

# ***Charm production at high ET***

**The top-charm frontier at the LHC**

**CERN TH, January 14-17 2014**

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# Heavy quark multiplicities in gluon jets,

Mangano, Nason, PLB285 (1992) 160-166

$$a = -\frac{1}{4} \left[ 1 + \frac{2C_A}{3\pi b} \left( 1 - \frac{C_F}{C_A} \right) \right]$$

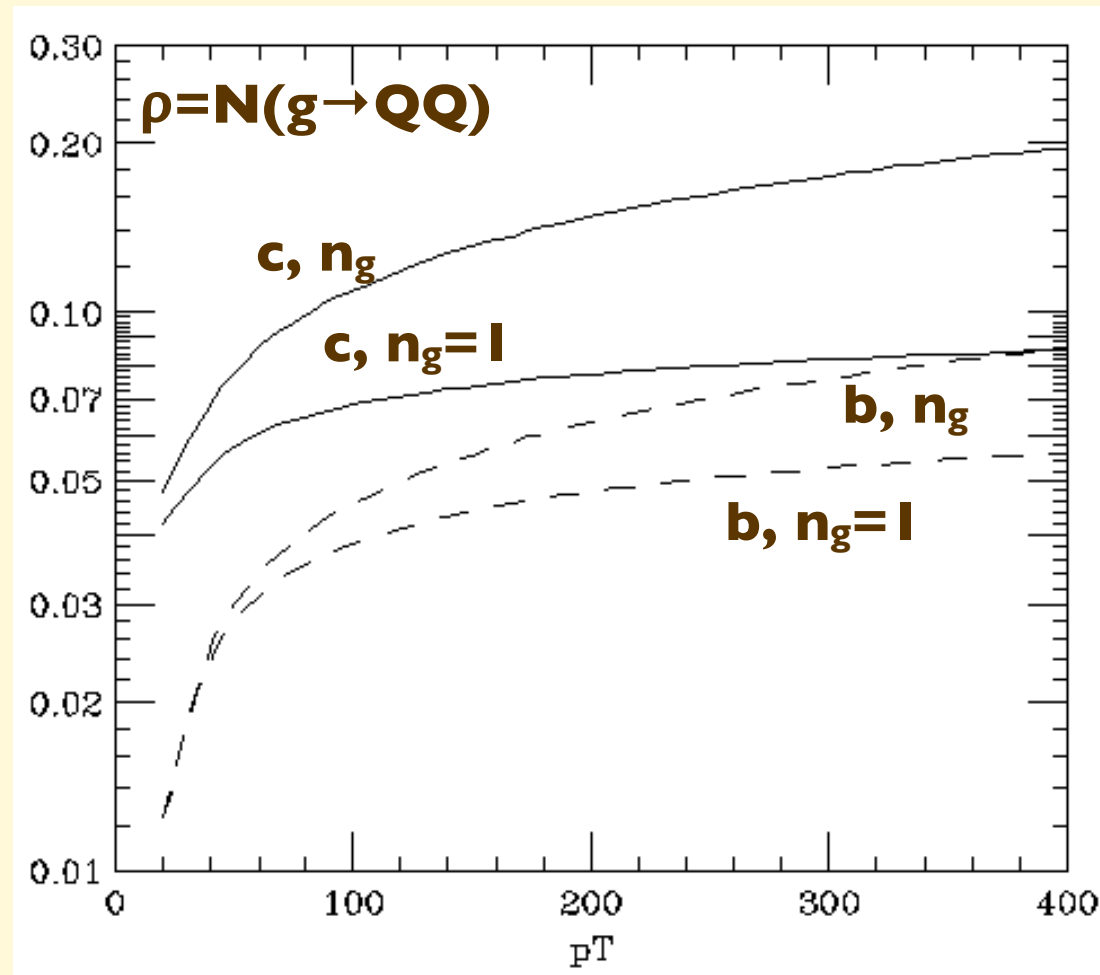
$$b = \frac{11C_A - 2N_F}{12\pi}$$

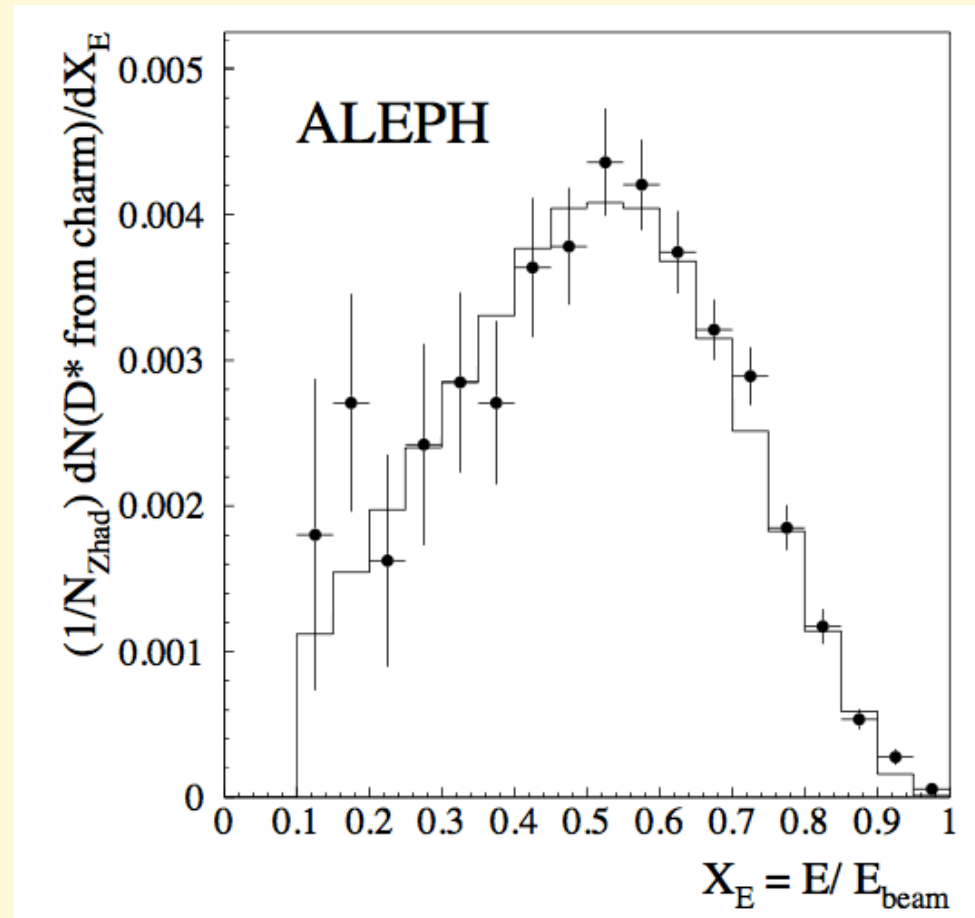
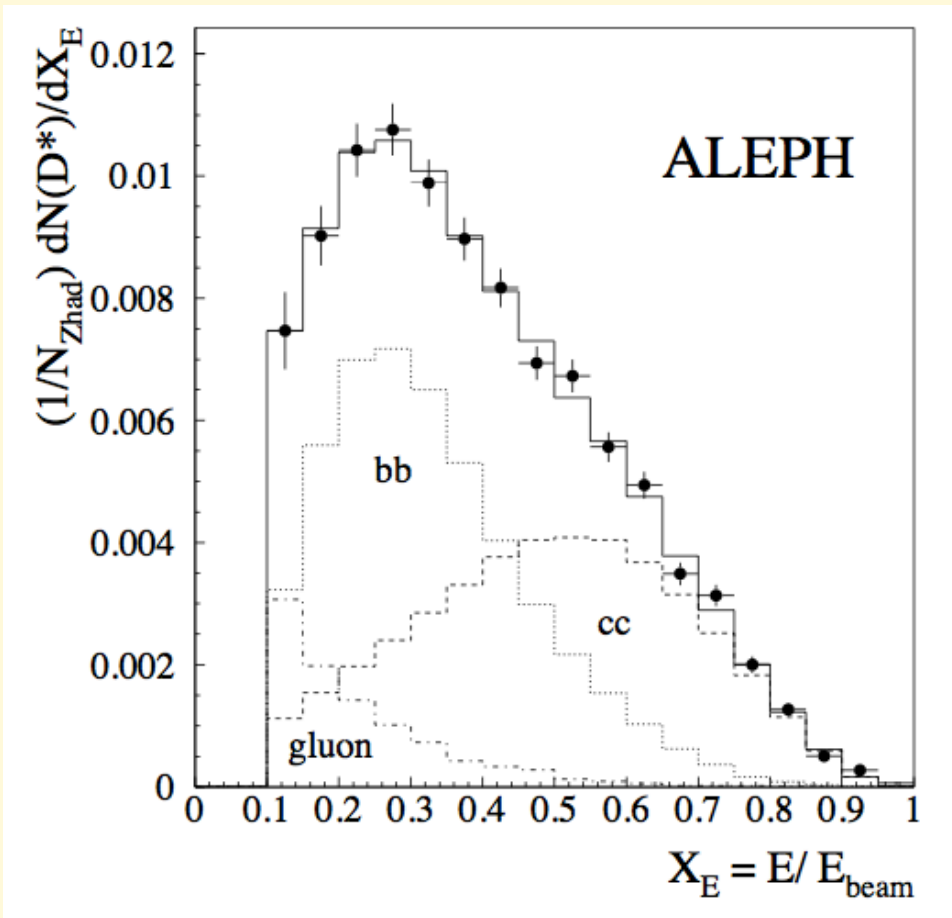
$$\rho = \frac{1}{6\pi} \int_{4m^2}^{Q^2} \frac{dK^2}{K^2} \alpha_s(K^2) \left( 1 + \frac{2m^2}{K^2} \right) \sqrt{1 - \frac{4m^2}{K^2}}$$

