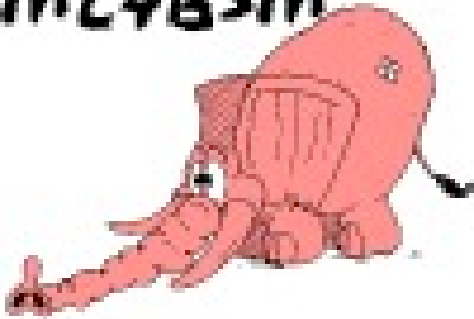


MC4BSM-3: Welcome and Orientation

MC4BSM



G. Azuelos, U. Montreal/ATLAS

C. Grojean, CERN/Theory

J. Hubisz, ANL/Theory

B. Kersevan, Jozef Stefan/ATLAS

J. Lykken, FNAL/Theory

F. Maltoni, UCL/Theory

K. Matchev, U. Florida/Theory

F. Moortgat, ETH Zurich/CMS

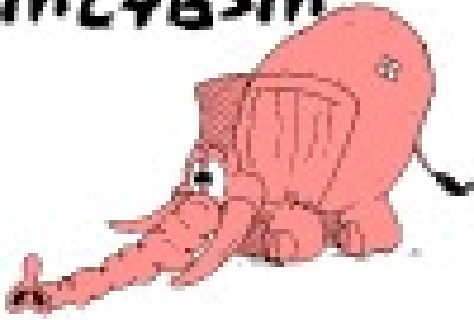
S. Mrenna, FNAL/Theory

M. Perelstein, Cornell/Theory

P. Skands, CERN+FNAL/Theory

J. Wells, CERN+Michigan/Theory

MC4BSM



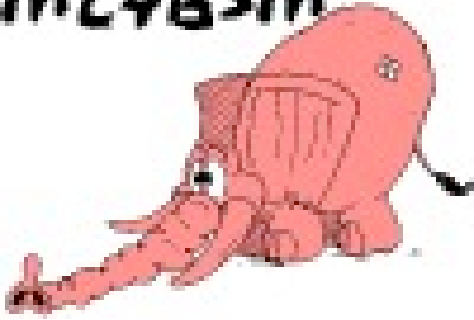
WORKSHOP INFO

- **LAPTOPS** need to be registered with CERN:

[https://network.cern.ch/sc/fcgi/
sc.fcgi?Action=VisitorRegistration](https://network.cern.ch/sc/fcgi/sc.fcgi?Action=VisitorRegistration)

- **COFFEE BREAKS** in TH common room (3rd floor)
- **RECEPTION** tonight at
- After-workshop event **TOURS** Wed 12 Mar
 - **CMS** at 10:00
 - **ATLAS** at 14:45
- A few spaces still left! People interested in joining either should contact **Peter Skands** in the coffee breaks

MC4BSM



BSM Physics Overview

“A (brief and not unweighted) random walk through the theory landscape”

