

CERN Accelerator School
Chavannes de Bogis, Switzerland
7th February 2014

Beam-Beam Interactions

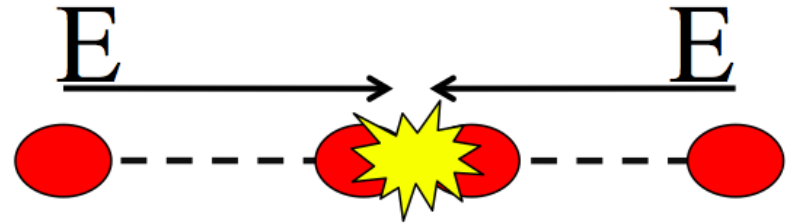
Tatiana Pieloni (BE-ABP-ICE)

Thanks to the Beam-beam Team:
D. Banfi, J. Barranco, X. Buffat and W. Herr



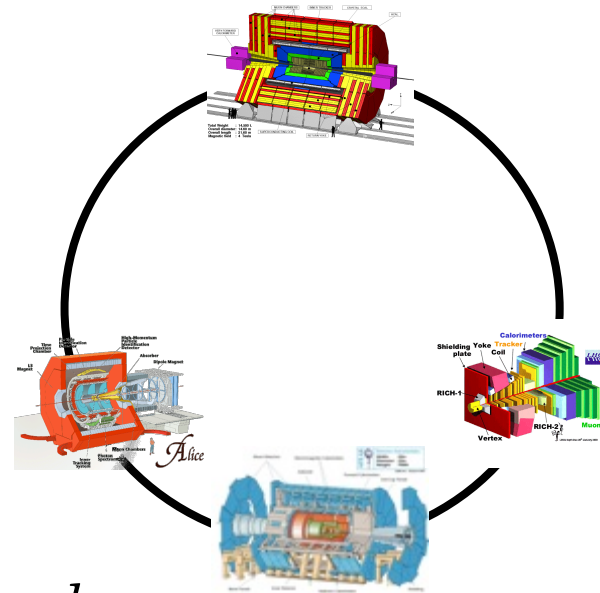
Hadron Circular Colliders

$$E^* \approx 2 \times E$$



$$N_{event}/s = L \cdot \sigma_{event}$$

$$L \propto \frac{N_p^2}{\sigma_x \sigma_y} \cdot n_b \cdot f_{rev}$$



Bunch intensity:

$$N_p = 1.15 - 1.65 \cdot 10^{11} \text{ ppb}$$

Transverse Beam size:

$$\sigma_{x,y} = 16 - 30 \mu m$$

Number of bunches

$$1370 - 2808$$

Revolution frequency

$$11 \text{ kHz}$$

When do we have beam-beam effects?

➤ They occur when two beams get closer and collide

➤ Two types

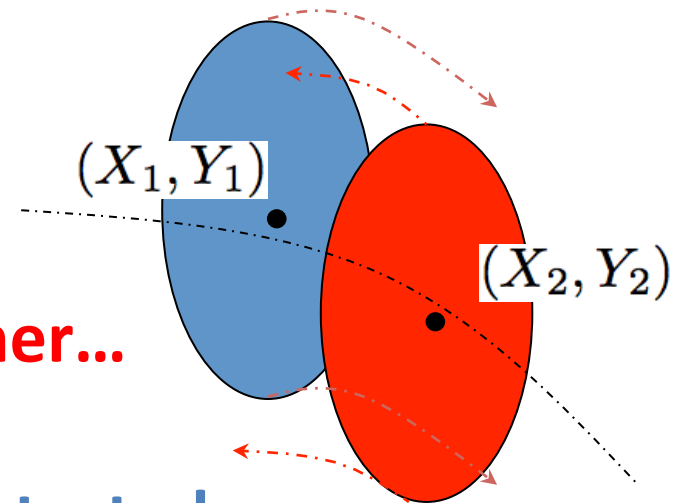
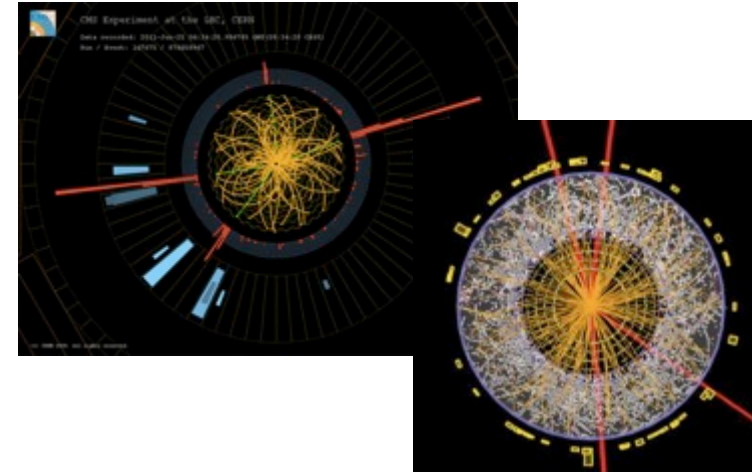
➤ High energy collisions between two particles (**wanted**)

➤ Distortions of beam by electromagnetic forces (**unwanted**)

➤ **Unfortunately: usually both go together...**

➤ 0.001% (or less) of particles collide

➤ 99.999% (or more) of particles are distorted

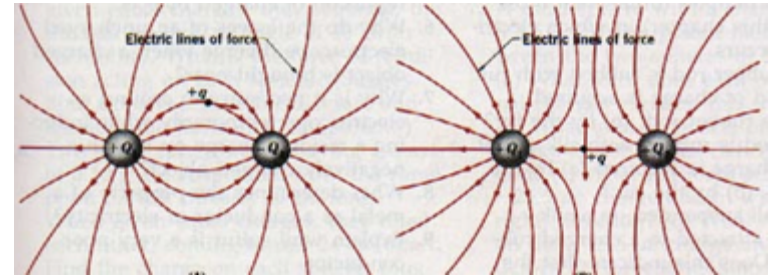


Beam-beam effects: overview

- **Circular Colliders:** interaction occurs at every turn
 - Many effects and problems
 - Try to understand some of them
- Overview of effects (single particle and multi-particle effects)
- Qualitative and physical picture of effects
- Observations from the LHC
- Mathematical derivations and more info in References or at
Beam-beam webpage <http://lhc-beam-beam.web.cern.ch/lhc-beam-beam/>
And CAS Proceedings

Beams EM potential

- Beam is a collection of charges
- Beam is an electromagnetic potential for other charges

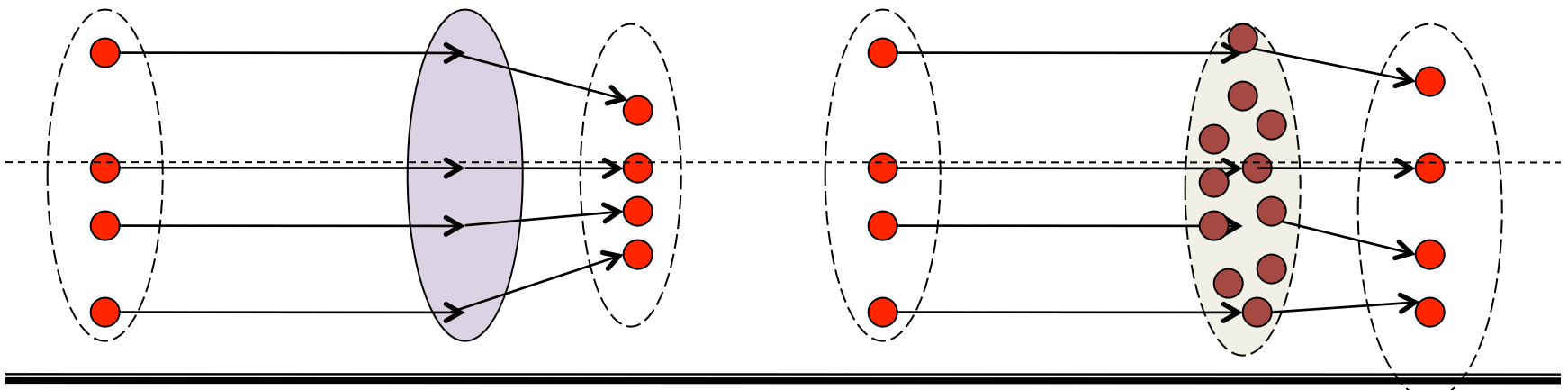


Force on itself (**space charge**) and opposing beam (**beam-beam effects**)

Single particle motion and whole bunch motion **distorted**

Focusing quadrupole

Opposite Beam



A beam acts on particles like an electromagnetic lens, but...

Beam-beam Mathematics

General approach in electromagnetic problems Reference[5] already applied to beam-beam interactions in Reference[1,3, 4]

$$\Delta U = -\frac{1}{\epsilon_0} \rho(x, y, z)$$

Derive potential from Poisson equation for charges with distribution ρ

Solution of Poisson equation

$$U(x, y, z, \sigma_x, \sigma_y, \sigma_z) = \frac{1}{4\pi\epsilon_0} \int \int \int \frac{\rho(x_0, y_0, z_0) dx_0 dy_0 dz_0}{\sqrt{(x - x_0)^2 + (y - y_0)^2 + (z - z_0)^2}}$$

$$\vec{E} = -\nabla U(x, y, z, \sigma_x, \sigma_y, \sigma_z)$$

Then compute the fields

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

From Lorentz force one calculates the force acting on test particle with charge q

Making some assumptions we can simplify the problem and derive analytical formula for the force...

Round Gaussian distributions:

Gaussian distribution for charges:

Round beams:

Very relativistic, Force has only radial component :

$$\sigma_x = \sigma_y = \sigma$$

$$\beta \approx 1 \quad r^2 = x^2 + y^2$$

$$F \propto \frac{N_p}{\sigma} \cdot \frac{1}{r} \cdot \left[1 - e^{-\frac{r^2}{2\sigma^2}} \right]$$

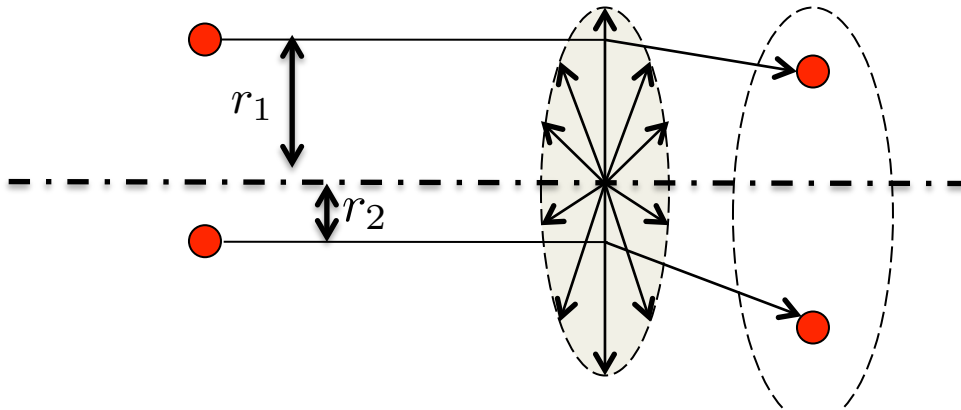
Beam-beam Force

$$\Delta r' = \frac{1}{mc\beta\gamma} \int F_r(r, s, t) dt$$

$$\Delta r' = -\frac{N_p r_0}{r} \cdot \frac{r}{r^2} \left[1 - e^{-\frac{r^2}{2\sigma^2}} \right]$$

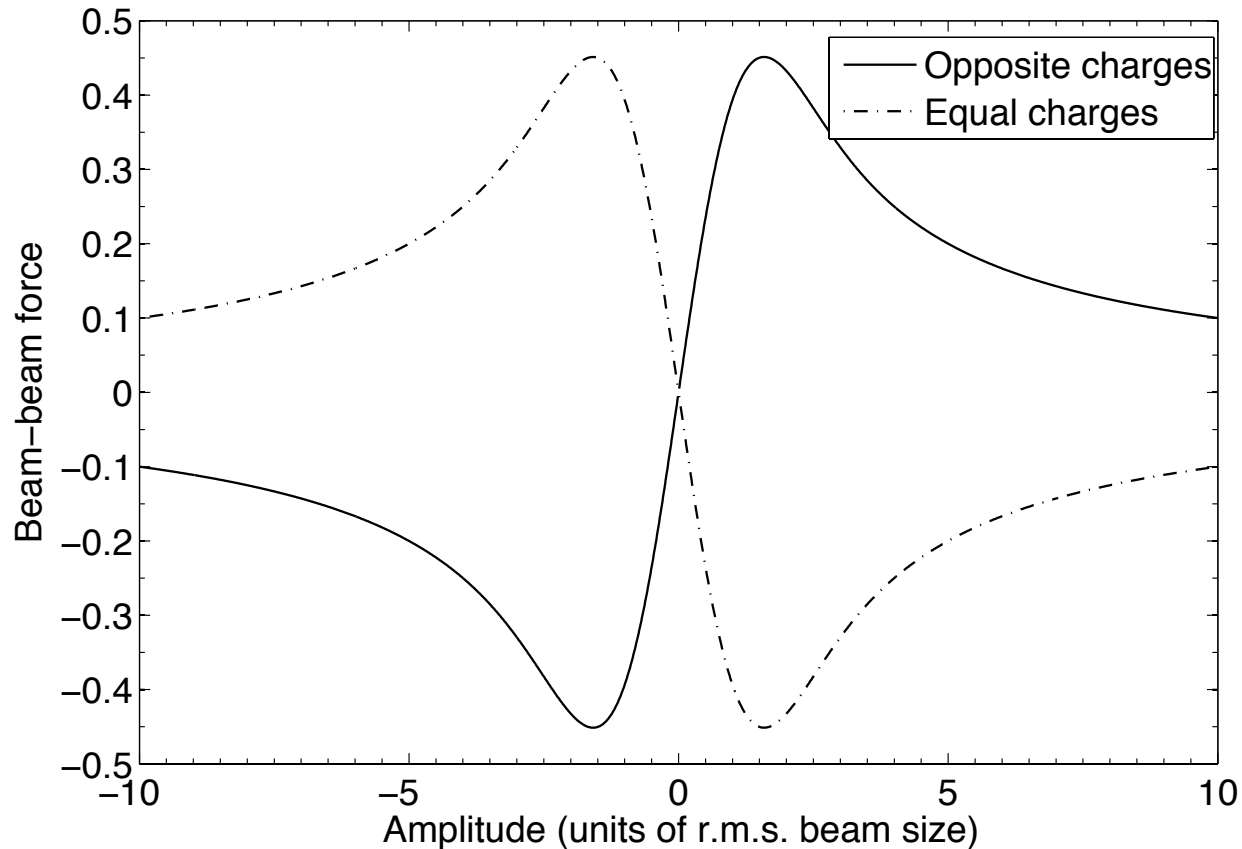
Beam-beam kick obtained
integrating the force over the
collision (i.e. time of passage)

Only radial component in
relativistic case



How does this force looks
like?

Beam-beam Force



$$F_r(r) = \pm \frac{ne^2(1 + \beta_{rel}^2)}{2\pi\epsilon_0} \frac{1}{r} \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right) \right]$$

Why do we care?

Pushing for luminosity means stronger beam-beam effects

$$\mathcal{L} \propto \frac{N_p^2}{\sigma_x \sigma_y} \cdot n_b$$

$$F \propto \frac{N_p}{\sigma} \cdot \frac{1}{r} \cdot \left[1 - e^{-\frac{r^2}{2\sigma^2}} \right]$$

Physics fill lasts for many hours 10h – 24h

Strongest non-linearity in a collider YOU CANNOT AVOID!

La Une | Mercredi 4 juillet 2012 | Dernière mise à jour 18:08

Mon journal numérique | Abonnements | Météo: Genève 20°

Tribune de Genève

Recherche

Home | Emploi | Auto | Petites Annonces

CARNET NOIR
L'acteur de télévision Andy Griffith est mort à 86 ans

FACEBOOK
CHAMP-DOLLON
Pour s'être plaint sur Facebook, un gardien est puni

FORMULE 1
La pilote Maria De Villota grièvement blessée

GENÈVE SUISSE MONDE ÉCONOMIE BOURSE SPORTS CULTURE PEOPLE VIVRE AUTO HIGH-TECH SAVOIRS SERVICES

PHYSIQUE

Une nouvelle particule a été découverte



Une nouvelle particule a été découverte par des chercheurs du CERN lancés sur la trace du boson de Higgs. Plus...

Mis à jour il y a 2 minutes

Bourse

SMI 6'191.43 -0.04%
Stoxx50 2'437.70 -0.21%
DJIA 12'943.82 +0.56%

Les plus lus

- 1 Le Conseil d'Etat a fait valser ses cadres. Une première
- 2 Les fontaines de Dubai pleurent Whitney Houston
- 3 Champ-Dollon: pour s'être plaint sur Facebook, un gardien est puni
- 4 La pilote Maria De Villota grièvement blessée
- 5 Sentier des Toblerones: Apple censure la ville de Gland

Genève au fil du temps



Aux origines du casino de Genève, Le Kursaal fut l'une des attractions de la ville pendant 150 ans.

[Voir nos galeries photo](#)

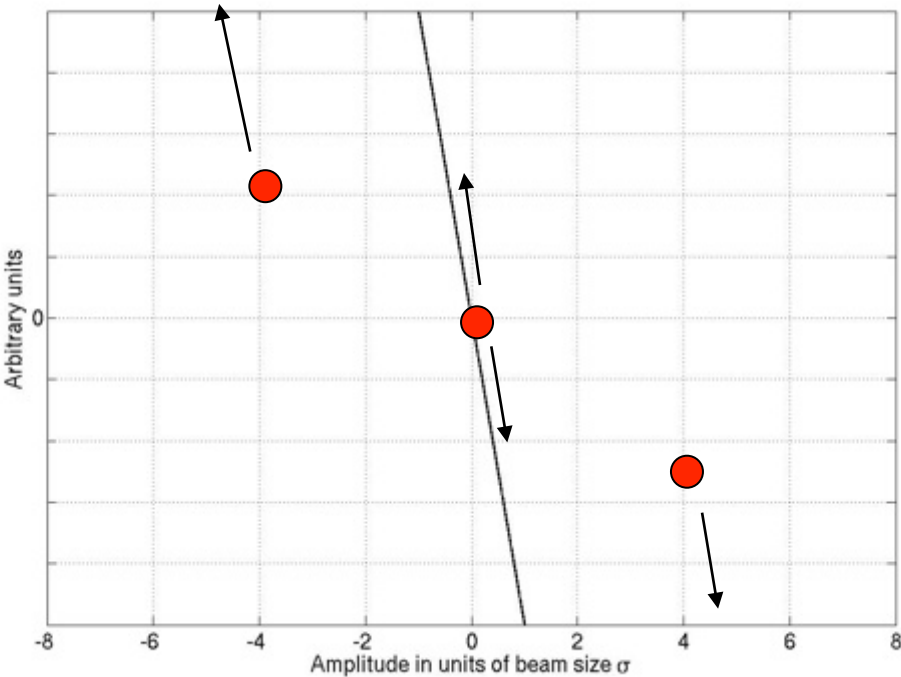
Cet été, bien informé rime avec mobilité

Two main questions:

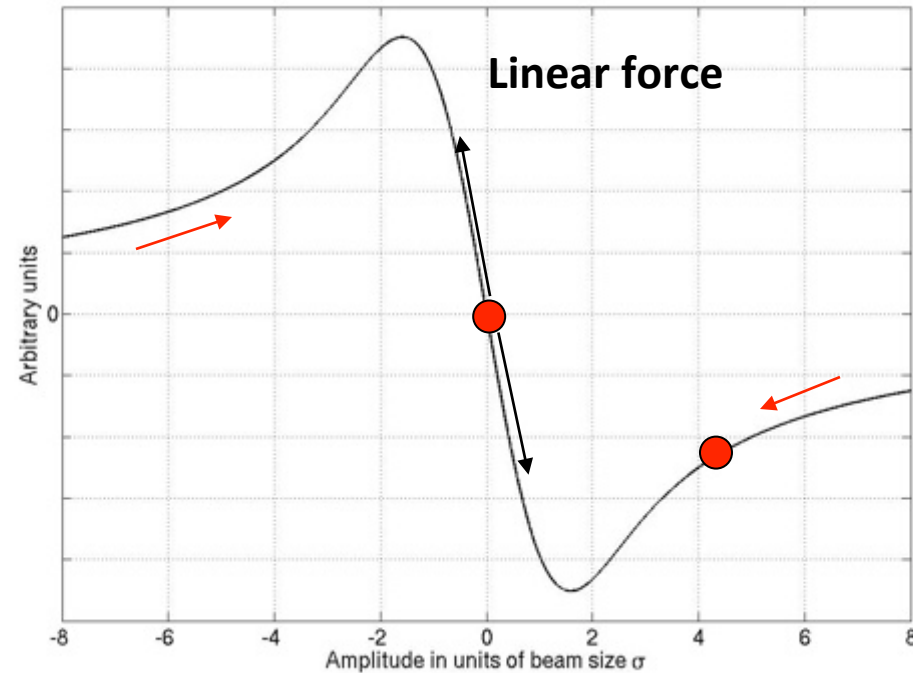
What happens to a single particle?
What happens to the whole beam?

Beam-Beam Force: single particle...

Lattice defocusing quadrupole



Beam-beam force



$$F = -k \cdot r$$

$$F \propto \frac{N_p}{\sigma} \cdot \frac{1}{r} \cdot \left[1 - e^{-\frac{r^2}{2\sigma^2}} \right]$$

For small amplitudes: linear force

For large amplitude: very non-linear

The beam will act as a strong non-linear electromagnetic lens!

