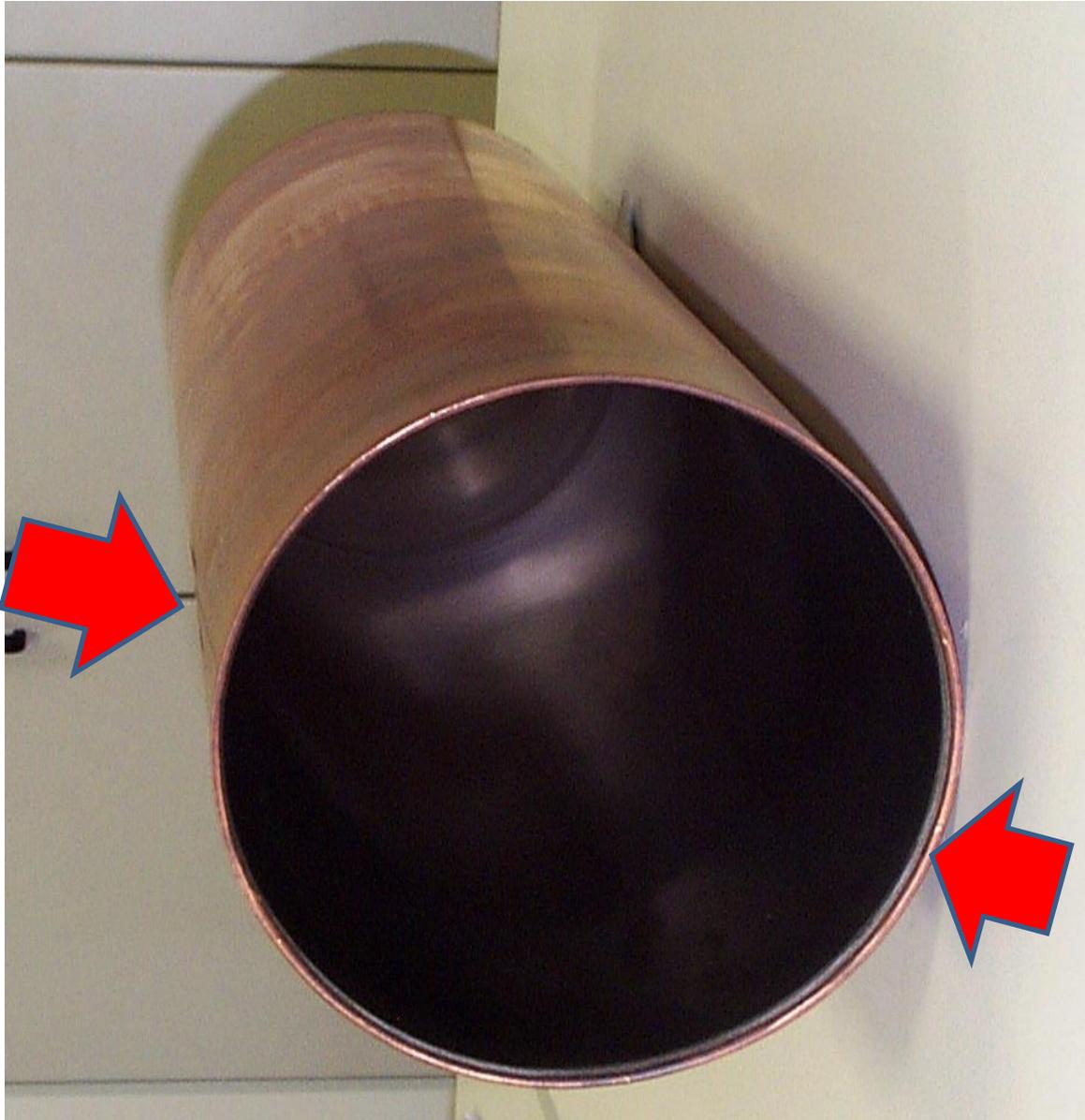
Crazy Idea n. 2

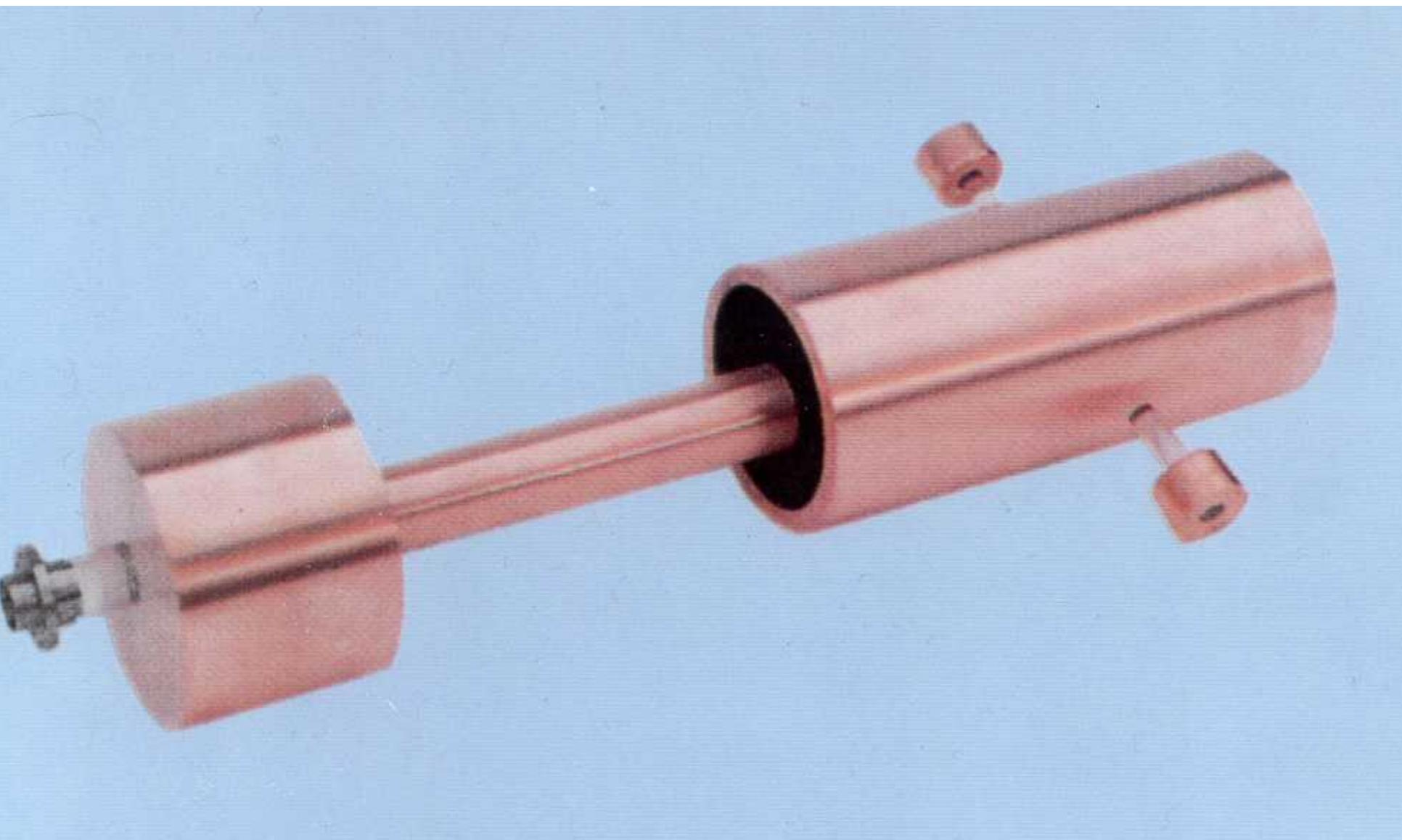


Take the most simple half wave cavity



A totally **Seamless** Nb Clad Cu Crab cavity is possible





Crazy Idea n. 3

Sputtering Nb onto plates or into tubes and then **spinning**

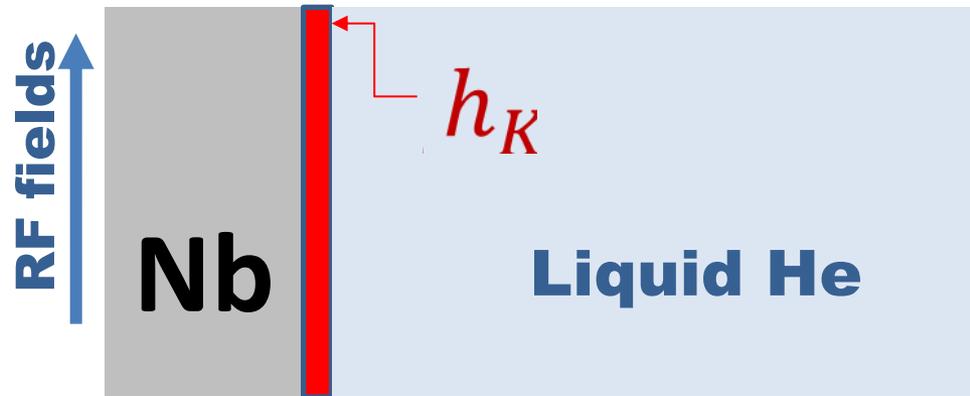




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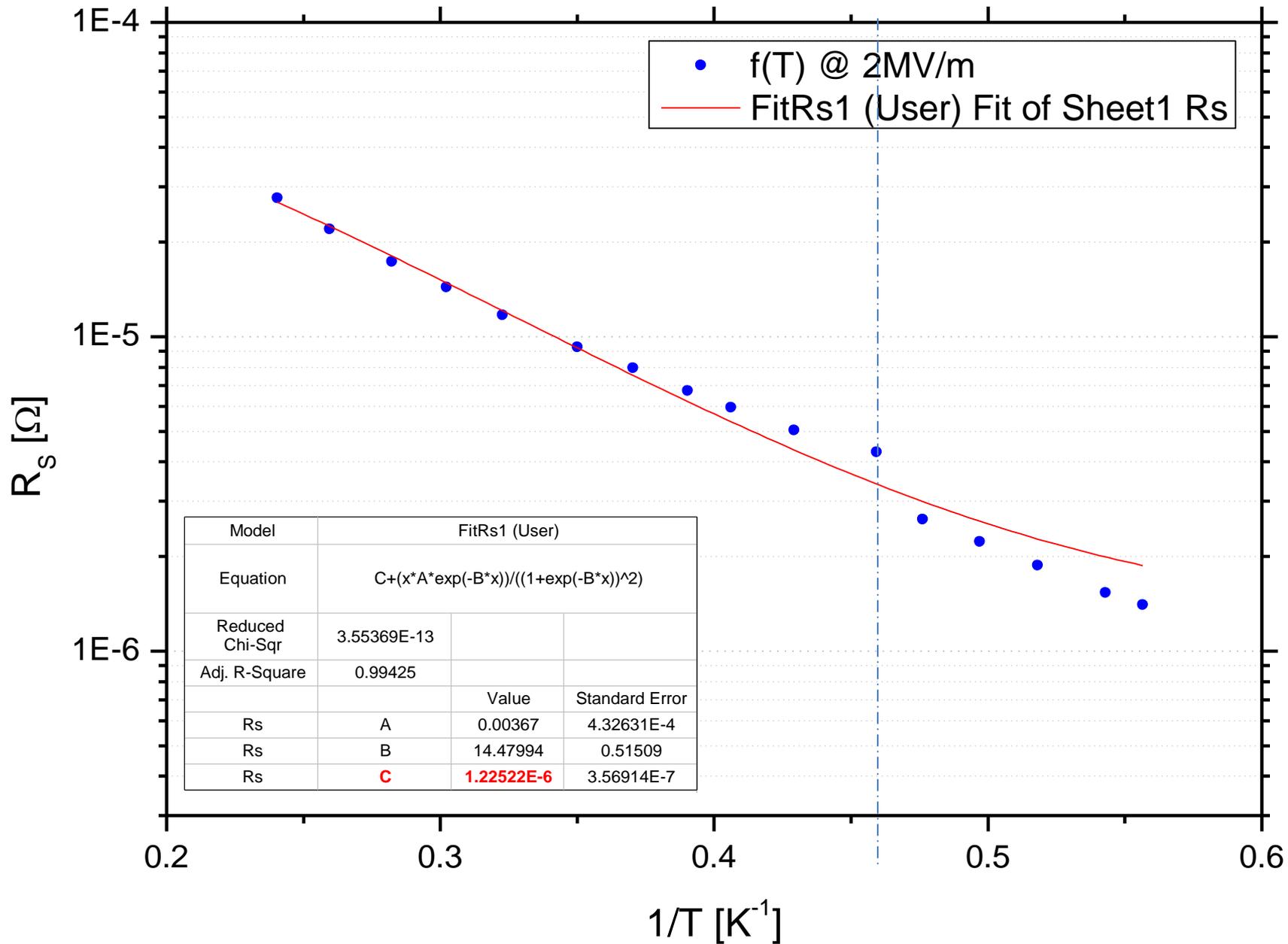
Kapitza conductance



$$h_{K_0} = \lim_{\Delta T_S \rightarrow 0} \frac{q}{\Delta T_S}$$

This quantity has a strong T^n temperature dependence with n varying between 2 and 4

R_s Nb 122 After ATM Annealing



**Effect of high temperature heat treatments on the quality factor of
 a large-grain superconducting radio-frequency niobium cavity**

P. Dhakal,¹ G. Ciovati,¹ G. R. Myneni,^{1,*} K. E. Gray,² N. Groll,² P. Maheshwari,³
 D. M. McRae,⁴ R. Pike,⁵ T. Proslir,² F. Stevie,³ R. P. Walsh,⁴ Q. Yang,⁶ and J. Zasadzinski⁷