$$\times n_g(Q^2, K^2), \quad (1.1)$$

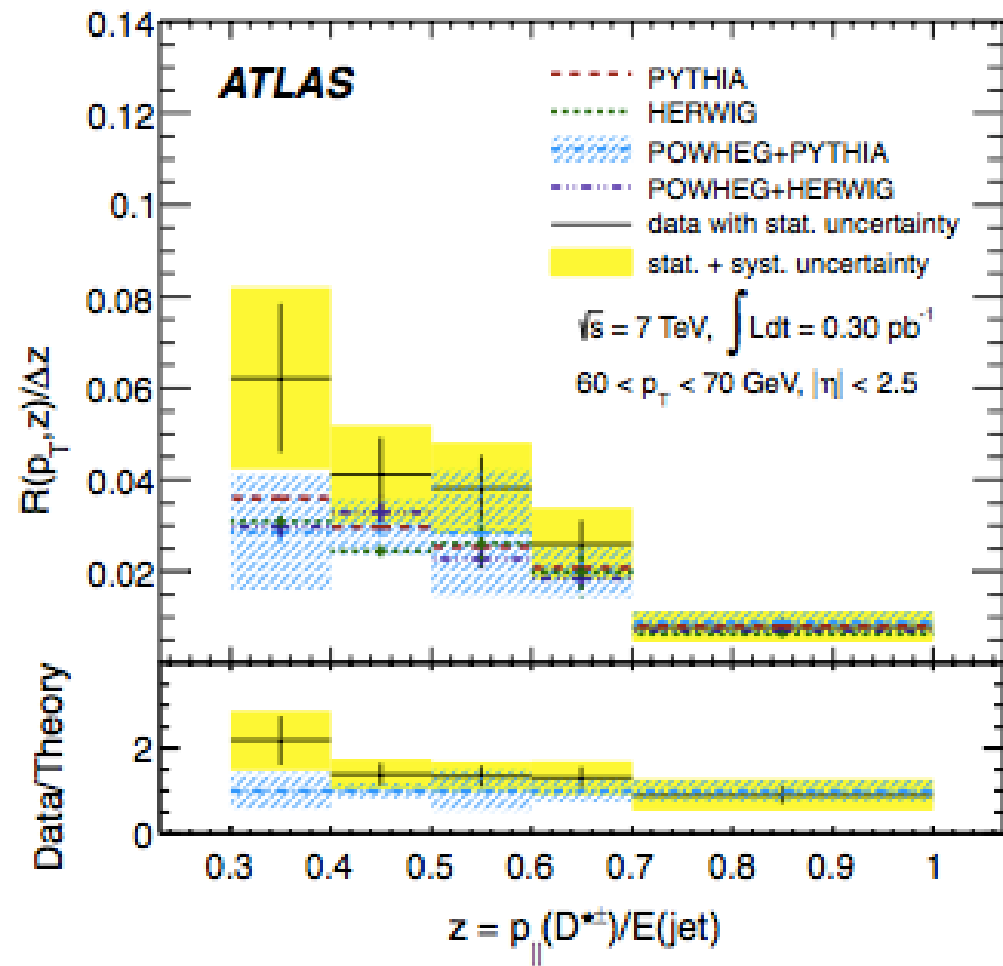
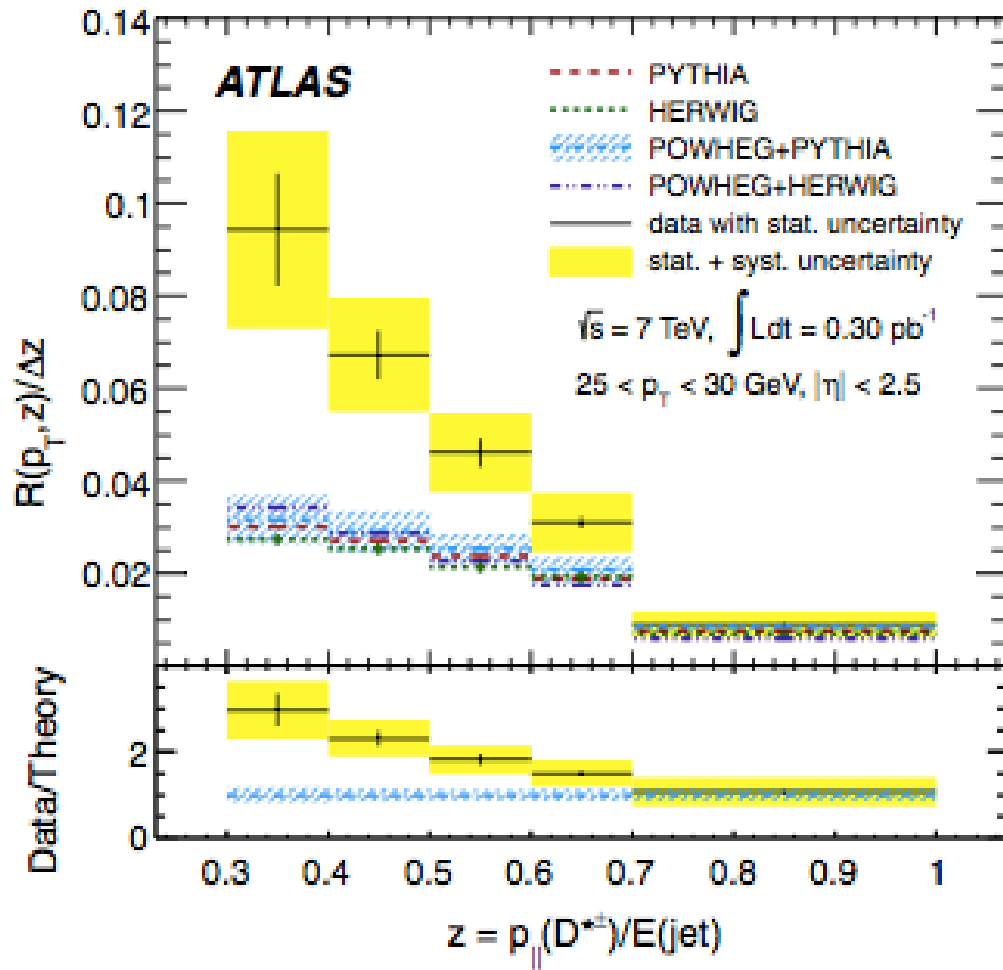
$$n_g(Q^2, K^2) = \left( \frac{\log(Q^2/\Lambda^2)}{\log(K^2/\Lambda^2)} \right)^a$$

$$\times \cosh \left( \sqrt{\frac{2C_A}{\pi b}} \left( \sqrt{\log \frac{Q^2}{\Lambda^2}} - \sqrt{\log \frac{K^2}{\Lambda^2}} \right) \right)$$

