

On the onset of the ridge structures and the universal geometrical scaling of the elliptic flow

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Non orthodox view of ridge and elliptic flow

Scaling law (geometrical scaling) for the elliptic flow

Onset of the ridge structures in pp, pA and AA as interactions among strings, above a critical string density (string percolation)

$$\frac{1}{N_A} \frac{dN_{ch}}{dp_T^2} = \frac{1}{Q_0^2} F(\tau), \quad \tau = \frac{p_T^2}{(Q_s^A)^2}$$

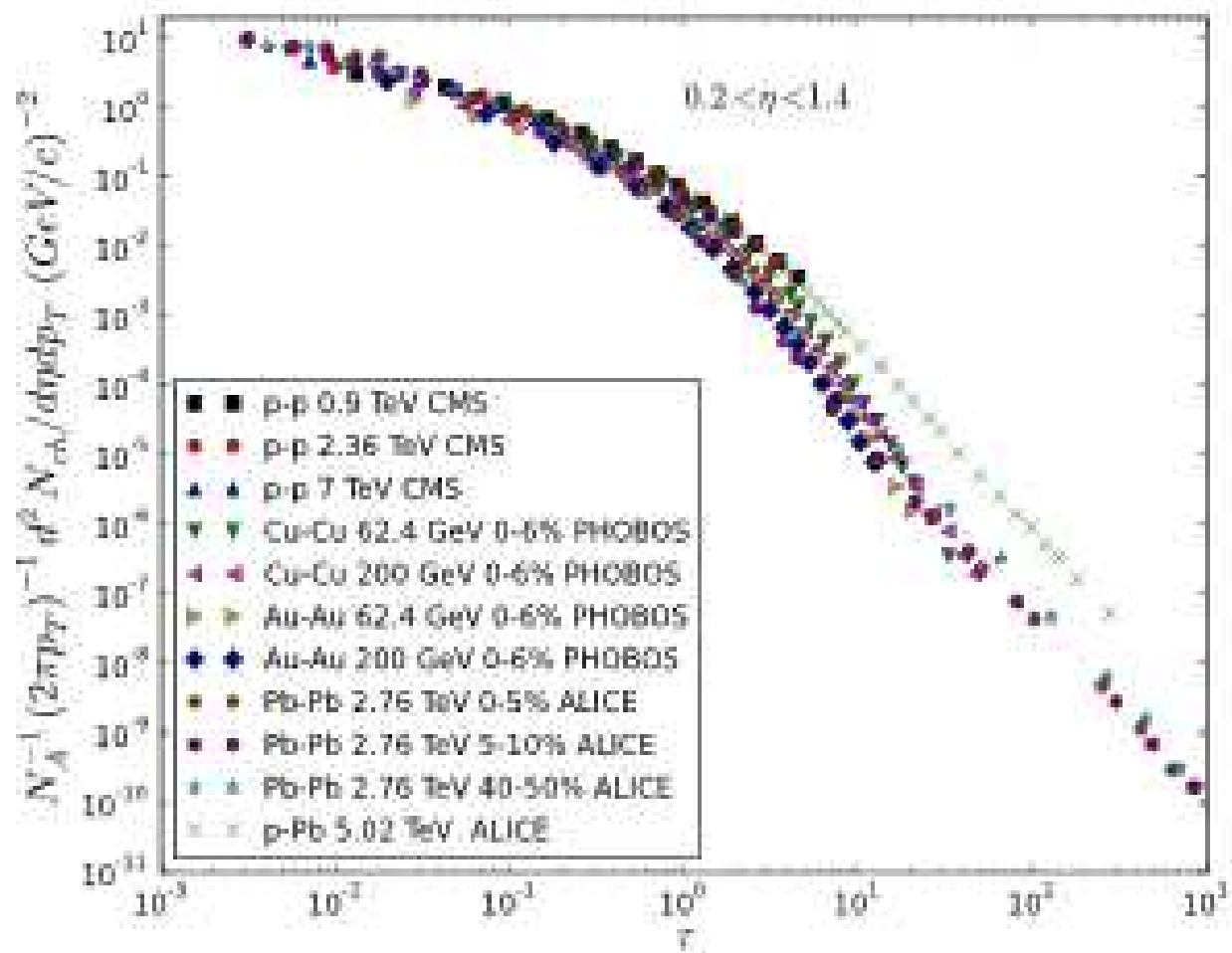
$$(Q_s^A)^2 = (Q_s^p)^2 A^{\alpha(s)/2} N_A^{1/6}$$

$$\alpha(s) = \frac{1}{3} \left(1 - \frac{1}{1 + \ln(\sqrt{s/s_0} + 1)} \right)$$

and

$$(Q_s^p)^2 = Q_0^2 \left(\frac{W}{p_T} \right)^\lambda,$$

with $Q_0 = 1 \text{ GeV}$, $W = \sqrt{s} \times 10^{-3}$ and $\lambda = 0.27$.