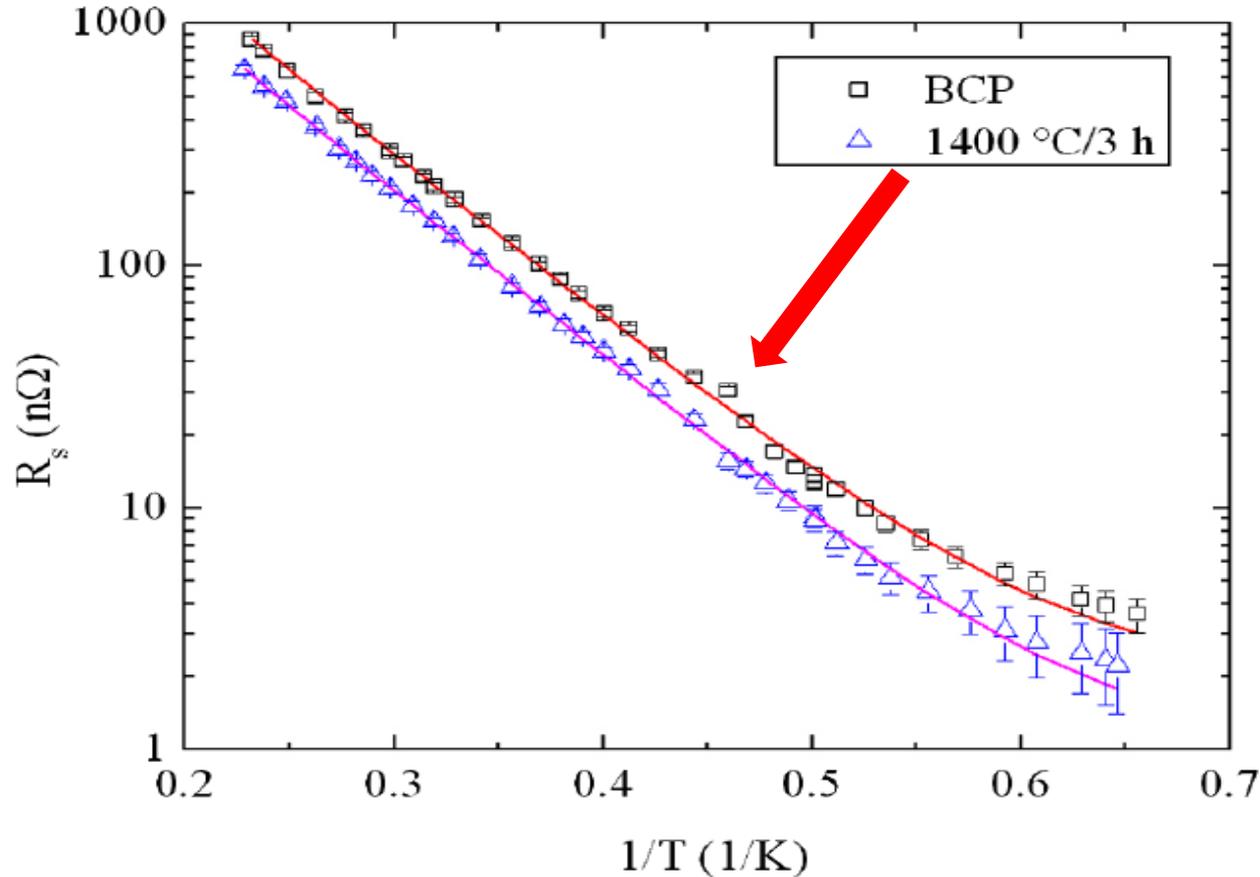
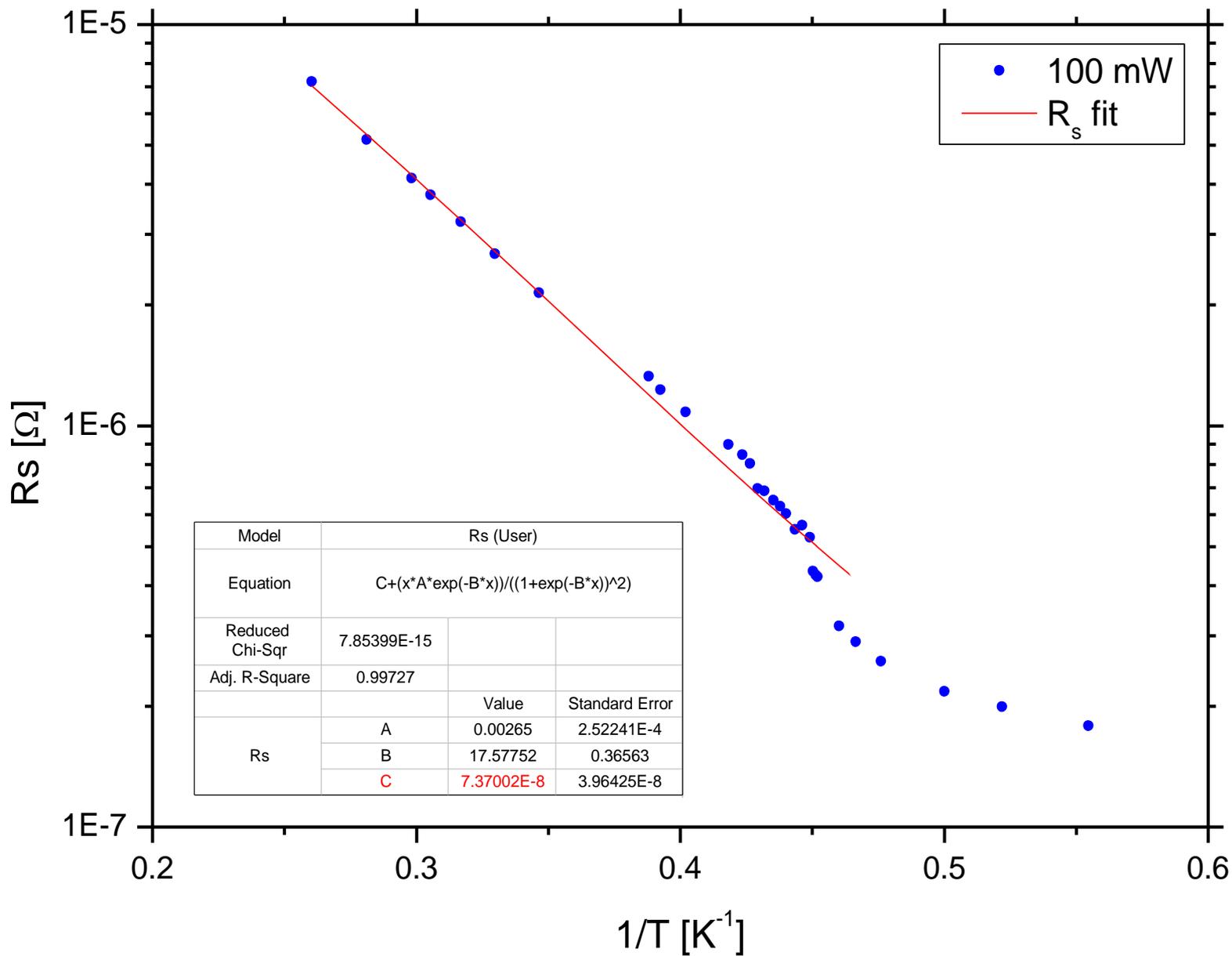
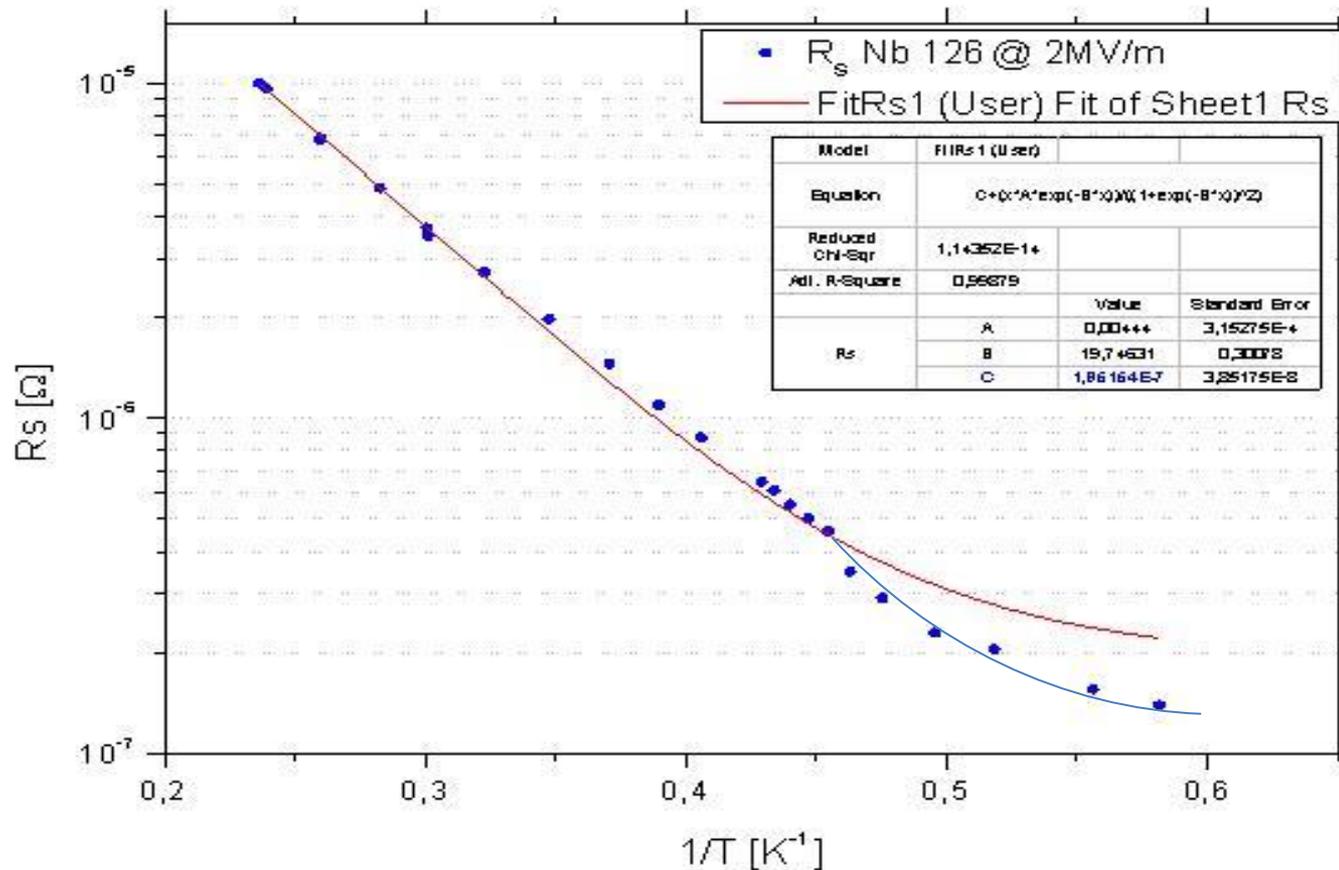


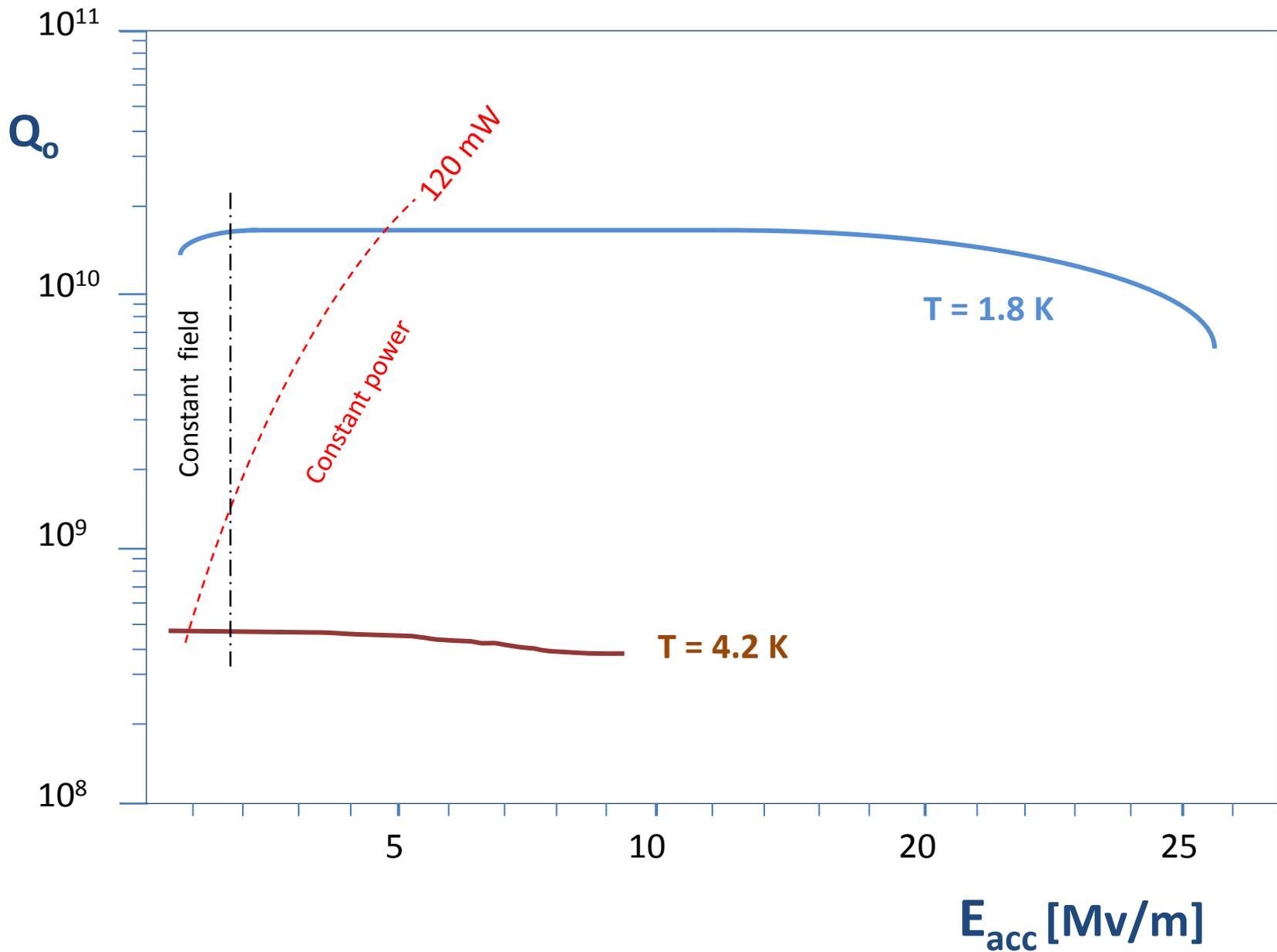
FIG. 9. R_s vs $1/T$ measured after BCP and after HT at 1400°C . Solid lines are least-square fits with $R_s(T) = R_{\text{BCS}}(T) + R_{\text{res}}$. The values of the fit parameters are $\Delta/kT_c = 1.87 \pm 0.02$, $\ell = (303 \pm 85)$ nm, $R_{\text{res}} = (2.0 \pm 0.3)$ n Ω after BCP and $\Delta/kT_c = 1.90 \pm 0.01$, $\ell = (76 \pm 17)$ nm, $R_{\text{res}} = (1.0 \pm 0.2)$ n Ω after HT at 1400°C .