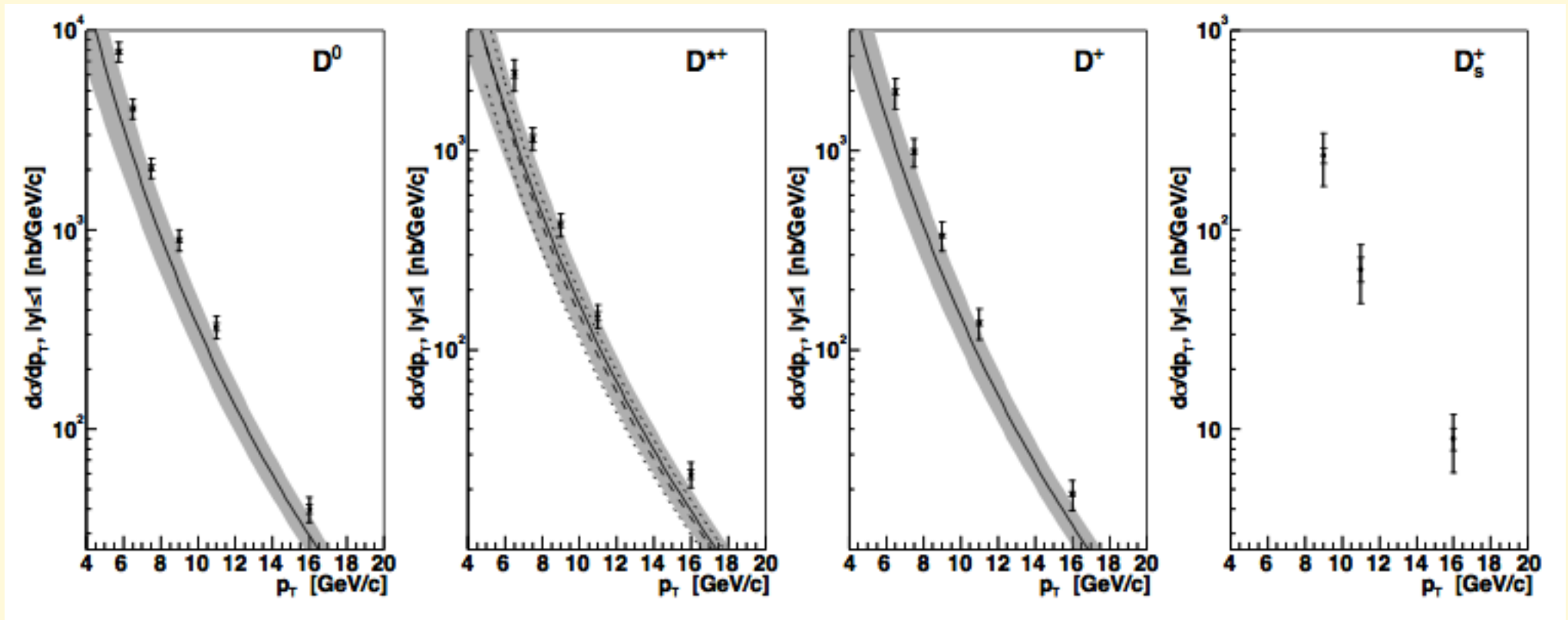


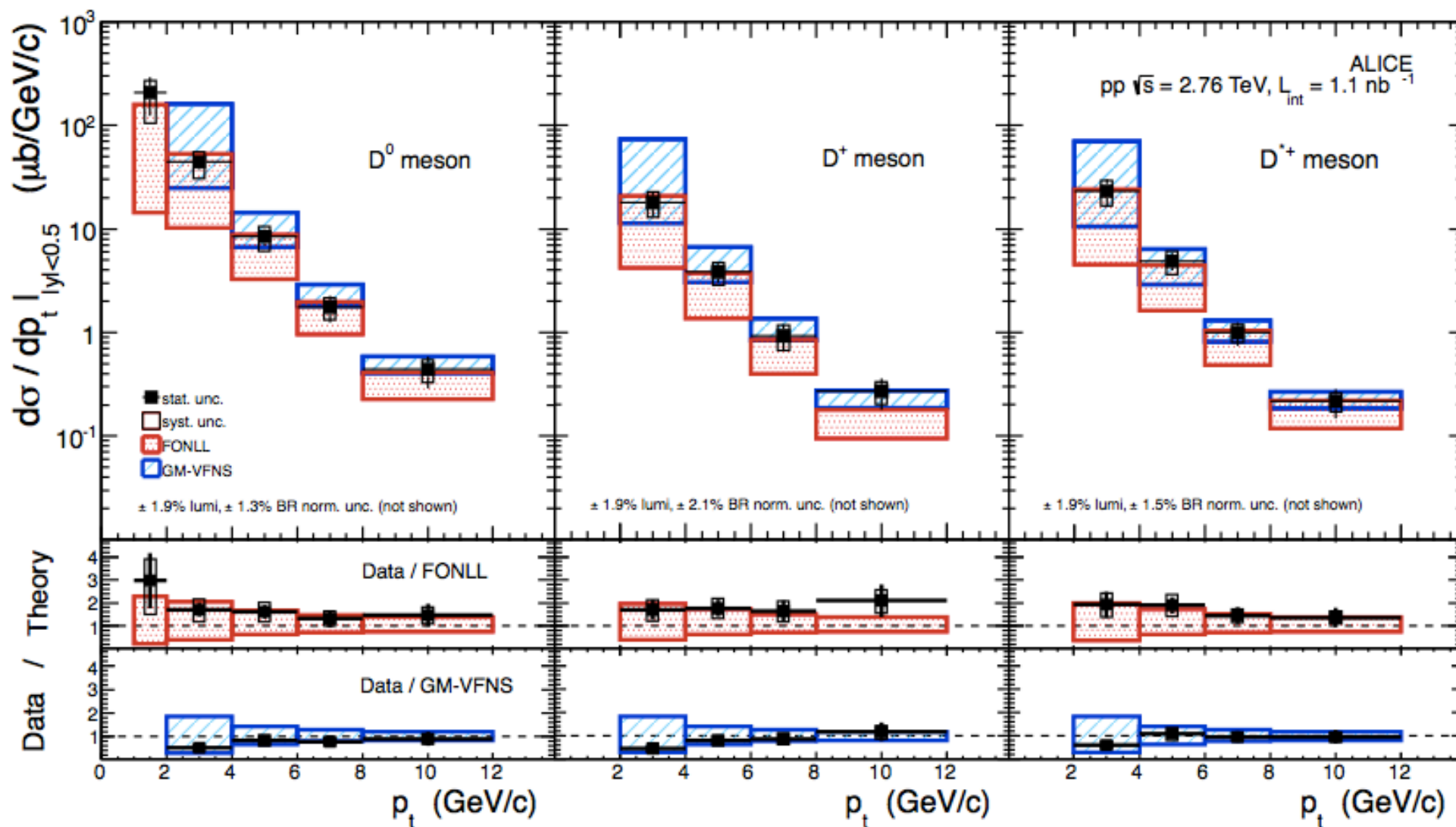


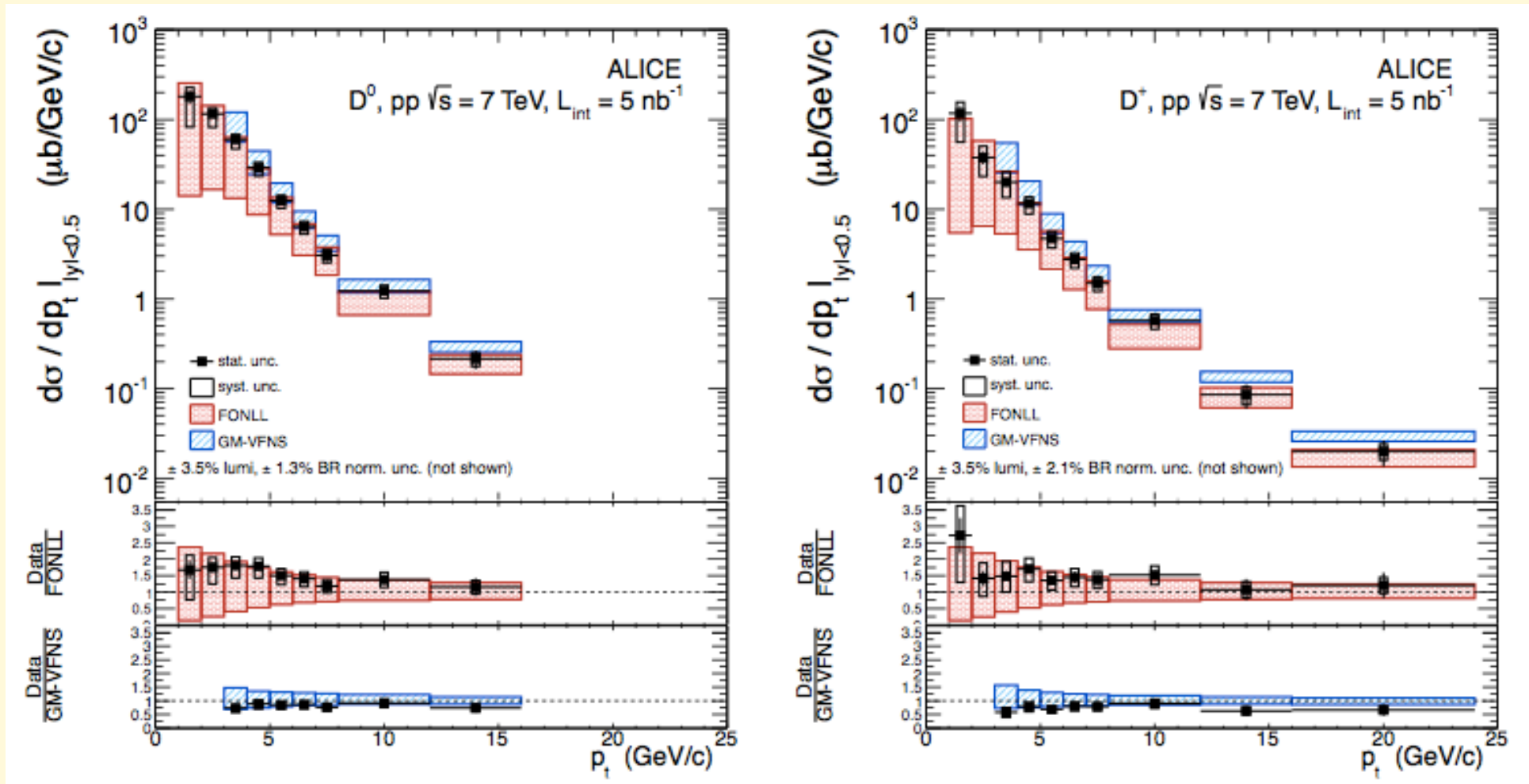
ALEPH, arXiv:hep-ex/9909032



CDF, Phys.Rev.Lett. 91 (2003) 241804