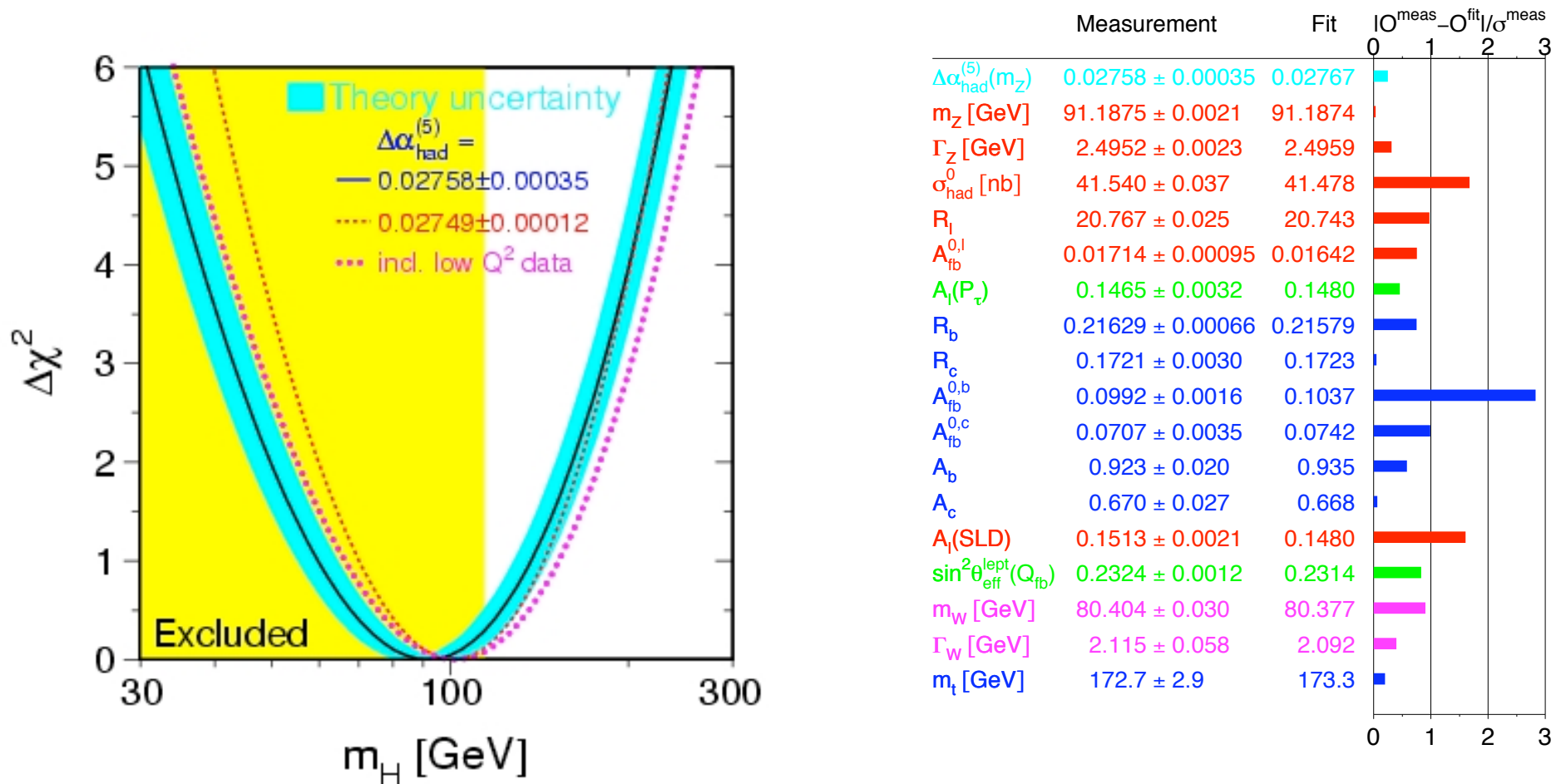
M. Perelstein, Cornell

MC4BSM-3 Workshop, CERN, March 10 2008

Introduction

- **Standard Model:** Electroweak gauge symmetry $SU(2) \times U(1)$ is **fundamental**, but **spontaneously broken** at low energies down to e&m $U(1)$
- Uncovering the **mechanism** of electroweak symmetry breaking (**EWSB**) is the central question for the LHC
- The Standard Model explanation of EWSB: **Higgs phenomenon**
- Postulate a new particle - the **Higgs boson** - of spin 0
- Vacuum is filled with **Higgs condensate**, which breaks the symmetry

Is the Higgs Really There?



- Standard Model with a **light Higgs** provides a good fit to all data, indirect determination of H mass:

$$M_H < 186 \text{ GeV} \quad (95\% \text{ c.l.})$$

Light Higgs → New Physics

- No **elementary spin-0** particles are known to exist: scalar mass is **unstable** with respect to radiative corrections
- If SM is valid up to some UV scale Λ where **new physics** occurs and divergent loop integrals are cut off, we expect

$$\mu^2(M_{\text{ew}}) = \mu^2(\Lambda) + c_1 \frac{1}{16\pi^2} \Lambda^2 + c_2 \frac{1}{16\pi^2} \log\left(\frac{\Lambda}{M_{\text{ew}}}\right) + \text{finite}$$

with $c_1 \sim 1$

- Unless the 1st and 2nd terms cancel each other precisely (“fine-tuning”), we should get $\mu \sim \Lambda/(4\pi)$
- New physics must occur at **~ 1 TeV scale!**

Supersymmetry

- Supersymmetry, with superpartners around 1 TeV, would do the job
- The minimal SUSY model - the **MSSM** - is the best-studied by far of the BSM models: successful gauge unification, etc.
- However, **non-observation** of the Higgs at LEP2 presents a significant problem for the MSSM
- At tree level, there is a **firm upper bound** on the mass of the lightest of the two CP-even Higgs bosons:

$$m(h^0) < M_Z$$

- **Experimentally,** $m(h^0) > 114 \text{ GeV}$
- Either the MSSM is **wrong**, or **loop corrections** to $m(h^0)$ are **large** (25%)

