Predictions for a 3.55 keV photon line from dark matter decay to axions in the Milky Way

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- Sharp peak in the cool core region of Perseus
- Morphology inconsistent with direct decay of dark matter to photons

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 Dark matter decays to an axion like particle (ALP) which mixes with the photon in astrophysical magnetic fields (c.f. M Rummel's talk)

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• Fits the observed morphology of the 3.55 keV flux

ALPs

$$\mathcal{L} = rac{1}{2} \partial_{\mu} a \partial^{\mu} a - rac{1}{2} m_a^2 a^2 + rac{a}{M} \mathbf{E} \cdot \mathbf{B}$$

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$$\mathcal{L} = \frac{1}{2} \partial_{\mu} a \partial^{\mu} a - \frac{1}{2} m_{a}^{2} a^{2} + \frac{a}{M} \mathbf{E} \cdot \mathbf{B}$$
$$\left(\omega + \begin{pmatrix} \Delta_{\gamma} & 0 & \Delta_{\gamma ax} \\ 0 & \Delta_{\gamma} & \Delta_{\gamma ay} \\ \Delta_{\gamma ax} & \Delta_{\gamma ay} & \Delta_{a} \end{pmatrix} - i \partial_{z} \end{pmatrix} \begin{pmatrix} \mid \gamma_{x} \rangle \\ \mid \gamma_{y} \rangle \\ \mid a \rangle \end{pmatrix} = 0$$

•
$$\Delta_{\gamma} = rac{-\omega_{
m pl}^2}{2\omega}$$

• Plasma frequency:
$$\omega_{pl} = \left(4\pi \alpha \frac{n_e}{m_e}\right)^{\frac{1}{2}}$$

•
$$\Delta_{a}=rac{-m_{a}^{2}}{\omega}$$
 (Here we take $m_{a}=0$)

• Mixing:
$$\Delta_{\gamma a i} = rac{B_i}{2M}$$

Over a distance R of $R/L \gg 1$ domains, with **B** randomised between each domain, we can approximate:

$$P\simeq 6.9\times 10^{-7}\left(\frac{L}{1\,\mathrm{kpc}}\frac{R}{30\,\mathrm{kpc}}\right)\left(\frac{B_{\perp}}{1\,\mu\mathrm{G}}\frac{10^{13}\,\mathrm{GeV}}{M}\right)^{2},$$

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- For small conversion probabilities, P (a → γ) increases with the domain size L.
- Conversion suppressed by high electron densities.
- We use a full discretized simulation for propogation in the Milky Way.

Dark Matter Lifetime

• To reproduce observed flux with direct dark matter decay to photons:

 $\tau_{\rm direct}\sim 5\times 10^{27}\,{\rm s}$

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• Detailed simulation of the central region of Coma (Angus *et al*, 2013; c.f. A Powell's talk) taking $M = 10^{13} \,\text{GeV}$ gives:

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 Therefore to reproduce the observed flux with DM → a → γ we require:

$$au_{
m ALP} \sim 5 imes 10^{24} \, {
m s} \left(rac{10^{13} \, {
m GeV}}{M}
ight)^2$$

What does $DM \rightarrow a \rightarrow \gamma$ predict in the Milky Way?

NFW Profile:

$$\rho_{NFW}(r) = rac{
ho_s}{\left(rac{r}{r_s}
ight) \left(1 + rac{r}{r_s}
ight)^2},$$

where $r_s = 20 \text{ kpc}$ and $\rho_{NFW}(R_{\odot}) = 0.4 \text{ GeV cm}^{-3}$ (Parameters from Fermi-LAT collaboration, 1305.5597) • Recent model by Jansson and Farrar (2012) based on 40,000 extragalactic Faraday RMs

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- The random field has a coherence length of less than 100 pc, and so does not contribute significantly to $a \rightarrow \gamma$ conversion.
- Regular component, composed of disc, halo and X fields, has structure on the scale of the Milky Way.

Milky Way Magnetic Field



Average Transverse Magnetic Field/µG

The average regular transverse magnetic field experienced by an ALP on a path starting 20 kpc from the Earth and ending at the Earth.