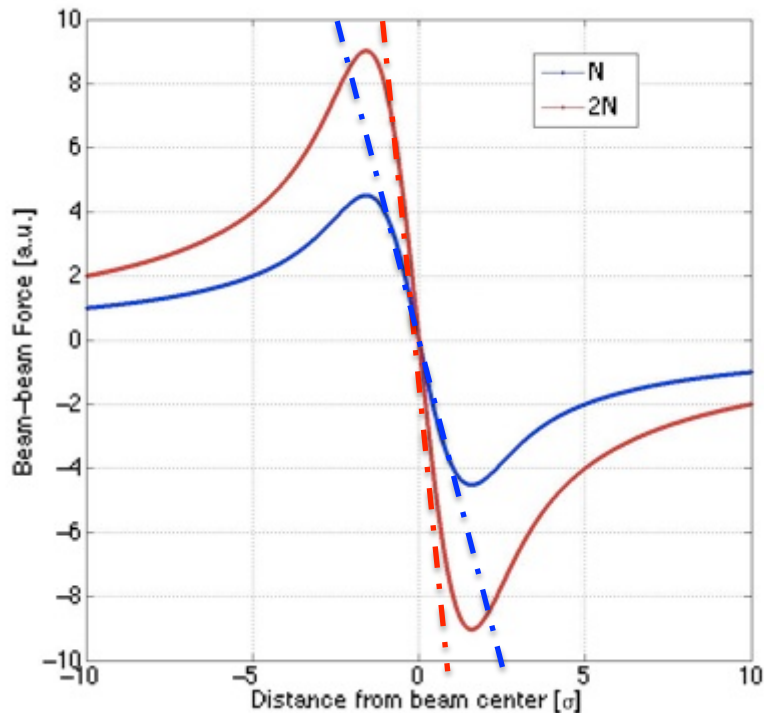
Can we quantify the beam-beam strenght?

Quantifies the strength of the force
but does NOT reflect the nonlinear
nature of the force

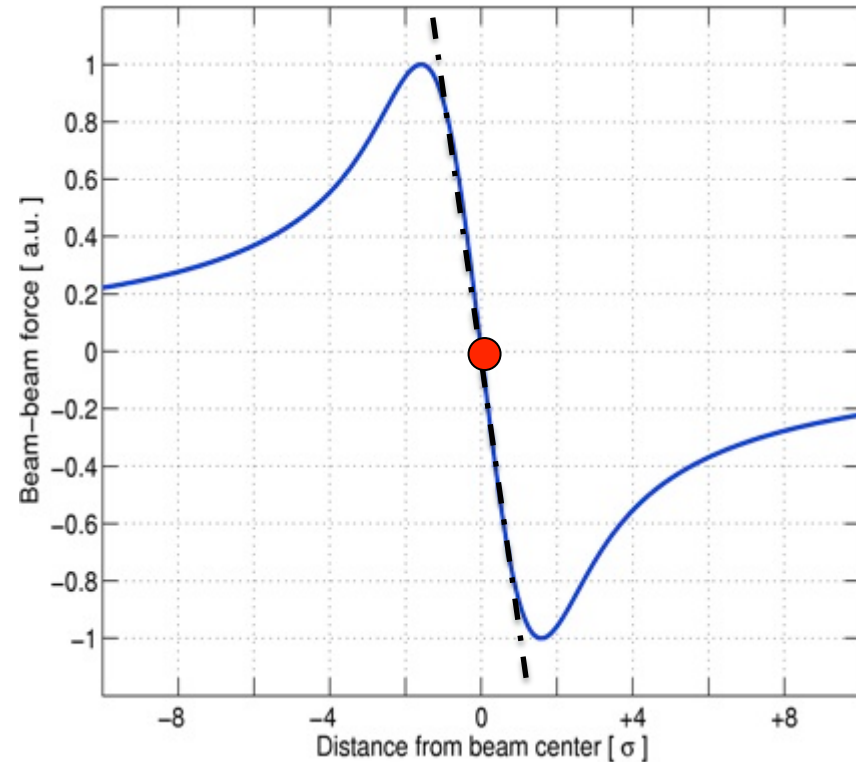
For small amplitudes: linear force

$$F \propto -\xi \cdot r$$

The slope of the force gives you
the beam-beam parameter ξ



Beam-beam force



$$\Delta r' = -\frac{N_p r_0}{r} \cdot \frac{r}{r^2} \cdot \left[1 - e^{-\frac{r^2}{2\sigma^2}} \right]$$

$$\Delta r' = \frac{2N_p r_0}{\gamma} \cdot \frac{1}{r} \cdot \left[1 - \left(1 - \frac{r^2}{2\sigma^2} + \dots \right) \right]$$

Colliders:

For round beams:

$$\xi = \frac{\beta^*}{4\pi} \cdot \frac{\delta(\Delta r')}{\delta r} = \frac{Nr_0\beta^*}{4\pi\gamma\sigma^2}$$

For non-round beams:

$$\xi_{x,y} = \frac{Nr_0\beta_{x,y}^*}{2\pi\gamma\sigma_{x,y}(\sigma_x + \sigma_y)}$$

Examples:

Parameters	LHC nominal	LHC 2012
Intensity $N_{p,e}/\text{bunch}$	$1.15 \cdot 10^{11}$	$1.6 \cdot 10^{11}$
Energy GeV	7000	4000
Beam emittance	$3.75 \mu\text{mrad}$	$2.2\text{-}2.5 \mu\text{mrad}$
Crossing angle (μrad)	285	290
$\beta_{x,y}^*$ (m)	1.25-0.05	0.60-0.60
Luminosity	$1 \cdot 10^{34}$	$7.6 \cdot 10^{33}$
ξ_{bb}	0.0034	0.006

Linear Tune shift

For small amplitudes beam-beam can be approximated as linear force as a quadrupole

$$F \propto -\xi \cdot r$$

Focal length:

$$\frac{1}{f} = \frac{\Delta x'}{x} = \frac{Nr_0}{\gamma\sigma^2} = \frac{\xi \cdot 4\pi}{\beta^*}$$

Beam-beam matrix:

$$\begin{pmatrix} 1 & 0 \\ -\frac{\xi \cdot 4\pi}{\beta^*} & 1 \end{pmatrix}$$

Perturbed one turn matrix with perturbed tune ΔQ and beta function at the IP β^* :

$$\begin{aligned} & \begin{pmatrix} \cos(2\pi(Q + \Delta Q)) & \beta^* \sin(2\pi(Q + \Delta Q)) \\ -\frac{1}{\beta^*} \sin(2\pi(Q + \Delta Q)) & \cos(2\pi(Q + \Delta Q)) \end{pmatrix} \\ &= \begin{pmatrix} 1 & 0 \\ -\frac{1}{2f} & 1 \end{pmatrix} \cdot \begin{pmatrix} \cos(2\pi Q) & \beta_0^* \sin(2\pi Q) \\ -\frac{1}{\beta_0^*} \sin(2\pi Q) & \cos(2\pi Q) \end{pmatrix} \cdot \begin{pmatrix} 1 & 0 \\ -\frac{1}{2f} & 1 \end{pmatrix} \end{aligned}$$

Linear tune

Solving the one turn matrix one can derive the tune shift ΔQ and the perturbed beta function at the IP β^* :

Tune is changed

$$\cos(2\pi(Q + \Delta Q)) = \cos(2\pi Q) - \frac{\beta_0^* \cdot 4\pi\xi}{\beta^*} \sin(2\pi Q)$$

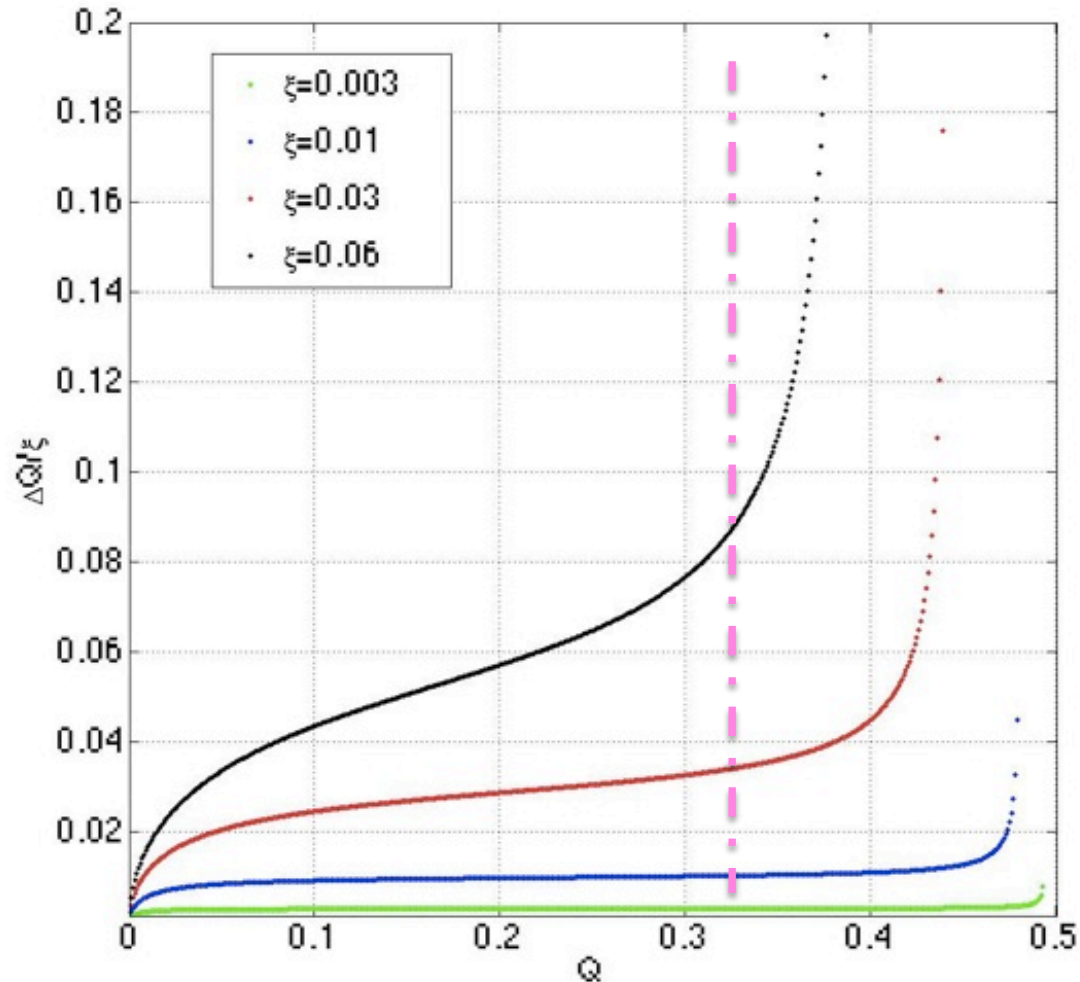
β -function is changed:

$$\frac{\beta^*}{\beta_0^*} = \frac{\sin(2\pi Q)}{\sin(2\pi(Q + \Delta Q))}$$

...how do they change?

Tune dependence of tune shift and dynamic beta

Tune shift as a function of tune



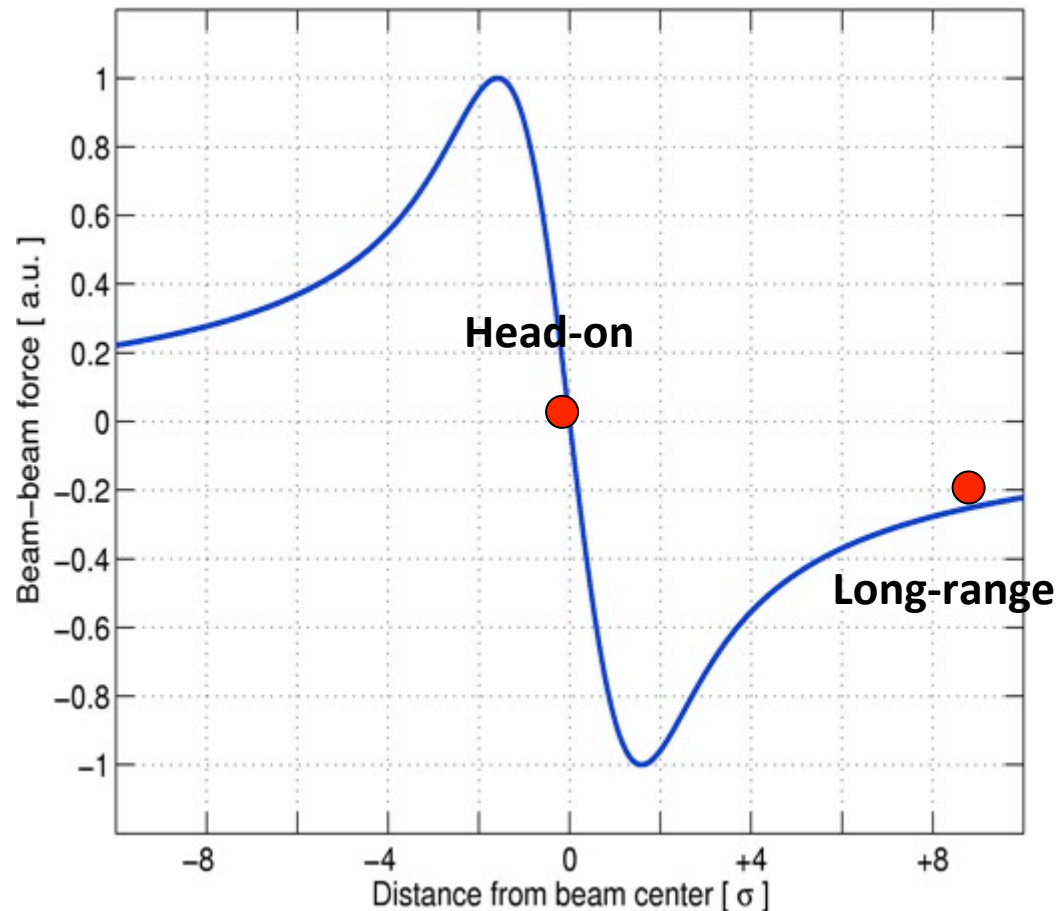
Larger ξ



Strongest variation with Q

Head-on and Long-range interactions

Beam-beam force



$$L \propto \frac{N_p^2}{\sigma_x \sigma_y} \cdot n_b \cdot f_{rev}$$

Other beam passing in the center force: **HEAD-ON** beam-beam interaction

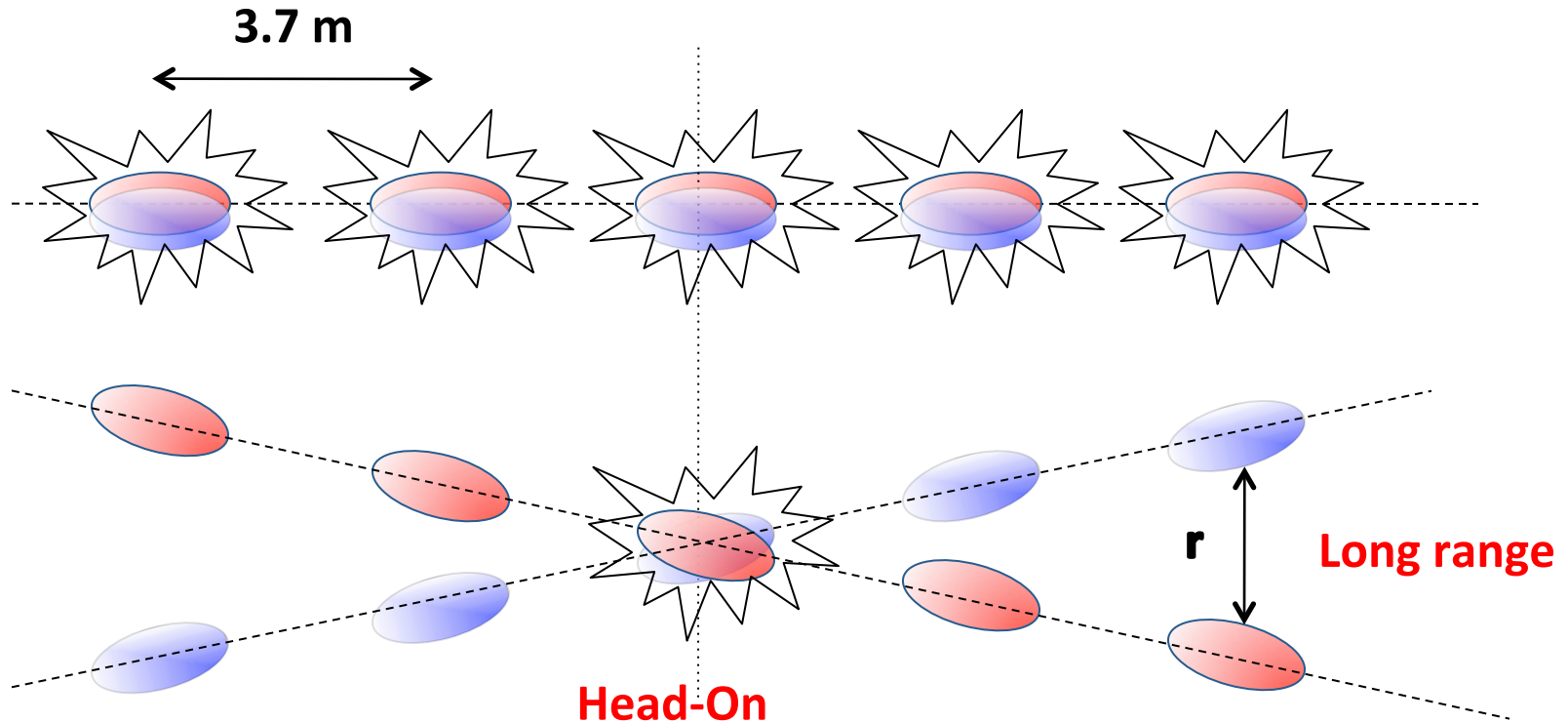
Other beam passing at an offset of the force: **LONG-RANGE** beam-beam interaction

Multiple bunch Complications

MANY INTERACTIONS

$$\mathcal{L} \propto \frac{N_p^2}{\sigma_x \sigma_y} \cdot n_b$$

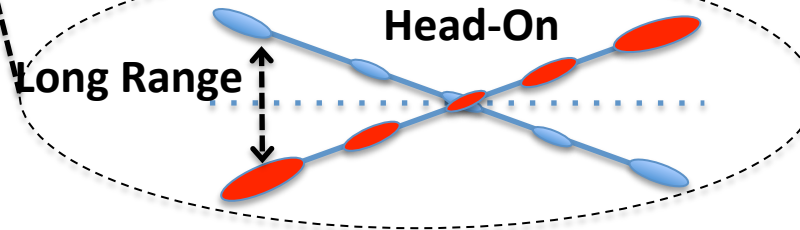
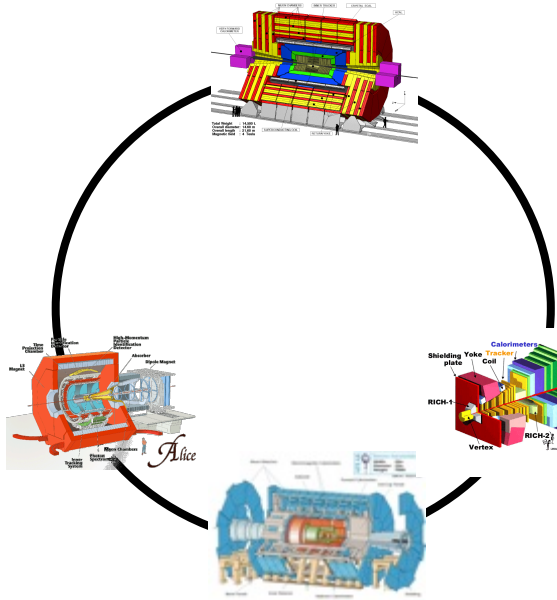
Num. of bunches : $n_b = 2808$



For 25ns case 124 BBIs per turn: 4 HO and 120 LR

LHC, KEKB... colliders

- Crossing angle operation
- High number of bunches in train structures



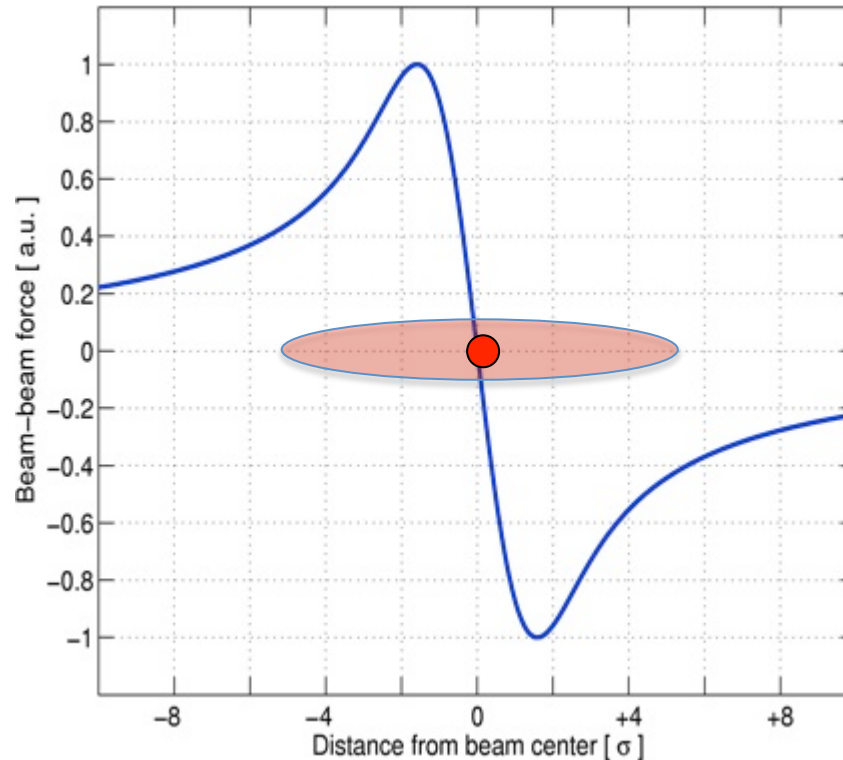
72 bunches

|||||

	SppS	Tevatron	RHIC	LHC
Number Bunches	6	36	109	2808
LR interactions	9	70	0	120/40
Head-on interactions	3	2	2	4