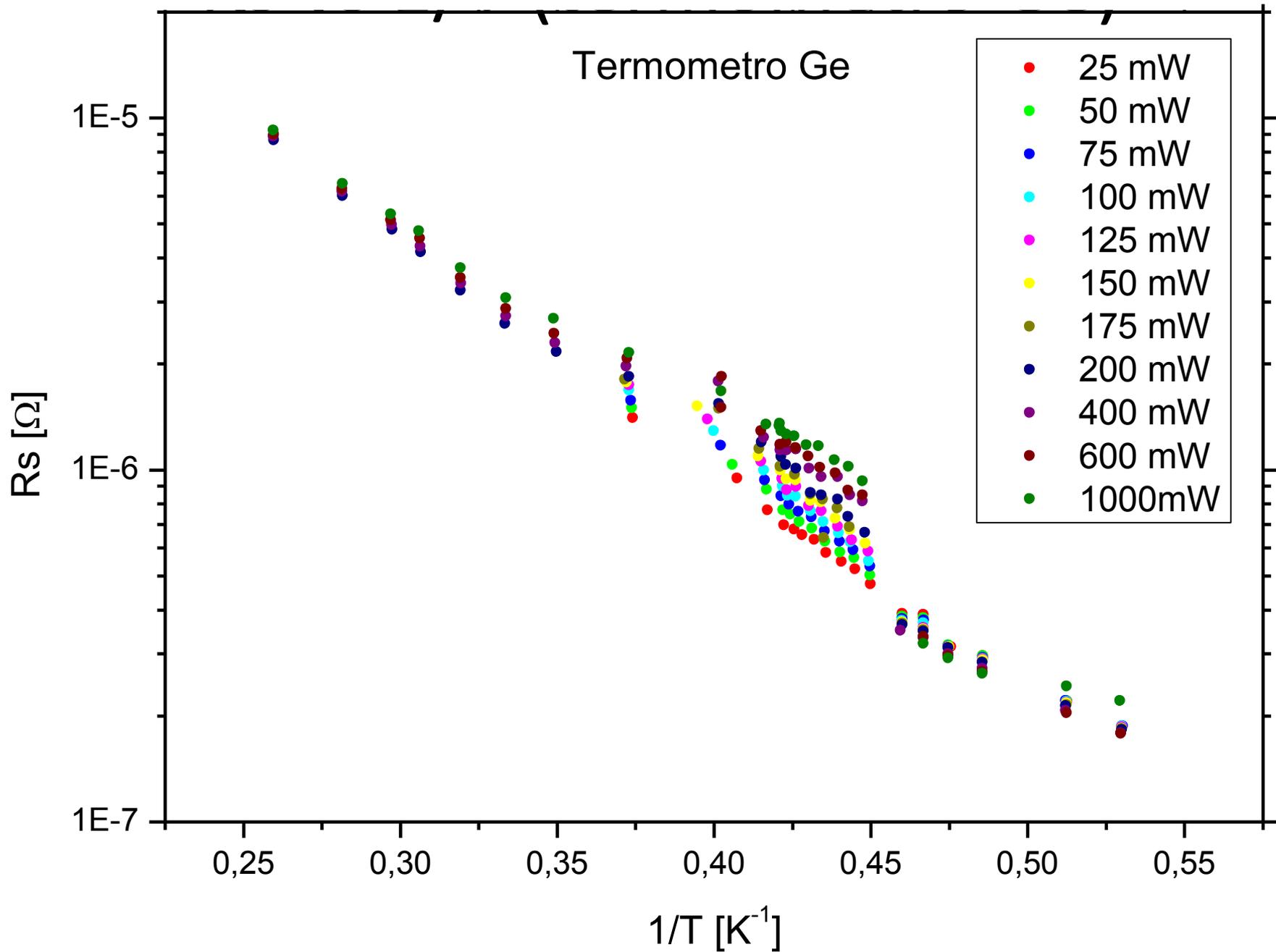
R_s Nb 122 After 3rd UHV Annealing



R_{RES} depends on the status of Liq- He instead than on Nb material?







$$R_{BCS}(T_0) = \frac{A\omega^2}{T_0} \exp\left[-\frac{sT_c}{2T_0}\right]$$

$$R_{BCS}(T_0 + \Delta T) \approx \frac{A\omega^2}{T_0} \exp\left[-\frac{sT_c}{2(T_0 + \Delta T)}\right]$$

$$R_{BCS}(T_0 + \Delta T) \approx \frac{A\omega^2}{T_0} \exp\left[-\frac{sT_c}{2T_0} \left(1 - \frac{\Delta T}{T_0}\right)\right]$$

$$R_{BCS}(T_0 + \Delta T) \approx \frac{A\omega^2}{T_0} \exp\left[-\frac{sT_c}{2T_0}\left(1 - \frac{\Delta T}{T_0}\right)\right]$$

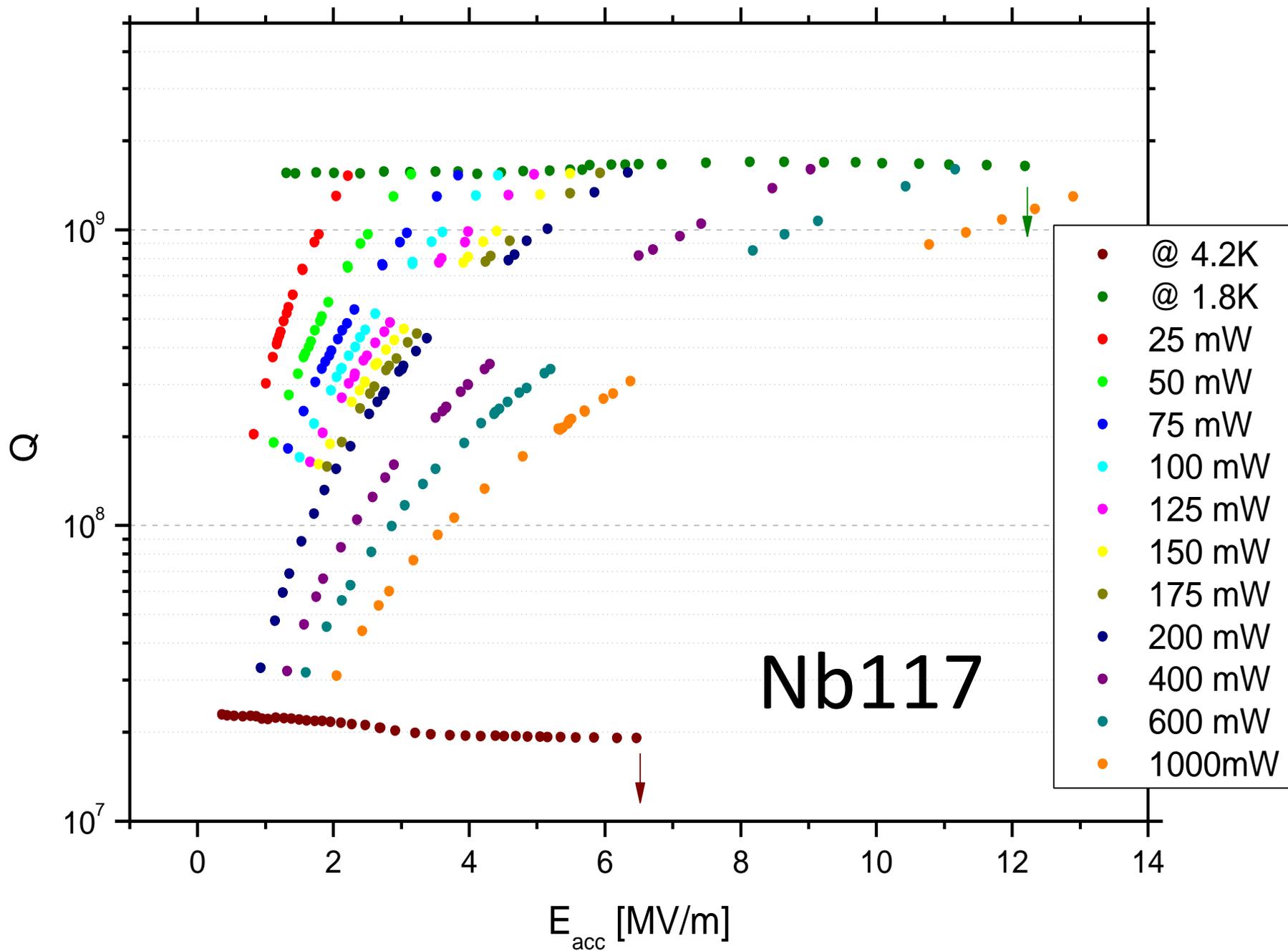
$$R_{BCS}(T_0 + \Delta T) \approx \frac{A\omega^2}{T_0} \exp\left[-\frac{sT_c}{2T_0}\right] * \exp\left[-\frac{sT_c}{2} \frac{\Delta T}{T_0^2}\right]$$

$$R_{BCS}(T_0 + \Delta T) \approx R_{BCS}(T_0) * \exp\left[-\frac{sT_c}{2} \frac{\Delta T}{T_0^2}\right]$$

$$R_S(T) \approx R_{BCS}(T) + R_{RES}$$

?????

$$R_S(T_0, W) \approx R_{BCS}(T_0) * \exp\left[-\frac{sT_C}{2} \frac{\Delta T}{T_0^2}\right] + R_{RES}$$



the Q-factor decreases linearly with W ,
but at a certain point it becomes
almost constant!