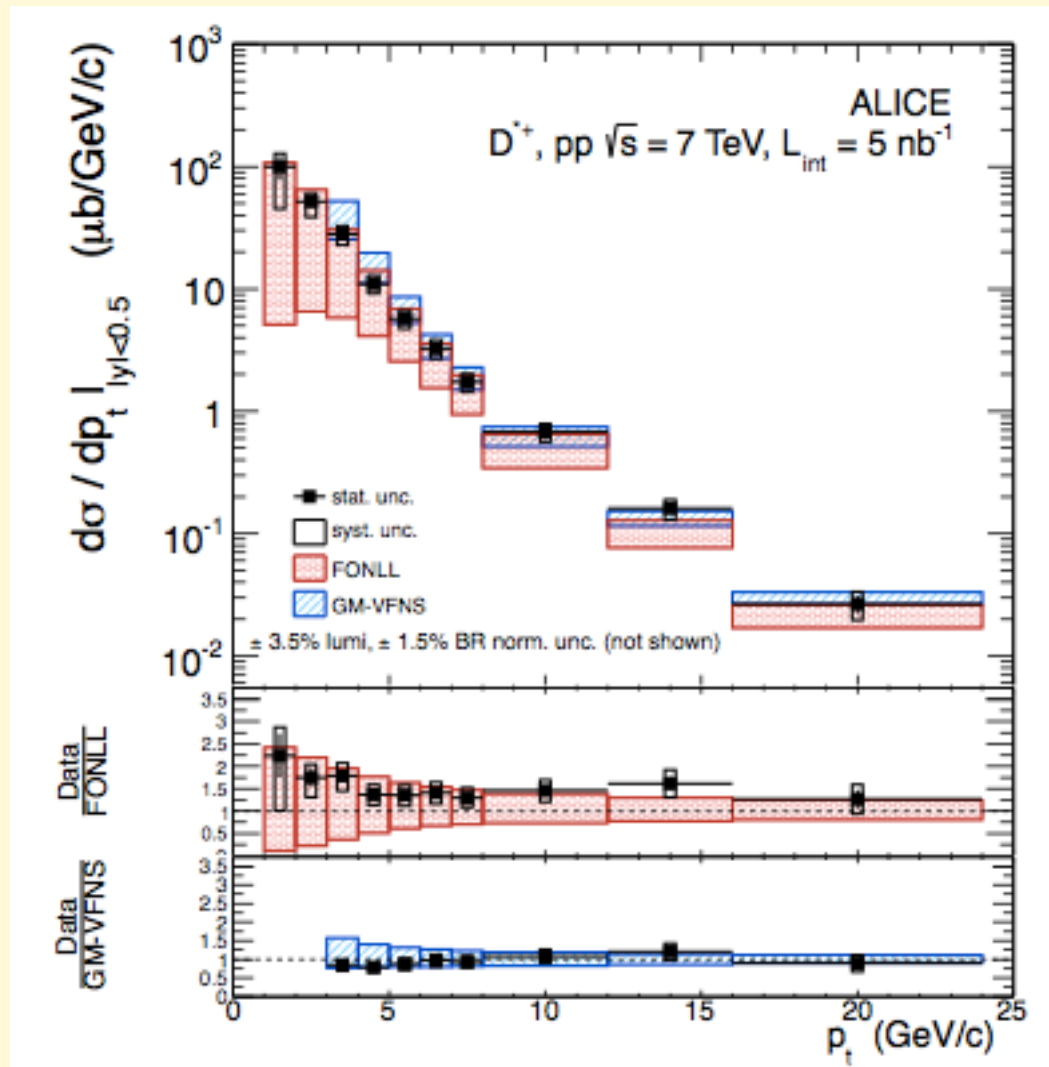




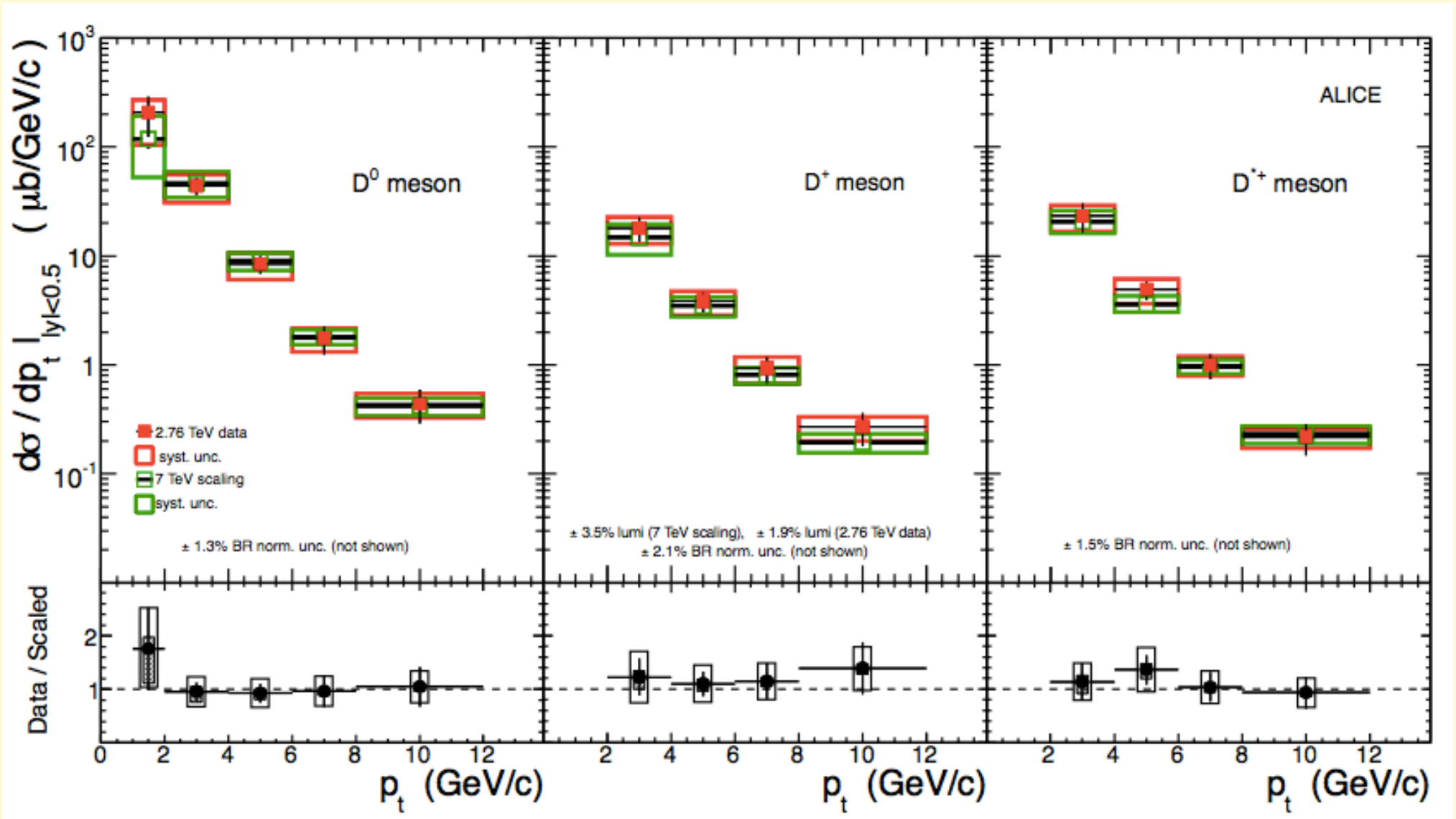


**FONLL**, Cacciari et al, JHEP 05 (1998) 007, JHEP 10 (2012) 137

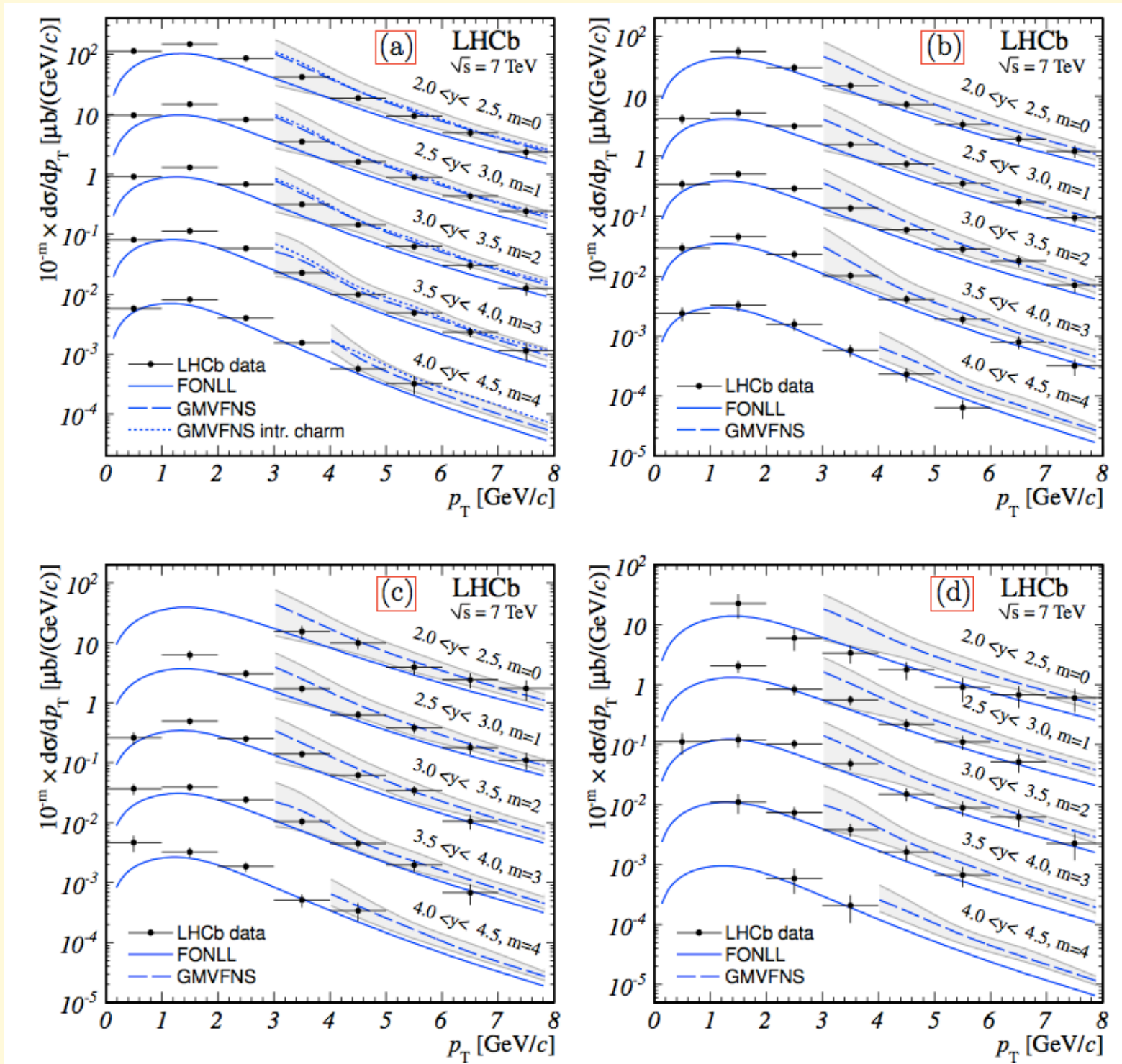