A beam is a collection of particles

Beam-beam force

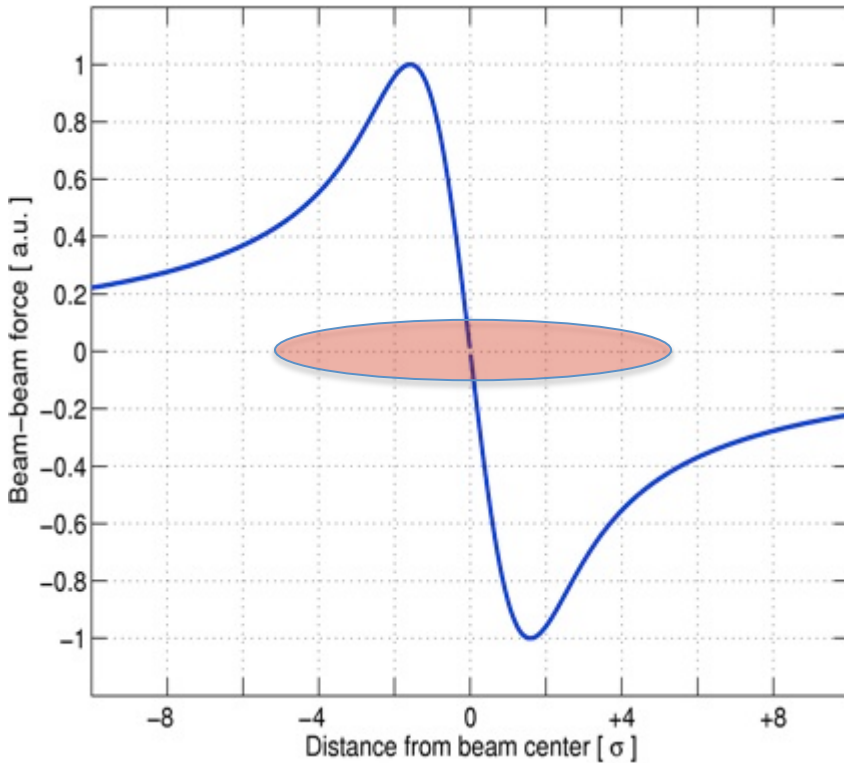


**Beam 2 passing in the center of force produce by Beam 1
Particles of Beam 2 will experience different ranges of the beam-beam forces**

**Tune shift as a function of amplitude (detuning with amplitude or
tune spread)**

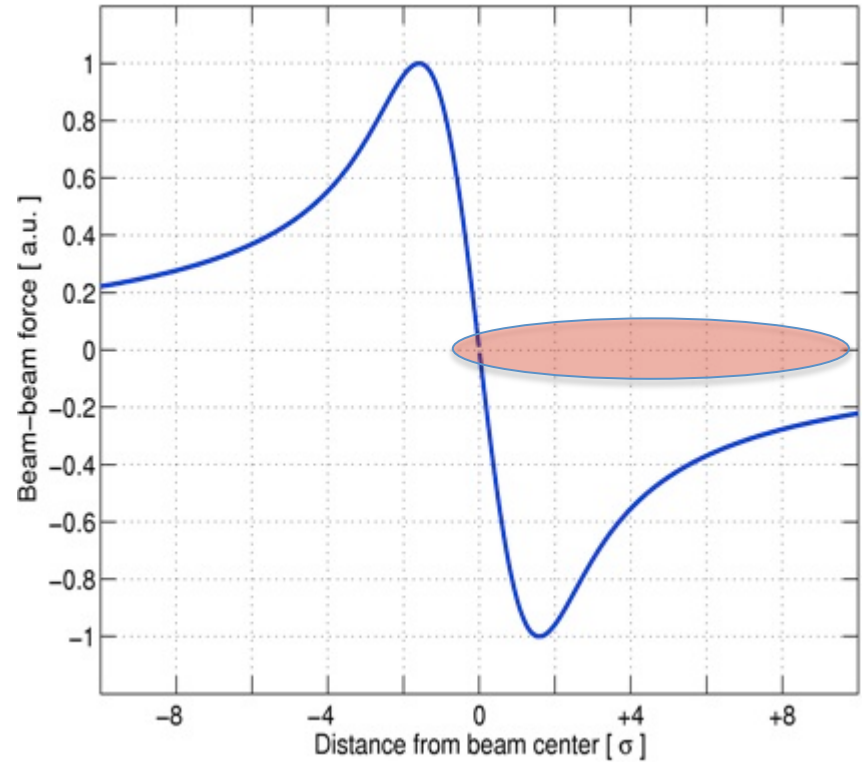
A beam will experience all the force range

Beam-beam force



Second beam passing in the center
HEAD-ON beam-beam interaction

Beam-beam force

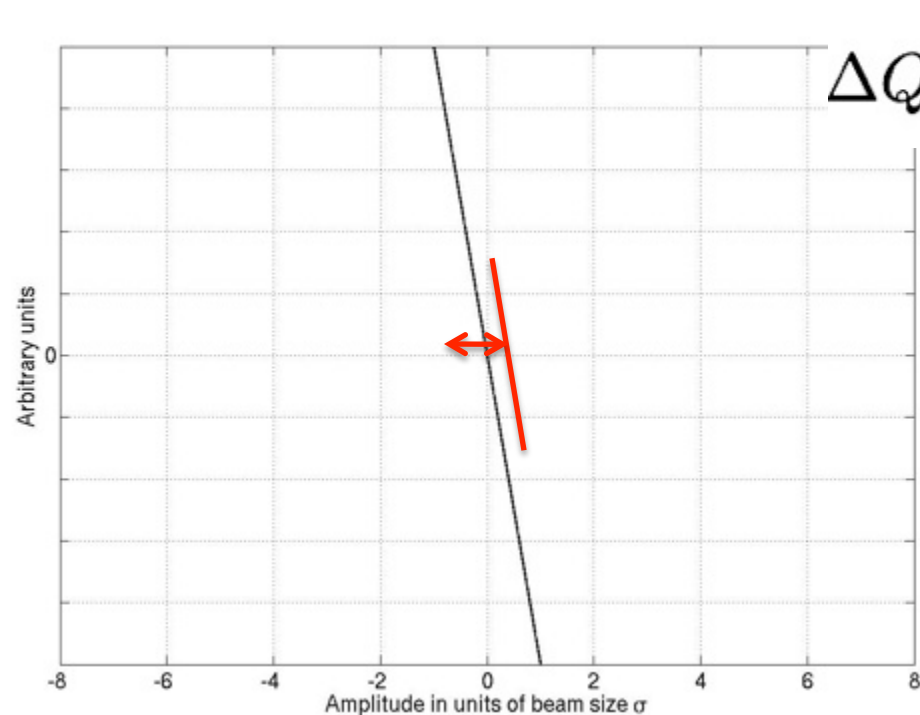


Second beam displaced offset
LONG-RANGE beam-beam interaction

Different particles will see different force

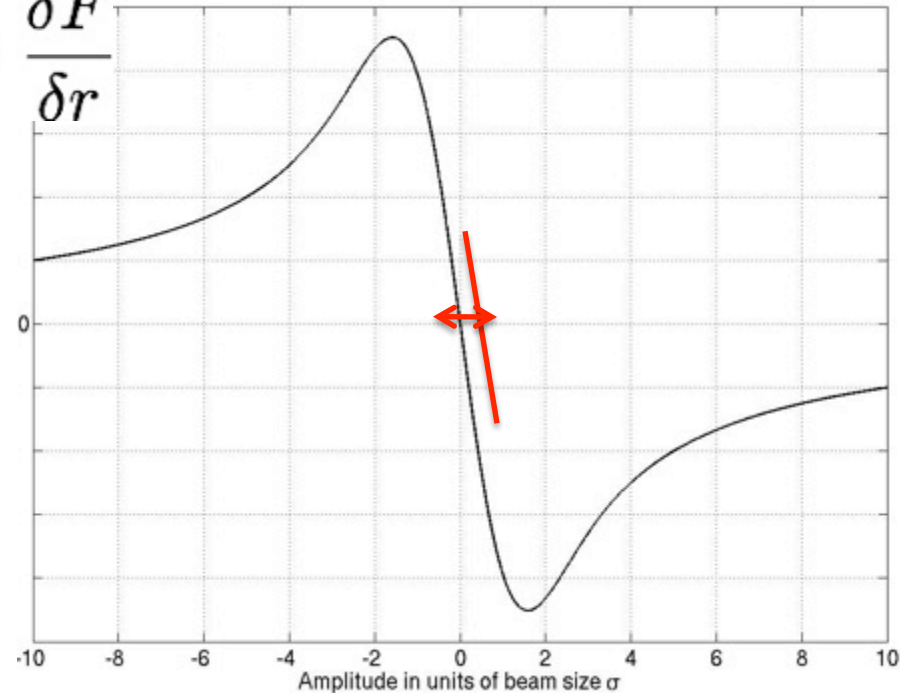
Detuning with Amplitude for head-on

Instantaneous tune shift of test particle when it crosses the other beam is related to the derivative of the force with respect to the amplitude



$$\Delta Q_{quad} = const$$

$$\Delta Q \propto \frac{\delta F}{\delta r}$$



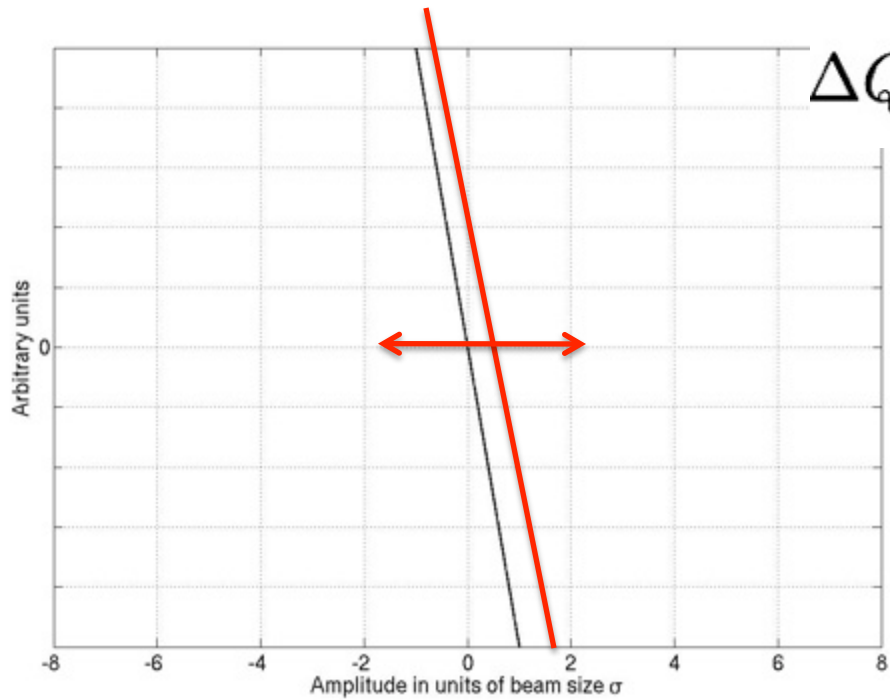
$$\Delta Q_{bb} \approx const$$

For small amplitude test particle
linear tune shift

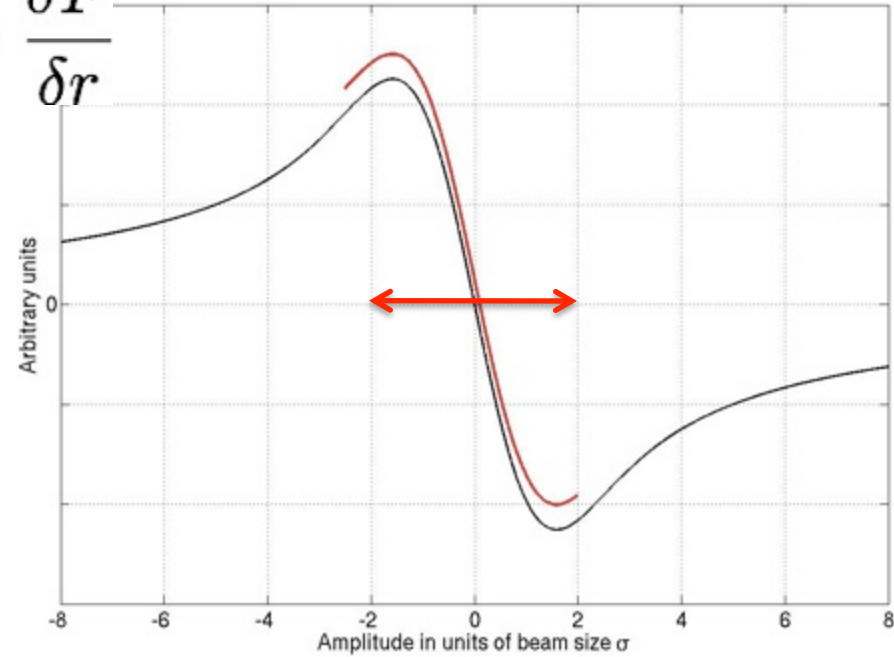
$$\lim_{r \rightarrow 0} \Delta Q(r) = -\frac{Nr_0\beta^*}{4\pi\gamma\sigma^2} = \xi$$

Detuning with Amplitude for head-on

Beam with many particles this results in a tune spread



$$\Delta Q \propto \frac{\delta F}{\delta r}$$



$$\Delta Q_{quad} = const$$

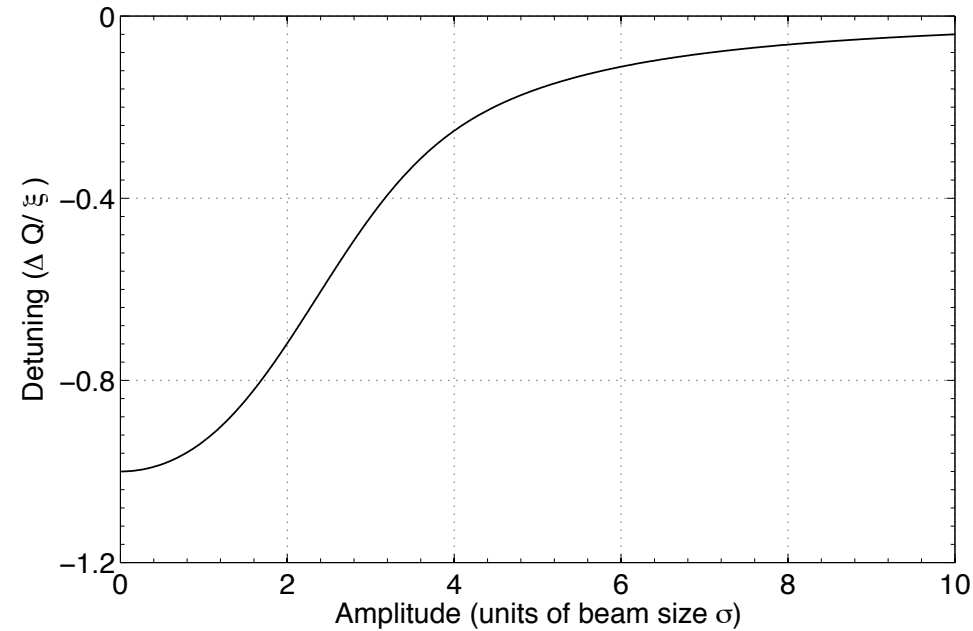
$$\Delta Q_{bb} \neq const$$

$$\Delta Q(x) = \frac{Nr_0\beta}{4\pi\gamma\sigma^2} \cdot \frac{1}{(\frac{x}{2})^2} \cdot (\exp -(\frac{x}{2})^2 I_0(\frac{x}{2})^2 - 1)$$

Mathematical derivation in Ref [3] using Hamiltonian formalism and in Ref [4] using Lie Algebra

Head-on detuning with amplitude and footprints

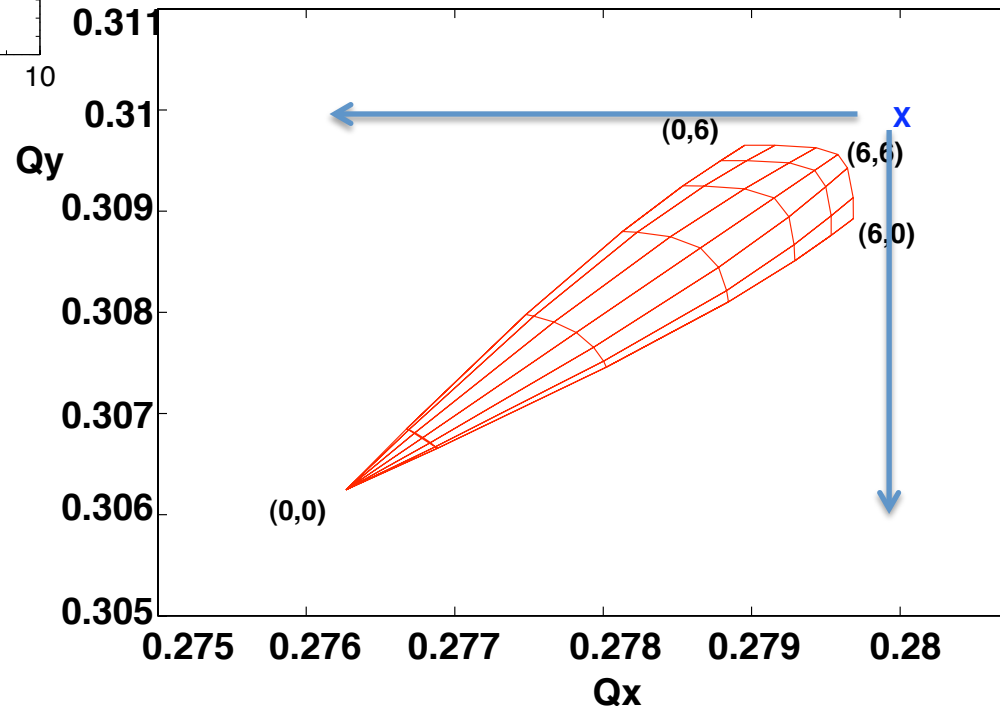
1-D plot of detuning with amplitude



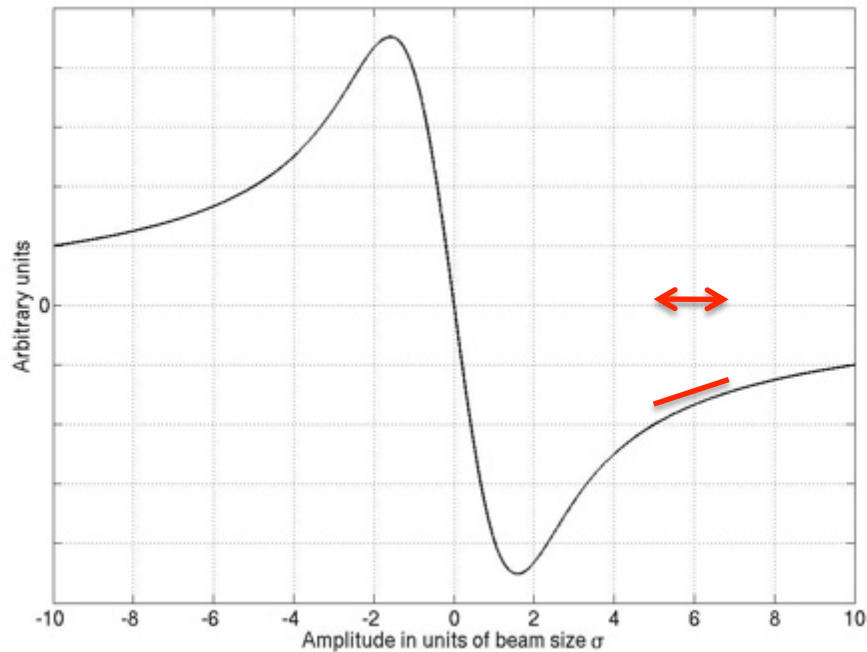
FOOTPRINT
2-D mapping of the detuning with
amplitude of particles

And in the other plane?
THE SAME DERIVATION
same tune spread

Tune footprint for head-on collision



And for long-range interactions?

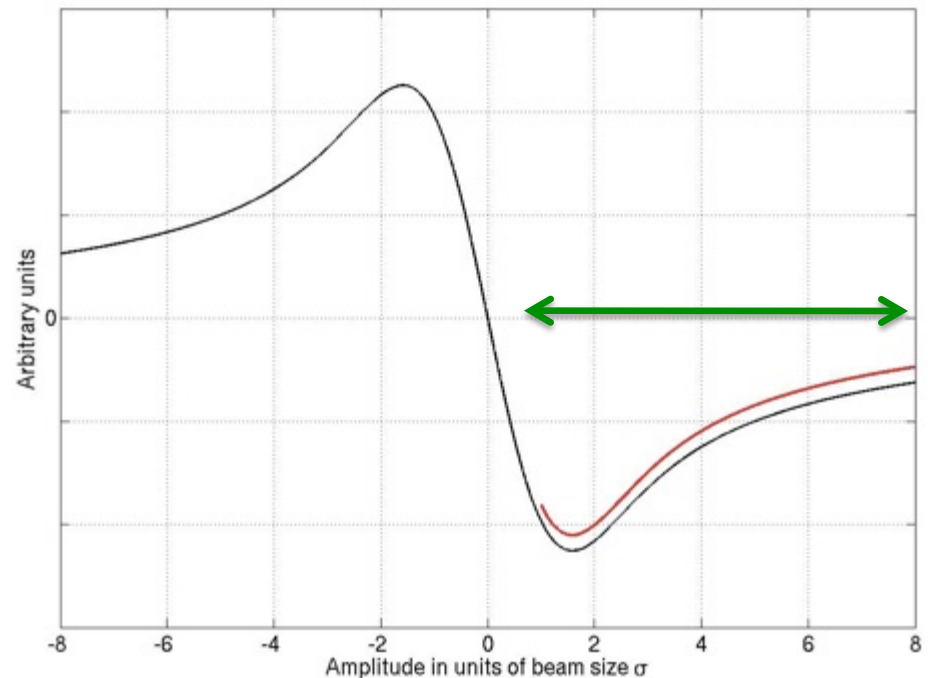


Second beam centered at d (i.e. 6σ)

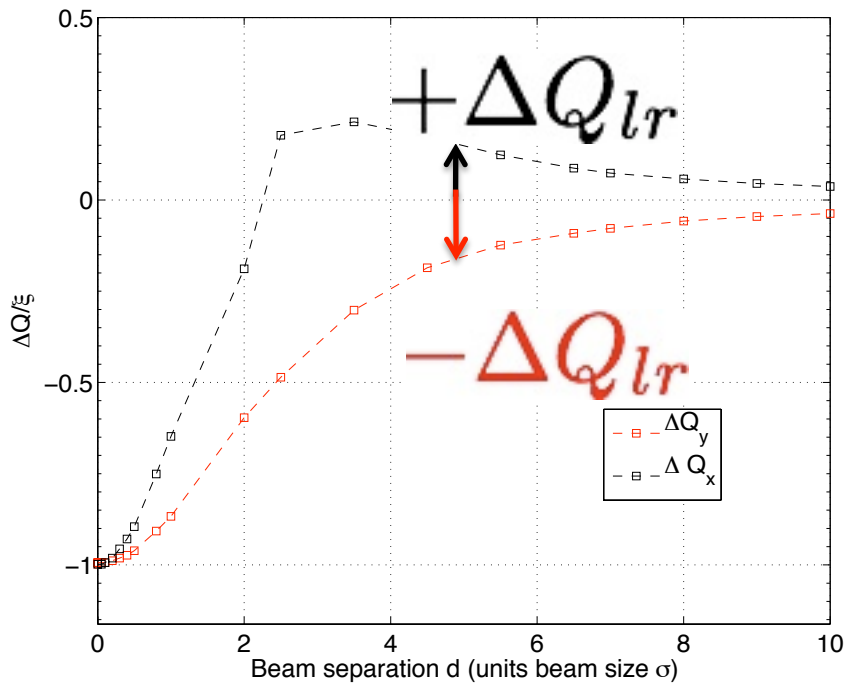
- Small amplitude particles **positive tune shifts**
- Large amplitude can go to **negative tune shifts**

Long range tune shift scaling for distances $d > 6\sigma$

$$\Delta Q_{lr} \propto -\frac{N}{d^2}$$



Long-range footprints



Separation in vertical plane!