**The critical power where the losses
change slope do correspond to the
He boiling nucleation?**

Q-SLOPE ANALYSIS OF NIOBIUM SC RF CAVITIES

K.Saito[#], KEK, 1-1 Oho, Tsukuba-shi, Ibaraki-ken, Japan

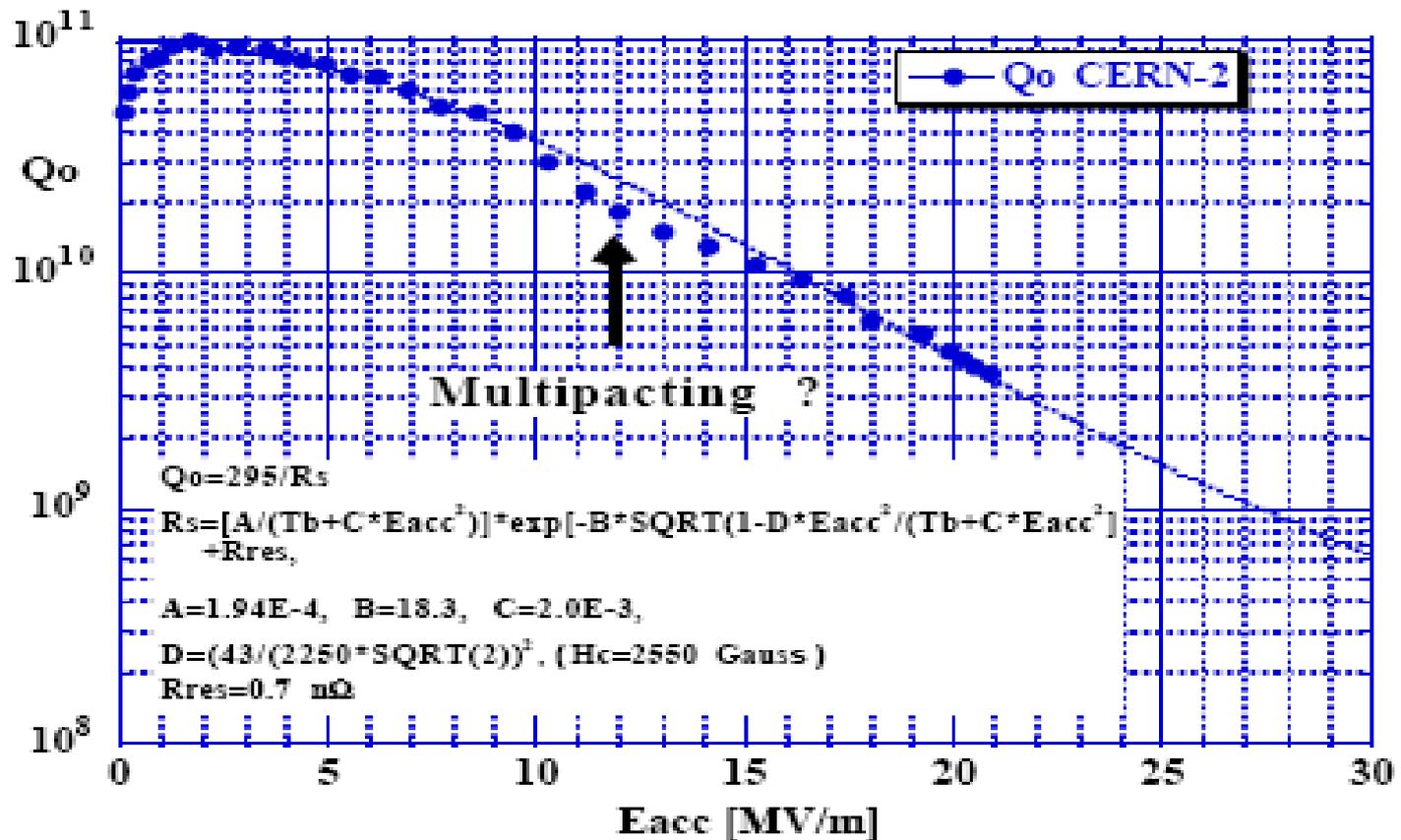


Figure 20: Q₀-E_{acc} excitation curve fitting by the combined model for the 1500MHz niobium film coated cavity at CERN.

Is it possible that

**He-II will have memory of the
boiling nucleation of He-I ?**

EXPERIMENTAL COMPARISON AT KEK OF HIGH GRADIENT PERFORMANCE OF DIFFERENT SINGLE CELL SUPERCONDUCTING CAVITY DESIGNS

F. Furuta^a, K. Saito^a, T. Saeki^a, H. Inoue^a, Y. Morozumi^a, T. Higo^a, Y. Higashi^a, H. Matsumoto^a, S. Kazakov^a, H. Yamaoka^a, K. Ueno^a, Y. Kobayashi^a, R. S. Orr^a and J. Sekutowicz^b

^aKEK High Energy Accelerator Research Organization, 1-1 Oho, Tsukuba 305-0801, Japan