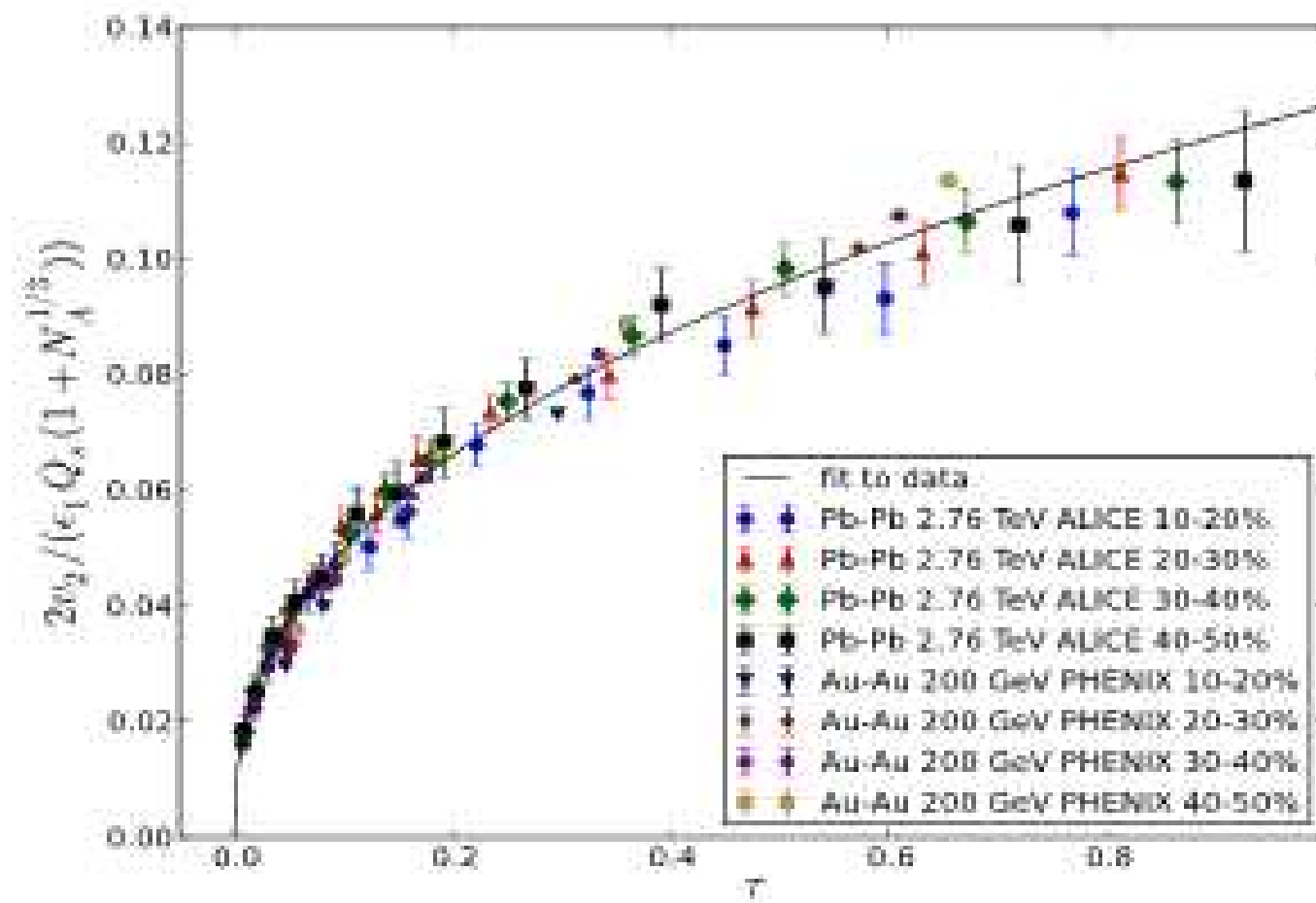


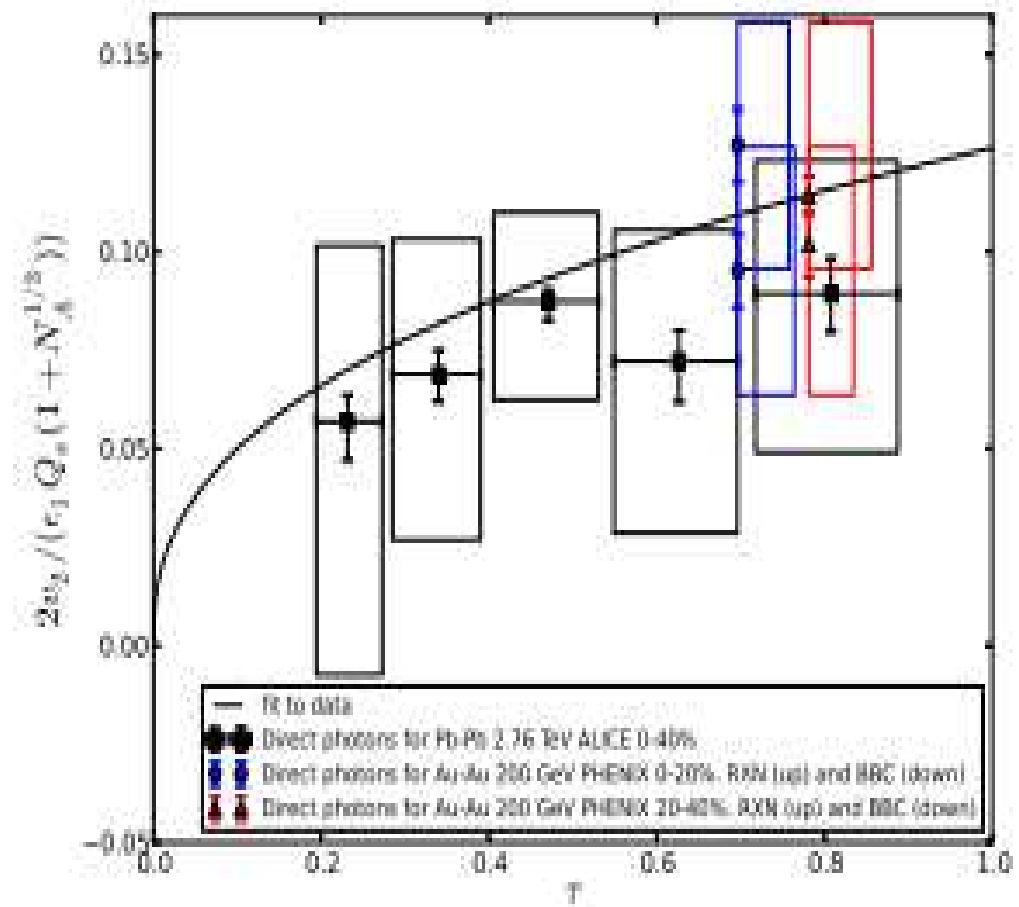
$$\frac{v_2(p_T)}{\epsilon_1 Q_s^A L} = f(\tau) \quad \text{or} \quad v_2 k_n = \epsilon_1 \tau \varphi(\tau) \quad k_n = \frac{\lambda_{mf} p}{L}$$

$$\epsilon_1 = \frac{2}{\pi} \int_0^{\pi/2} d\varphi \cos 2\varphi \frac{R^2 - R_\varphi^2}{R^2}, \quad R_\varphi = \frac{R_A \sin(\varphi - \alpha)}{\sin \varphi}$$

$$\alpha = \arcsin\left(\frac{b}{2R_A} \sin \varphi\right), \quad R^2 = \langle R_\varphi^2 \rangle = \frac{2}{\pi} \int_0^{\pi/2} d\varphi R_\varphi^2$$

$$\mathbf{L} = (1 + N_A^{1/3})/2,$$





$$(Q_{s\varphi}^A)^2 \equiv \frac{L}{\lambda_{mf p}} \frac{1}{R_\varphi^2} = \frac{1}{k_n R_\varphi^2} = \frac{Q_s^A L}{R_\varphi^2}.$$

$$v_2 = \frac{\int_0^{\pi/2} d\varphi \cos 2\varphi \frac{dN}{dp_T^2 d\varphi}}{\frac{1}{2\pi} \frac{dN}{dp_T^2}} = \frac{\int_0^{\pi/2} d\varphi \cos 2\varphi F(\tau_\varphi)}{\frac{1}{2\pi} F(\tau)}$$