**GM-VFNS**, Kniehl et al, Phys. Rev. D71 (2005) 014018, EPJC72 (2012) 2082



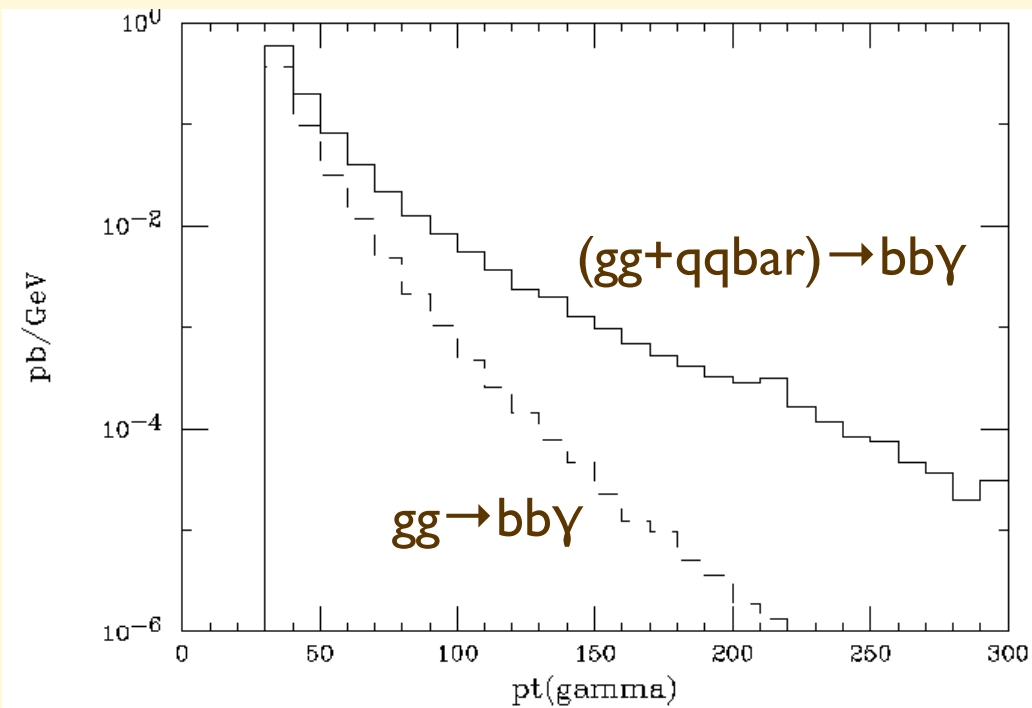
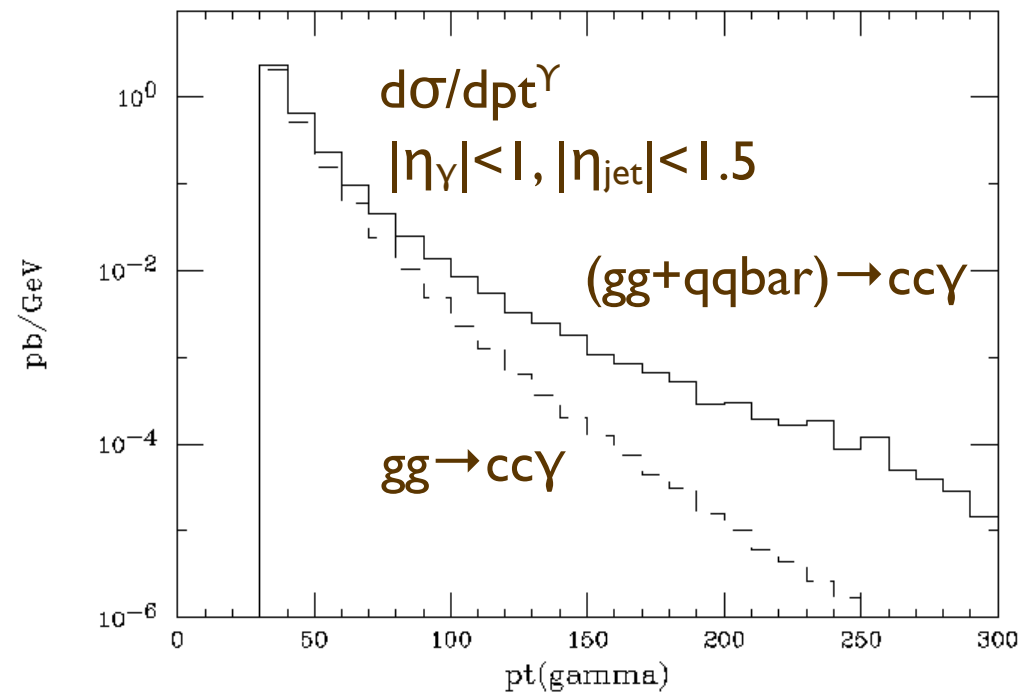




