

Heavy Tetraquark States and Quarkonium Hybrids

Wei Chen

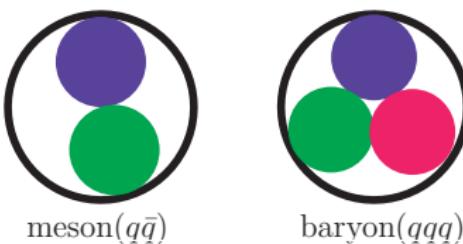
Subatomic Physics Institute, University of Saskatchewan (Canada)

IHEP, Protvino, Russia, June 23 – 27, 2014

Contents

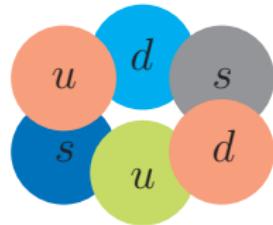
- ① Exotic Hadron States
- ② QCD Sum Rule
- ③ Heavy Tetraquark States
- ④ Quarkonium Hybrid States
- ⑤ Summary

Exotic Hadrons

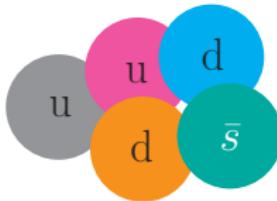


- **Quark model** is quite successful in the classification of hadrons:
 J^{PC} quantum numbers and flavour quantum numbers;
- However, the hadron structures are more complicated in **QCD**. It may allow for hadrons which lie outside the naive quark model;
- Mesons with exotic J^{PC} quantum numbers:
 $J^{PC} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, \dots$
- Baryons with exotic flavour quantum numbers:
dibaryon, pentaquark, ...

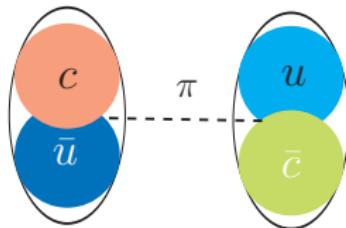
Different configurations of exotic hadron in QCD:



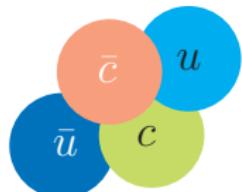
dibaryon



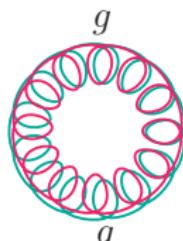
pentaquark



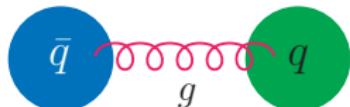
molecule



tetraquark



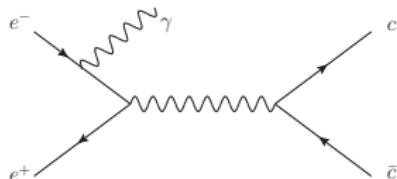
glueball



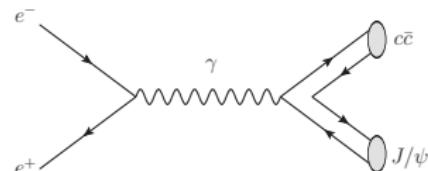
hybrid

Production mechanisms of XYZ states

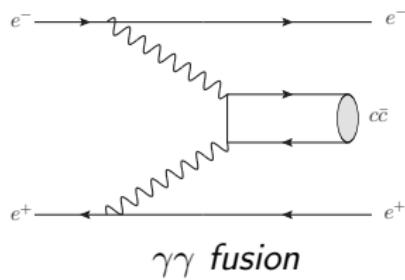
Since 2003, many new **charmonium-like states** are discovered experimentally:



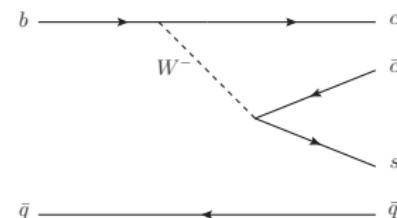
Initial State Radiation (ISR)