ALPs: Photoelectric Absorption

In the Milky Way, photoelectric absorption can be significant. Damping parameter: $\Gamma = \sigma_{\text{eff}} (n_{HI} + 2n_{H2})$

$$H = \begin{pmatrix} \Delta_{\gamma} & 0 & \Delta_{\gamma ax} \\ 0 & \Delta_{\gamma} & \Delta_{\gamma ay} \\ \Delta_{\gamma ax} & \Delta_{\gamma ay} & \Delta_{a} \end{pmatrix} - \begin{pmatrix} i\frac{\Gamma}{2} & 0 & 0 \\ 0 & i\frac{\Gamma}{2} & 0 \\ 0 & 0 & 0 \end{pmatrix} = M - iD,$$
$$\rho = \begin{pmatrix} |\gamma_x\rangle \\ |\gamma_y\rangle \\ |a\rangle \end{pmatrix} \otimes (|\gamma_x\rangle & |\gamma_y\rangle & |a\rangle)^*$$
$$\rho(z) = e^{-iHz}\rho(0)e^{iH^{\dagger}z}.$$

For $\omega = 3.55 \, {\rm keV}$, we find the effect is negligible.

$P(a ightarrow \gamma)$ in the Milky Way



Log₁₀[Conversion Probability]

The ALP to photon conversion probability for an ALP with $M = 10^{13} \text{ GeV}$ starting at 30 kpc from the Earth and propagating to Earth.

For direct decay $DM \rightarrow \gamma$:

$$F = \frac{1}{\tau_{\text{direct}}} \frac{1}{4\pi} \int_{r=0}^{\infty} \frac{\rho_{DM}(r,\theta,\phi)}{m_{DM}} dr \, \text{sr}^{-1},$$

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For $DM \rightarrow a \rightarrow \gamma$:

$$F = \frac{1}{\tau_{\rm ALP}} \frac{1}{4\pi} \int_{r=0}^{\infty} \frac{\rho_{DM}(r,\theta,\phi)}{m_{DM}} P(r,\theta,\phi) \, dr \, {\rm sr}^{-1},$$

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- Evaluate the integral as a sum for the first 30 spherical shells.
- For ALPs produced greater than 30 kpc from the Earth, we take $P(r > 30 \text{ kpc}, \theta, \phi) = P(30 \text{ kpc}, \theta, \phi)$ and integrate ρ_{DM} from r = 30 kpc to $r = \infty$ numerically.

Expected Flux for $DM \rightarrow a \rightarrow \gamma$



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Expected Flux for $DM \rightarrow a \rightarrow \gamma$ and $P(a \rightarrow \gamma)$

Log10 [Expected Flux/cm⁻²s⁻¹sr⁻¹]





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Expected Flux for $DM \rightarrow \gamma$



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Expected Flux





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Milky Way Results

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- The corresponding maximum count rate on ASTRO-H's Soft X-ray Spectrometer is $\sim 4\times 10^{-8}\,{\rm s}^{-1}.$

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- The maximal flux in the $DM\to a\to\gamma$ scenario is $\sim 2\times 10^{-4}\,{\rm cm^{-2}s^{-1}sr^{-1}}$
- The corresponding maximum count rate on ASTRO-H's Soft X-ray Spectrometer is $\sim 4\times 10^{-8}\,{\rm s}^{-1}.$
- Detection with ASTRO-H would be impossible, in contrast to the direct decay case.

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- There is no sign of field reversals among the spirals.
- The vertical scale height is at least 1 kpc.
- M31 is near edge on (inclination angle = 77.5°), so ALPs originating from dark matter decay pass through a large coherent transverse magnetic field on their way to Earth.

Single domain small angle approximation for the conversion probability for a 3.55 keV ALP created at the centre of M31 and propagating to Earth:

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B_{\perp} \sim 5 \,\mu\text{G}

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Single domain small angle approximation for the conversion probability for a 3.55 keV ALP created at the centre of M31 and propagating to Earth: $B_{\perp} \sim 5 \,\mu\text{G}$ $L \sim R \sim 20 \,\text{kpc}$

$$P_{a
ightarrow\gamma,M31}\sim 2.3 imes 10^{-4}\left(rac{10^{13}\,{
m GeV}}{M}
ight)^2$$

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- This is two orders of magnitude higher than typical conversion probabilities for the Milky Way, and only smaller by a factor of four than typical conversion probabilites in galaxy clusters.
- In the DM → a → γ scenario the observed signal strength from M31 can be comparable to that from clusters, consistent with the results of Boyarsky et al.
- M31 is an unusually favourable galaxy for observing the 3.55 keV line.