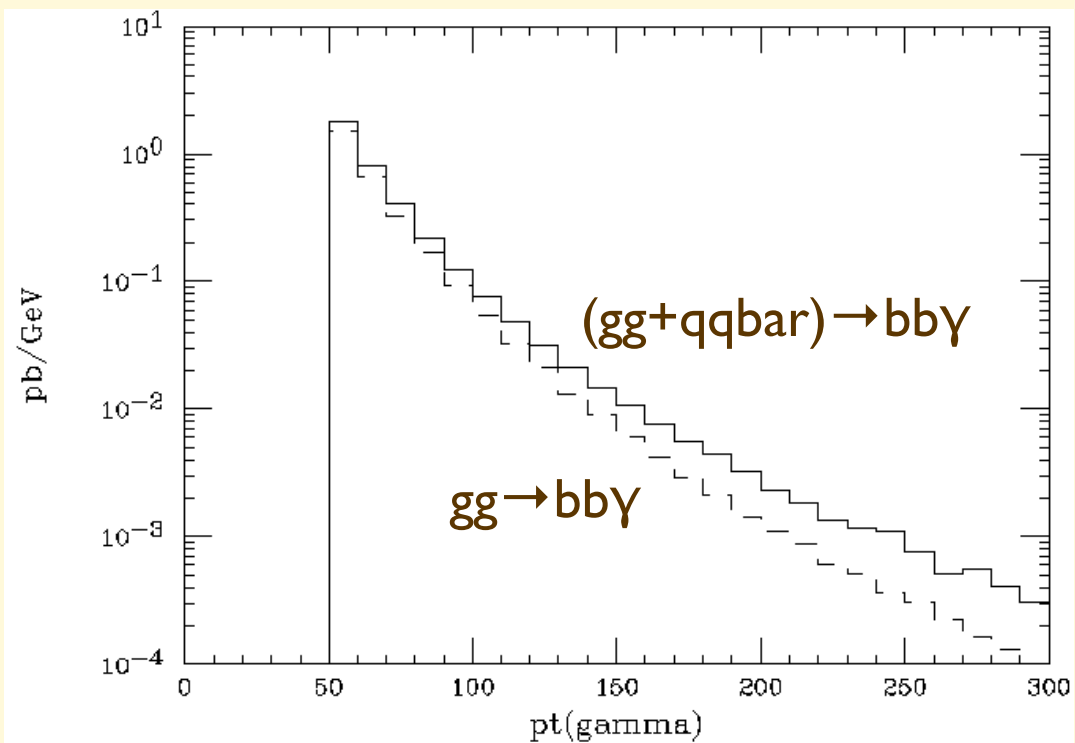
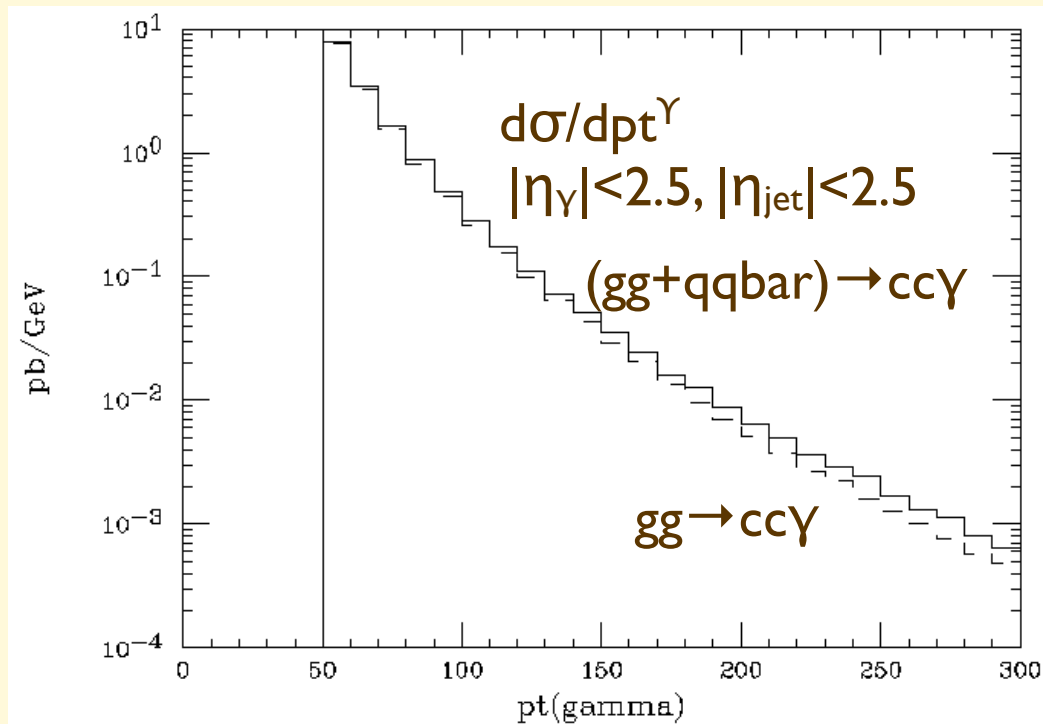
2.76 TeV data vs rescaling of 7 TeV data using FONLL 7 TeV/2.76 TeV ratio



# Tevatron

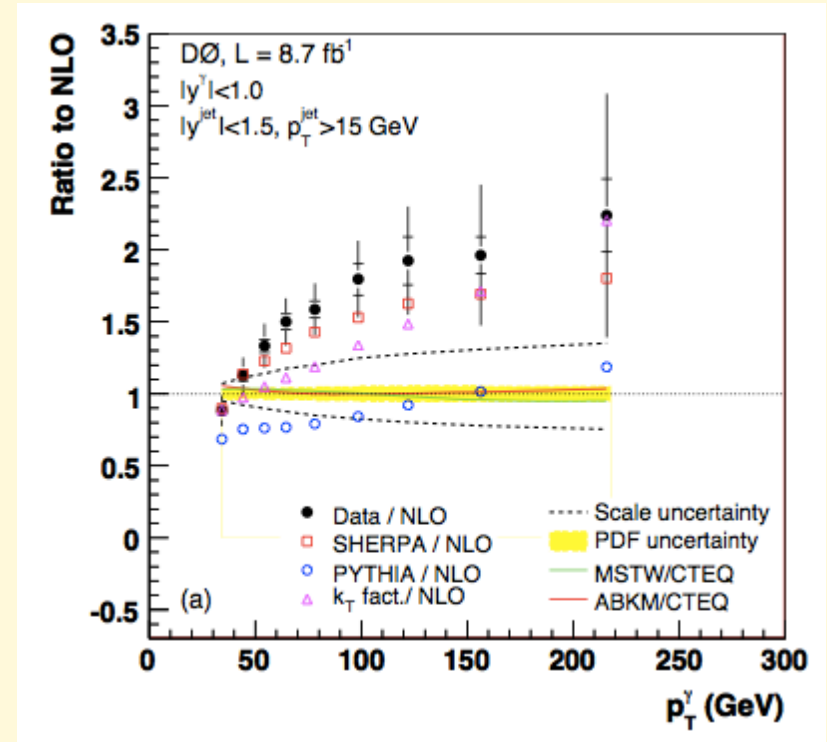
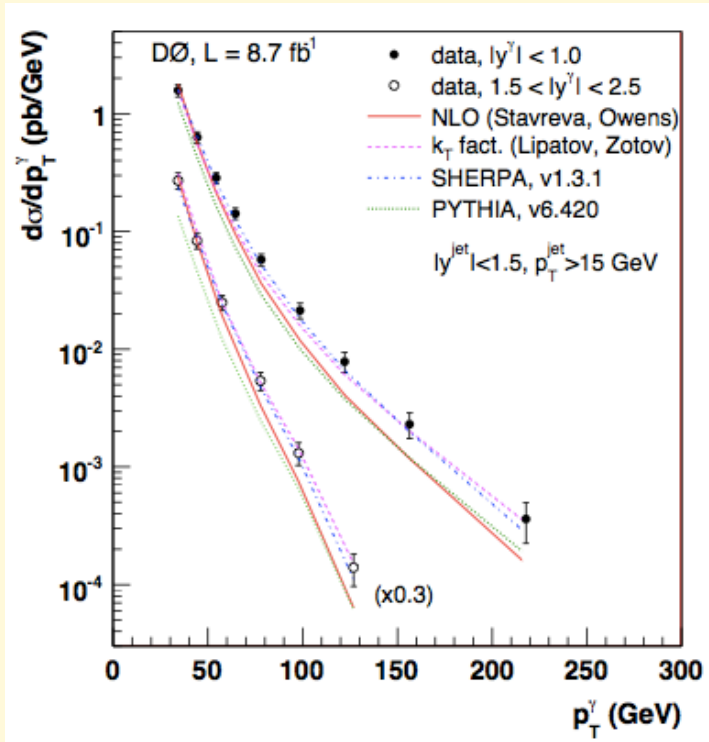