Double charmonium production



$\gamma\gamma$ fusion

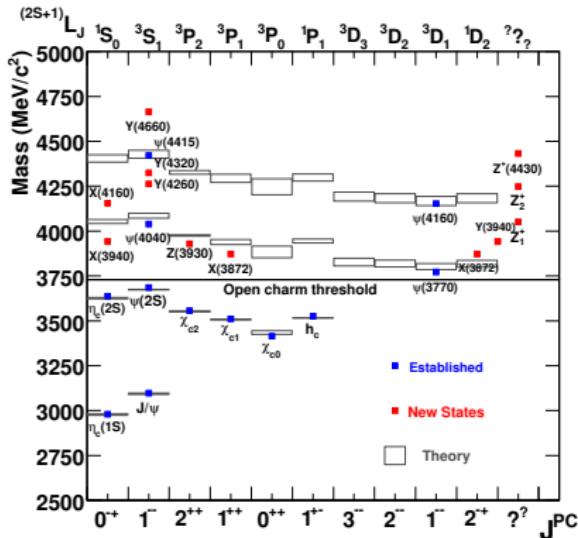


B decay

State	M (MeV)	Γ (MeV)	J^{PC}	Process (decay mode)	Experiment (# σ)	1 st observation
$X(3823)$	3823.1 ± 1.9	< 24	? $?^-$	$B \rightarrow K + (x_c c\gamma)$	Belle (3.8)	Belle 2013
$X(3872)$	3871.68 ± 0.17	< 1.2	1 $^{++}$	$B \rightarrow K + (J/\psi \pi^+ \pi^-)$	Belle (12.8), BABAR (8.6)	Belle 2003
				$p\bar{p} \rightarrow (J/\psi \pi^+ \pi^-) + \dots$	CDF (np), D \emptyset (5.2)	
				$B \rightarrow K + (J/\psi \pi^+ \pi^- \pi^0)$	Belle (4.3), BABAR (4.0)	
				$B \rightarrow K + (D^0 \bar{D}^0 \pi^0)$	Belle (6.4), BABAR (4.9)	
				$B \rightarrow K + (J/\psi \gamma)$	Belle (4.0), BABAR (3.6)	
				$B \rightarrow K + (\psi(2S) \gamma)$	BABAR (3.5), Belle (0.4)	
				$pp \rightarrow (J/\psi \pi^+ \pi^-) + \dots$	LHCb	
$X(3915)$	3917.5 ± 1.9	20 ± 5	0 $^{++}$	$B \rightarrow K + (J/\psi \omega)$	Belle (8.1), BABAR (19)	Belle 2004
				$e^+ e^- \rightarrow e^+ e^- + (J/\psi \omega)$	Belle (7.7), BABAR (7.6)	
$x_{c2}(2P)$	3927.2 ± 2.6	24 ± 6	2 $^{++}$	$e^+ e^- \rightarrow e^+ e^- + (D\bar{D})$	Belle (5.3), BABAR (5.8)	Belle 2005
$X(3940)$	3942^{+9}_{-8}	37^{+27}_{-17}	? $?^+$	$e^+ e^- \rightarrow J/\psi + (D^* \bar{D})$	Belle (6.0)	Belle 2007
				$e^+ e^- \rightarrow J/\psi + (\dots)$	Belle (5.0)	
$G(3900)$	3943 ± 21	52 ± 11	1 $^{--}$	$e^+ e^- \rightarrow \gamma + (D\bar{D})$	BABAR (np), Belle (np)	BABAR 2007
$Y(4008)$	4008^{+121}_{-49}	226 ± 97	1 $^{--}$	$e^+ e^- \rightarrow \gamma + (J/\psi \pi^+ \pi^-)$	Belle (7.4)	Belle 2007
$Y(4140)$	4144.5 ± 2.6	15^{+11}_{-7}	? $?^+$	$B \rightarrow K + (J/\psi \phi)$	CDF (5.0), CMS (>5)	CDF 2009
$X(4160)$	4156^{+29}_{-25}	139^{+113}_{-65}	? $?^+$	$e^+ e^- \rightarrow J/\psi + (D^* \bar{D}^*)$	Belle (5.5)	Belle 2007