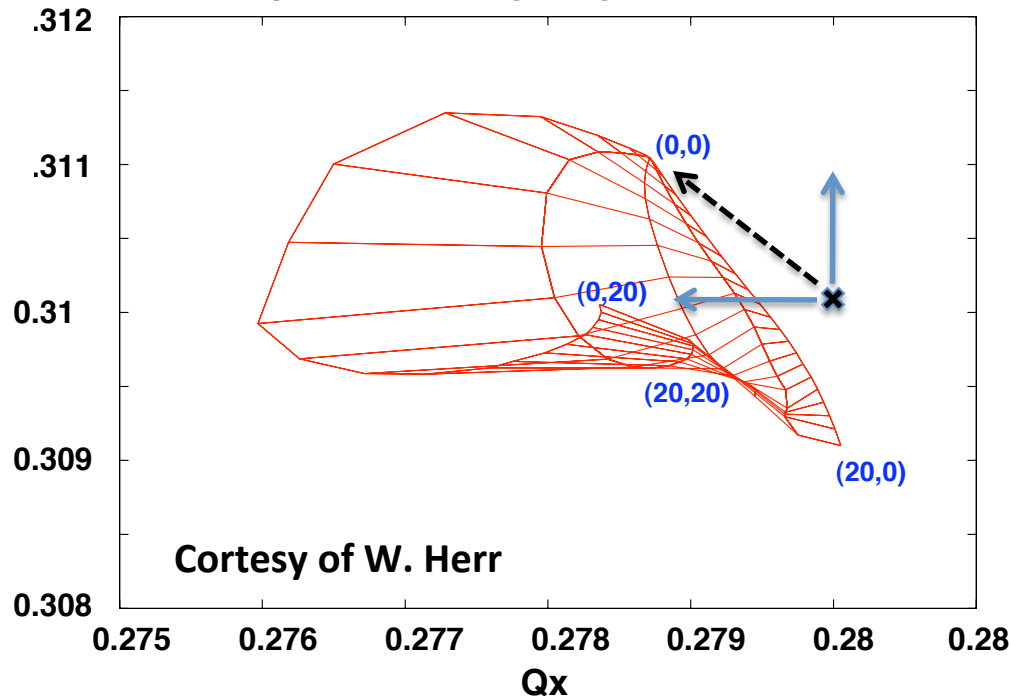
And in horizontal plane?

The test particle is centered with the opposite beam

tune spread more like for head-on at large amplitudes

The picture is more complicated now the LARGE amplitude particles see the second beam and have larger tune shift

footprint from long range interactions

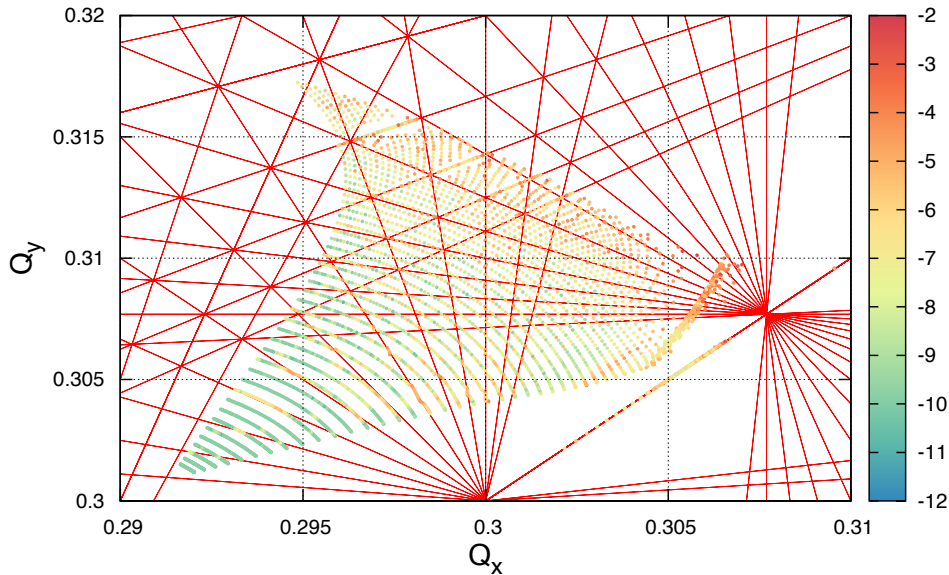


Beam-beam tune shift and spread

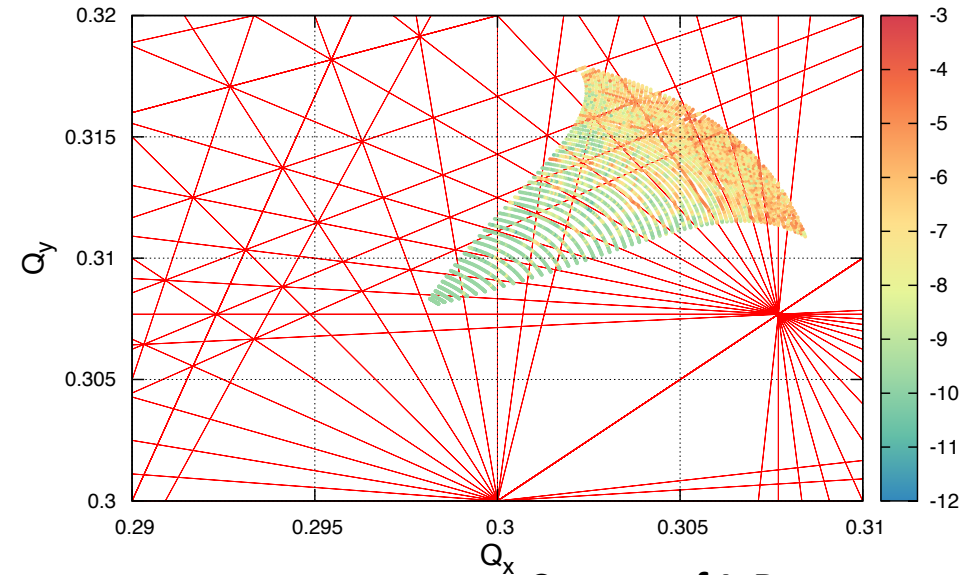
When long-range effects become important footprint wings appear (large amplitude particles deviate)

Aim to reduce the area as much as possible and avoid chaotic motion!

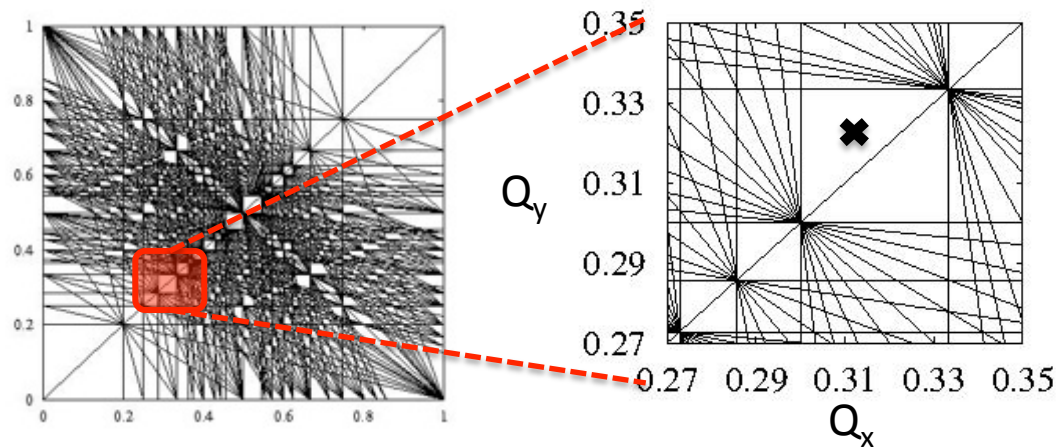
296mrad



390mrad



Courtesy of J. Barranco

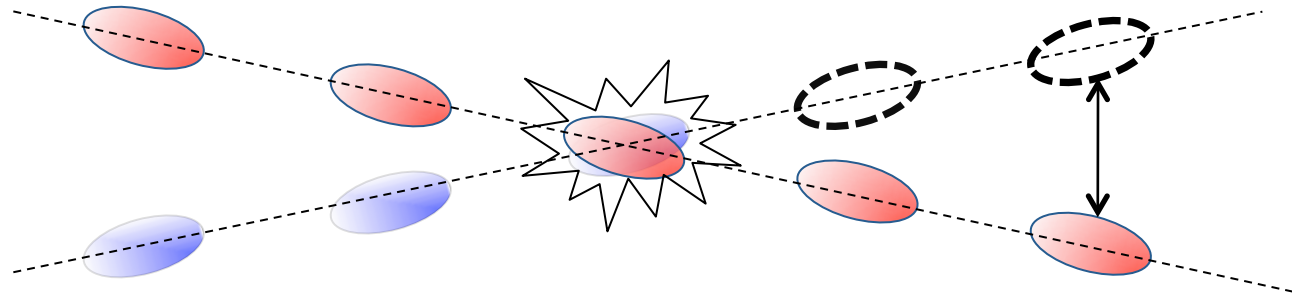


Complications

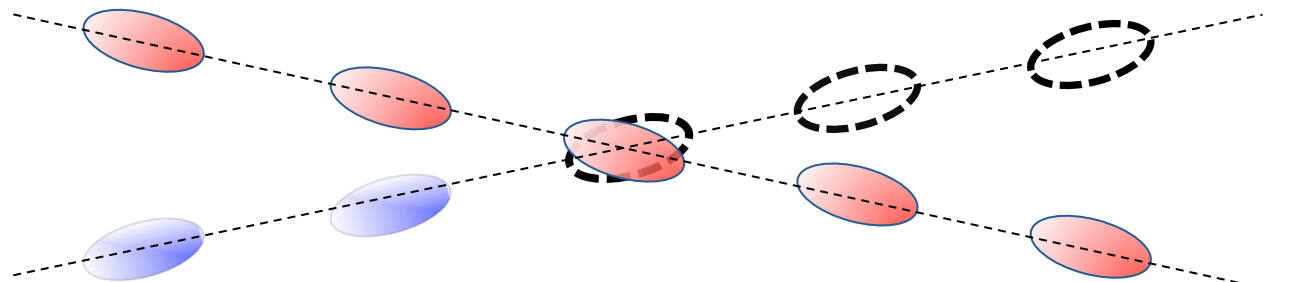
PACMAN and SUPER PACMAN bunches



Pacman:
miss long range BBI



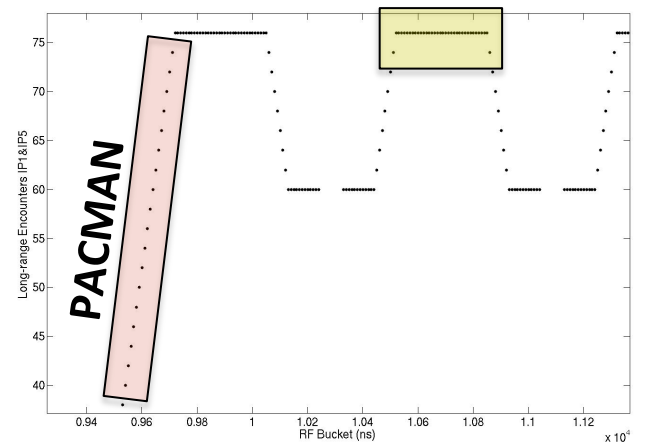
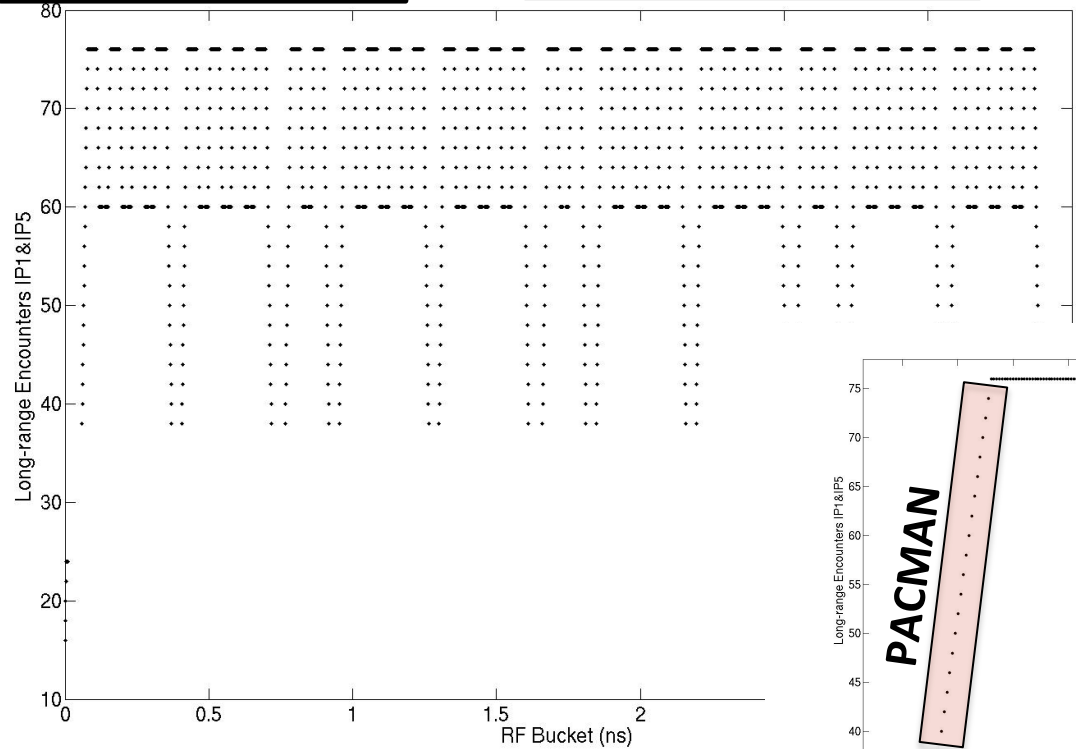
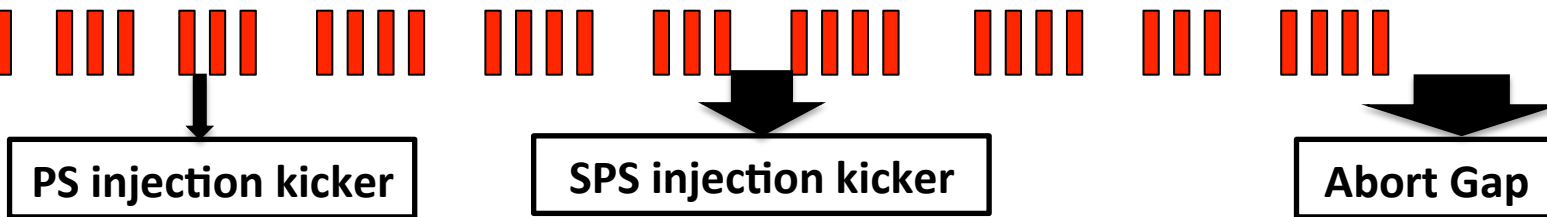
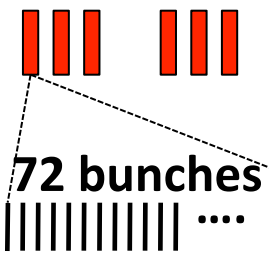
Super Pacman:
miss head-on BBI



IP2 and IP8 depending on filling scheme

Different bunch families: Pacman and Super Pacman

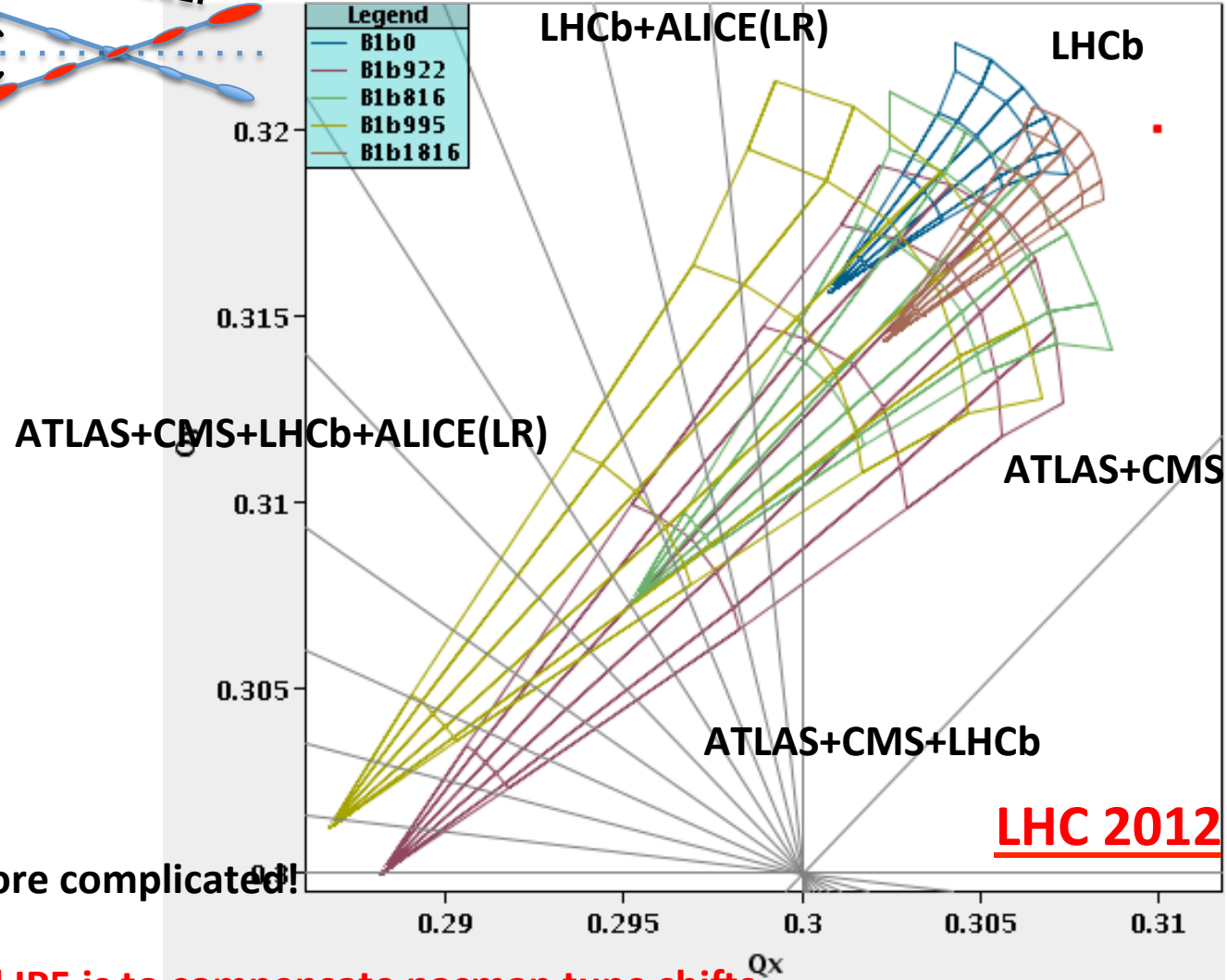
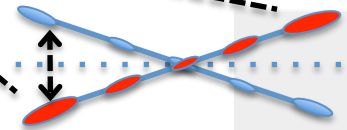
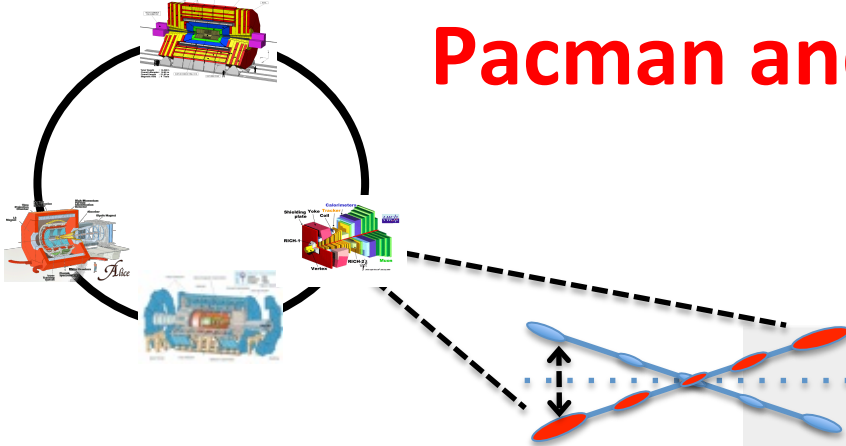
LHC Complications: filling schemes



Nominal Bunches

Pacman Bunches: different number of long-range interactions

Pacman and Super-pacman



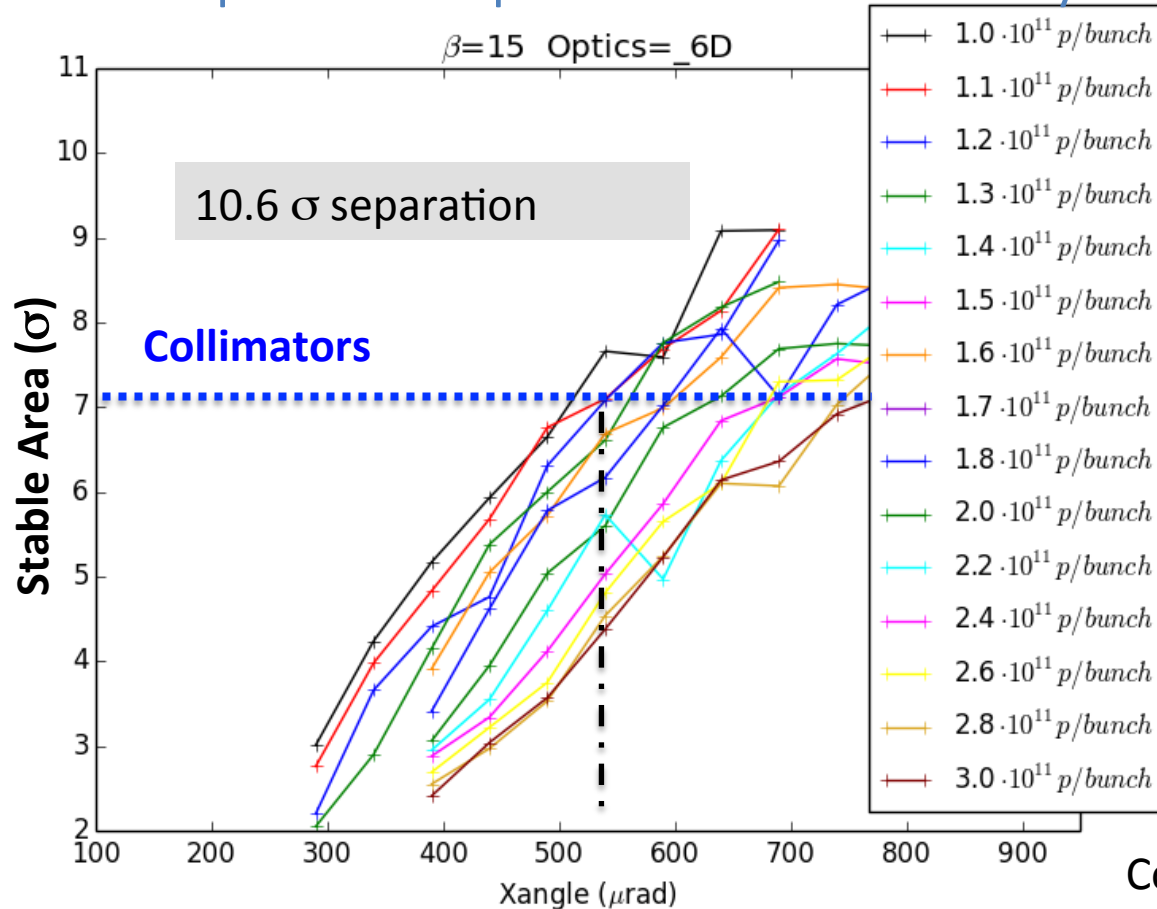
...operationally it is even more complicated!
 ...intensities, emittances...

