

Radiative Generation of Quark Masses and Mixing Angles in the 2HDM

[arXiv:1403.2382]

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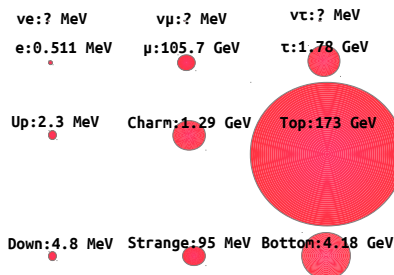


Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)



Introduction

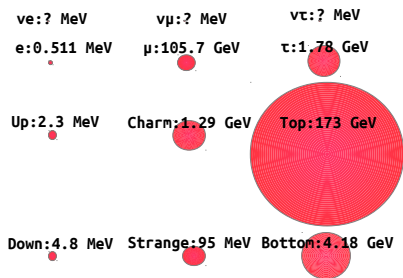
- **Motivation:** Hierarchy for masses and mixing angles in the Standard Model. \Rightarrow **New Physics?**



$$V_{CKM} = \begin{pmatrix} 0.974 & 0.225 & 0.0035 \\ 0.225 & 0.973 & 0.041 \\ 0.0087 & 0.04 & 0.999 \end{pmatrix}$$

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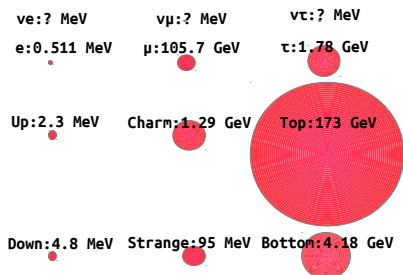


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- ▶ **Our goal:** Reproduce masses and mixing angles.
- ▶ **How?:** 2HDM (~~extra symmetries~~)

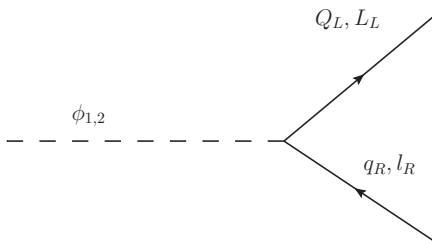
The Two-Higgs Doublet Model

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- ▶ Standard Model + one extra Higgs doublet.

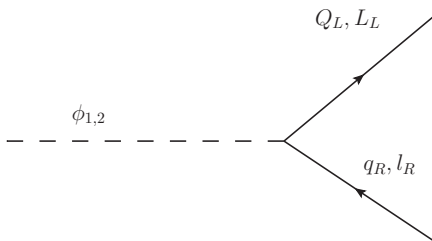
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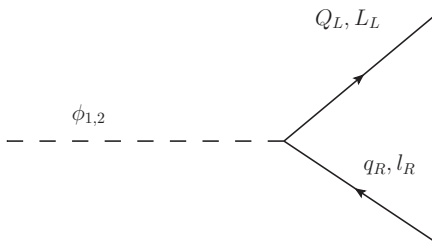
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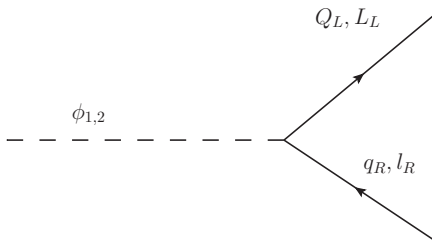
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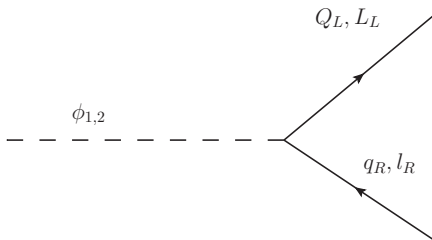
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or LFV processes.

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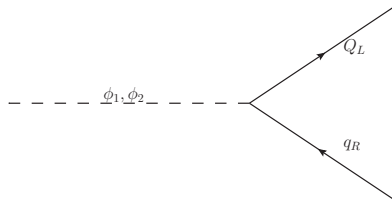


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- ▶ Decoupling limit \checkmark SM vacuum.