# LHC

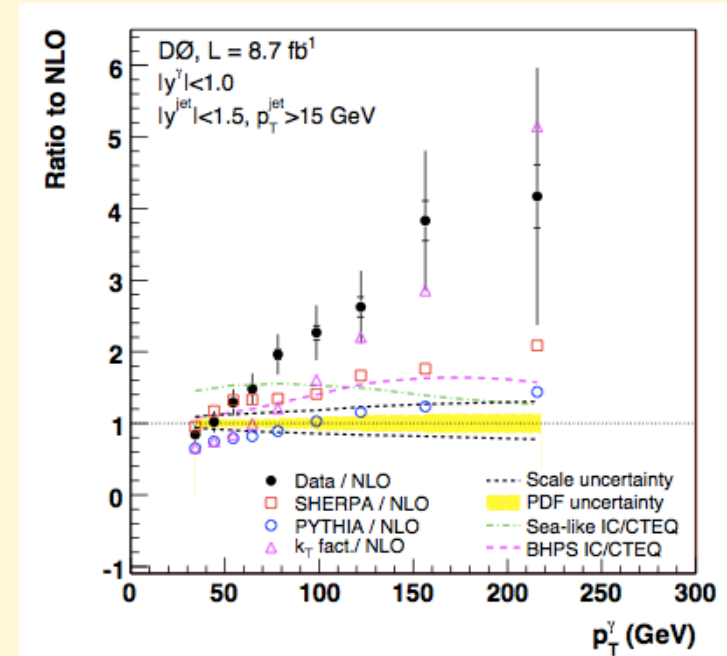
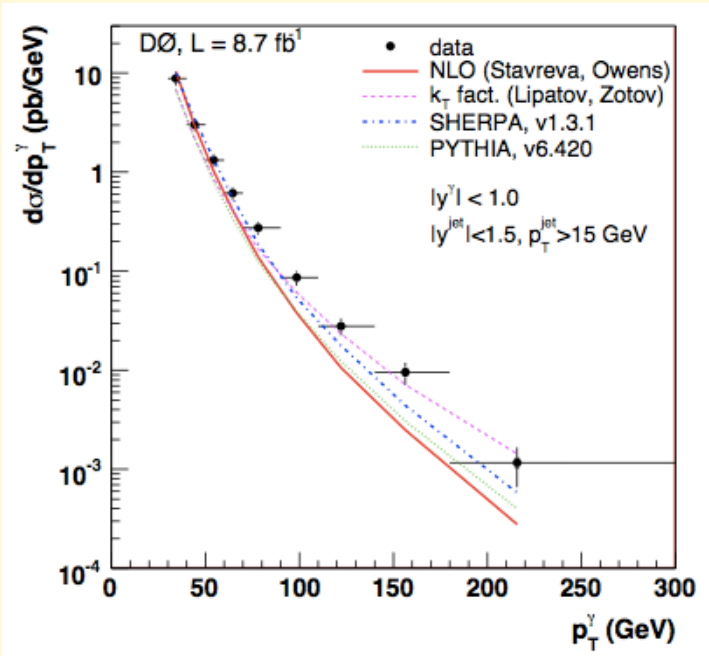


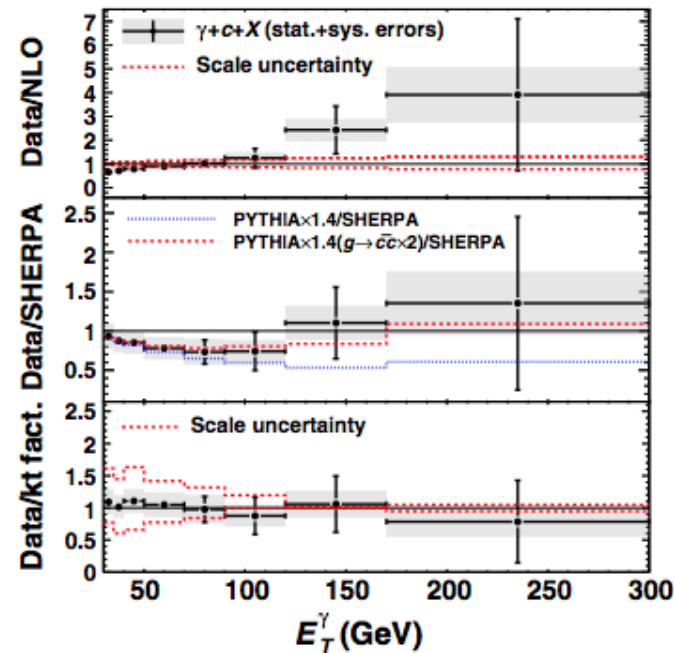
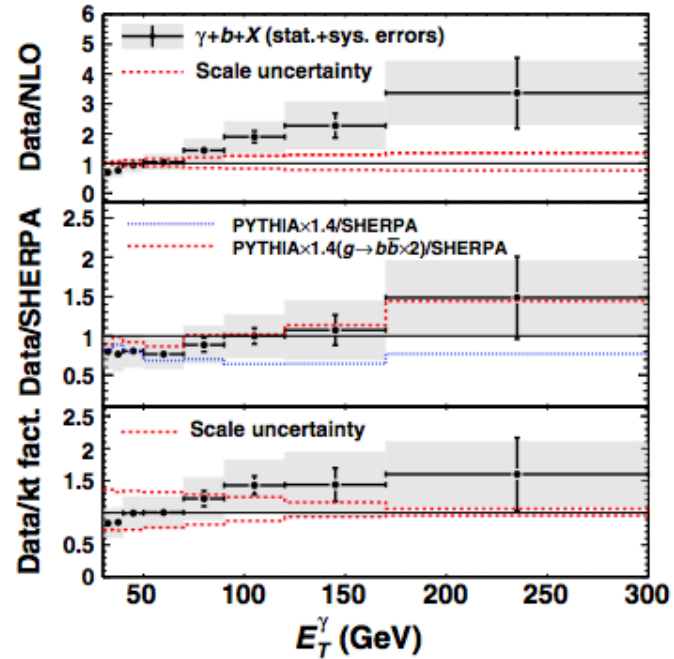
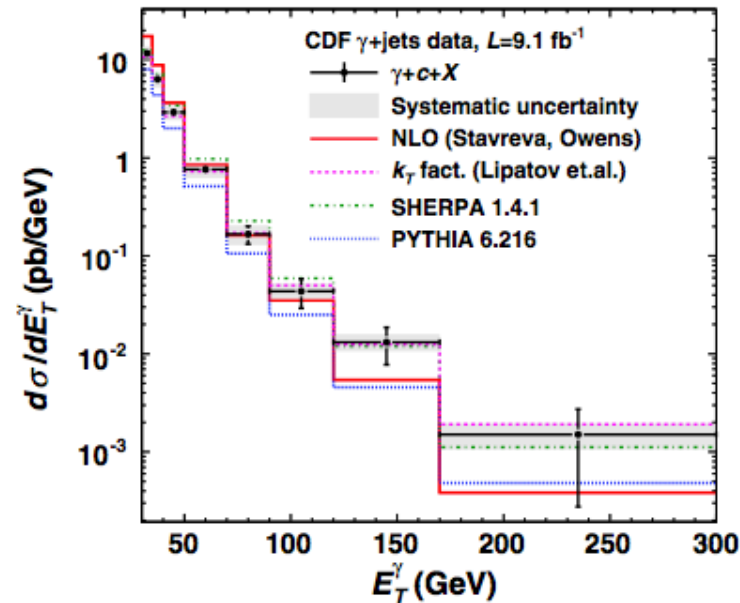
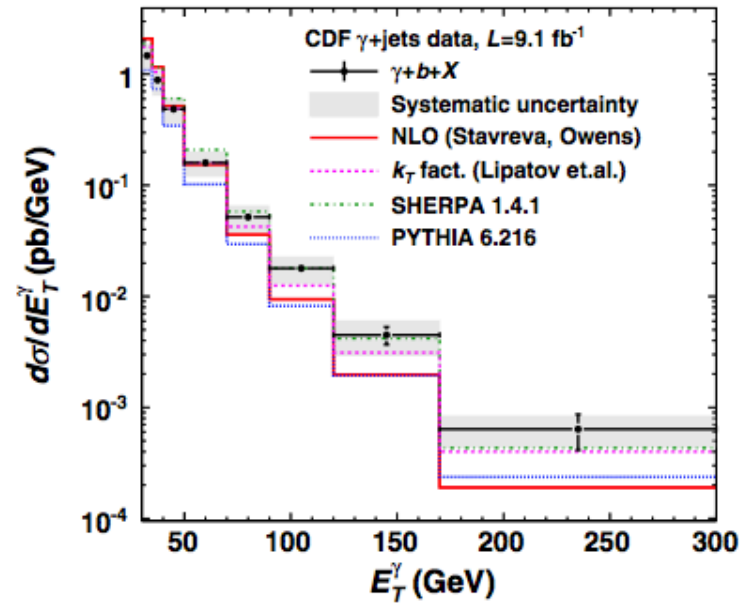
# D0 results