Supersymmetry and Tuning

- Loop corrections to the Higgs potential change **the Higgs vev**, not just its mass!
- Parameters can always be adjusted to reproduce the known Higgs vev...
- BUT, if the loop contribution to the vev is **BIG**, it would need to be cancelled precisely by other terms \Rightarrow classic example of **fine-tuning!**
- So: need loops to change Higgs **mass by a lot** while **not** changing the Higgs **vev by a lot** \Rightarrow **difficult!**
- Fine-tuning of **at least 1%** is needed to make things work
- If SUSY is realized, it may very well be a **non-minimal** version (e.g. extra scalars coupled to the Higgs sector)

Non-Minimal Supersymmetry

- Models with one extra singlet studied extensively (NMSSM)

$$W = \lambda S H_1 H_2 + \frac{\kappa}{3} S^3$$

- Typically large λ is required to raise m_h ; λ hits Landau pole before the GUT scale, **extra structure** is required to preserve coupling unification (e.g. “Fat Higgs”)
- Both scalars and **inos** are affected: singlino admixture in neutralinos!
- Example of a realistic simulation in NMSSM: Cornell 2007 LHC Olympics Black Box
- Used Madgraph/Madevent; took **~hour** to “add” the model
- MG/ME \Rightarrow (BRIDGE) \Rightarrow LHA file \Rightarrow PYTHIA \Rightarrow (PGS): realistic events **within hours!**
- More complex models are easily possible - **flexibility** is key!!

Quantum Gravity at TeV

- At Planck scale, SM has to be embedded into a theory with quantum gravity - string theory?
- It is believed that that theory must be **finite** - all divergences cut off at M_{Pl}
- If $M_{\text{Pl}} \sim 1 \text{ TeV}$, there is **no** hierarchy problem!
- **ADD model**: SM on a 4D brane inside higher-D space, with extra dimensions compactified with

$$R \sim M_{\text{Pl}}^{-1} \left(\frac{M_{\text{Pl},4}}{M_{\text{Pl}}} \right)^{2/n} \gg M_{\text{Pl}}^{-1}$$

- At $E < M_{\text{Pl}}$, model-independent **missing energy** signature due to graviton emission into the extra dimensions
- If two partons collide at super-plankian energies $E \gg M_{\text{Pl}}$, a **black hole** must form!

[see talk by Godang]

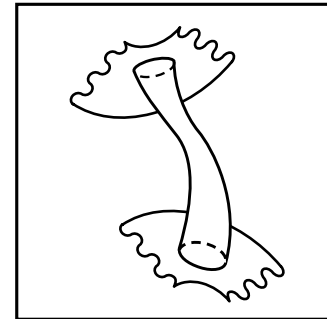
String Theory at TeV

- Given existing constraints on M_{Pl} , it seems pretty **unlikely** that the LHC will probe the region $E \gg M_{\text{Pl}}$

- In any (weakly coupled) string theory, **Regge excitations** of SM particles lie below Planck scale

$$M_n = \sqrt{n}M_S, \quad M_S \ll M_{\text{Pl}}$$

- Reggeons appear as s-channel **resonances** in SM scattering processes!



- Easy to see, more realistic target than BHs
- But need to be **distinguished** from Zprimes etc.
- Excited Reggeons have **spin > 2** \Rightarrow at present not handled by the general-purpose MC generators!