State	M (MeV)	Γ (MeV)	$J^P C$	Process (decay mode)	Experiment (# σ)	1 st observation
$Y(4260)$	4263^{+8}_{-9}	95 ± 14	1^{--}	$e^+ e^- \rightarrow \gamma + (J/\psi \pi^+ \pi^-)$	$B\!\!AB\!\!AR$ (8.0), CLEO (5.4)	$B\!\!AB\!\!AR$ 2005
					Belle (15)	
				$e^+ e^- \rightarrow (J/\psi \pi^+ \pi^-)$	CLEO (11)	
				$e^+ e^- \rightarrow (J/\psi \pi^0 \pi^0)$	CLEO (5.1)	
$Y(4274)$	$4274.4^{+8.4}_{-6.7}$	32^{+22}_{-15}	$?^+$	$B \rightarrow K + (J/\psi \phi)$	CDF (3.1)	CDF 2010
$X(4350)$	$4350.6^{+4.6}_{-5.1}$	$13.3^{+18.4}_{-10.0}$	$0/2^{++}$	$e^+ e^- \rightarrow e^+ e^- (J/\psi \phi)$	Belle (3.2)	Belle 2009
$Y(4360)$	4361 ± 13	74 ± 18	1^{--}	$e^+ e^- \rightarrow \gamma + (\psi(2S) \pi^+ \pi^-)$	$B\!\!AB\!\!AR$ (np), Belle (8.0)	$B\!\!AB\!\!AR$ 2007
$X(4630)$	4634^{+9}_{-11}	92^{+41}_{-32}	1^{--}	$e^+ e^- \rightarrow \gamma (\Lambda_c^+ \Lambda_c^-)$	Belle (8.2)	Belle 2007
$Y(4660)$	4664 ± 12	48 ± 15	1^{--}	$e^+ e^- \rightarrow \gamma + (\psi(2S) \pi^+ \pi^-)$	Belle (5.8)	Belle 2007
$Z_c^+(3900)$	3898 ± 5	51 ± 19	$1?^-$	$Y(4260) \rightarrow \pi^- + (J/\psi \pi^+)$	BESIII (np), Belle (5.2)	BESIII 2013
$Z_1^+(4050)$	4051^{+24}_{-43}	82^{+51}_{-55}	?	$B \rightarrow K + (\chi_{c1}(1P) \pi^+)$	Belle (5.0), $B\!\!AB\!\!AR$ (1.1)	Belle 2008
$Z_2^+(4250)$	4248^{+185}_{-45}	177^{+321}_{-72}	?	$B \rightarrow K + (\chi_{c1}(1P) \pi^+)$	Belle (5.0), $B\!\!AB\!\!AR$ (2.0)	Belle 2008
$Z^+(4430)$	4443^{+24}_{-18}	107^{+113}_{-71}	?	$B \rightarrow K + (\psi(2S) \pi^+)$	Belle (6.4), $B\!\!AB\!\!AR$ (2.4)	Belle 2007
$Y_b(10888)$	10888.4 ± 3.0	$30.7^{+8.9}_{-7.7}$	1^{--}	$e^+ e^- \rightarrow (\Upsilon(nS) \pi^+ \pi^-)$	Belle (2.0)	Belle 2010
$Z_b^+(10610)$	10607.2 ± 2.0	18.4 ± 2.4	1^{+-}	$\Upsilon(5S) \rightarrow \pi^- + (\Upsilon/h_b \pi^+)$,	Belle (16)	Belle 2011
$Z_b^+(10650)$	10652.2 ± 1.5	11.5 ± 2.2	1^{+-}	$\Upsilon(5S) \rightarrow \pi^- + (\Upsilon/h_b \pi^+)$,	Belle (16)	Belle 2011

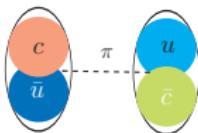
Charmonium Spectrum



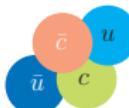
- They contain a $c\bar{c}$ pair and are above the $D\bar{D}$ threshold;
- Some are very narrow or even charged;
- Absence of the open-charm decay channels;
- Their masses are not compatible with the prediction for $c\bar{c}$ states;
- They are good candidates for exotic mesons!

Possible Interpretations

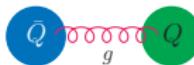
Molecule: loosely bound states composed of a pair of mesons, probably bound by the long-rang color-singlet pion exchange



Tetraquark: bound states of four quarks, bound by colored-force between quarks, decay through rearrangement, some are charged



Hybrid: composed of a pair of quarks and one excited gluonic field



In this talk, I will focus on the tetraquark states and quarkonium hybrids.

QCD Sum Rule

Two-point correlation function:

$$\Pi_{\mu\nu}(q^2) = i \int d^4x e^{iq\cdot x} \langle 0 | T[J_\mu(x) J_\nu^\dagger(0)] | 0 \rangle,$$

- Hadron level: described by the dispersion relation

$$\Pi(q^2) = \frac{(q^2)^N}{\pi} \int \frac{\text{Im}\Pi(s)}{s^N(s - q^2 - i\epsilon)} ds + \sum_{n=0}^{N-1} b_n(q^2)^n,$$

$$\begin{aligned} \rho(s) &= \frac{1}{\pi} \text{Im}\Pi(s) = \sum_n \delta(s - m_n^2) \langle 0 | J | n \rangle \langle n | J^\dagger | 0 \rangle \\ &= f_X^2 \delta(s - m_X^2) + \text{continuum}, \end{aligned}$$

- Borel transform:

$$\Pi(M_B^2) = \int_0^\infty ds \rho(s) e^{-s/M_B^2},$$

QCD Sum Rule

- Quark-gluon level: evaluated via operator product expansion(OPE)

$$\rho(s) = \rho^{\text{pert}}(s) + \rho^{\langle\bar{q}q\rangle}(s) + \rho^{\langle GG \rangle}(s) + \rho^{\langle\bar{q}q\rangle^2}(s) + \rho^{\langle\bar{q}Gq\rangle}(s) + \dots,$$

- Sum rules: quark-hadron duality

$$\mathcal{L}_k(s_0, M_B^2) = \int_{4m_Q^2}^{s_0} ds e^{-s/M_B^2} \rho(s) s^k = f_X^2 m_X^{2k} e^{-m_X^2/M_B^2},$$