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Some of the flux from the Milky Way halo in the DM → a → γ scenario will be significantly lower than for the case of direct dark matter decay to photons, and will be unobservable with ASTRO-H.

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- Solution The flux from the Milky Way halo in the DM → a → γ scenario will be significantly lower than for the case of direct dark matter decay to photons, and will be unobservable with ASTRO-H.
- ALP to photon conversion probabilities for M31 are two orders of magnitude larger than in the Milky Way, and are comparable with the conversion probabilities in clusters.
- A non-observation of the 3.55 keV line from the Milky Way with ASTRO-H will not rule out a dark matter origin of the signal.

ALPs: Single Domain

$$\tan (2\theta) = 5.0 \times 10^{-3} \times \left(\frac{10^{-3} \,\mathrm{cm}^{-3}}{n_e}\right) \left(\frac{B_{\perp}}{1 \,\mu\mathrm{G}}\right) \left(\frac{\omega}{3.55 \,\mathrm{keV}}\right) \left(\frac{10^{13} \,\mathrm{GeV}}{M}\right)$$
$$\Delta = 0.015 \times \left(\frac{n_e}{10^{-3} \,\mathrm{cm}^{-3}}\right) \left(\frac{3.55 \,\mathrm{keV}}{\omega}\right) \left(\frac{L}{1 \,\mathrm{kpc}}\right).$$

$$P(a \rightarrow \gamma) = \sin^2(2\theta) \sin^2\left(\frac{\Delta}{\cos 2\theta}\right)$$

•
$$P(a \rightarrow \gamma) \propto \frac{B_{\perp}^2}{M^2}$$

- For small conversion probabilities, P (a → γ) increases with the domain size L.
- Conversion suppressed by high electron densities.

Milky Way Electron Density

- Gomez and Cox (2001) model, sum of thick and thin disk components
- Thick disk component:

$$n_{
m thick} = 1.77 imes 10^{-2} \, {
m cm}^{-3} imes rac{{
m sech}^2 \left(rac{
ho}{15.4 \, {
m kpc}}
ight)}{{
m sech}^2 \left(rac{R_\odot}{15.4 \, {
m kpc}}
ight)} {
m sech}^2 \left(rac{z}{1.10 \, {
m kpc}}
ight)$$

• Thin disk component:

$$n_{
m thin} = 1.07 imes 10^{-2} \, {
m cm}^{-3} imes rac{{
m sech}^2 \left(rac{
ho}{3.6 \, {
m kpc}}
ight)}{{
m sech}^2 \left(rac{R_\odot}{3.6 \, {
m kpc}}
ight)} {
m sech}^2 \left(rac{z}{0.04 \, {
m kpc}}
ight)$$

• Minimum n_e of $10^{-7} \, \mathrm{cm}^{-3}$

Photo-electric absorption

- Scattering cross sections in Morrison and McCammon (1983), given as an effective cross section per hydrogen atom: $\Gamma = \sigma_{\text{eff}} (n_{HI} + 2n_{H2})$
- For a 3.55 keV photon, $\sigma_{\rm eff} \sim 10^{-23}\,{\rm cm}^2.$
- Neutral hydrogen distributions from Misiriotis et al (2006):

$$n_{HI} = \begin{cases} 0.32 \,\mathrm{cm}^{-3} \mathrm{exp} \left(-\frac{\rho}{18.24 \,\mathrm{kpc}} - \frac{|z|}{0.52 \,\mathrm{kpc}} \right), & \text{if } \rho \ge 2.75 \,\mathrm{kpc} \\ 0, & \text{otherwise} \end{cases}$$

$$n_{H2} = 4.06 \,\mathrm{cm}^{-3} \mathrm{exp} \left(-\frac{
ho}{2.57 \,\mathrm{kpc}} - \frac{|z|}{0.08 \,\mathrm{kpc}}
ight)$$

• Discretize the field:

$$\rho_k = e^{-iH_k\delta z} \rho_{k-1} e^{iH_k^\dagger \delta z}$$

- Simulate the propagation of ALPs created at 1 kpc, 2 kpc, and 30 kpc from the Earth. Magnetic field has structure on the scale of the whole Milky Way must simulate each path seperately.
- Assume ALPs created over 30 kpc from the Earth propagate without conversion up to this point.
- 20000 points in the sky, 4000 domains for each path