QCD Redux: Composite Higgs, Technicolor, and Their Cousins

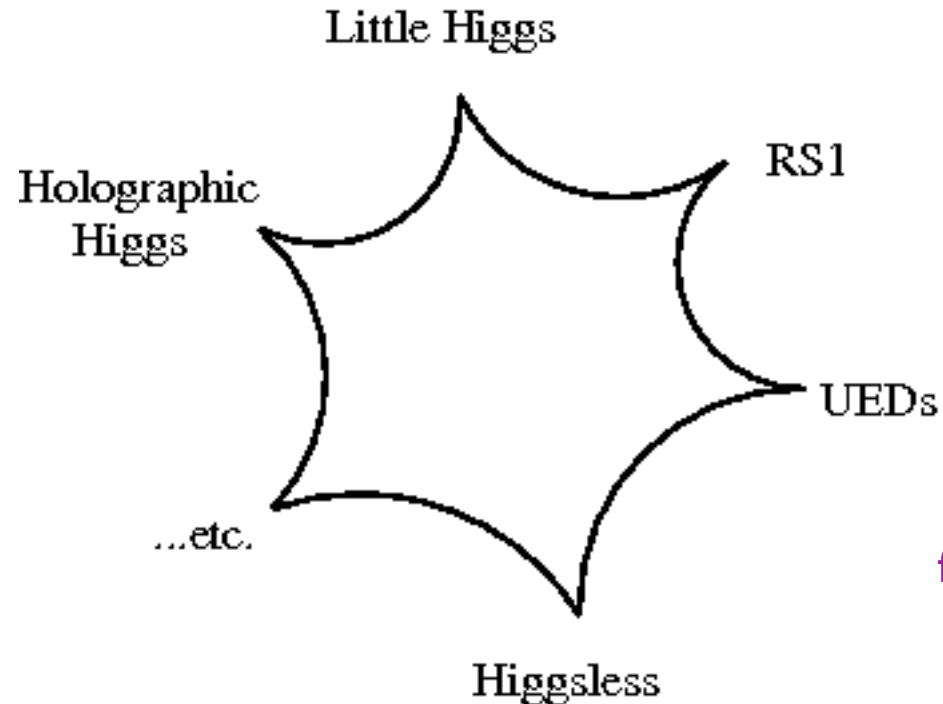


figure credit: Ian Low

- All these models involve **new strong dynamics at TeV**, a la QCD confinement (but with interesting twists)
- They can be “continuously deformed” into each other (“**little M-theory**”)

Composite Higgs

- Many spin-0 particles exist in nature - **mesons**
- They are **composite**, made of spin-1/2 quarks, bound by QCD strong force
- Above the QCD confinement scale, the good degrees of freedom are quarks \Rightarrow **no** hierarchy problem!
- Can the Higgs be a meson bound by a **new strong force**?
- Old idea, but difficult to build models - **non-perturbative** physics!
- New insight: **AdS/CFT duality** \Rightarrow some strongly coupled 4D models are “dual” to weakly coupled, calculable models with an extra dimension!
- Setup: **Randall-Sundrum** (RS) 5D model

Warped (RS) Extra Dimension

- Original model had the SM **on the TeV brane**
- It was subsequently realized that models with SM gauge fields and fermions **in the “bulk”** are more interesting:

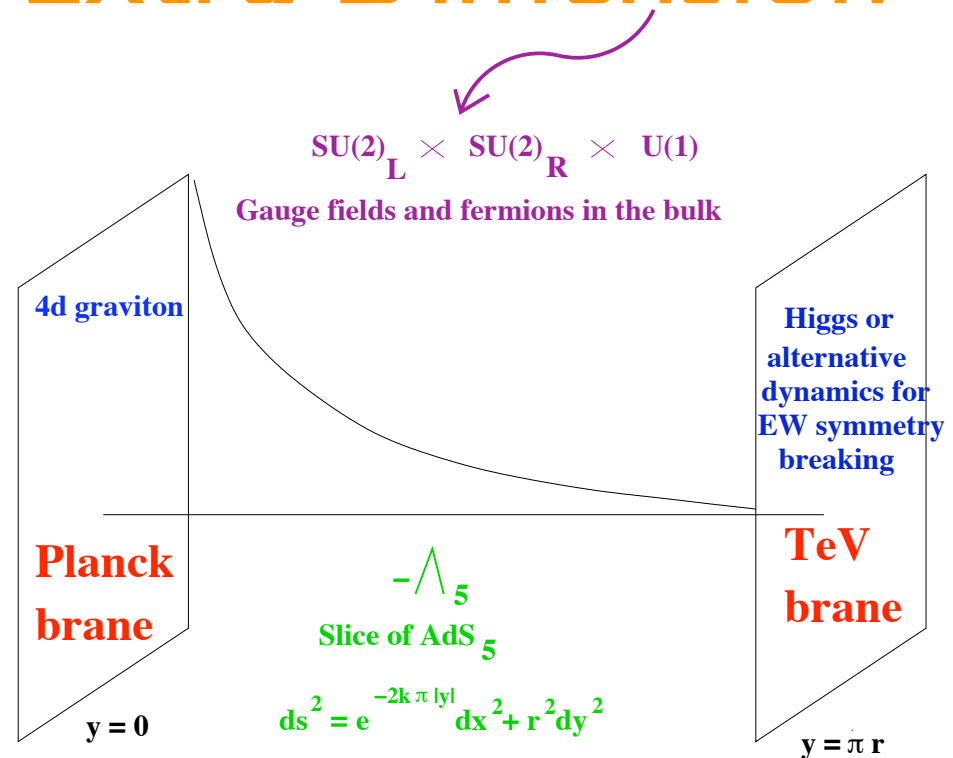


figure credit: G. Servant

- natural solution to **fermion mass hierarchy** problem
- natural suppression of **flavor-changing neutral currents**
- possibility of **gauge coupling unification**, as in the MSSM

NB: collider pheno is **different** from the “SM on the brane” model, e.g. small KK-graviton coupling to light quarks/leptons!