- Hadron mass:

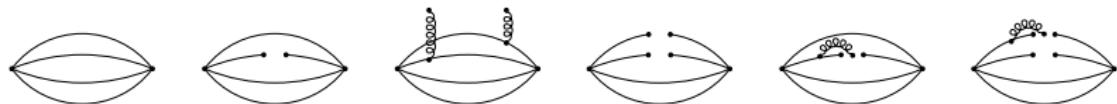
$$m_X(s_0, M_B^2) = \sqrt{\frac{\mathcal{L}_1(s_0, M_B^2)}{\mathcal{L}_0(s_0, M_B^2)}}.$$

We use this equation to predict the hadron mass in our analysis!

Tetraquark Sum Rule Analysis

Spectral density $\rho(s)$ at the quark-gluon level:

$$\rho(s) = \rho^{\text{pert}}(s) + \rho^{\langle \bar{q}q \rangle}(s) + \rho^{\langle GG \rangle}(s) + \rho^{\langle \bar{q}q \rangle^2}(s) + \rho^{\langle \bar{q}Gq \rangle}(s),$$

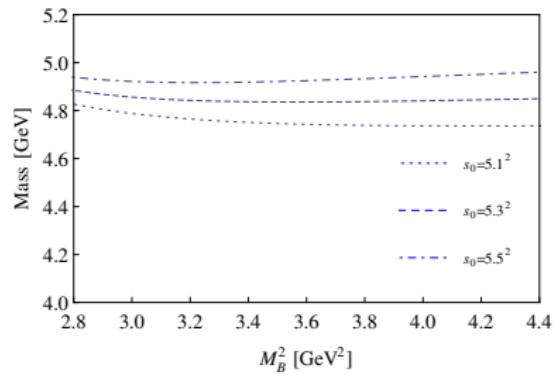
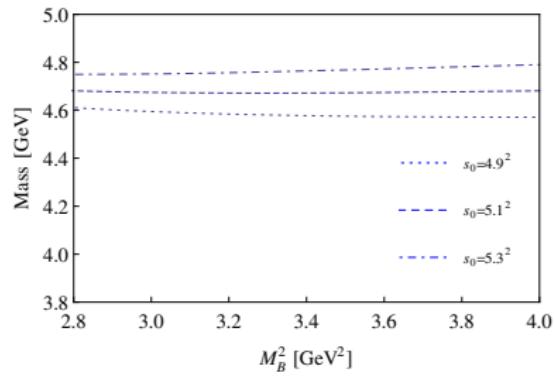
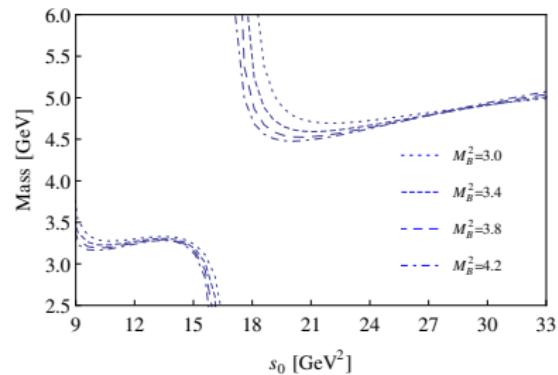
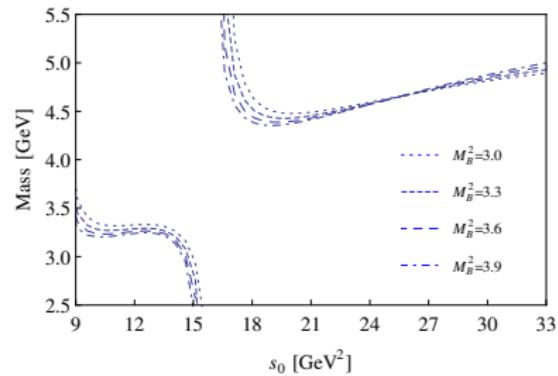


Tetraquark mass:

$$m_X(s_0, M_B^2) = \sqrt{\frac{\mathcal{L}_1(s_0, M_B^2)}{\mathcal{L}_0(s_0, M_B^2)}},$$

- We first fix the value of **threshold parameter s_0** ;
- A **Borel window** is obtained by studying the **pole contribution** and **OPE convergence**;
- Stabilities of the mass sum rules are also required.