The Quark Sector

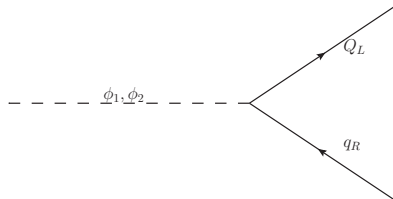
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(rank-1):



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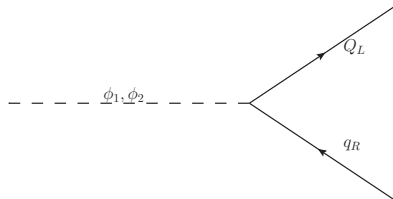


$$Y_u^{(1)}|_{\text{tree}} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & y_u^{(1)} \end{pmatrix}, \quad Y_d^{(1)}|_{\text{tree}} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & \epsilon y_d^{(1)} \\ 0 & 0 & y_d^{(1)} \end{pmatrix}$$

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- ▶ $Y_u^{(1)}|_{\text{tree}} \ \& \ Y_d^{(1)}|_{\text{tree}} \Rightarrow \begin{cases} m_t \\ m_b \end{cases} \quad \text{@ tree level}$

$$Y_u^{(2)}|_{\text{tree}} = U_L^\dagger \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & y_u^{(2)} \end{pmatrix} U_R, \quad Y_d^{(2)}|_{\text{tree}} = D_L^\dagger \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & y_d^{(2)} \end{pmatrix} D_R$$

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- $Y_{u,d}^{(2)}|_{\text{tree}}$ just depend on $U_{L3i}, U_{R3i}, D_{L3i}, D_{R3i}$. Parametrize:

$$\begin{aligned} (U_L)_{31} &= e^{i\rho_{uL}} \sin \theta_{uL} \sin \omega_{uL}, \\ (U_L)_{32} &= e^{i\xi_{uL}} \sin \theta_{uL} \cos \omega_{uL}, \\ (U_L)_{33} &= \cos \theta_{uL}, \end{aligned}$$

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- ▶ Neglect phases.
- ▶ Assume for simplicity $y_u^{(1)}, \mathbf{y}_u^{(2)} \gg y_d^{(1)}, \mathbf{y}_d^{(2)}$

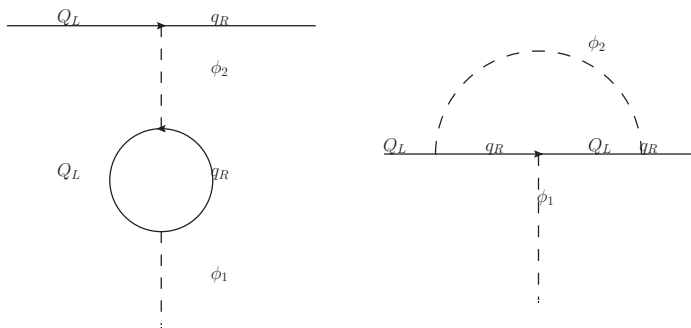
- ▶ 1-loop from β function:

$$Y_{u,d}^{(1)}|_{1\text{-loop}} \simeq Y_{u,d}^{(1)}|_{\text{tree}} + \frac{1}{16\pi^2} \beta_{u,d}^{(1)} \log \frac{\Lambda}{M_{\phi 2}}$$

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- ▶ 1-loop diagram (generate 2nd mass):



Quark Masses

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Mass hierarchy for the quarks:

$$\frac{y_c}{y_t} \simeq \left(\frac{1}{16\pi^2} \log \frac{\Lambda}{M_H} \right) \frac{3}{4} (y_u^{(2)})^2 \times \text{mixing angles}$$