Alternating crossing IP1 and IP5 is to compensate pacman tune shifts

Particles stability

Dynamic Aperture: area in amplitude space with stable motion

Stable area of particles depends on beam intensity and crossing angle



$$d_{lr} \propto \sqrt{\frac{\beta^* \alpha^2 \gamma}{\epsilon_n}}$$

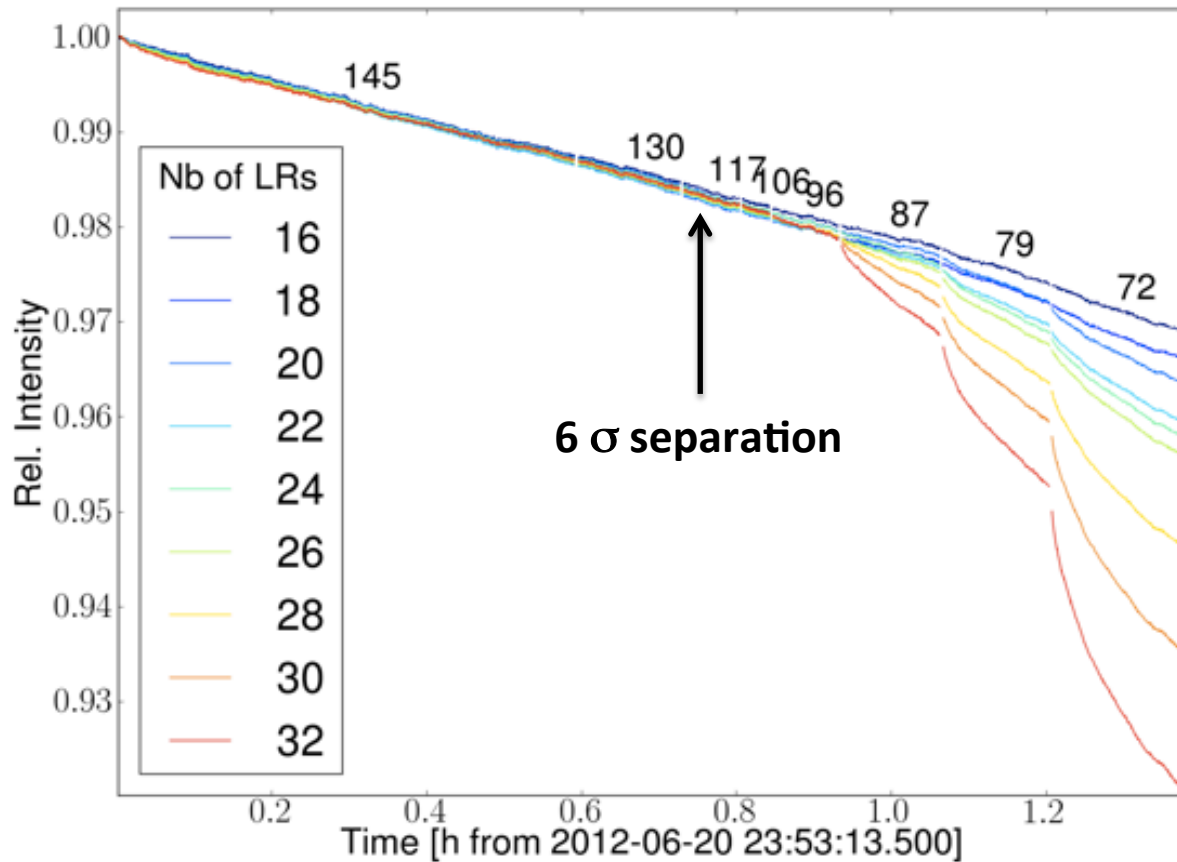
$$F_{bb} \propto N_p$$

Cortesy D. Banfi

Stable area depends on beam-beam interactions therefore the choice of running parameters (crossing angles, β^* , intensity, working points) is the result of careful study of different effects!

DO we see the effects of LR in the LHC?

Courtesy X. Buffat



Small crossing angle
= small separation

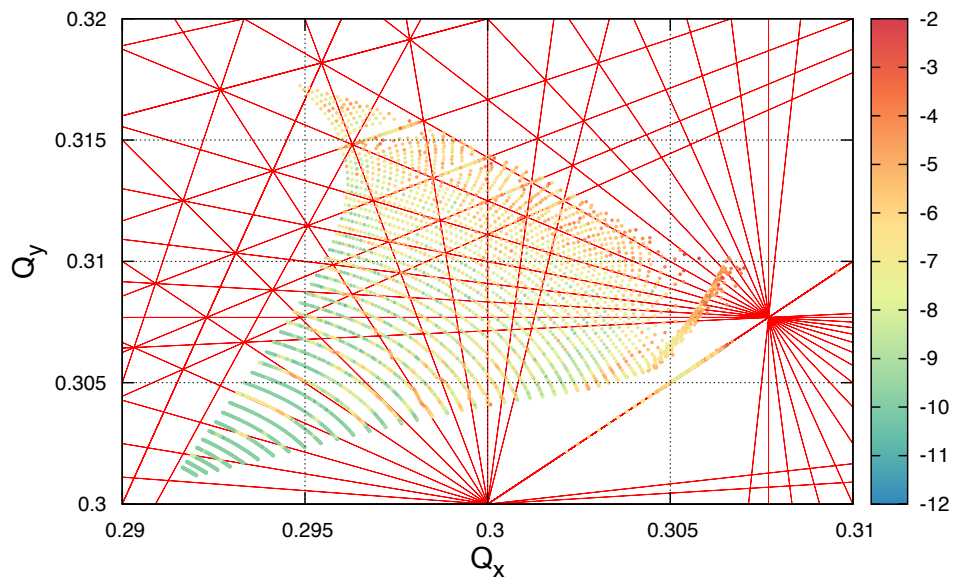
If separation of long range too
small particles become unstable
and get lost

Bunches start losing large amplitude particles
Particle losses follow number of Long range interactions
Nominal LHC will have twice the number of interactions

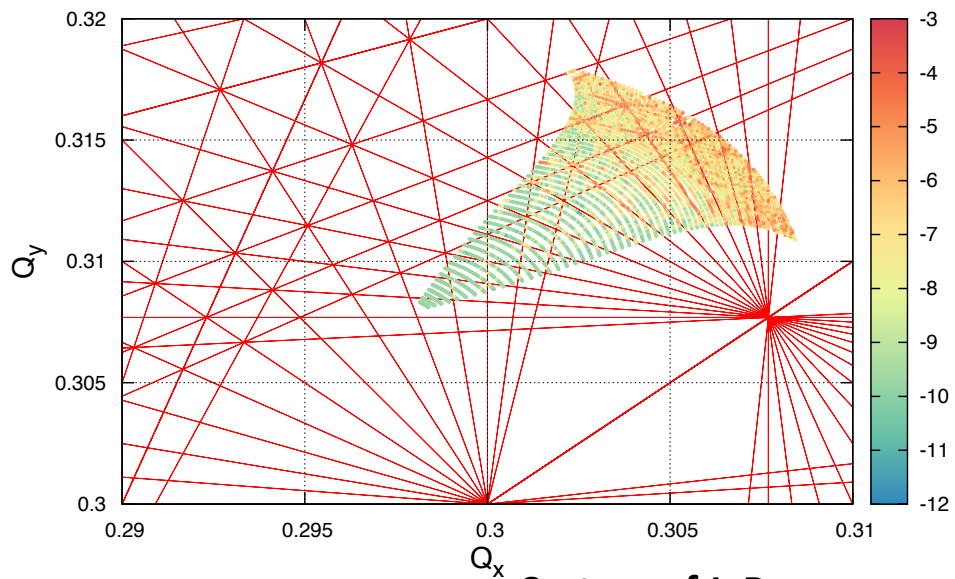
Beam-beam tune shift and spread

When long-range effects become important footprint wings appear
Aim to reduce the area as much as possible and avoid chaotic motion!

296mrad



390mrad



Courtesy of J. Barranco

Orbit Effects

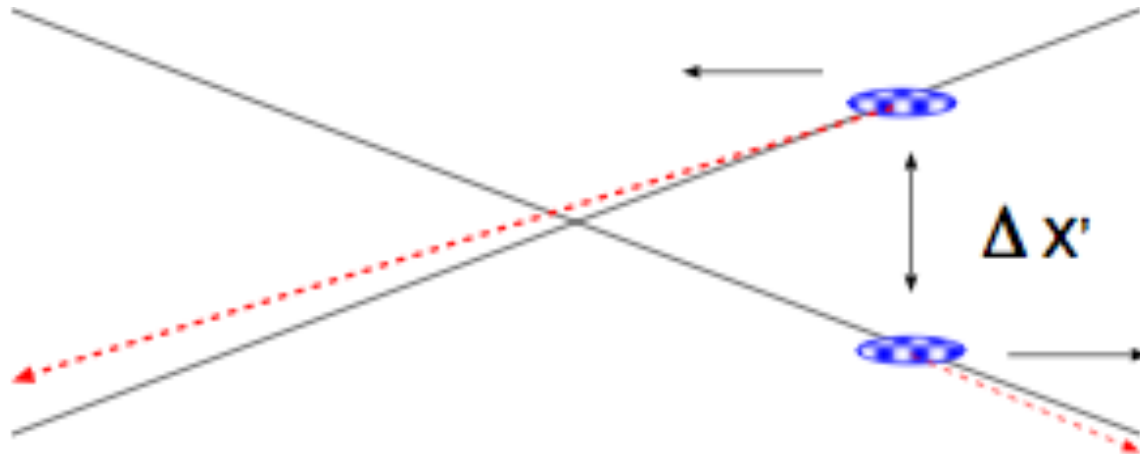
Long Range Beam-beam interactions lead to orbit effects

Long range kick
$$\Delta x'(x + d, y, r) = -\frac{2Nr_0}{\gamma} \frac{(x + d)}{r^2} \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right)\right]$$

For well separated beams $d \gg \sigma$

The force has an amplitude independent contribution: **ORBIT KICK**

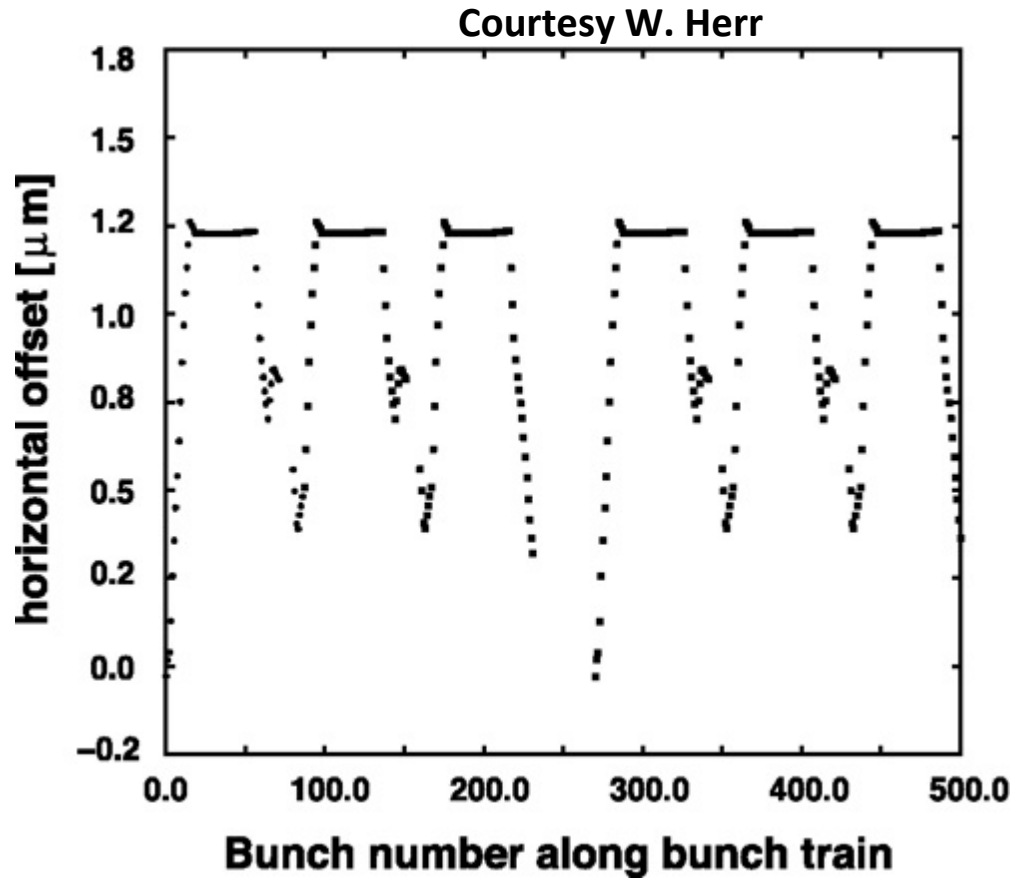
$$\Delta x' = \frac{const}{d} \left[1 - \frac{x}{d} + O\left(\frac{x^2}{d^2}\right) + \dots\right]$$



Orbit can be corrected but we should remember PACMAN effects

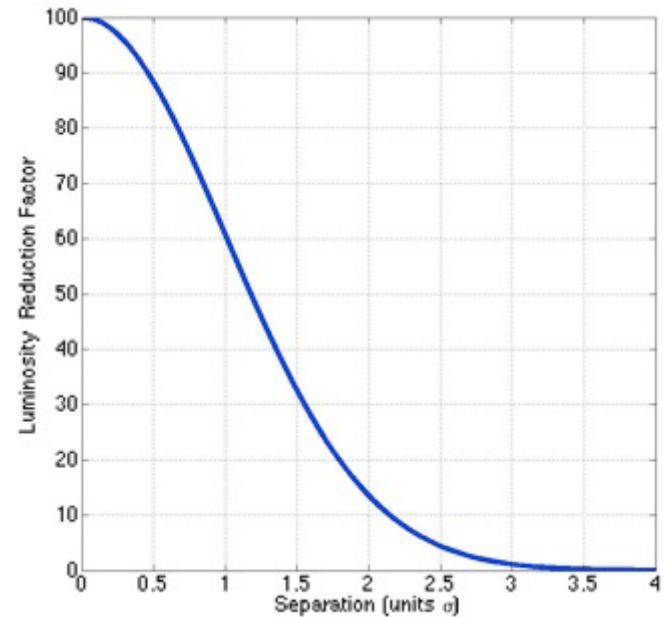
LHC orbit effects

Orbit effects different due to Pacman effects and the many long-range add up giving a non negligible effect



$$L = L_0 \cdot e^{-\frac{d^2}{4\sigma_x^2}}$$

↑
↓
d = 0 - 0.4 units of beam size



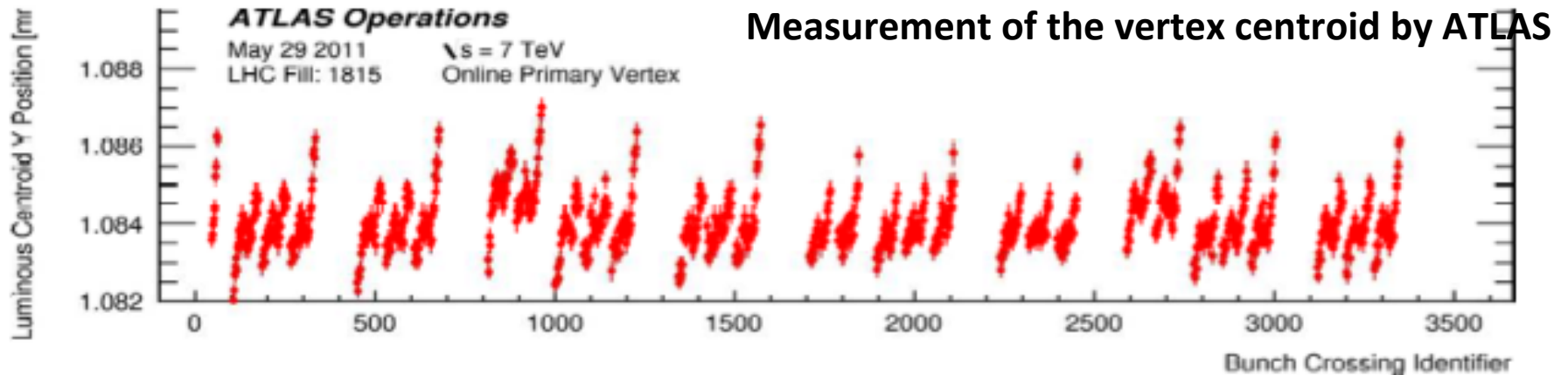
Long range orbit effect

Long range interactions leads to orbit offsets at the experiment a direct consequence is deterioration of the luminosity

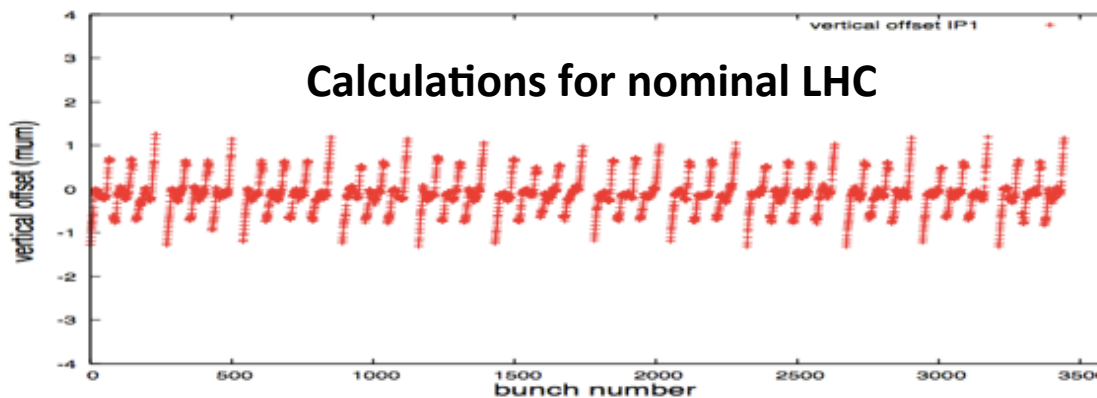
2011-07-05

file:///afs/cern.ch/user/z/zwe/Desktop/PNG/bcid_vs_posY_pm_posYErr.png

#1

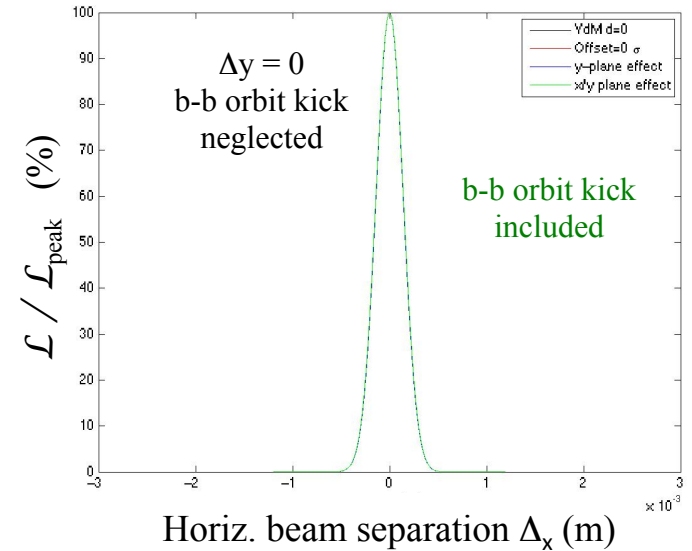
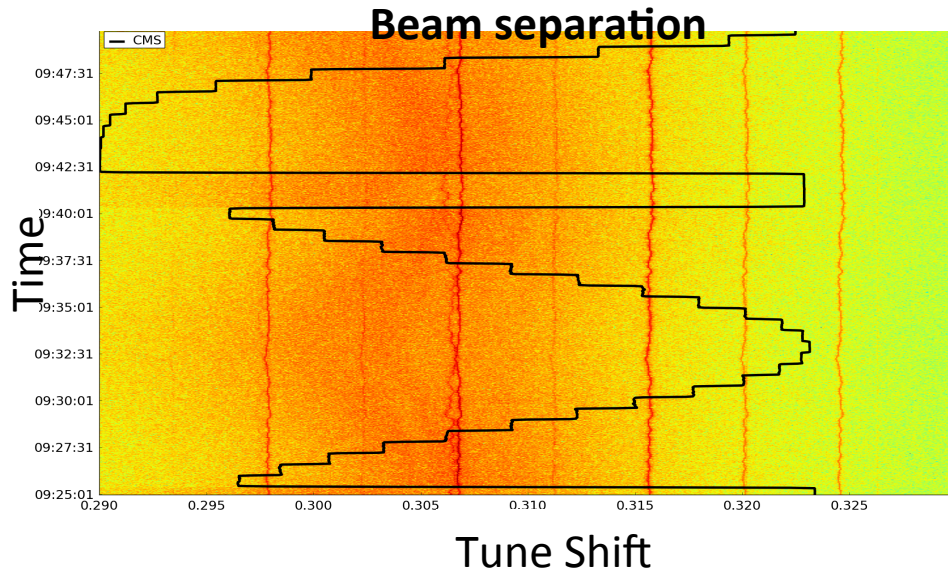


Courtesy W. Kozanecki



Effect is already visible with reduced number of interactions

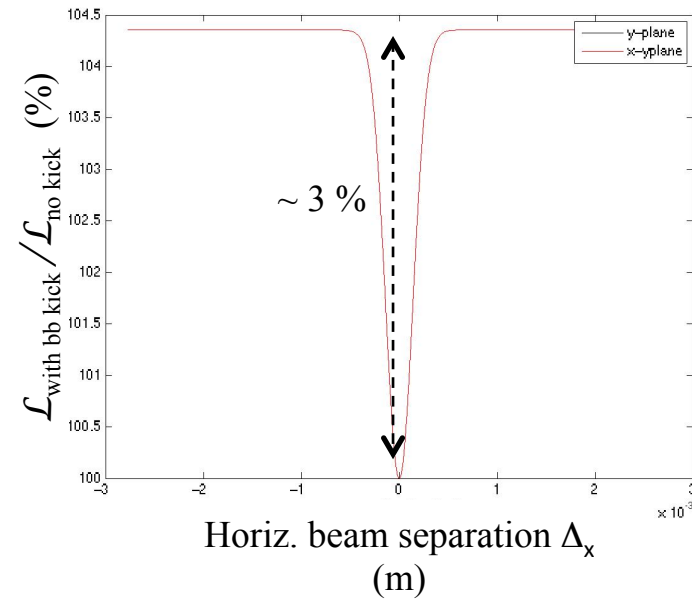
Van der Meer Scans: tune shift and orbit



Van der Meer scan is a scan of the beam-beam force

- Tune shift
- Orbit effect
- Dynamic beta

Beam-beam effects need to be taken into account for precision measurements!