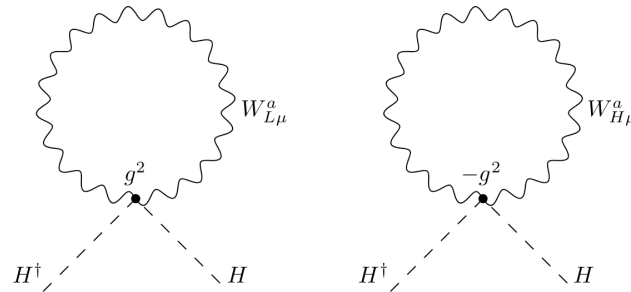
[see talks by **Randall, Perez, Lee**]

Universal Extra Dimensions and KK Parity

- UED is a model with **flat extra dimension(s)** of $R^{-1} \sim 1 \text{ TeV}$, all SM fields (including Higgs) propagate in the bulk
- Natural discrete symmetry, **KK parity**: first-level KK states are odd, lightest odd state (**LKP**) is stable, typically heavy photon
- Interesting **pheno**: faking SUSY at the LHC, dark matter
[see talk by [Przysniezniak](#)]
- Disadvantage 1: strongly coupled at $\Lambda \sim 10 \text{ TeV}$, needs a “**UV completion**” e.g. to address gauge coupling unification
- Disadvantage 2: $m_h \sim 0.01\Lambda$ requires 1% fine-tuning (“**little hierarchy**”)
- An interesting alternative that addresses these issues: a **warped** RS-style extra dimension with KK parity
[[Agashe](#), [Falkowski](#), [Low](#), [Servant](#), [12/07](#)]

Gauge-Higgs Unification

- A zero-mass photon does not require fine-tuning - mass is protected by **gauge symmetry**
- In a **5D** theory, the gauge field $A_M(x) \rightarrow A_\mu(x), A_5(x)$
- If the 5th dimension is infinite, A_5 is **naturally** massless!
- After **compactification**, $m(A_5) \sim 1/R \Rightarrow$ good if $1/R \sim M_W$
- Higgs mass quadratic divergences are **anceled** by KK modes:



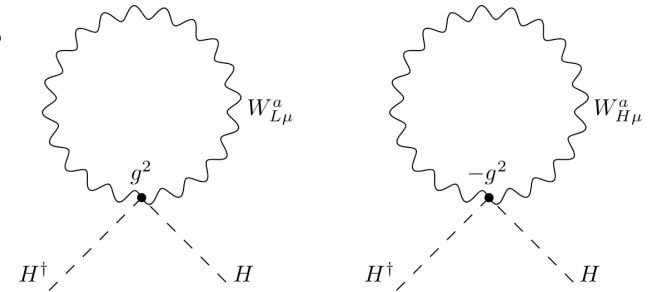
- A realistic GHU implementation, using a **warped** extra dimension, predicts $m_h < 140 \text{ GeV}$ and KK states at **2 TeV**

[Agashe, Contino, Pomarol, 12/04]

Little Higgs

- Quadratic divergence cancellation by **same-spin states** can also occur in a purely 4D theory - Little Higgs

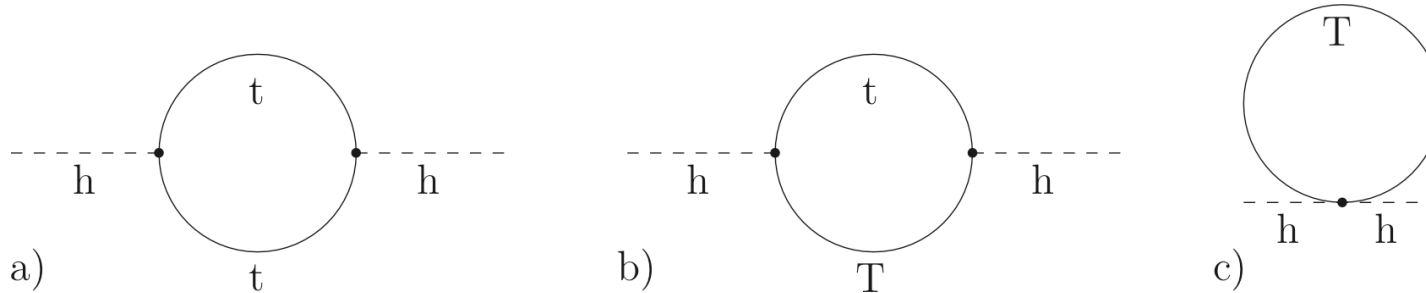
[LH \leftrightarrow effective theory of the first two KK modes in **GHU!**]