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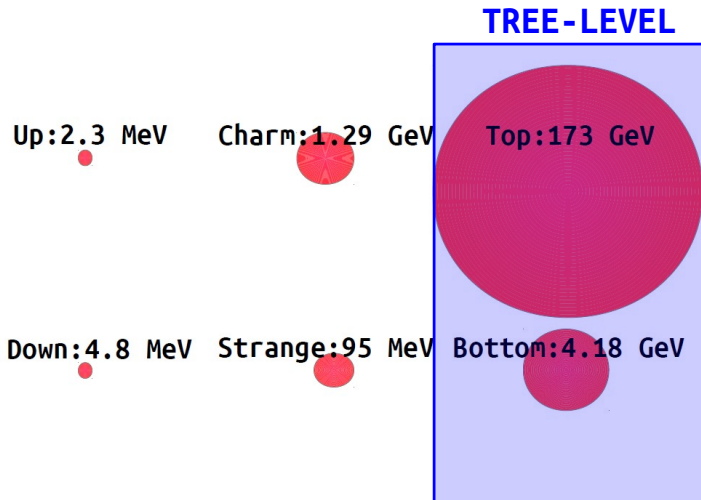
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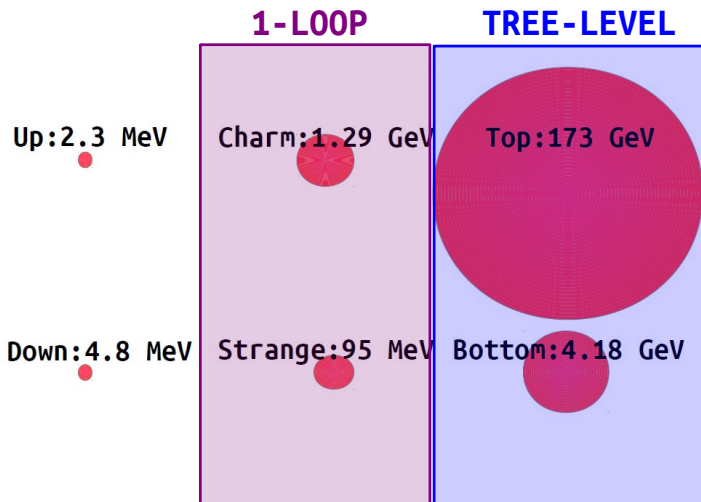
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- ▶ First generation massless.

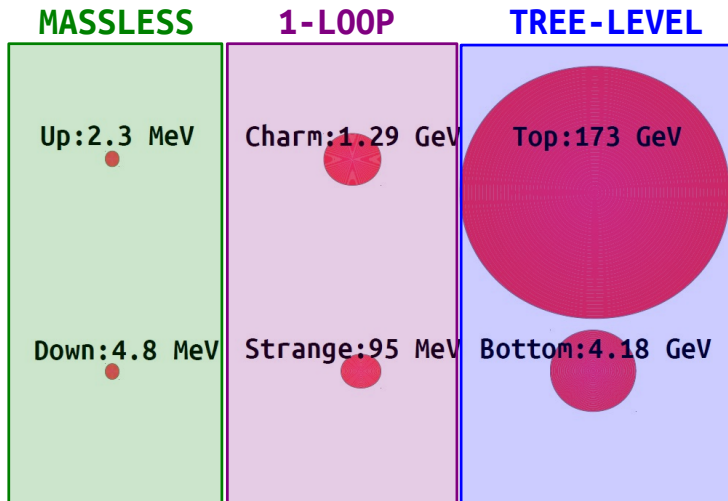
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$$V_{us} \simeq -V_{cd} \simeq \frac{3 \sin \theta_{d_L} \cos \theta_{u_L} \sin(\omega_{d_L} - \omega_{u_L})}{N_d}$$

$$N_d = [9 \sin^2 \theta_{d_L} \cos^2 \theta_{u_L} + 4 \cos^2 \theta_{d_L} \sin^2 \theta_{u_L} - 3 \sin 2\theta_{d_L} \sin 2\theta_{u_L} \cos(\omega_{d_L} - \omega_{u_L})]^{1/2}$$

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- ▶ Remaining off-diagonal terms @ 1st order in perturbation theory:

$$V_{ub} \simeq \left(\frac{1}{16\pi^2} \log \frac{\Lambda}{M_H} \right) \frac{3y_u^{(1)} y_u^{(2)} y_d^{(2)}}{y_d^{(1)}} \times \text{mixing angles}$$

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Tree-level

Quark Mixing Angles

0th Order

↓

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0th Order
1st Order Tree-level

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$$\frac{y_s}{y_b} \frac{V_{us}}{V_{ub}} \simeq \tan \theta_{d_R} ,$$
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$$\Rightarrow \theta_{u_R} \approx 0.16, \theta_{d_R} \approx 1.06$$

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- ▶ Degeneracy for other parameters:

- ✓ All masses and angles can be reproduced
(example: $y_u^{(2)} \approx 1.04$, $y_d^{(2)} \approx 0.02$, $\theta_{d_L} \approx 0.61$, $\theta_{u_L} \approx 0.51$,
 $\omega_{d_L} - \omega_{u_L} \approx 0.10$)