Coherent dipolar beam-beam modes

Coherent beam-beam effects arise from the forces which an exciting bunch exerts on a **whole test bunch** during collision

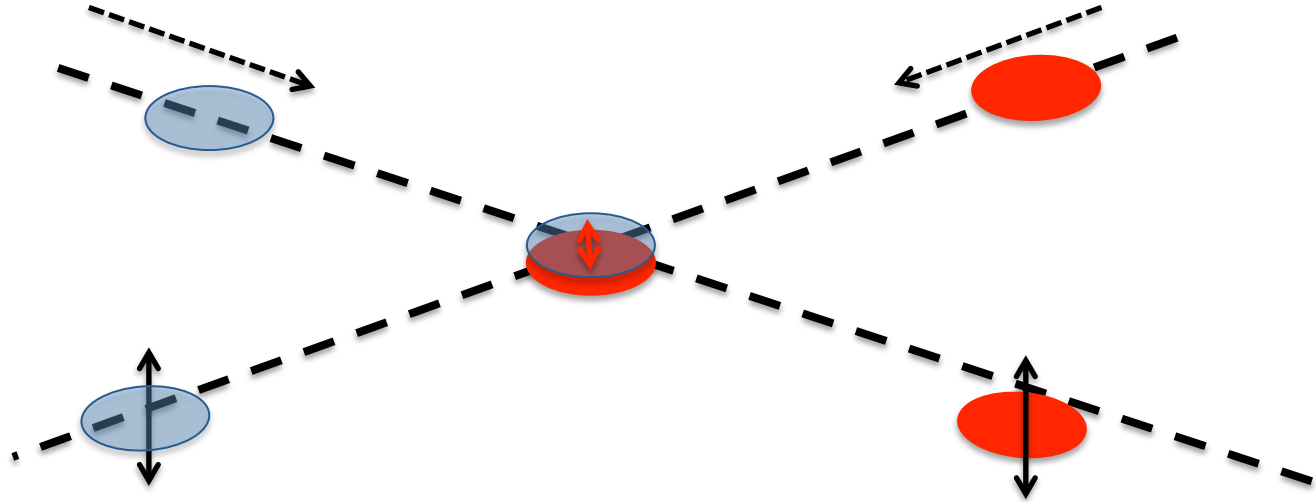
We study the **collective behaviour** of all particles of a bunch

Coherent motion requires an **organized behaviour** of all particles of the bunch

Coherent beam-beam force

- Beam distributions Ψ_1 and Ψ_2 mutually changed by interaction
- Interaction depends on distributions
 - Beam 1 Ψ_1 solution depends on beam 2 Ψ_2
 - Beam 2 Ψ_2 solution depends on beam 1 Ψ_1
- Need a **self-consistent** solution

Coherent beam-beam effects



- Whole bunch sees a kick as an entity (**coherent kick**)
- Coherent **kick seen by full bunch** different from single particle kick
- Requires **integration** of individual kick over particle distribution

$$\Delta r' = -\frac{N_p r_0}{r} \cdot \frac{r}{r^2} \cdot \left[1 - e^{-\frac{r^2}{4\sigma^2}} \right]$$

- Coherent kick of separated beams can excite coherent **dipolar oscillations**
- All bunches couple because each bunch “sees” many opposing bunches(LR): **many coherent modes** possible!

Coherent effects

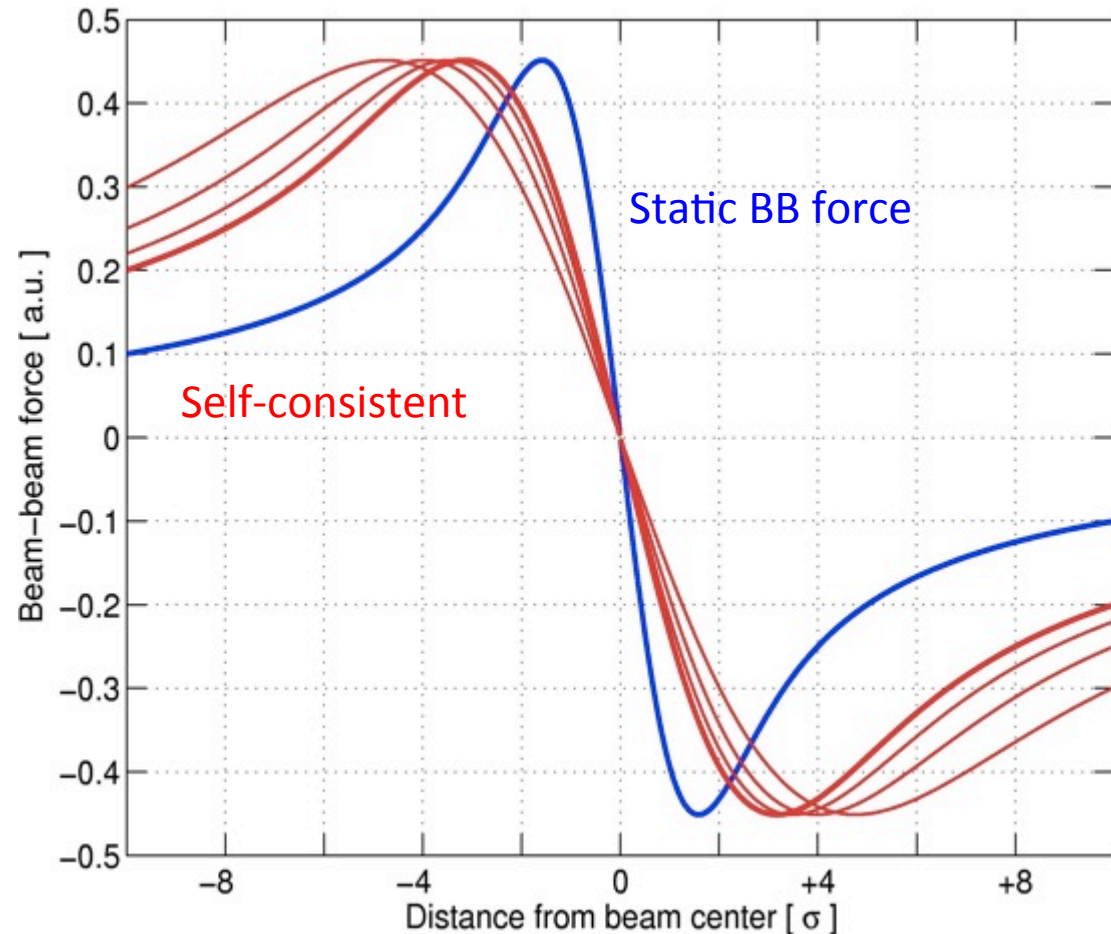
Self-consistent treatment needed

Perturbative methods

static source of distortion:
example magnet

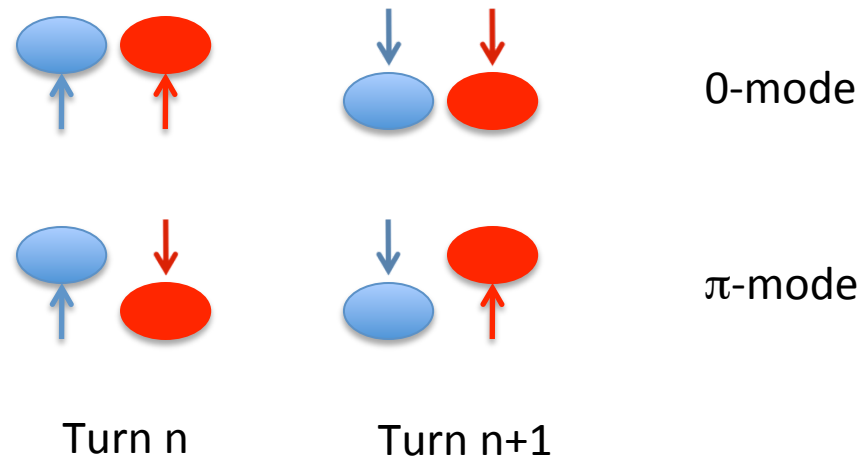
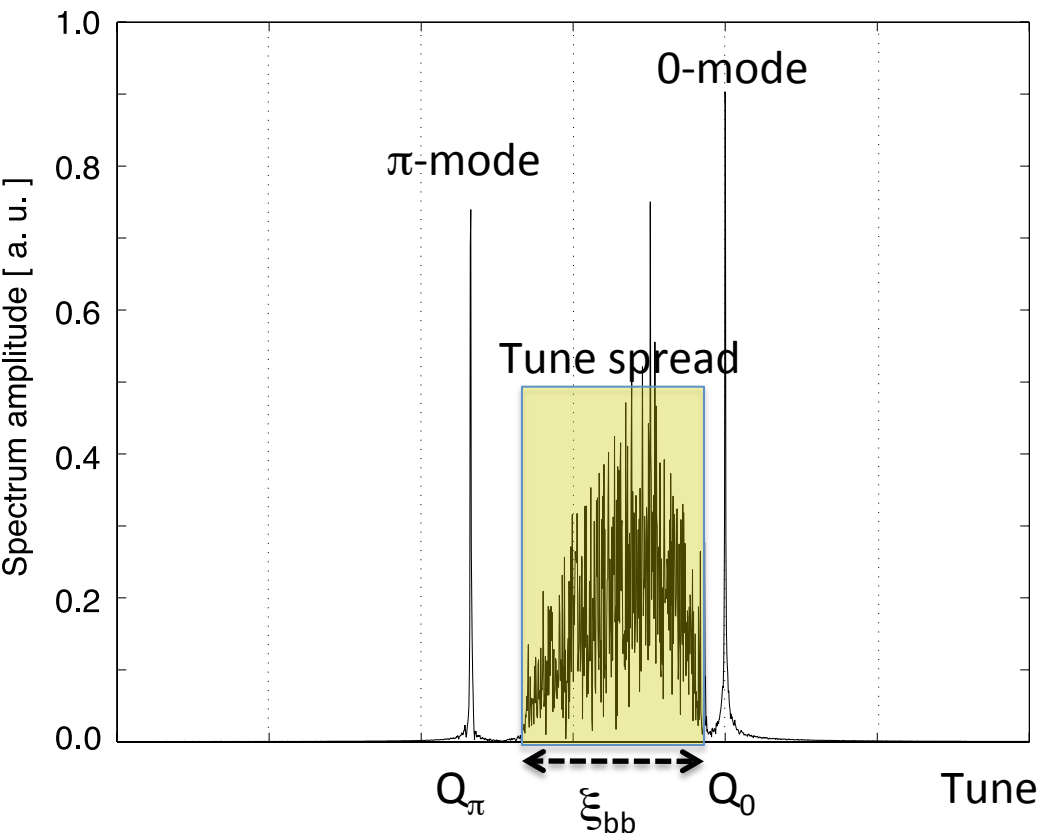
Self-consistent method

source of distortion changes
as a result of the distortion



For a complete understanding of BB effect a self-consistent treatment should be used

Simple case: one bunch per beam



MOVIE

0-mode at unperturbed tune Q_0

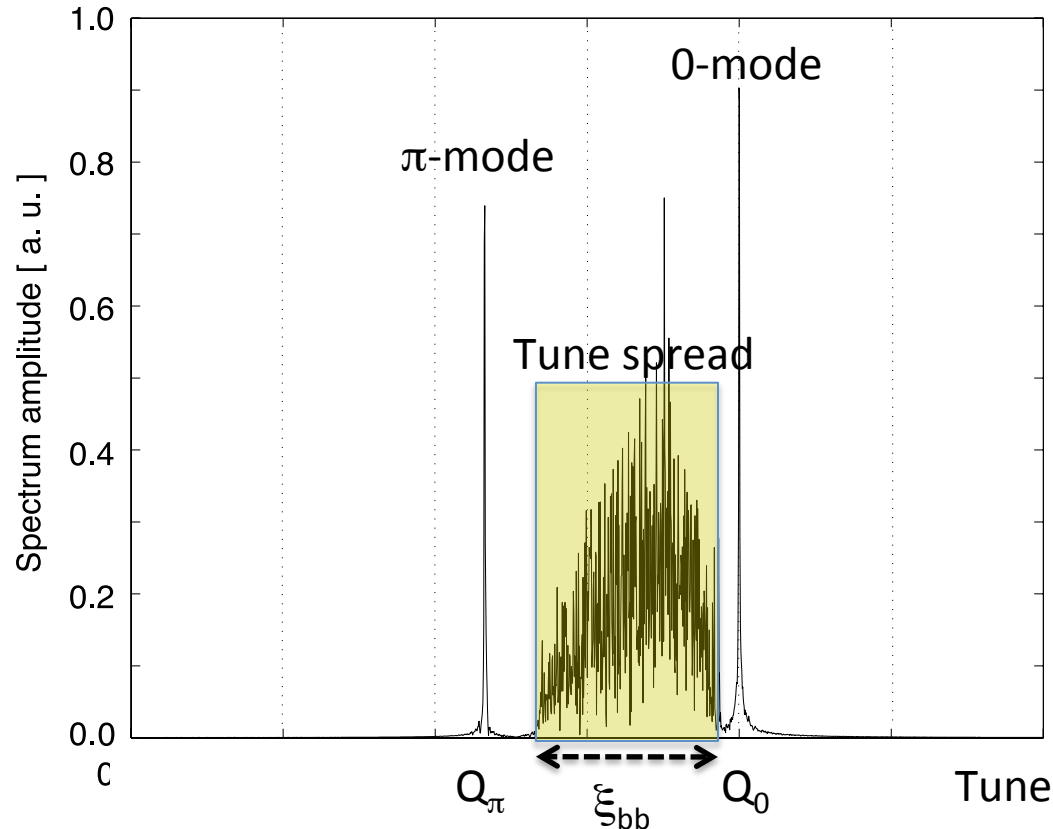
π -mode is shifted at $Q_\pi = 1.1-1.3 \xi_{bb}$

Incoherent tune spread range $[0, -\xi]$

$$\Delta Q = Y \cdot \xi$$

- Coherent mode: two bunches are “locked” in a coherent oscillation
- 0-mode is stable (mode with NO tune shift)
- π -mode can become unstable (mode with largest tune shift)

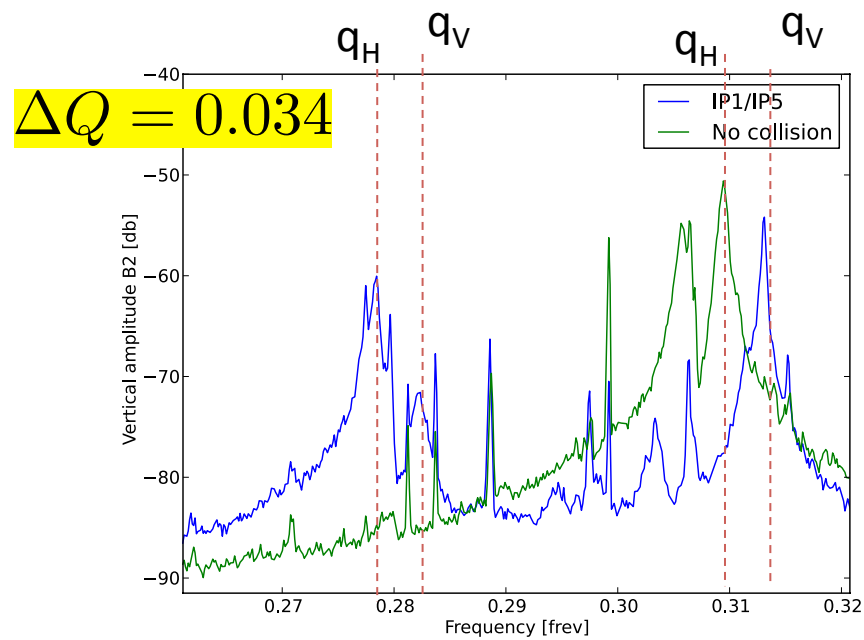
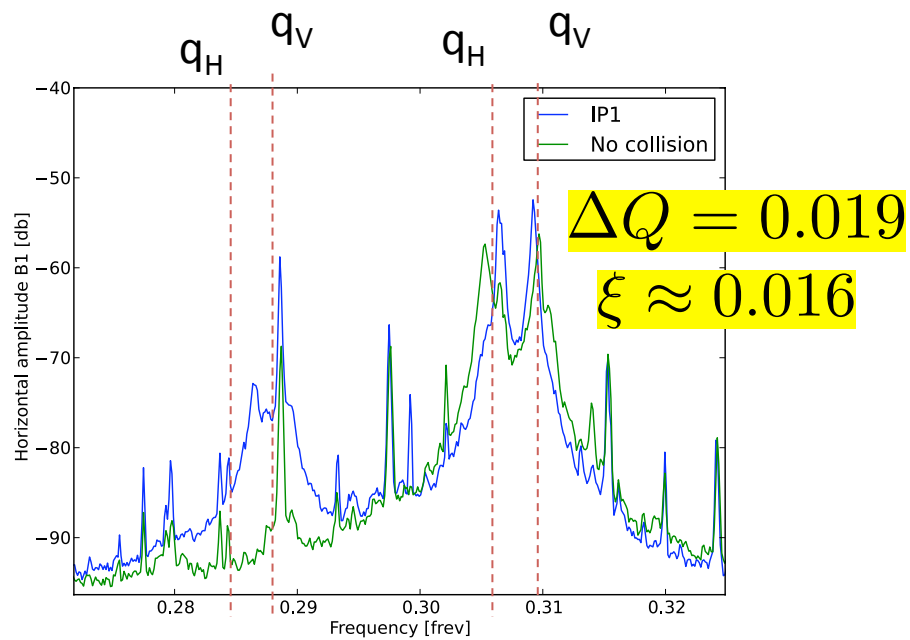
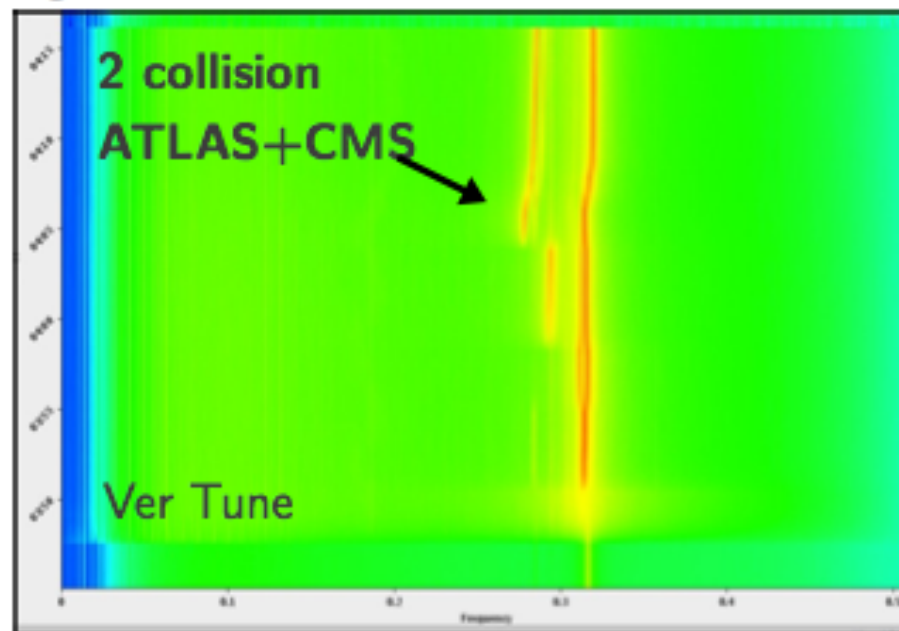
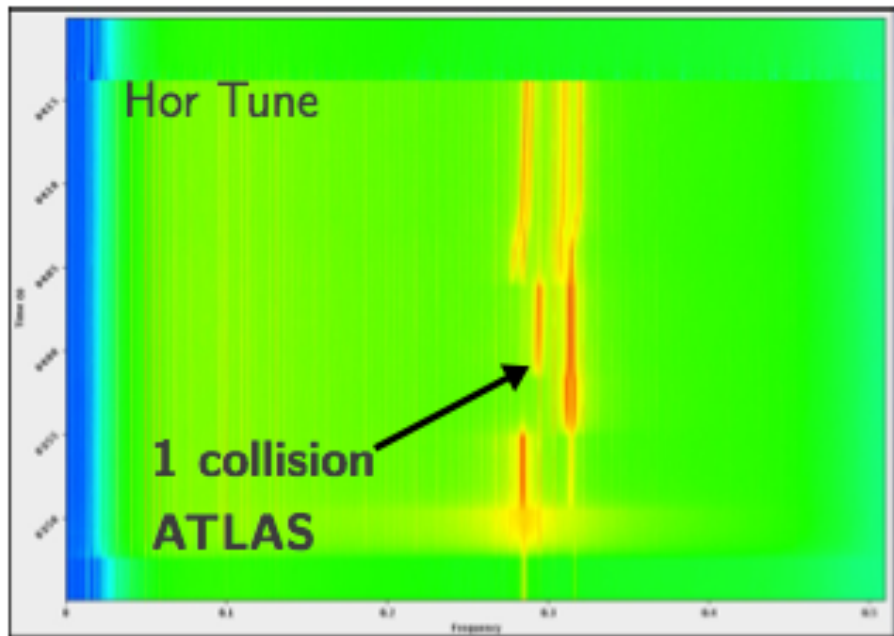
Simple case: one bunch per beam and Landau damping