^bDESY Deutsches Elektronen-Synchrotron, Notkestrasse 85, 22603 Hamburg, Germany

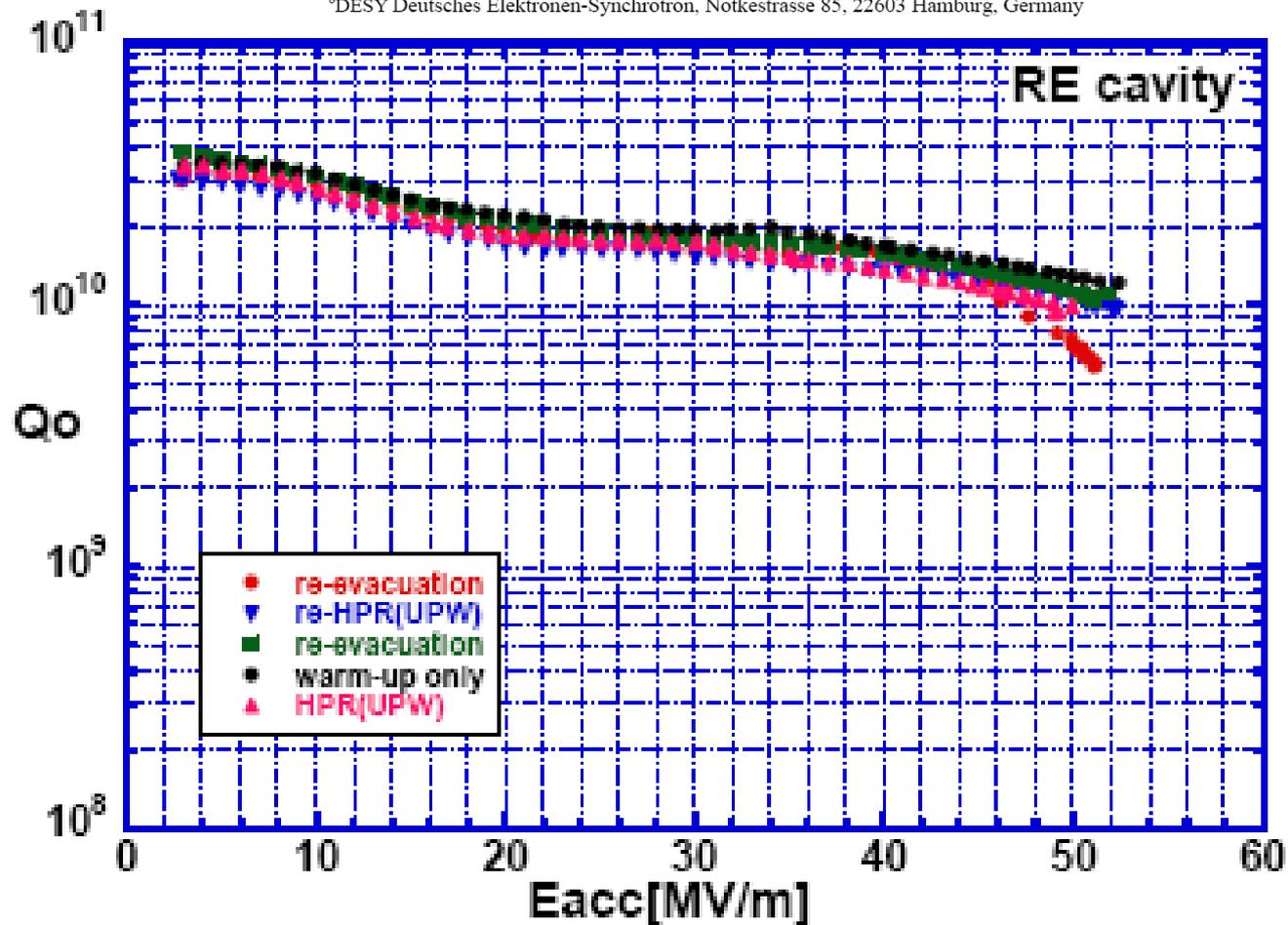
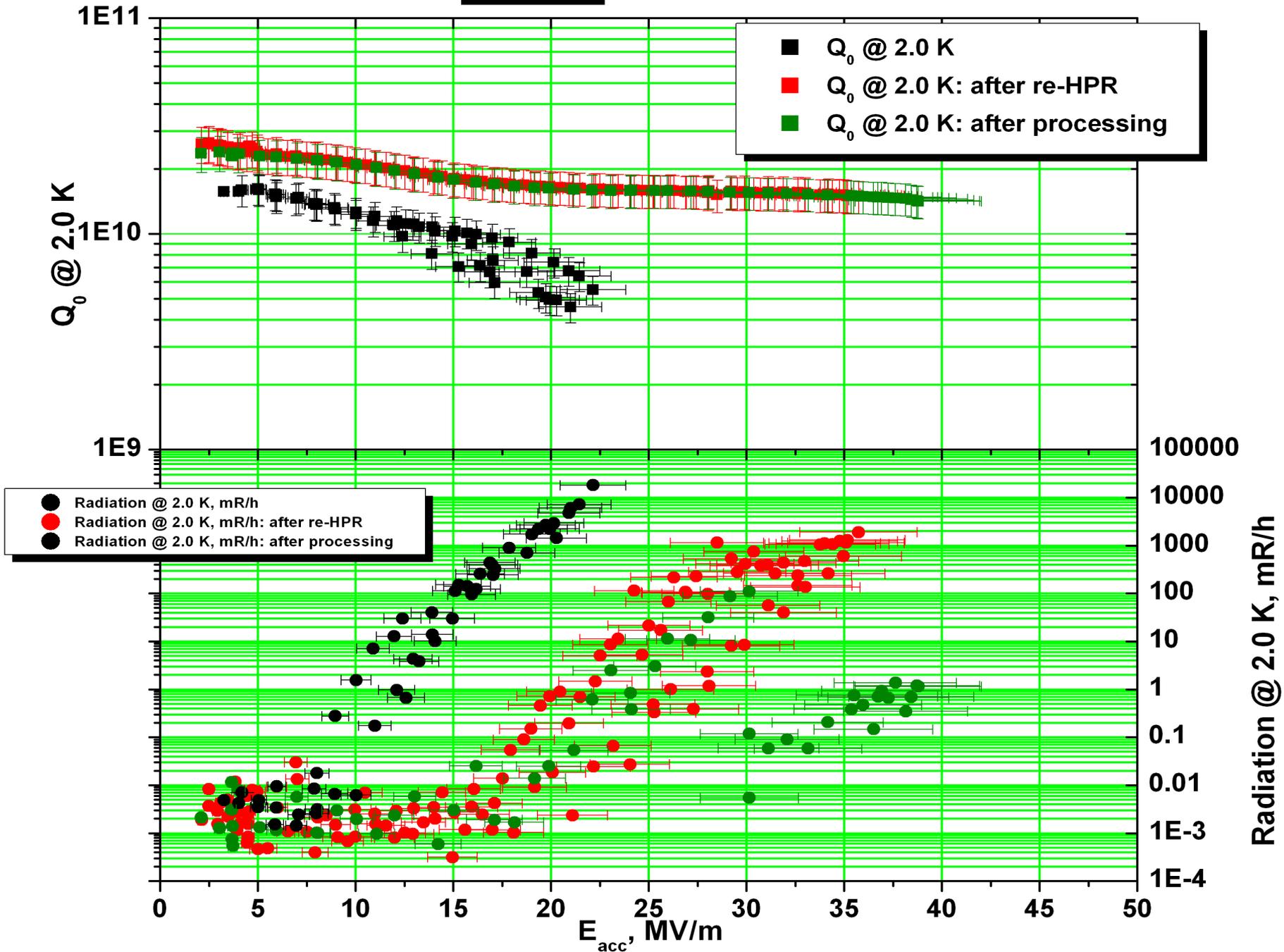
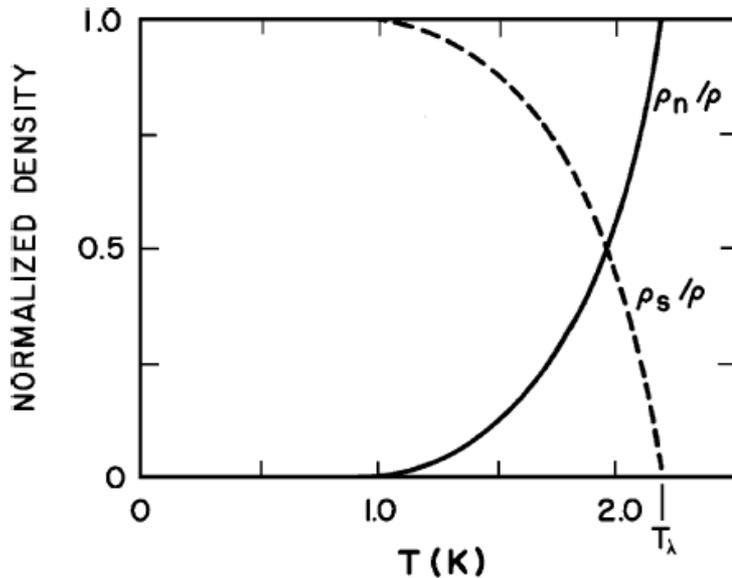


Figure 5: The reproducibility of high gradient.