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$\nu_e: ? \text{ MeV}$

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$\nu_\mu: ? \text{ MeV}$

$\mu: 105.7 \text{ GeV}$



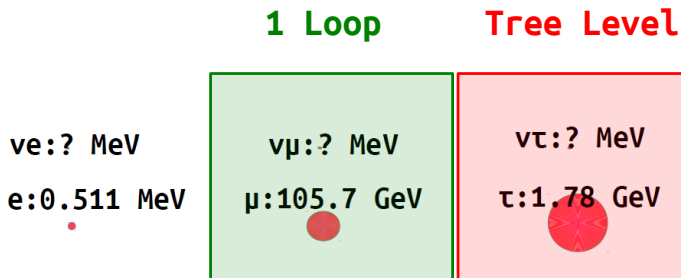
Tree Level

$\nu_\tau: ? \text{ MeV}$

$\tau: 1.78 \text{ GeV}$



Lepton Masses



Lepton Masses

Massless

ν_e :? MeV
e :0.511 MeV

1 Loop

ν_μ :? MeV
μ :105.7 GeV

Tree Level

ν_τ :? MeV
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PMNS Matrix

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Tree-Level

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0-Order **Tree-Level**

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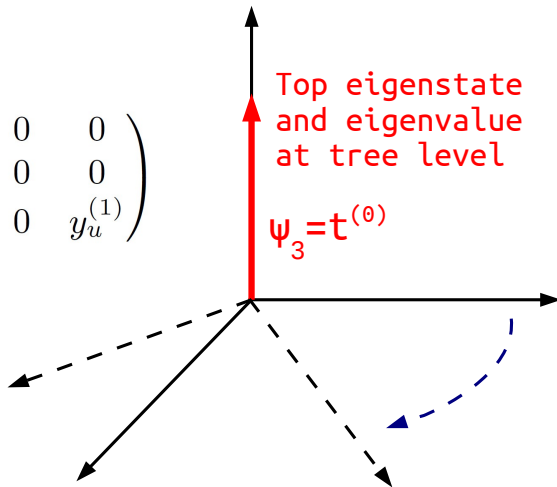
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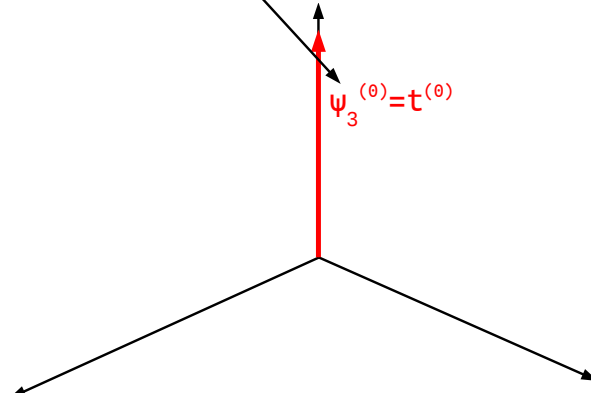
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Backup slides

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Backup slides

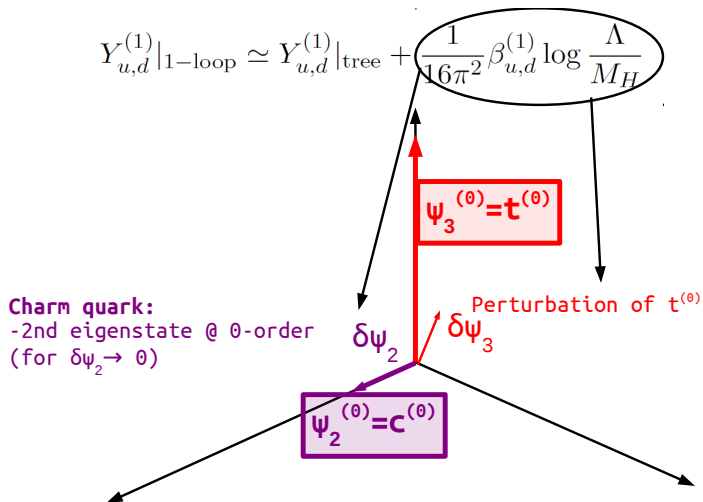
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$\psi_3^{(0)} = t^{(0)}$
 Perturbation of $t^{(0)}$
 $\delta\psi_2$
 $\delta\psi_3$

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- ▶ Eigenstates and eigenvalues of $\lambda W^{(n)}$ are eigenstates (0-order) and eigenvalues (1st-order) of H from the subspace $\perp |t^{(0)}\rangle$.