$\gamma+b$



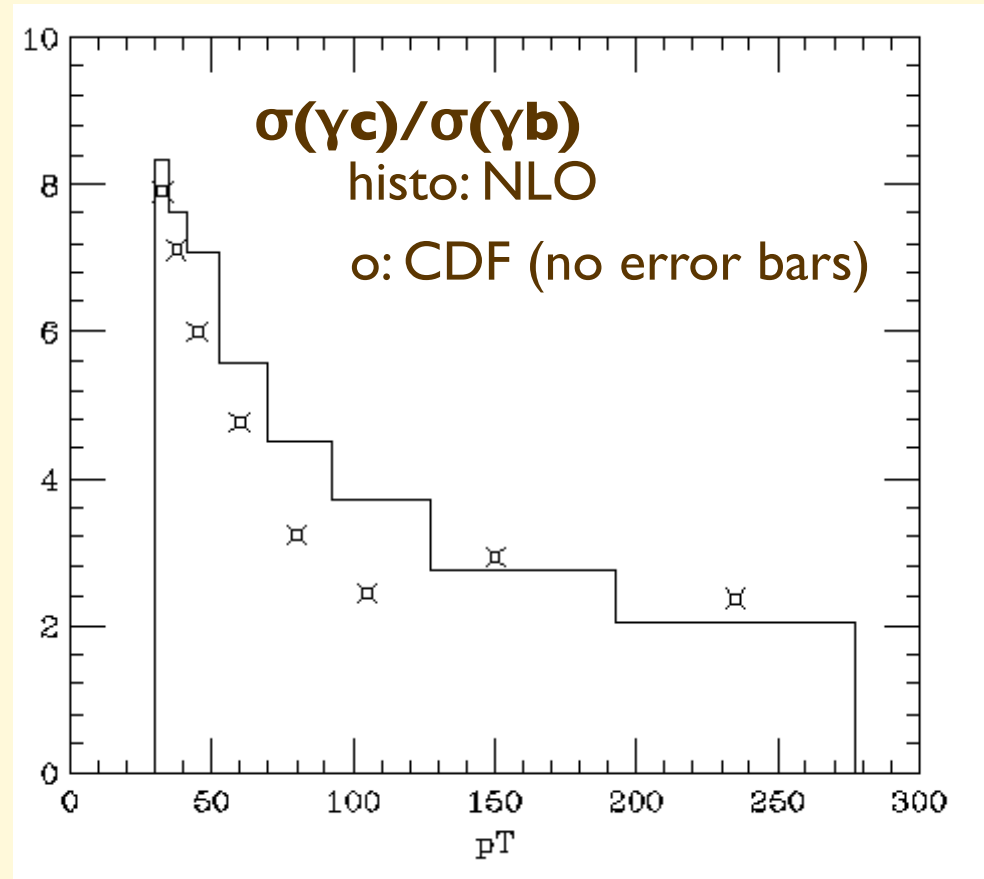
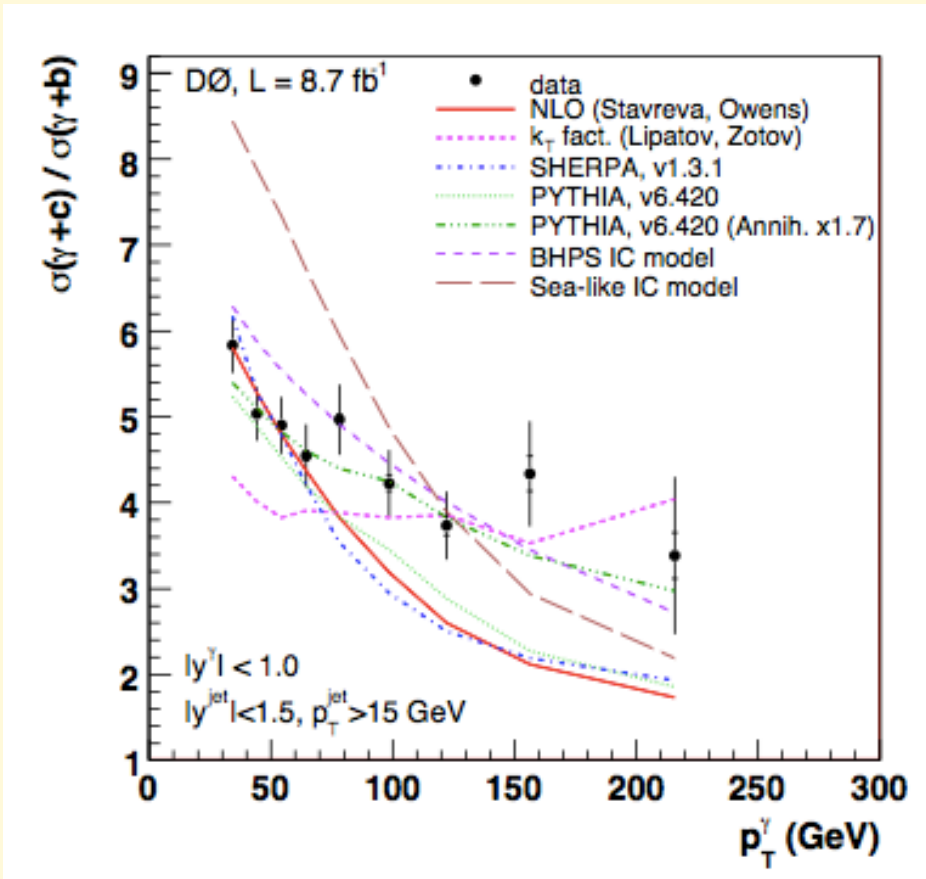
$\gamma+c$



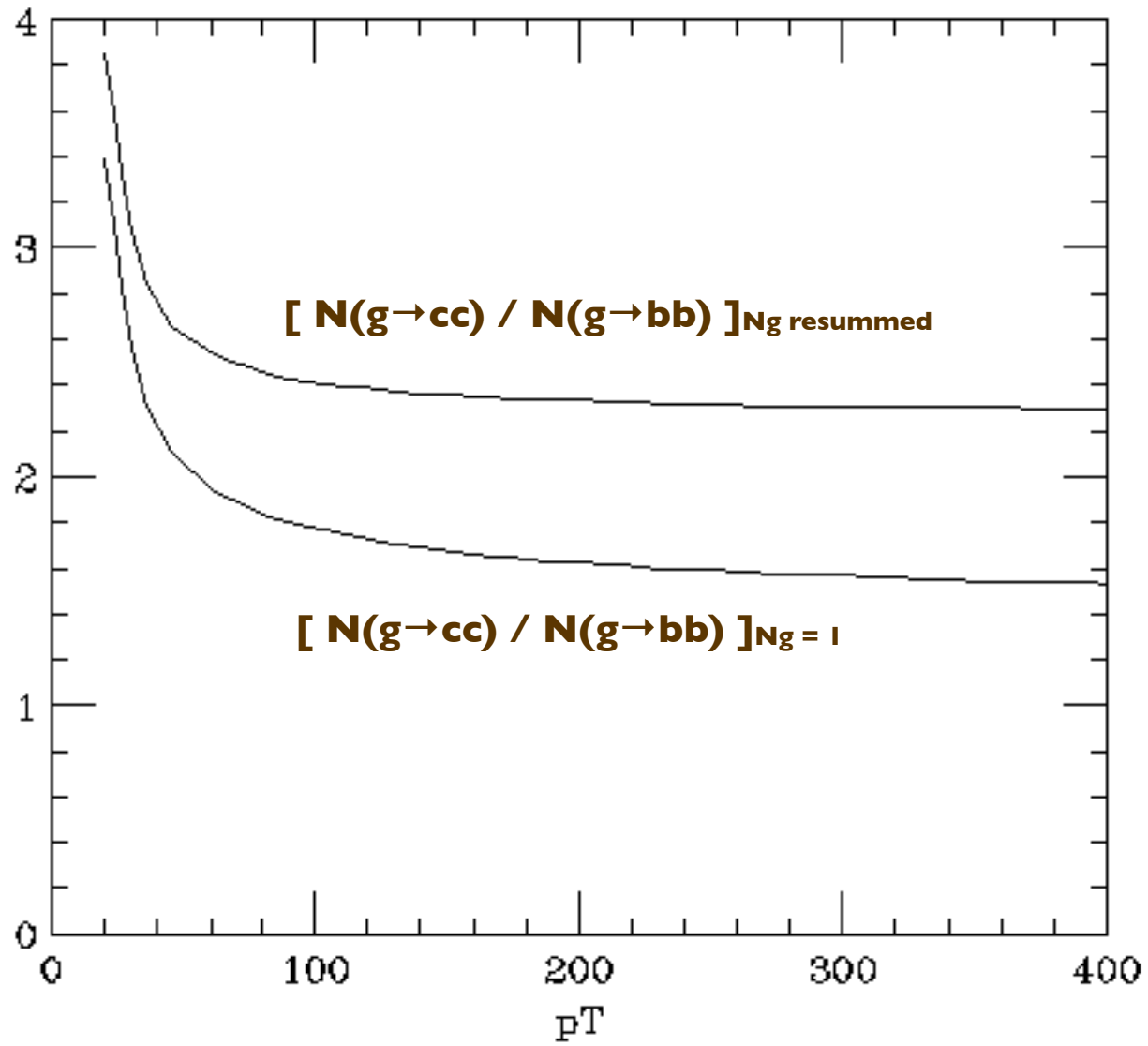


Similar trend in cdf and D0: an excess in both c and b.

However, if we look at the ratios c/b:



Are the CDF and D0 results consistent with each other?



Thus  $\gamma c$  production at large  $p_t$  at the LHC is more sensitive to the charm PDF than at the Tevatron, where gluon splitting has a major role at large  $p_t$ .