For $qc\bar{q}\bar{c}$ and $sc\bar{s}\bar{c}$ systems with $J^{PC} = 1^{-+}$



For $qc\bar{q}\bar{c}$ and $sc\bar{s}\bar{c}$ systems with $J^{PC} = 1^{--}$

Phys. Rev. D83, 034010 (2011)

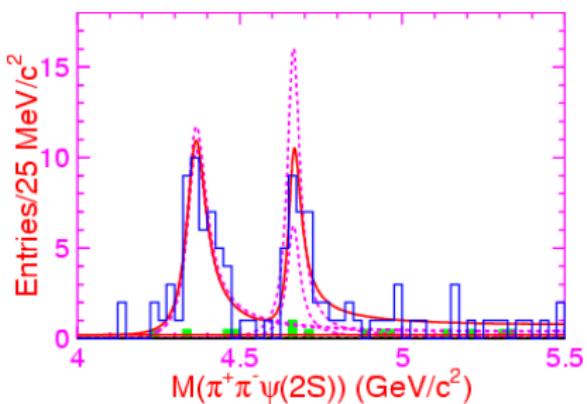
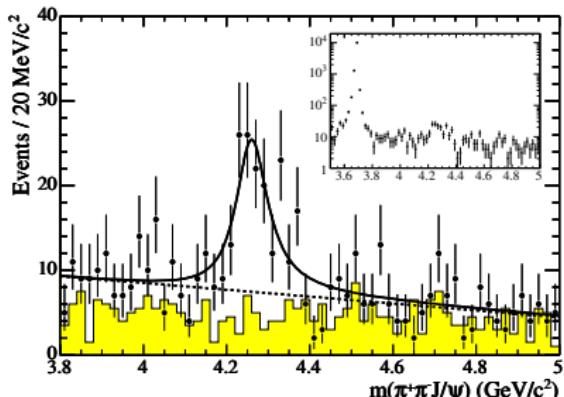
	Currents	$s_0(\text{GeV}^2)$	$[M_{\min}^2, M_{\max}^2](\text{GeV}^2)$	$m_X(\text{GeV})$	PC(%)
$qc\bar{q}\bar{c}$ system	$J_{1\mu}$	5.0^2	$2.9 \sim 3.6$	4.64 ± 0.09	44.1
	$J_{4\mu}$	5.0^2	$2.9 \sim 3.6$	4.61 ± 0.10	46.4
	$J_{7\mu}$	5.2^2	$2.9 \sim 4.1$	4.74 ± 0.10	47.3
$sc\bar{s}\bar{c}$ system	$J_{1\mu}$	5.4^2	$2.8 \sim 4.5$	4.92 ± 0.10	50.3
	$J_{2\mu}$	5.0^2	$2.8 \sim 3.5$	4.64 ± 0.09	48.6
	$J_{3\mu}$	4.9^2	$2.8 \sim 3.4$	4.52 ± 0.10	45.6
	$J_{4\mu}$	5.4^2	$2.8 \sim 4.5$	4.88 ± 0.10	51.7
	$J_{7\mu}$	5.3^2	$2.8 \sim 4.3$	4.86 ± 0.10	46.0
	$J_{8\mu}$	4.8^2	$2.8 \sim 3.1$	4.48 ± 0.10	43.2

The masses for $qc\bar{q}\bar{c}$ states are around **4.6 – 4.7 GeV**.

$Y(J^{PC} = 1^{--})$ family

Since 2005, *BABAR* and *Belle* discovered $Y(4260)$, $Y(4360)$ and $Y(4660)$ in the ISR process:

- $Y(4260)$: $e^+e^- \rightarrow \gamma_{ISR} J/\psi \pi^+ \pi^-$, PRL95, 142001(2005)
- $Y(4360)$, $Y(4660)$: $e^+e^- \rightarrow \gamma_{ISR} \psi(2S) \pi^+ \pi^-$, PRL99, 142002(2007)
- $Y(4660)$: $m = 4664 \pm 11 \pm 5$ MeV, $\Gamma = 48 \pm 15 \pm 3$ MeV



Our result supports $Y(4660)$ as a $qc\bar{q}\bar{c}$ tetraquark state!

For $qc\bar{q}\bar{c}$, $sc\bar{s}\bar{c}$, $qb\bar{q}\bar{b}$ and $sb\bar{s}\bar{b}$ systems with $J^{PC} = 1^{++}$

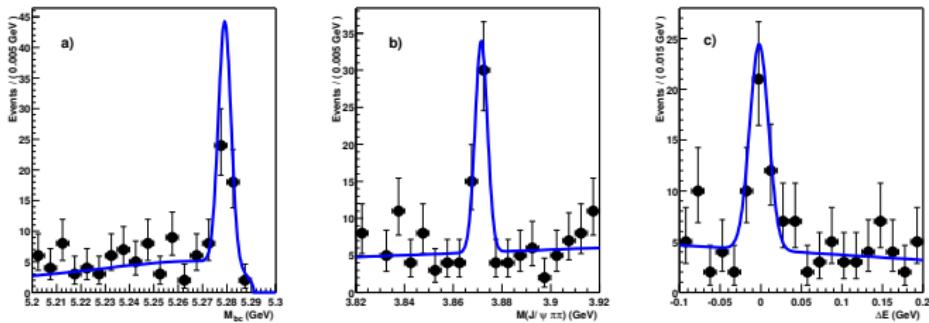
Phys. Rev. D83, 034010 (2011)