- In LH, Higgs is a **Goldstone boson** arising from a global symmetry breaking [a la **pions** in QCD]
- If the global symmetry is **exact**, $m_h = 0$ naturally!
- Goldstones only interact derivatively \Rightarrow need to **break** the global symmetry explicitly by gauge and Yukawa interactions
- Generically explicit breaking **reintroduces** quadratic divergences
- “Collective” breaking pattern in LH avoids quad. div. **at one loop**

EWSB in Littlest Higgs Model

- Higgs mass is dominated by **top and Top** loops:



- This contribution is log-divergent and **negative**:

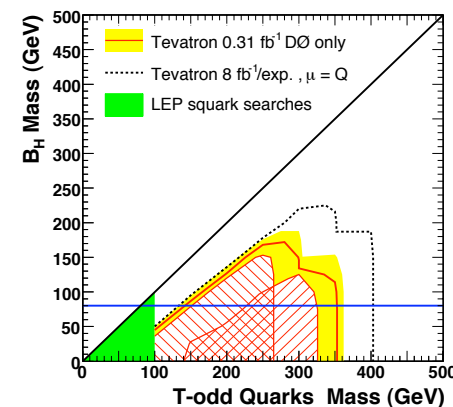
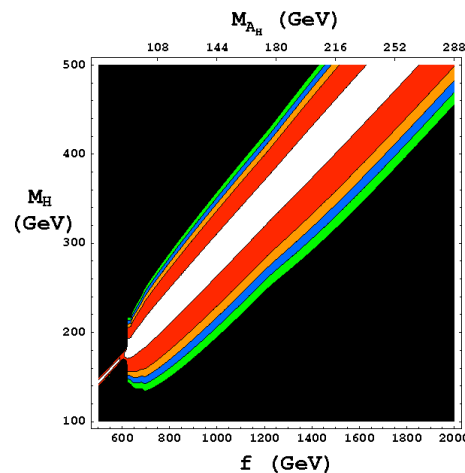
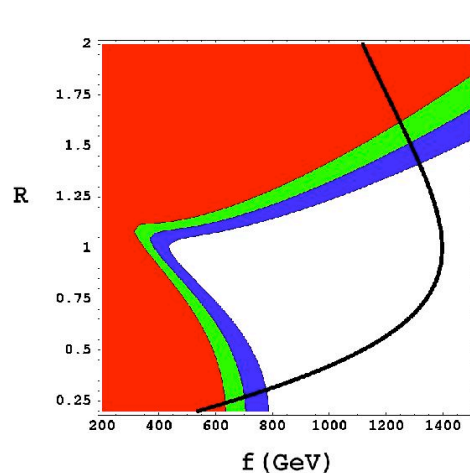
$$m_t^2(H) = -\frac{3\lambda_t^2 M_T^2}{8\pi^2} \log \frac{\Lambda^2}{M_T^2} .$$

- Quadratically divergent 2-loop contributions are of order $\frac{g^2 y_t^2 M_T^2}{16\pi^2}$, no log enhancement \Rightarrow **subdominant**
- 1-loop gauge contribution down by $g^2/y_t^2 \sim 0.25$
- EWSB is triggered **radiatively** - **simple mechanism!**

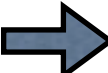
Little Higgs and T Parity

- LH models are weakly coupled at the TeV scale, **predictive!**
- The “first-generation” LH models strongly **disfavored** by precision electroweak data
- Best solution: introduce “**T Parity**”: new TeV-scale particles T-odd and only appear in loops in PEWVO [a la R parity of the MSSM]
- Littlest Higgs with T Parity (**LHT**) a benchmark model, many pheno studies: missing energy signatures at the LHC and the Tevatron, dark matter candidate, ...

[see talk by Figi]



What if There is No Higgs?