TB9RI028



1.8 K is very close to T_λ ,
so at 1.8K ρ_n is $\sim 34\%$!

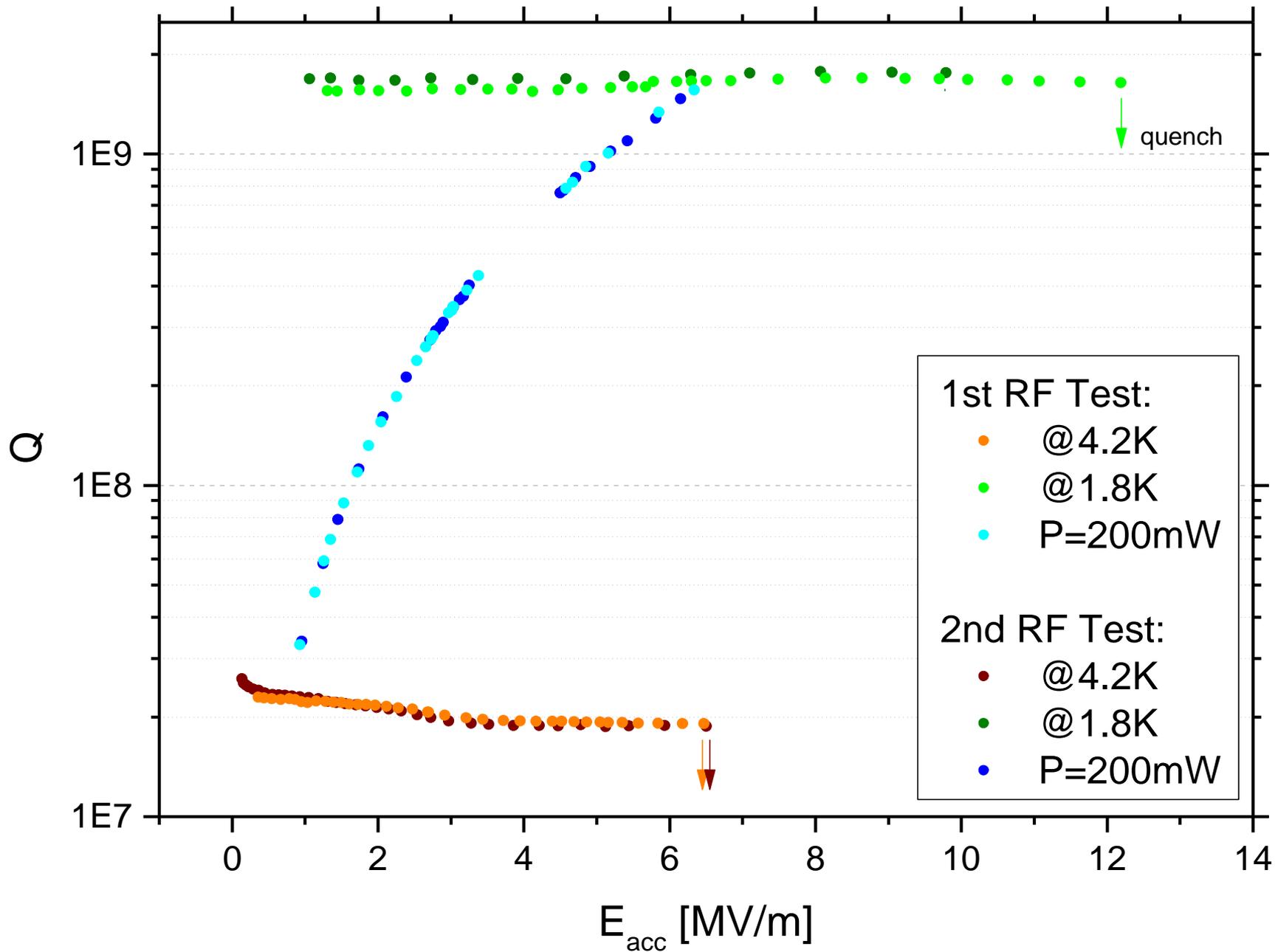


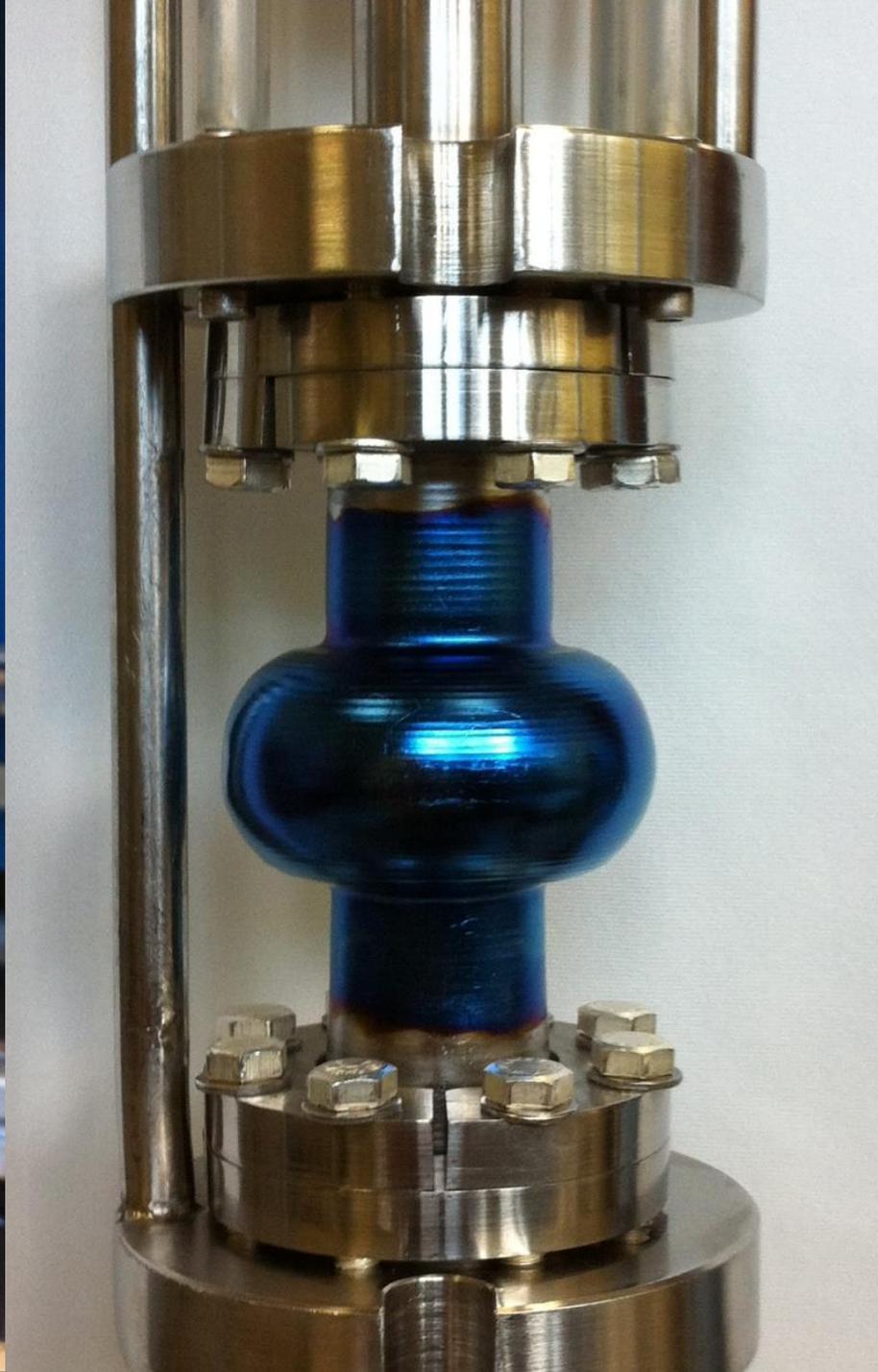
$$\rho = \rho_s + \rho_n$$

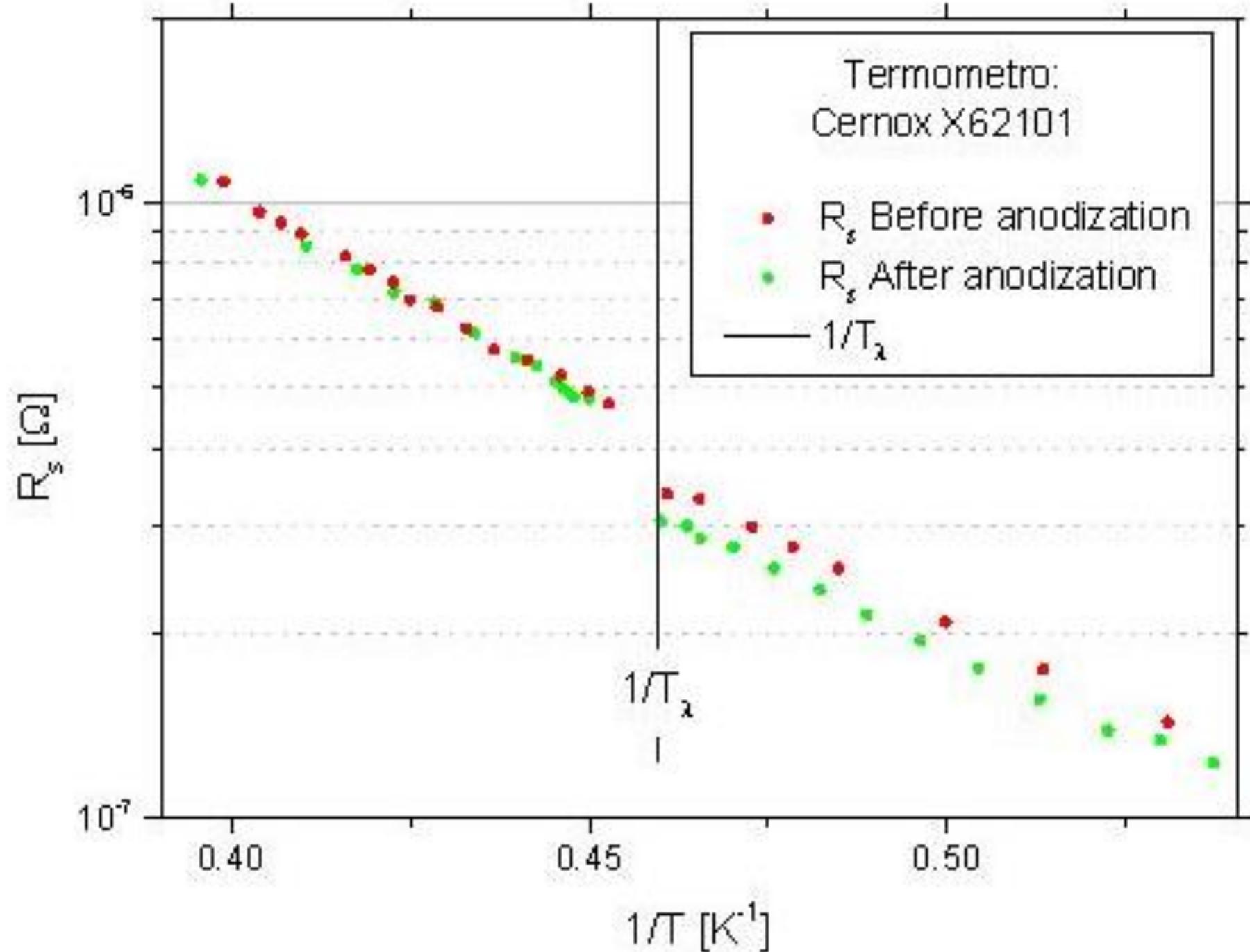
$$\frac{\rho_n}{\rho} = \left(\frac{T}{T_\lambda} \right)^{5.6} \quad \text{for } T \leq T_\lambda$$