	Currents	$s_0(\text{GeV}^2)$	$[M_{\min}^2, M_{\max}^2](\text{GeV}^2)$	$m_x(\text{GeV})$	PC(%)
$qc\bar{q}\bar{c}$ system	$J_{3\mu}$	4.6^2	$3.0 \sim 3.4$	4.19 ± 0.10	47.3
	$J_{4\mu}$	4.5^2	$3.0 \sim 3.3$	4.03 ± 0.11	46.8
$sc\bar{s}\bar{c}$ system	$J_{3\mu}$	4.6^2	$3.0 \sim 3.4$	4.22 ± 0.10	45.7
	$J_{4\mu}$	4.5^2	$3.0 \sim 3.3$	4.07 ± 0.10	44.4
$qb\bar{q}\bar{b}$ system	$J_{3\mu}$	10.9^2	$8.5 \sim 9.5$	10.32 ± 0.09	47.0
	$J_{4\mu}$	10.8^2	$8.5 \sim 9.2$	10.22 ± 0.11	44.6
	$J_{7\mu}$	10.7^2	$7.8 \sim 8.4$	10.14 ± 0.10	44.8
	$J_{8\mu}$	10.7^2	$7.8 \sim 8.4$	10.14 ± 0.09	44.8
$sb\bar{s}\bar{b}$ system	$J_{3\mu}$	10.9^2	$8.5 \sim 9.5$	10.34 ± 0.09	46.1
	$J_{4\mu}$	10.8^2	$8.5 \sim 9.1$	10.25 ± 0.10	43.3
	$J_{7\mu}$	10.8^2	$7.5 \sim 8.6$	10.24 ± 0.11	47.1
	$J_{8\mu}$	10.8^2	$7.5 \sim 8.6$	10.24 ± 0.10	47.1

The extracted mass **4.0 GeV** is slightly above the mass of **X(3872)**.

First New State: $X(3872)$

- In 2003, Belle discovered $X(3872)$ in $B^+ \rightarrow K^+ J/\psi \pi^+ \pi^-$ with $m = 3872.0 \pm 0.6 \pm 0.5 \text{ MeV}$, $\Gamma < 2.3 \text{ MeV}$, PRL91, 262001(2003)
- In 2013, LHCb (PRL110, 222001(2013)) determined the quantum numbers of $X(3872)$: $J^{PC} = 1^{++}$



Our results **don't preclude** the tetraquark interpretation of $X(3872)$.

For $qc\bar{q}\bar{c}$, $sc\bar{s}\bar{c}$, $qb\bar{q}\bar{b}$ and $sb\bar{s}\bar{b}$ systems with $J^{PC} = 1^{+-}$

Phys. Rev. D83, 034010 (2011)

	Currents	$s_0(\text{GeV}^2)$	$[M_{\min}^2, M_{\max}^2](\text{GeV}^2)$	$m_X(\text{GeV})$	PC(%)
$qc\bar{q}\bar{c}$ system	$J_{3\mu}$	4.6^2	$3.0 - 3.4$	4.16 ± 0.10	46.2
	$J_{4\mu}$	4.5^2	$3.0 - 3.3$	4.02 ± 0.09	44.6
	$J_{5\mu}$	4.5^2	$3.0 - 3.4$	4.00 ± 0.11	46.0
	$J_{6\mu}$	4.6^2	$3.0 - 3.4$	4.14 ± 0.09	47.0
$sc\bar{s}\bar{c}$ system	$J_{3\mu}$	4.7^2	$3.0 - 3.6$	4.24 ± 0.10	49.6
	$J_{4\mu}$	4.6^2	$3.0 - 3.5$	4.12 ± 0.11	47.3
	$J_{5\mu}$	4.5^2	$3.0 - 3.3$	4.03 ± 0.11	44.2
	$J_{6\mu}$	4.6^2	$3.0 - 3.4$	4.16 ± 0.11	46.0
$qb\bar{q}\bar{b}$ system	$J_{3\mu}$	10.6^2	$7.5 - 8.5$	10.08 ± 0.10	45.9
	$J_{4\mu}$	10.6^2	$7.5 - 8.5$	10.07 ± 0.10	46.2
	$J_{5\mu}$	10.6^2	$7.5 - 8.4$	10.05 ± 0.10	45.3
	$J_{6\mu}$	10.7^2	$7.5 - 8.7$	10.15 ± 0.10	47.6
$sb\bar{s}\bar{b}$ system	$J_{3\mu}$	10.6^2	$7.5 - 8.3$	10.11 ± 0.10	43.8
	$J_{4\mu}$	10.6^2	$7.5 - 8.4$	10.10 ± 0.10	44.1
	$J_{5\mu}$	10.6^2	$7.5 - 8.3$	10.08 ± 0.10	43.7
	$J_{6\mu}$	10.7^2	$7.5 - 8.5$	10.18 ± 0.10	46.5