- If physics at TeV scale is strongly coupled, a symmetry-breaking condensate can exist **without** a physical Higgs boson in the theory - **technicolor!**
- TC with QCD-like dynamics at TeV is **ruled out** by precision electroweak data
- Difficult to explore model space due to strong coupling
- New insight: **AdS/CFT duality**  some strongly coupled 4D models are “dual” to weakly coupled, calculable models with an extra dimension!
- 5D “**Higgsless**” models have been constructed, can pass precision electroweak tests with $\sim 1\%$ fine-tuning
- Renewed interest in **4D** TC model-building

[see talk by **Foadi**]

Conclusions and Outlook

- The mechanism which breaks electroweak symmetry remains a **fundamental, unsolved mystery**
- It must involve **new physics** at the TeV scale
- Several **theoretical ideas** for what new physics might be have been proposed: supersymmetry, dynamical symmetry breaking, extra dimensions, little higgs, ...
- True model is **unknown**: only indirect constraints and theoretical prejudice to guide us at this point...
- All models predict **discoveries** at the LHC - theory will be confronted with data soon!
- **Detailed predictions** will be needed for theoretical interpretation of the data

"NEW PHYSICS PIPELINE"

Build a Model

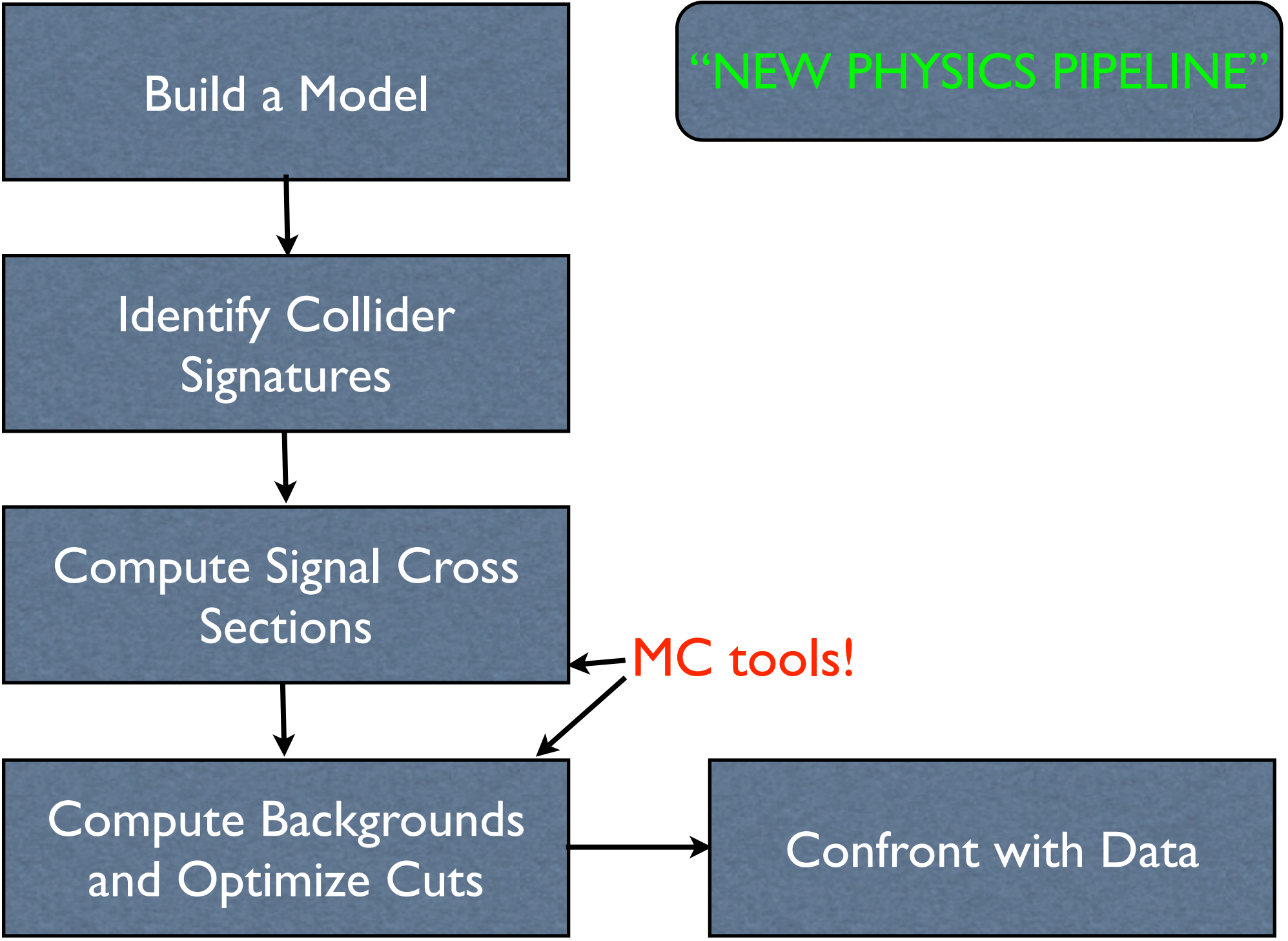
Identify Collider Signatures

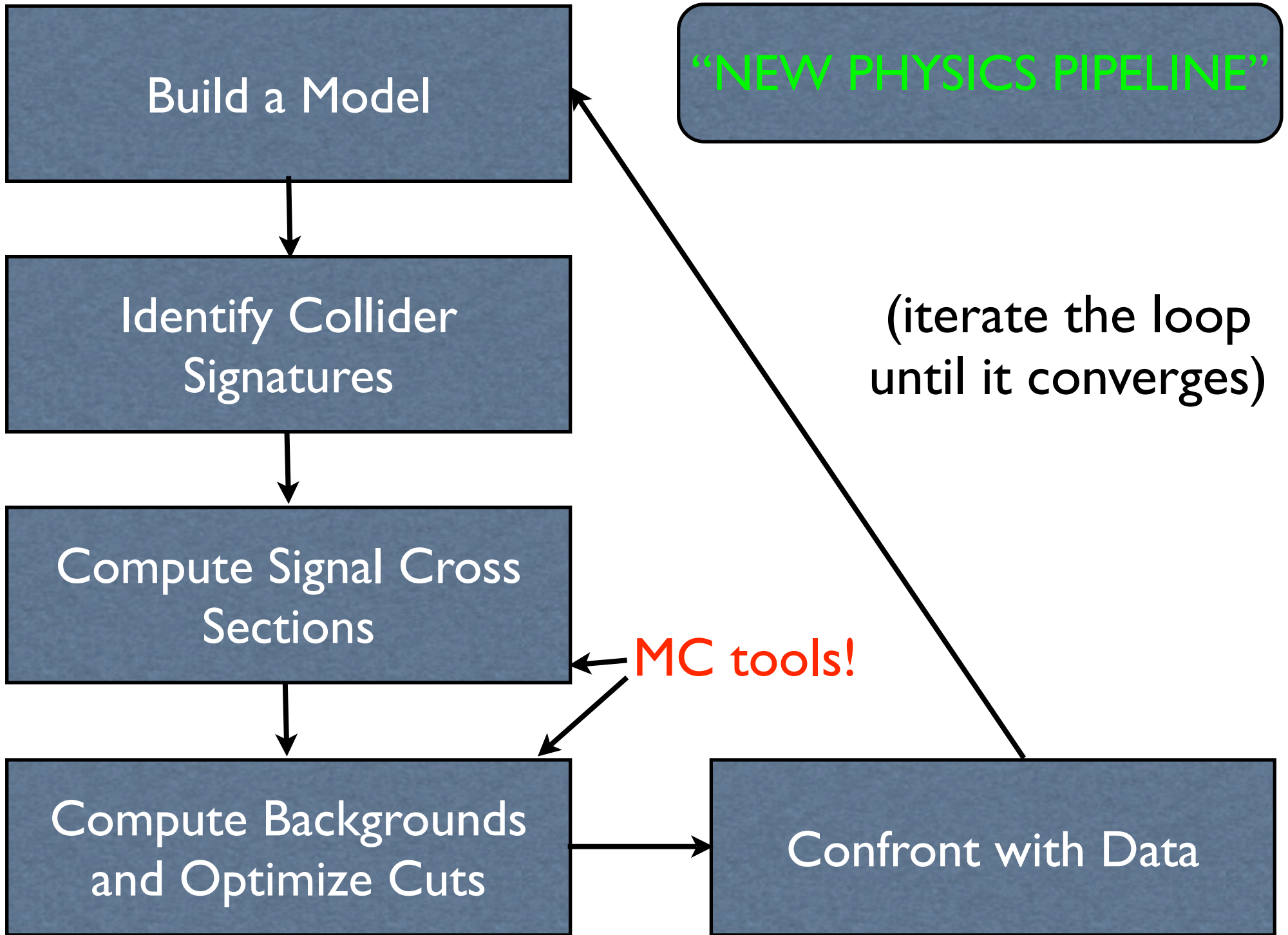
Compute Signal Cross Sections

Compute Backgrounds and Optimize Cuts

Confront with Data

MC tools!





- Since models will change “in real time”, **flexibility** is the key desired feature of the BSM MC tools
- In the past, general-purpose MC tools had a **small** number of models (or processes) hard-wired, with significant effort and expert-level coding required to add a new model/process
- Now: **any** new physics model can be **realistically** simulated within hours by a physicist with rudimentary software skills (e.g. myself) **as long as**:
 - all new particles are **spin ≤ 2**
 - couplings of “**known**” Lorentz structures only
 - no new long-lived colored states or exotic color rep’s
- Full implementation of such models ahead of data seems **unnecessary**; removing these limitations is more useful