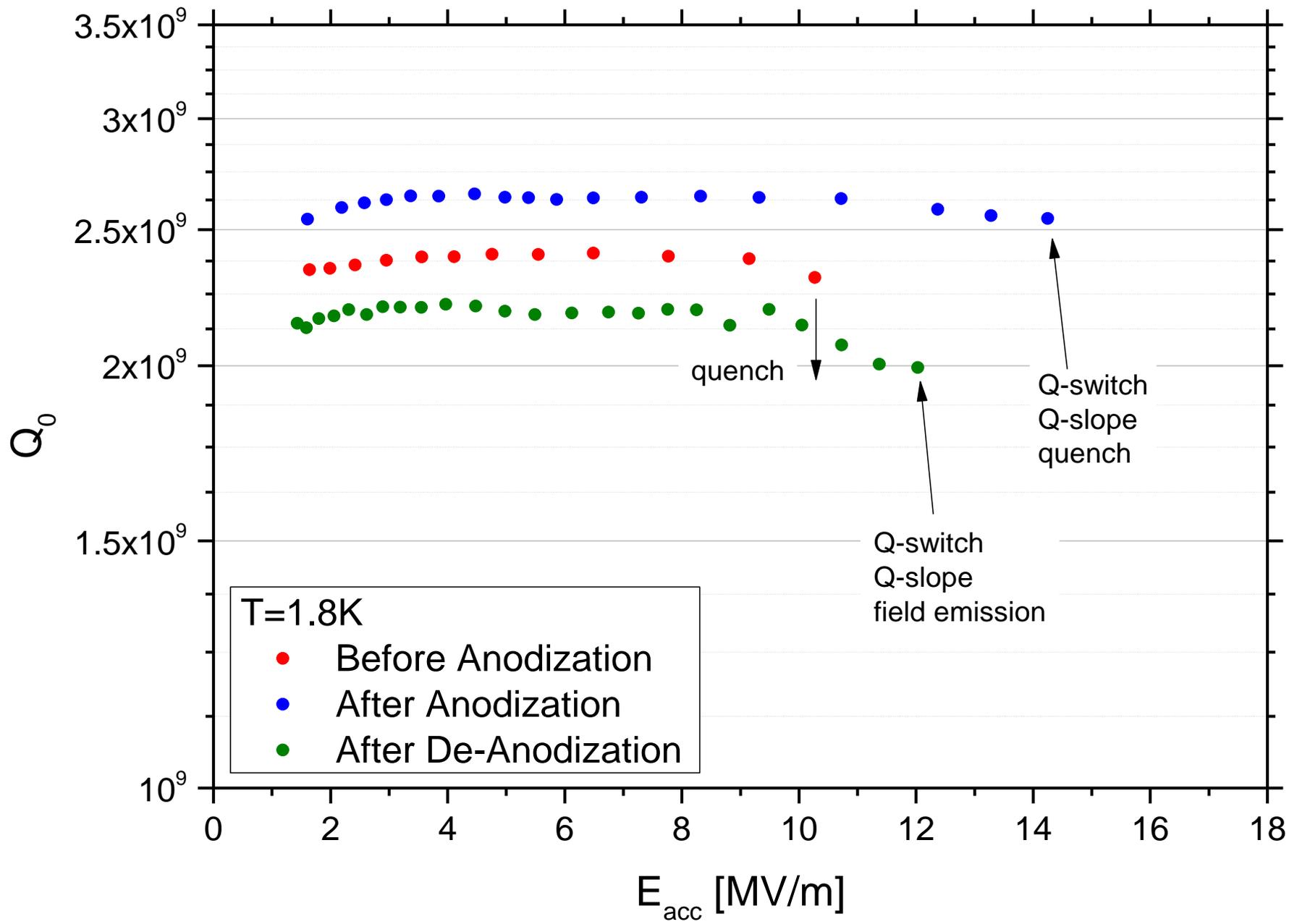


Measured twice



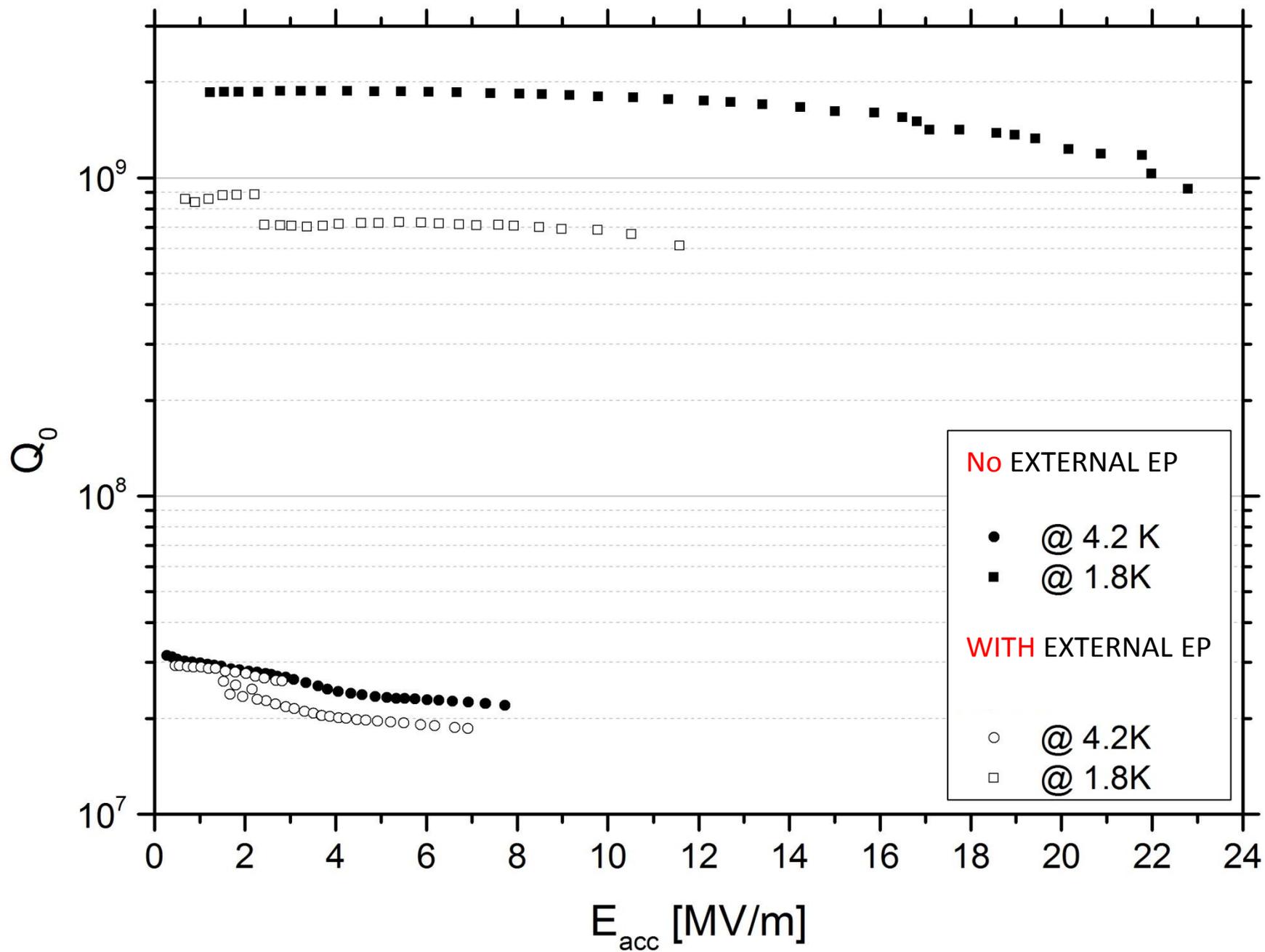






Why the **anodization is responsible of the Q increase?**

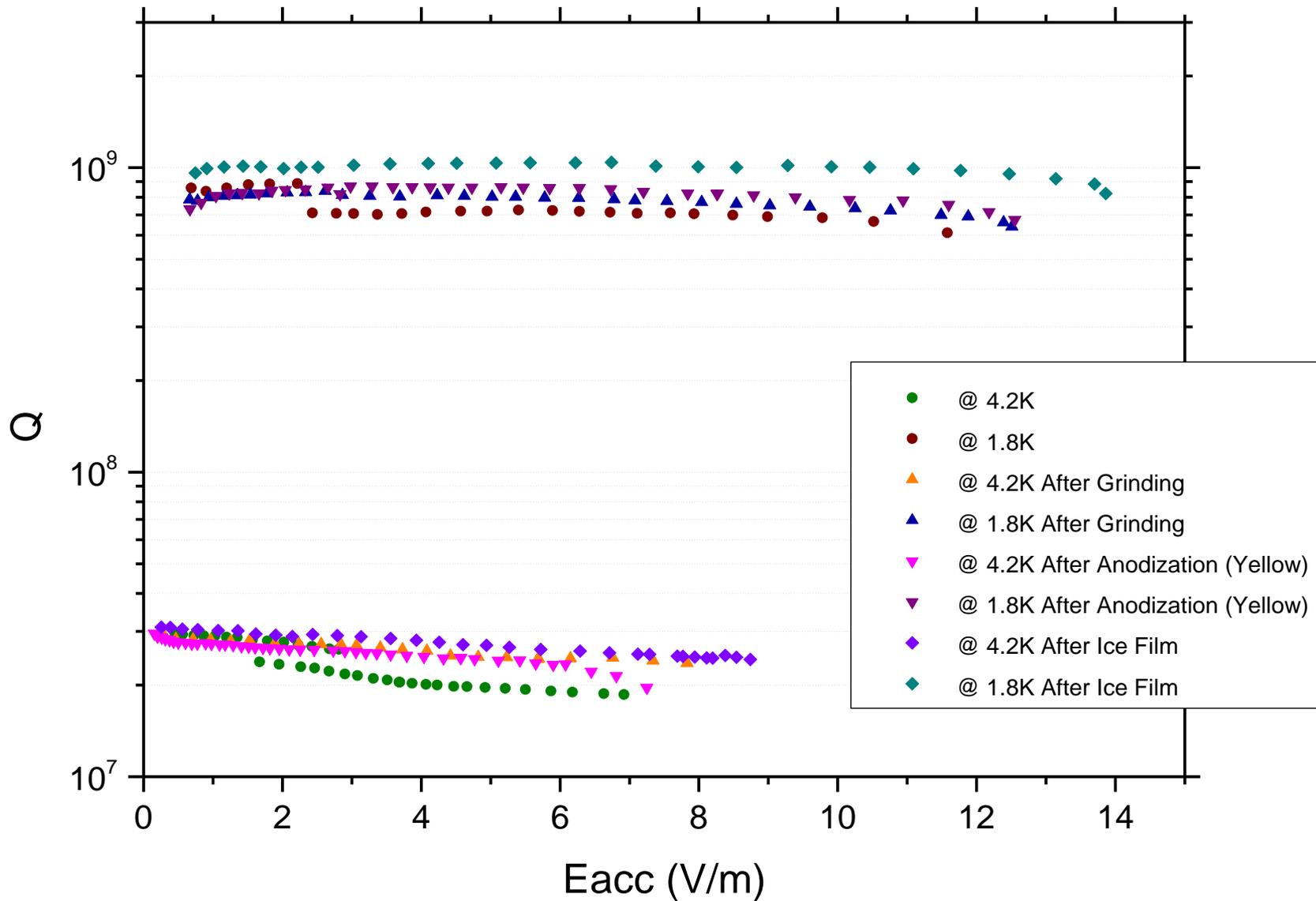
- **because of the **lower θ_D** ?**
- **because of the **boiling nucleation** on the external surface ?**
- **because the **oxide does not reflect** thermal phonons ?**
- **because of both ?**



If we mirror finish the cavity exterior surface, will this will behave as a Mirror for thermal phonons

Can water micro-cristallites
on the external surface of Nb
promote film boiling and then
positively affect cavity
performances?

Nb 127 with external EP



For years we have considered a cavity as an adiabatic system made by the RF fields + Nb, because the He bath has been considered as a stable and infinite reservoir at fixed temperature.

**Is it not the time now to
consider instead the
adiabatic system composed
by RF fields + Nb + Liquid
Helium ?**