In 2013, BESIII discovered two charged states $Z_c(3900)$ and $Z_c(4025)$:

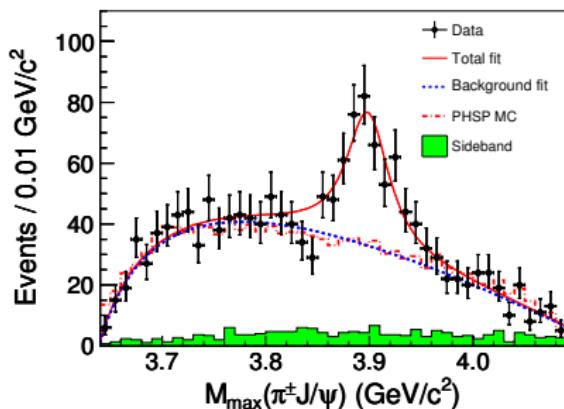
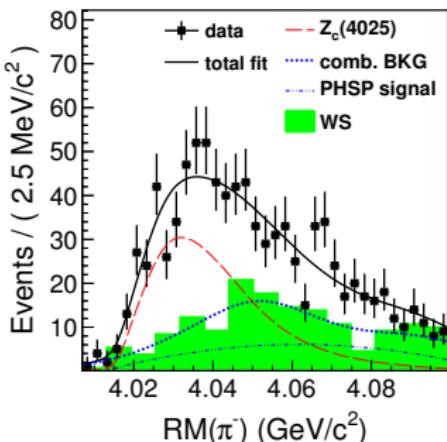
- $Z_c(3900)$: $e^+e^- \rightarrow \pi^+\pi^- J/\psi$, (PRL110, 252001(2013))

$$m = 3899.0 \pm 3.6 \pm 4.9 \text{ MeV}, \Gamma = 46 \pm 10 \pm 20 \text{ MeV}$$

- $Z_c(4025)$: $e^+e^- \rightarrow (D^*\bar{D}^*)^\pm\pi^\mp$, (PRL112, 132001(2014))

$$m = 4026.3 \pm 2.6 \pm 3.7 \text{ MeV}, \Gamma = 24.8 \pm 5.6 \pm 7.7 \text{ MeV}$$

They prefer $J^{PC} = 1^{+-}$



Our result supports $Z_c(3900)$ and $Z_c(4025)$ as $qc\bar{q}\bar{c}$ tetraquark states!

Hybrid Sum Rules

Two-point Correlation Function:

$$\Pi_{\mu\nu}(q^2) = i \int d^4x e^{iq \cdot x} \langle 0 | T[J_\mu(x) J_\nu^\dagger(0)] | 0 \rangle,$$

where $J_\mu(x)$ is the hybrid interpolating currents

$$J_\mu = g_s \bar{Q} \frac{\lambda^a}{2} \gamma^\nu G_{\mu\nu}^a Q, \quad J^{PC} = 1^{-+}, 0^{++},$$

$$J_\mu = g_s \bar{Q} \frac{\lambda^a}{2} \gamma^\nu \gamma_5 G_{\mu\nu}^a Q, \quad J^{PC} = 1^{+-}, 0^{--},$$

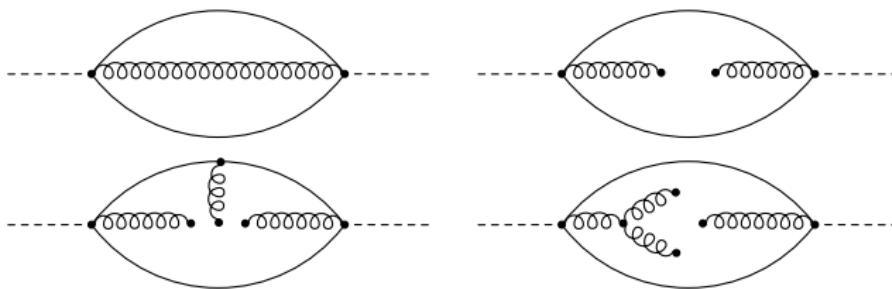
$$J_{\mu\nu} = g_s \bar{Q} \frac{\lambda^a}{2} \sigma_\mu^\alpha \gamma_5 G_{\alpha\nu}^a Q, \quad J^{PC} = 2^{-+}, 1^{++}, 1^{-+}, 0^{-+},$$

By replacing $G_{\mu\nu}^a$ with $\tilde{G}_{\mu\nu}^a = \frac{1}{2} \epsilon_{\mu\nu\alpha\beta} G^{\alpha\beta,a}$, we obtain the corresponding \tilde{J}_μ and $\tilde{J}_{\mu\nu}$ with opposite parities.