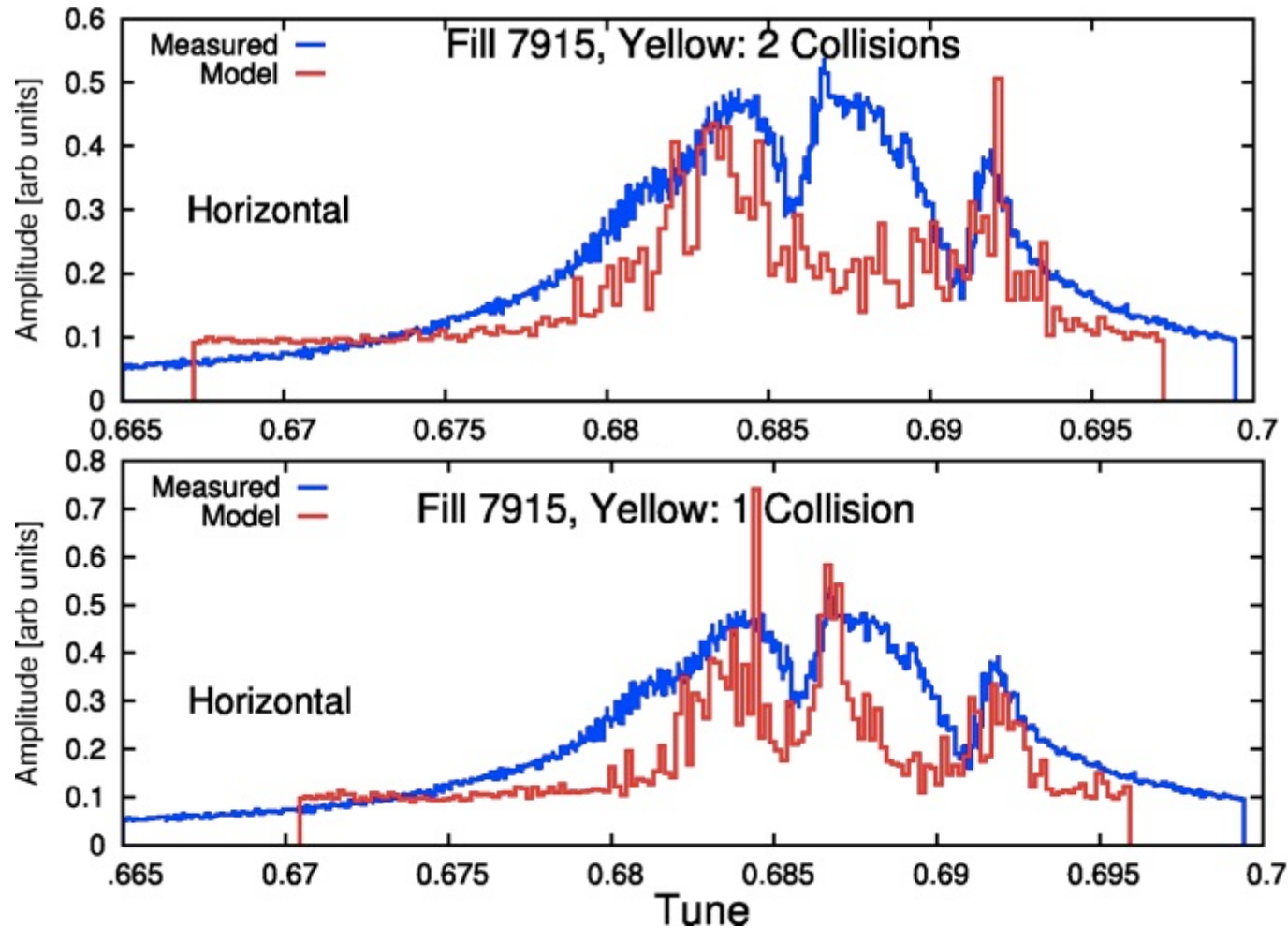
- Incoherent tune spread is the Landau damping region any mode with frequency laying in this range should not develop**
- π -mode has frequency out of tune spread (Y) so it is not damped!

Head-on beam-beam coherent mode: LHC

BBQ Signals



Beam-beam coherent modes and Landau Damping

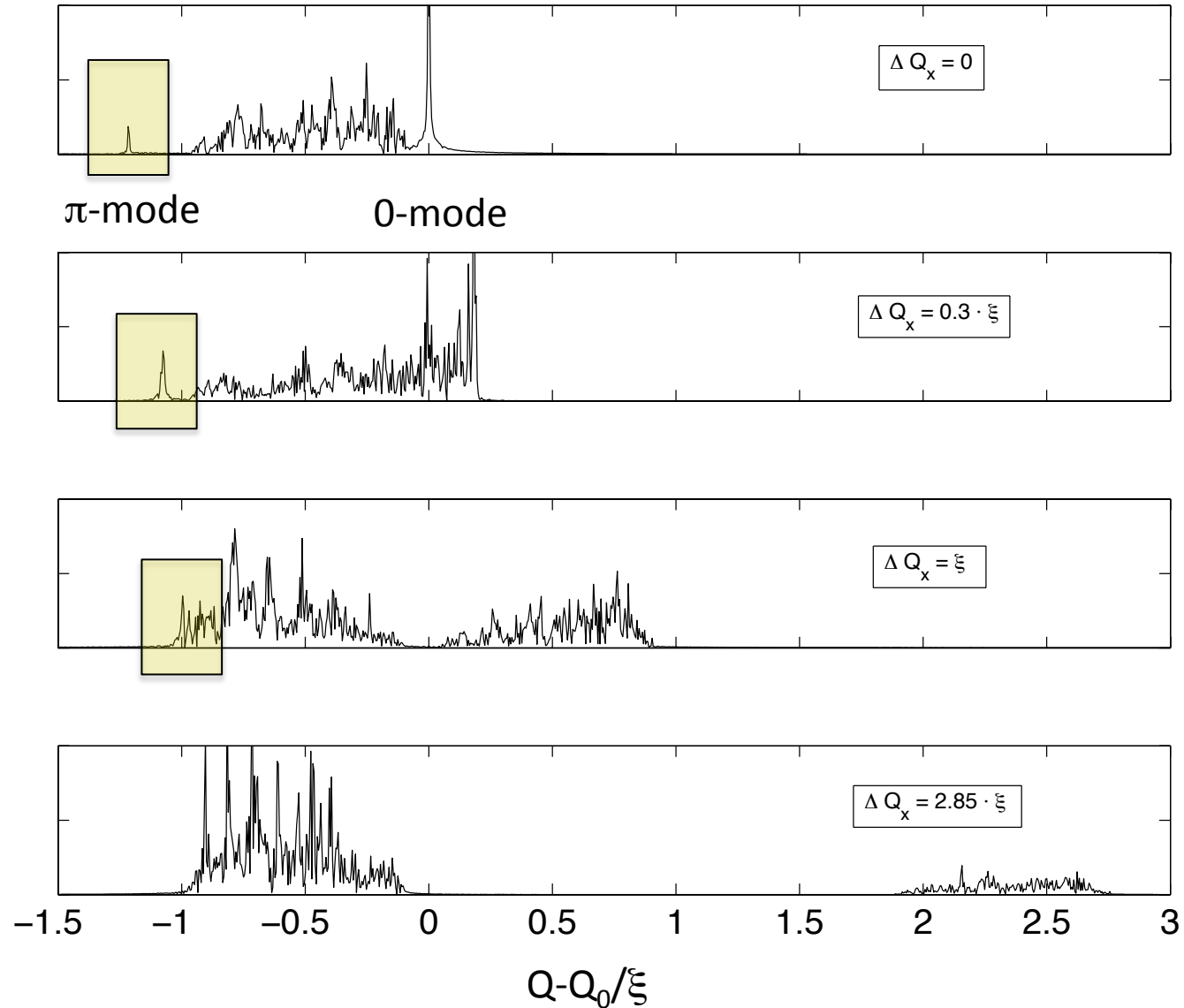


If you have an exciting force (i.e. impedance) at this frequency: fast instabilities could occur

Pacman effect on coherent modes

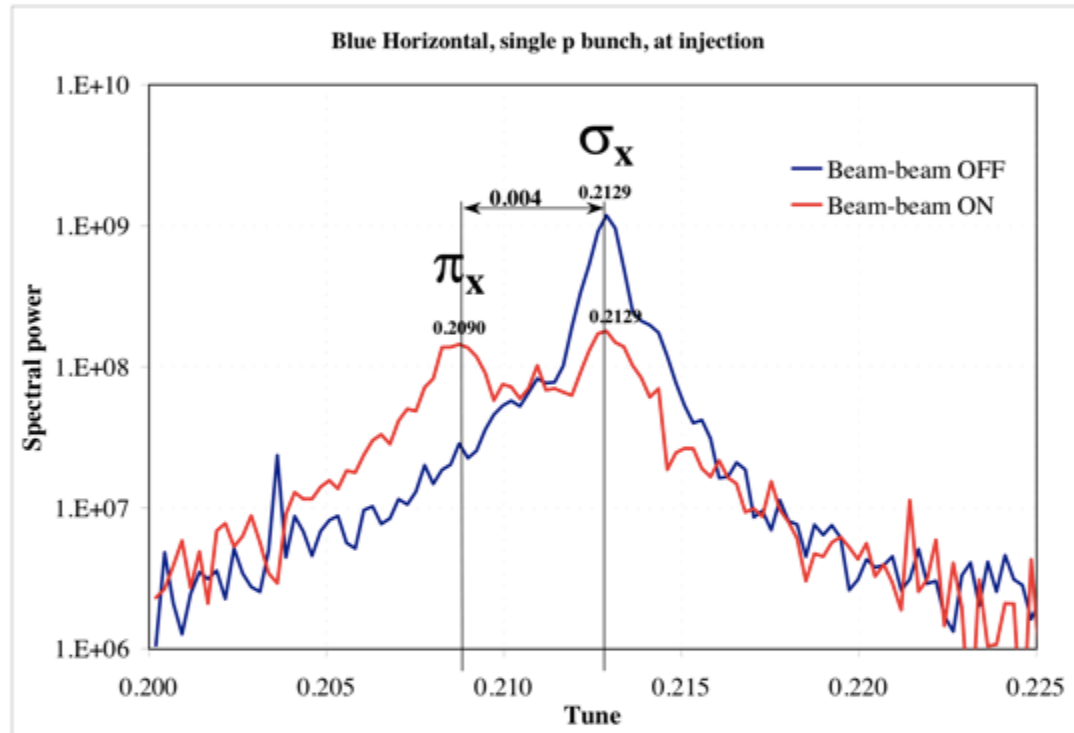
Single bunch diagnostic so important

Coherent modes and tune split



Tune split breaks symmetry and coherent mode cannot develop
LHC has used a TUNE-SPLIT in 2010 physics run!

Coherent modes at RHIC

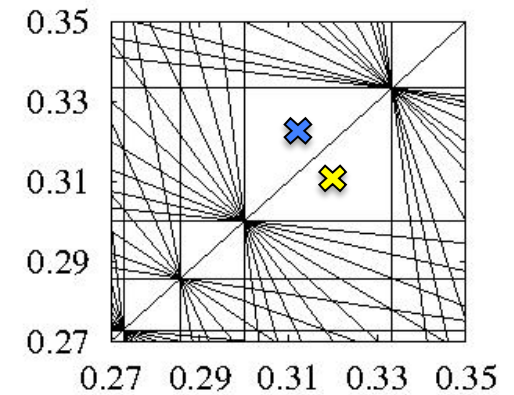
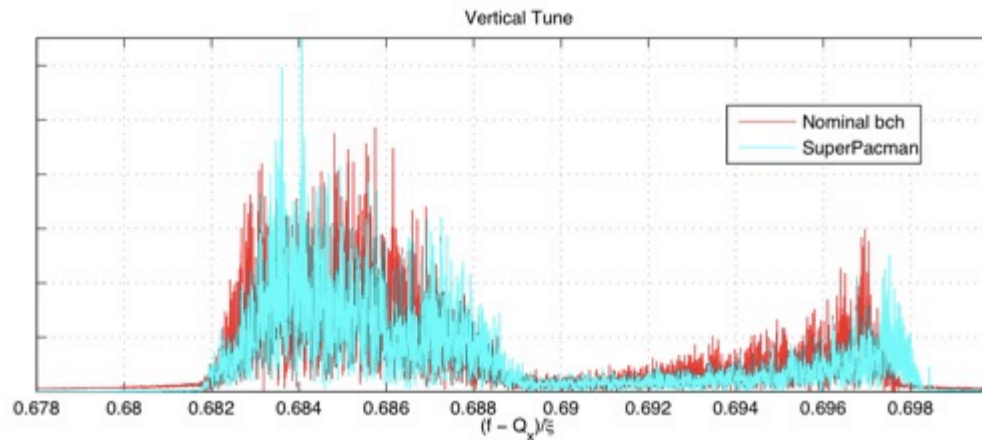
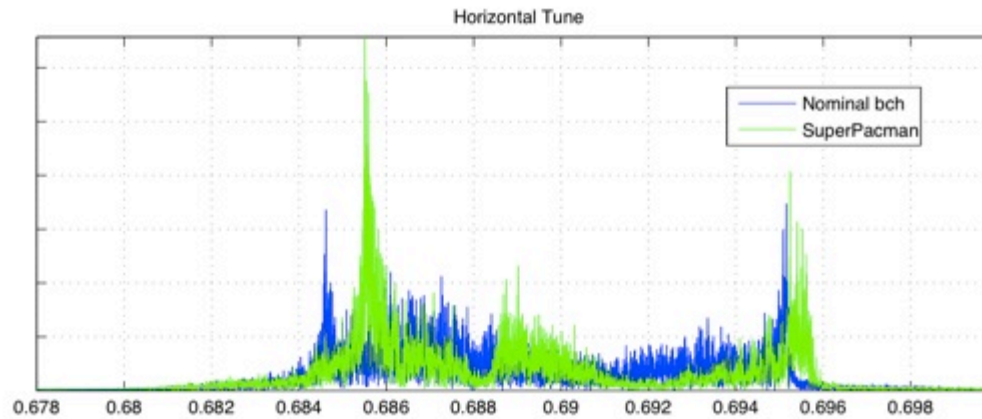


Courtesy W. Fischer (BNL)

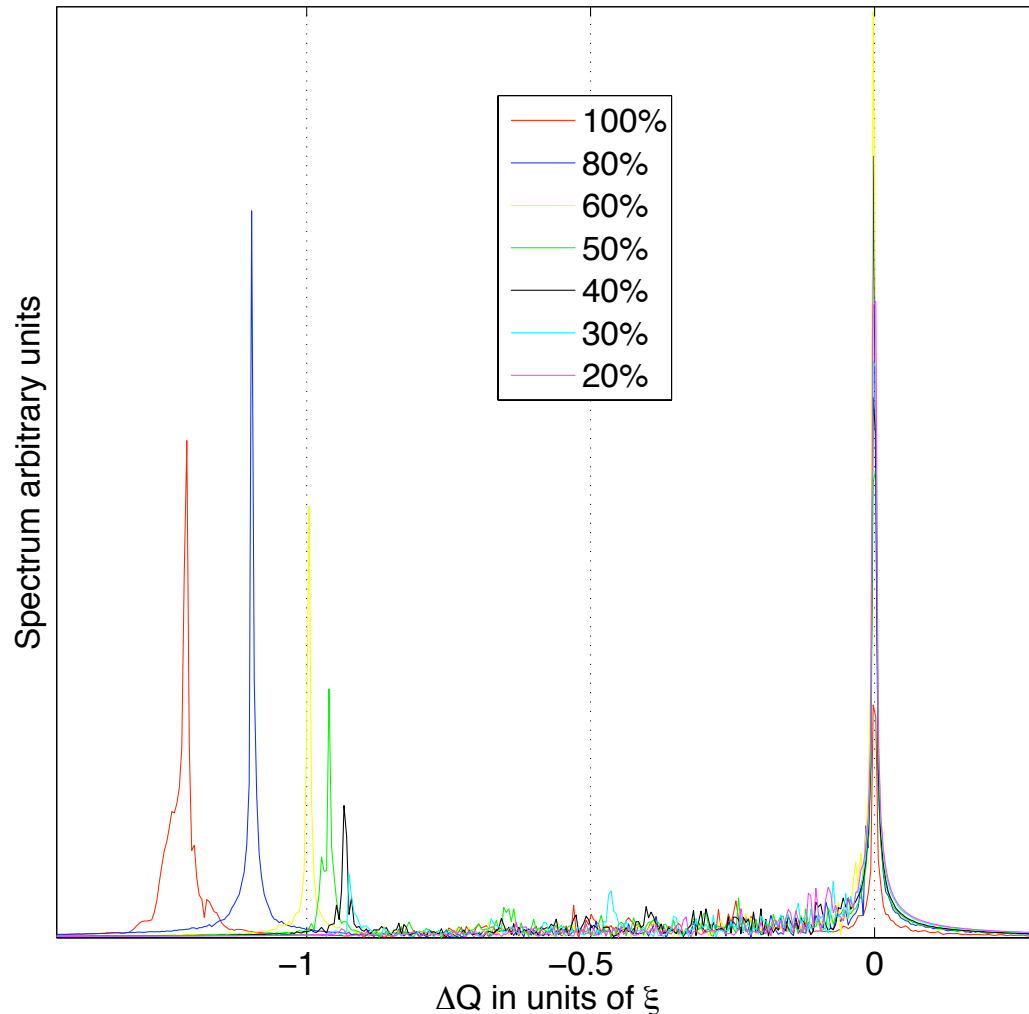
Tune spectra before collision and in collision two modes visible

Different tunes or intensities

RHIC running with “mirrored” tune for years to break coherent oscillations



Different bunch intensities



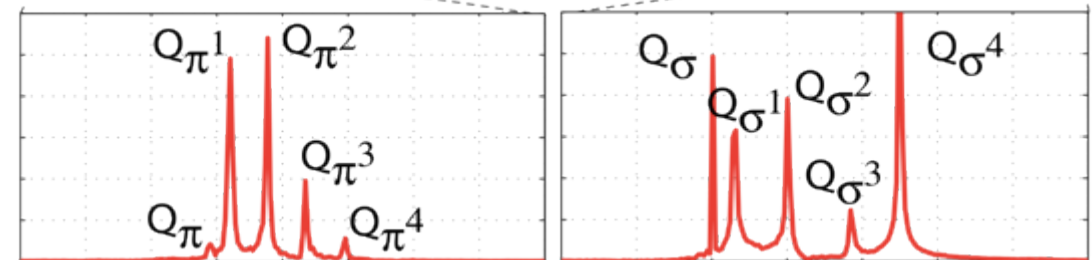
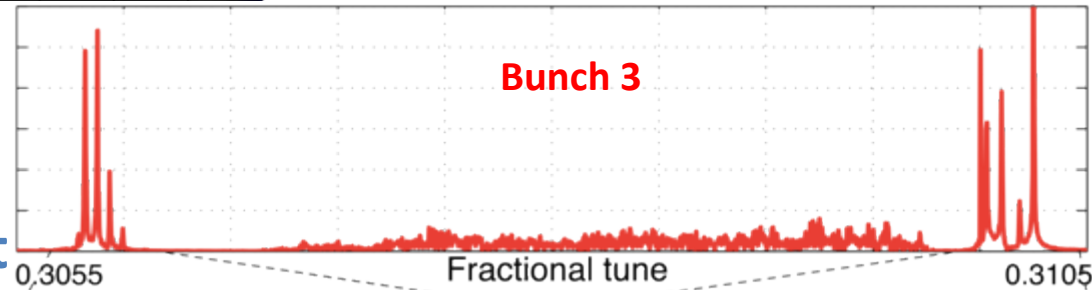
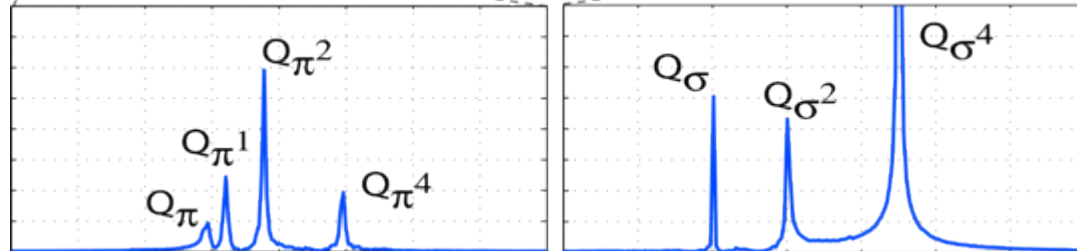
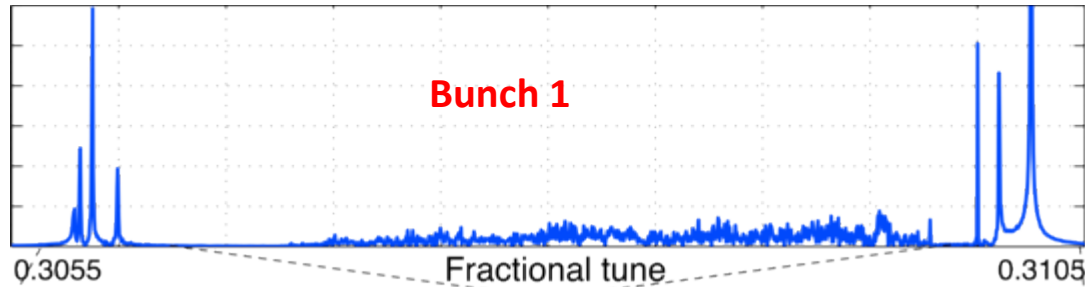
For two bunches colliding head-on in one IP the coherent mode disappears if intensity ratio between bunches is 55%

We assumed:

- equal emittances
- equal tunes
- NO PACMAN effects
(bunches will have different tunes)

For coherent modes the key is to break the symmetry in your coupled system...(tunes, intensities, collision patters...)