Sum Rule Analysis

- Quark-gluon level: $\rho(s)$ is calculated up to dimension six:

$$\rho(s) = \rho^{\text{pert}}(s) + \rho^{\langle GG \rangle}(s) + \rho^{\langle GGG \rangle} + \rho^{\langle jj \rangle},$$



- Fix on the value of threshold value s_0
- Determine the Borel window by studying OPE convergence and PC
- The Borel curves should be stable

Spectrum for $\bar{c}Gc$ charmonium hybrids

JHEP09(2013)019

J^{PC}	$s_0(\text{GeV}^2)$	$[M_{\min}^2, M_{\max}^2](\text{GeV}^2)$	$m_X(\text{GeV})$	PC(%)
1^{--}	15	$2.5 \sim 4.8$	3.36 ± 0.15	18.3
0^{-+}	16	$5.6 \sim 7.0$	3.61 ± 0.21	15.4
1^{-+}	17	$4.6 \sim 6.5$	3.70 ± 0.21	18.8
2^{-+}	18	$3.9 \sim 7.2$	4.04 ± 0.23	26.0
0^{+-}	20	$6.0 \sim 7.4$	4.09 ± 0.23	15.5
2^{++}	23	$3.9 \sim 7.5$	4.45 ± 0.27	21.5
1^{+-}	24	$2.5 \sim 8.4$	4.53 ± 0.23	33.2
1^{++}	30	$4.6 \sim 11.4$	5.06 ± 0.44	30.4
0^{++}	34	$5.6 \sim 14.6$	5.34 ± 0.45	36.3
0^{--}	35	$6.0 \sim 12.3$	5.51 ± 0.50	31.0

- Unstable channels are stabilized and the mass predictions are reliable!
- 1^{++} charmonium hybrid is much heavier than $X(3872)$ while 1^{--} hybrid is lighter than $Y(4260)$, which seem to preclude a pure hybrid interpretation for these states.

Spectrum for $\bar{b}Gb$ bottomonium hybrids

JHEP09(2013)019				
J^{PC}	$s_0(\text{GeV}^2)$	$[M_{\min}^2, M_{\max}^2](\text{GeV}^2)$	$m_X(\text{GeV})$	PC(%)
1^{--}	105	$11 \sim 17$	9.70 ± 0.12	17.2
0^{-+}	104	$14 \sim 16$	9.68 ± 0.29	17.3
1^{-+}	107	$13 \sim 19$	9.79 ± 0.22	20.4
2^{-+}	105	$12 \sim 19$	9.93 ± 0.21	21.7
0^{+-}	114	$14 \sim 19$	10.17 ± 0.22	17.6
2^{++}	120	$12 \sim 20$	10.64 ± 0.33	19.7
1^{+-}	123	$10 \sim 21$	10.70 ± 0.53	28.5
1^{++}	134	$13 \sim 27$	11.09 ± 0.60	27.7
0^{++}	137	$13 \sim 31$	11.20 ± 0.48	30.0
0^{--}	142	$14 \sim 25$	11.48 ± 0.75	24.1

We have confirmed the hybrid supermultiplet structure:

- **Lightest hybrid supermultiplet**: negative-parity states with $J^{PC} = 1^{--}, (0, 1, 2)^{-+}$;
- **Heavier hybrid supermultiplet**: positive-parity states with $J^{PC} = (0, 1)^{+-}, (0, 1, 2)^{++}$;
- **Heaviest 0^{--} hybrid** may suggest a highly excited gluonic structure.

Summary

Spectra of the quarkonium-like tetraquark states

- Construct all tetraquark interpolating currents in a systematic way;
- Calculate the two-point correlation functions;
- Obtain the spectra of $qc\bar{q}\bar{c}$ and $qb\bar{q}\bar{b}$ tetraquark states in QSR;
- **Support** the $qc\bar{q}\bar{c}$ tetraquark interpretation of $Y(4660)$;
- **Don't preclude** the tetraquark interpretation of $X(3872)$;
- **Support** $Z_c(3900)$, $Z_c(4025)$ states to be 1^{+-} $qc\bar{q}\bar{c}$ tetraquark states.

Spectra of the quarkonium hybrids

- Unstable channels are stabilized by the dimension six condensates;
- Predict the masses of the **exotic hybrid channels**: 1^{-+} , 0^{+-} , 0^{--} ;
- **Preclude** the pure hybrid interpretations of $X(3872)$ and $Y(4260)$;
- Confirm the **supermultiplet structures** of the hybrid spectrum.

THANK YOU VERY MUCH!