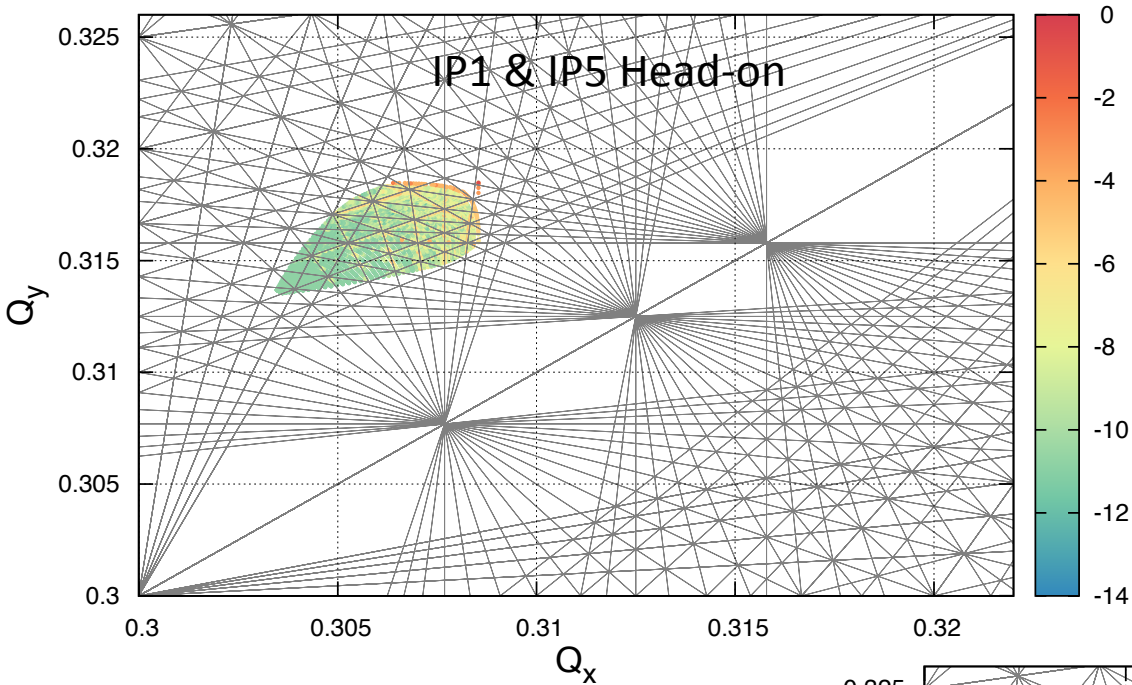
And Long range interactions?



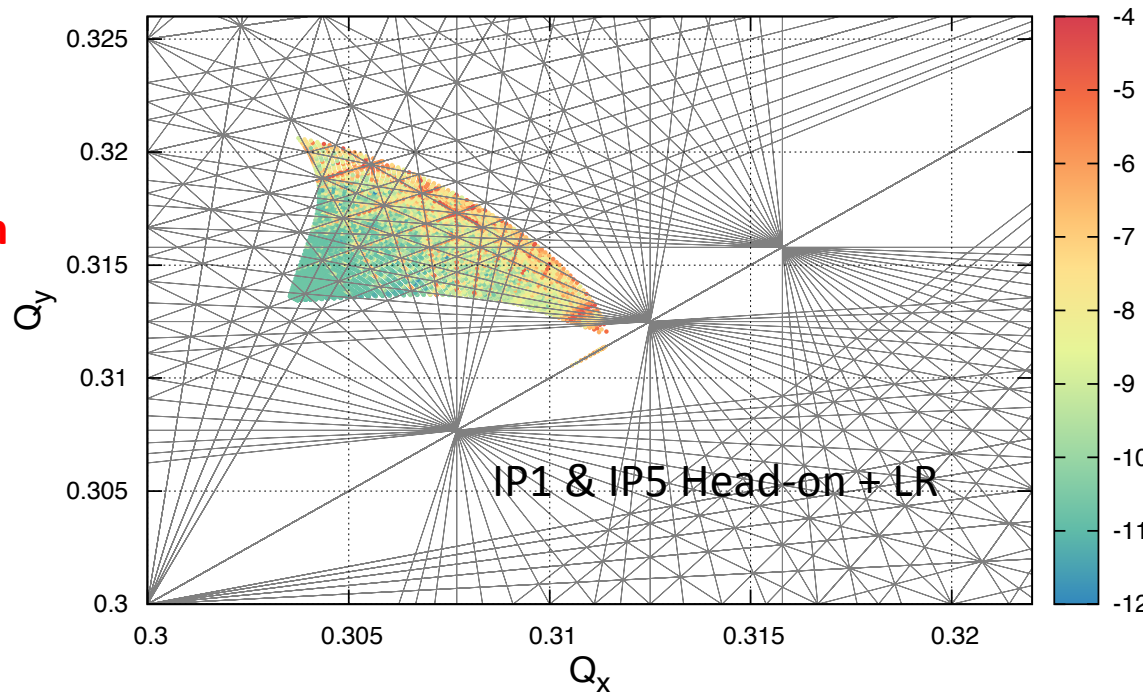
- Each bunch will have different number of modes and tune spectra
- No Landau damping of long-range coherent modes

Single bunch diagnostic can make the difference

Beam-Beam compensation ideas



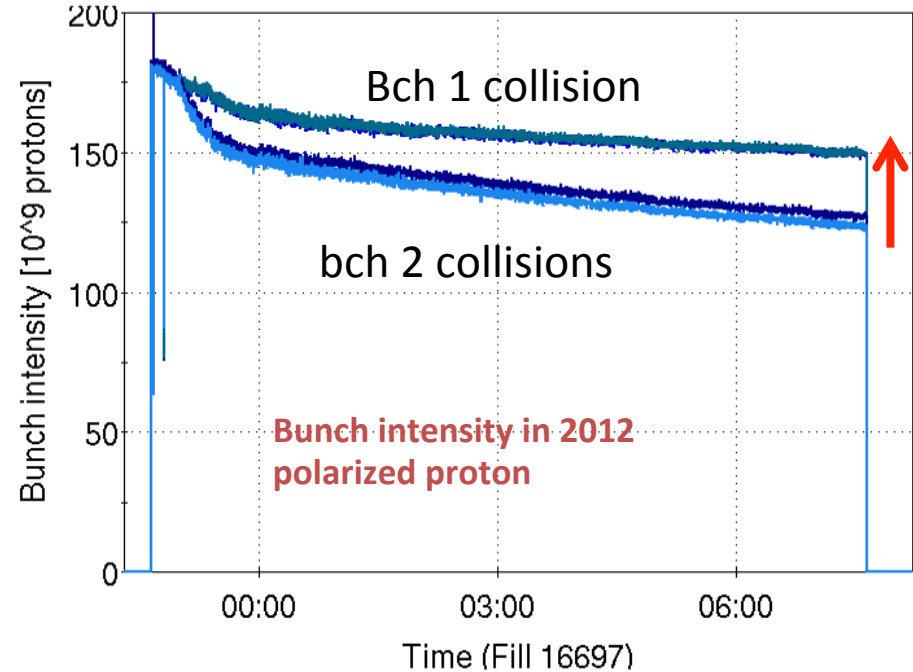
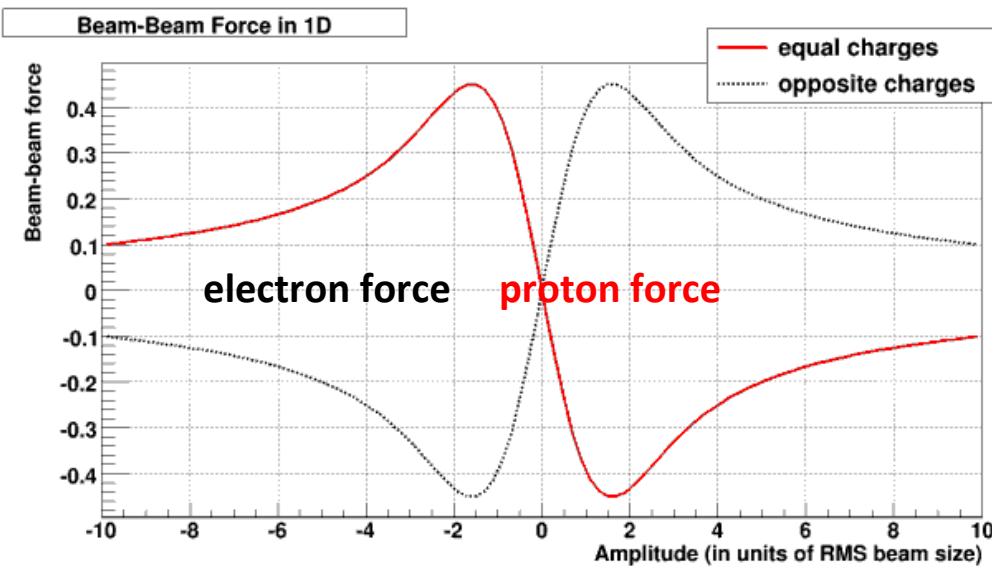
Head-on interaction shift down in tune diagram!
Long range interactions open wings!



Beam-beam compensations:

Head-on

- Linear e-lens, suppress head-on shift
- Non-linear e-lens, suppress tune spread

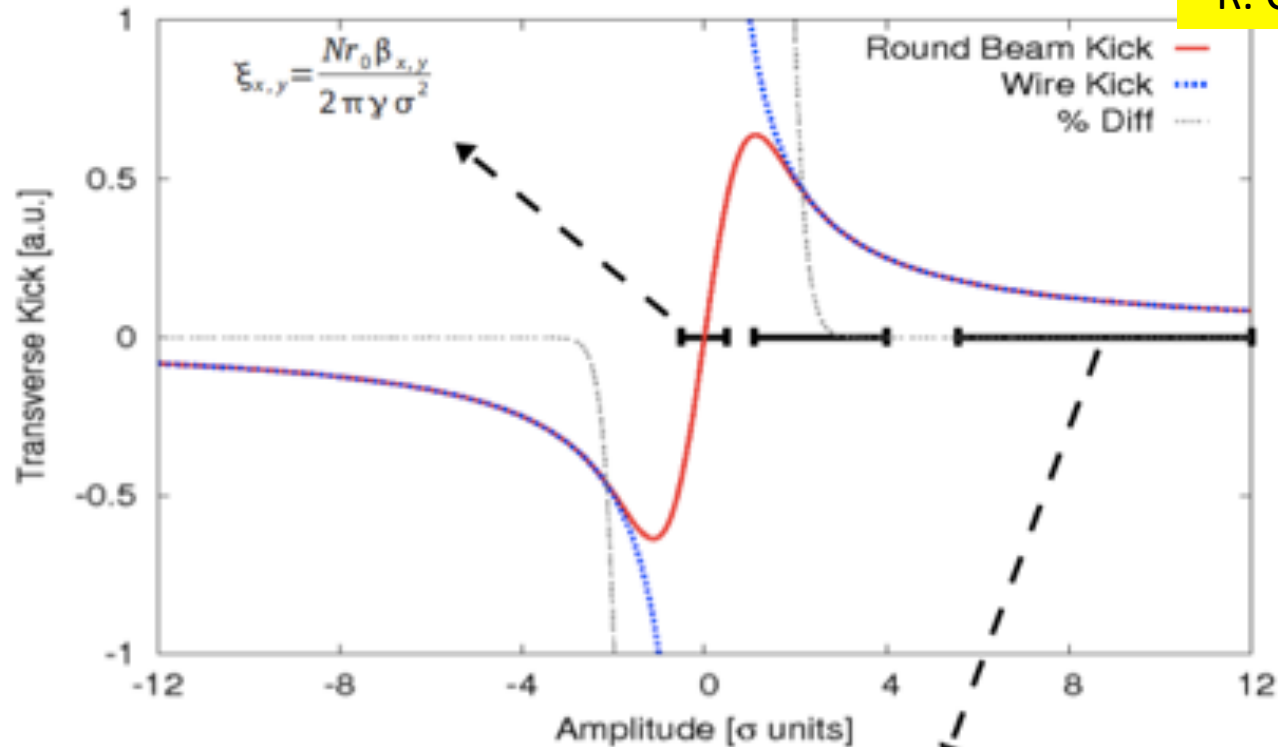


- Past experience: at Tevatron linear and non-linear e-lenses, also hollow...
- Present: test for half compensation at RHIC with non-linear e-lens

Beam-beam compensations: long-range

Beam-beam wire compensation

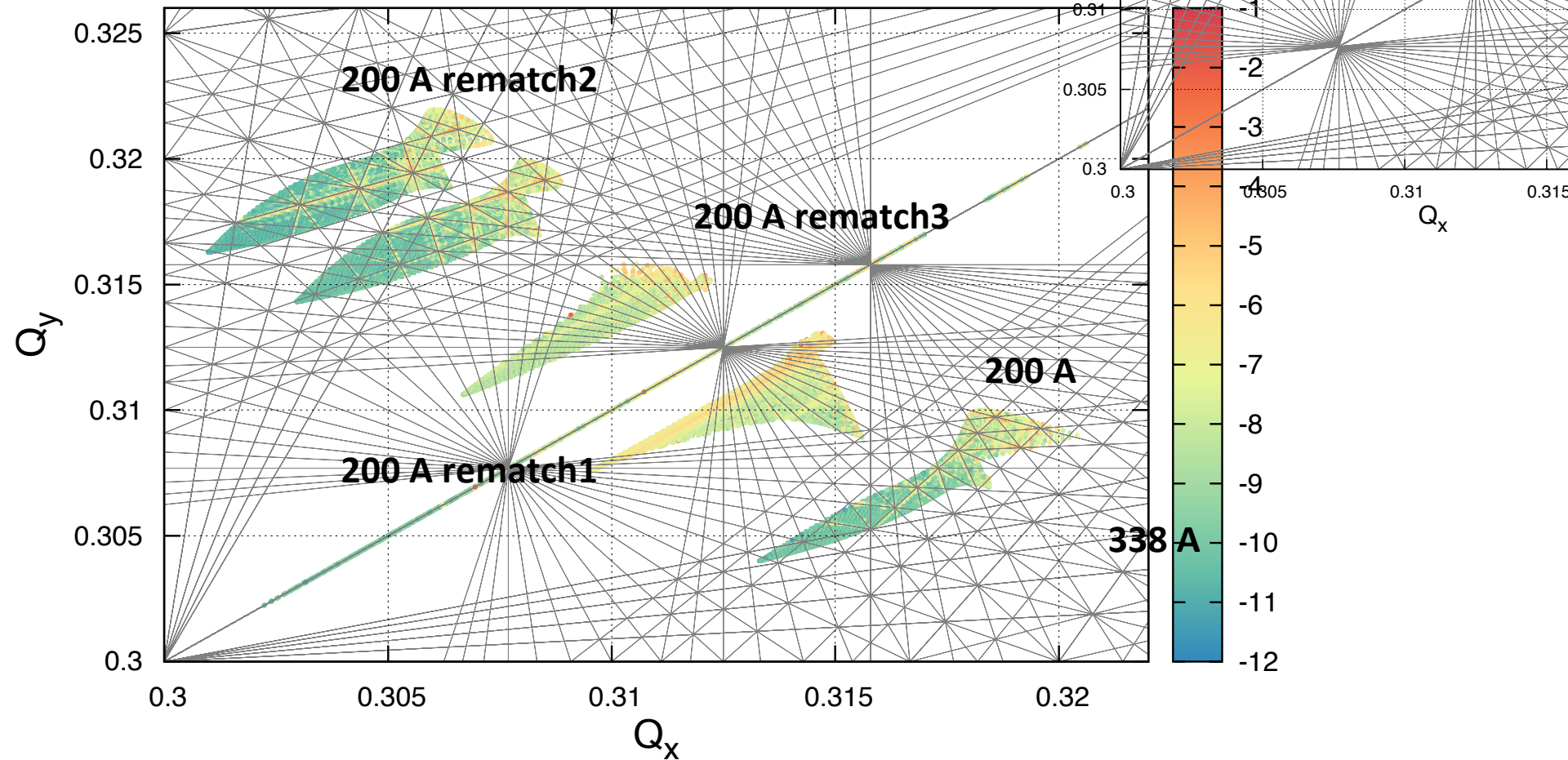
R. Calaga



$$\sigma \ll d: \quad \Delta x'(x, d) = -\frac{K}{d} \cdot \left(1 + \frac{x}{d} + \frac{x^2}{d^2} + \dots\right)$$

- Past experience: at RHIC several tests till 2009...
- Present: simulation studies on-going for possible use in HL-LHC...

BB Long range wire compensation



Many studies from F. Zimmerman and first proposal J. P. Koutchouk

Studies on-going for LHC Upgrade scenarios where crossing angle reduction is fundamental!

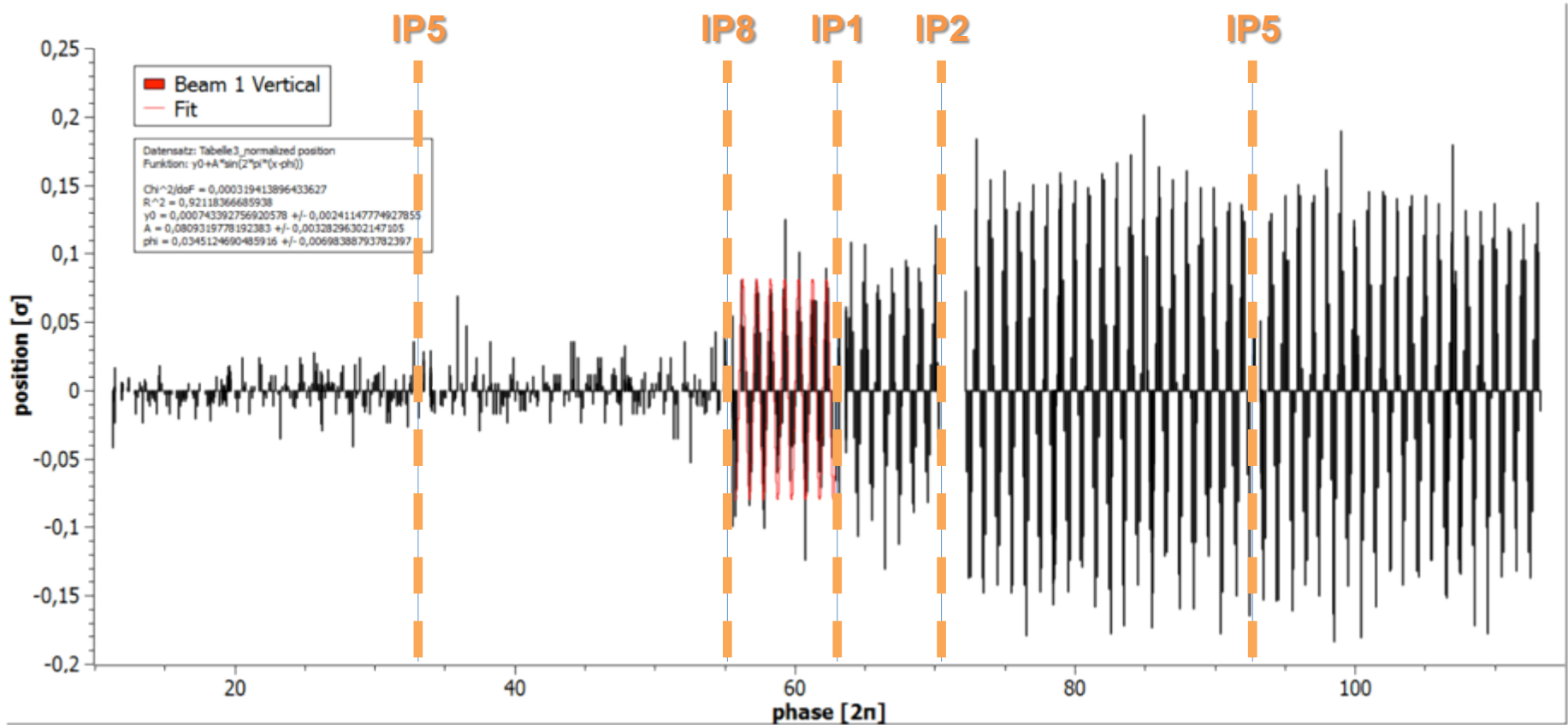
Crossing angle impact LHC nominal 15% HL-LHC 70-80%!

...not covered here...

- *Linear colliders special issues*
- *Asymmetric beams effects*
- *Coasting beams*
- *Beamstrahlung*
- *Synchrotron coupling*
- *Beam-beam experiments*
- *Beam-beam and other collective effects (impedance)*
- ...

Thank You!

Long range orbit effect observations:



Courtesy T. Baer

Vertical oscillation starts when one beam is ejected and dumped

References:

- [1] http://cern.ch/Werner.Herr/CAS2009/proceedings/bb_proc.pdf
- [2] V. Shiltsev et al, "Beam beam effects in the Tevatron", *Phys. Rev. ST Accel. Beams* 8, 101001 (2005)
- [3] Lyn Evans "The beam-beam interaction", CERN 84-15 (1984)
- [4] Alex Chao "Lie Algebra Techniques for Nonlinear Dynamics" SLAC-PUB-9574 (2002)
- [5] J. D. Jackson, "Classical Electrodynamics", John Wiley & Sons, NY, 1962.
- [6] H. Grote, F. Schmidt, L. H. A. Leunissen, "LHC Dynamic Aperture at Collision", LHC-Project-Note 197, (1999).
- [7] W. Herr, "Features and implications of different LHC crossing schemes", LHC-Project-Note 628, (2003).
- [8] A. Hofmann, "Beam-beam modes for two beams with unequal tunes", CERN-SL-99-039 (AP) (1999) p. 56.
- [9] Y. Alexahin, "On the Landau damping and decoherence of transverse dipole oscillations in colliding beams ", *Part. Acc.* 59, 43 (1996).

...much more on the LHC Beam-beam webpage:

<http://lhc-beam-beam.web.cern.ch/lhc-beam-beam/>