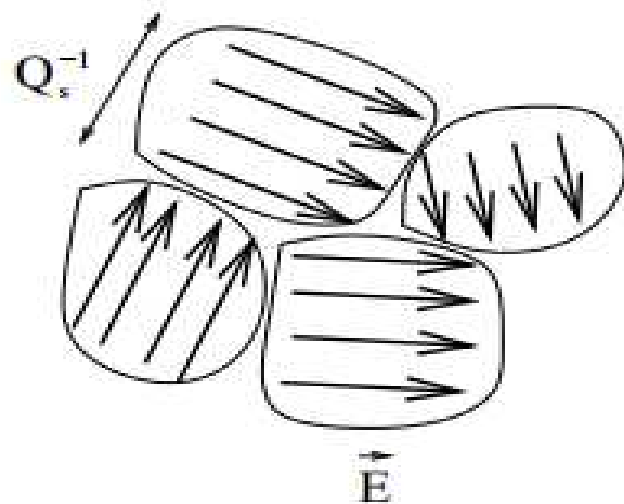
where

$$\tau_\varphi \equiv \frac{p_T^2}{(Q_{s\varphi}^A)^2} = \frac{p_T^2}{(Q_s^A)^2} \frac{R_\varphi^2}{L} Q_s^A = \tau \frac{R_\varphi^2}{L} Q_s^A$$

$$v_2 = \frac{2}{\pi} \int_0^{\pi/2} d\varphi \cos 2\varphi \frac{R^2 - R_\varphi^2}{R^2} \frac{1}{2\pi F(\tau)} \frac{dF}{d\tau} \tau Q_s^A L.$$

--Hydrodynamics, probably is not at the origin of encoding all the phi dependence in Q_s and only at most can preserve it

--It suggests models including gluon saturation (or string percolation) with a domain like structure (or cluster of strings structure)



--A parton interacts with the color fields of the cluster of strings(domain) in its way out, losing energy. This gives rise to the elliptic flow

--Two partons leaving the same cluster are correlated with a transverse correlation length($1/Q_s$), and therefore are collimated in azimuthal angle(near side ridge structure)

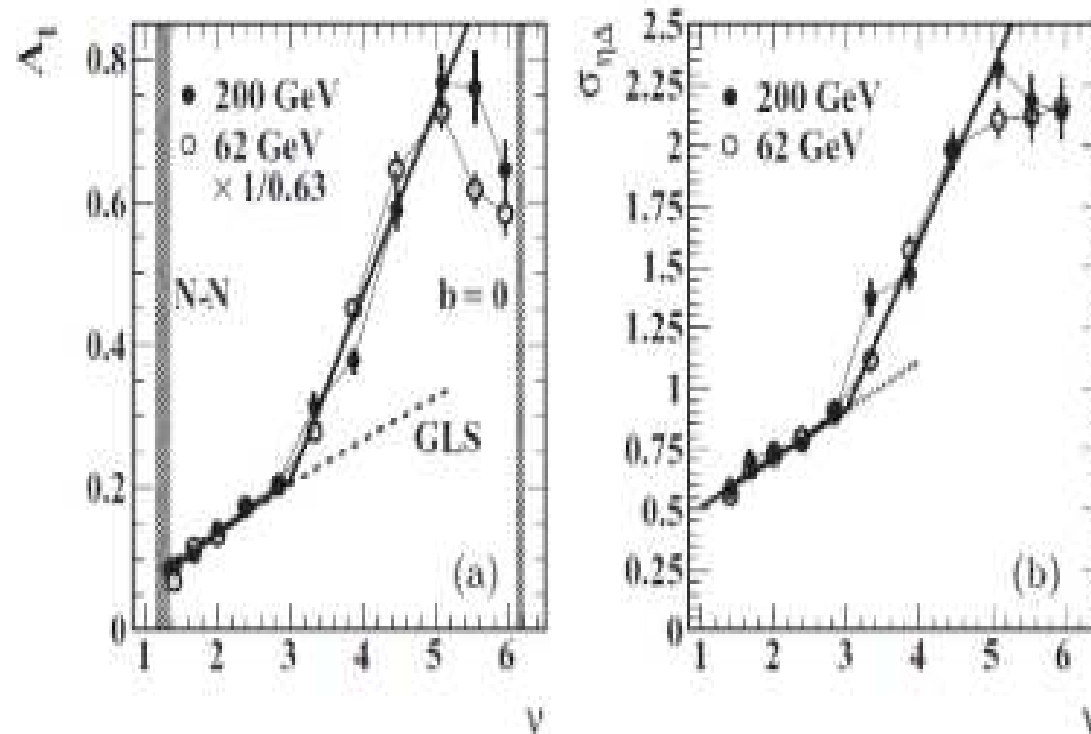
A.Kovner et al Phys Rev C83 034017 2011

M.A.Braun et al Nucl Phys A 906 14 2013

Results on 62 and 200 GeV Au-Au data from STAR

PHYSICAL REVIEW C 86, 064902 (2012)

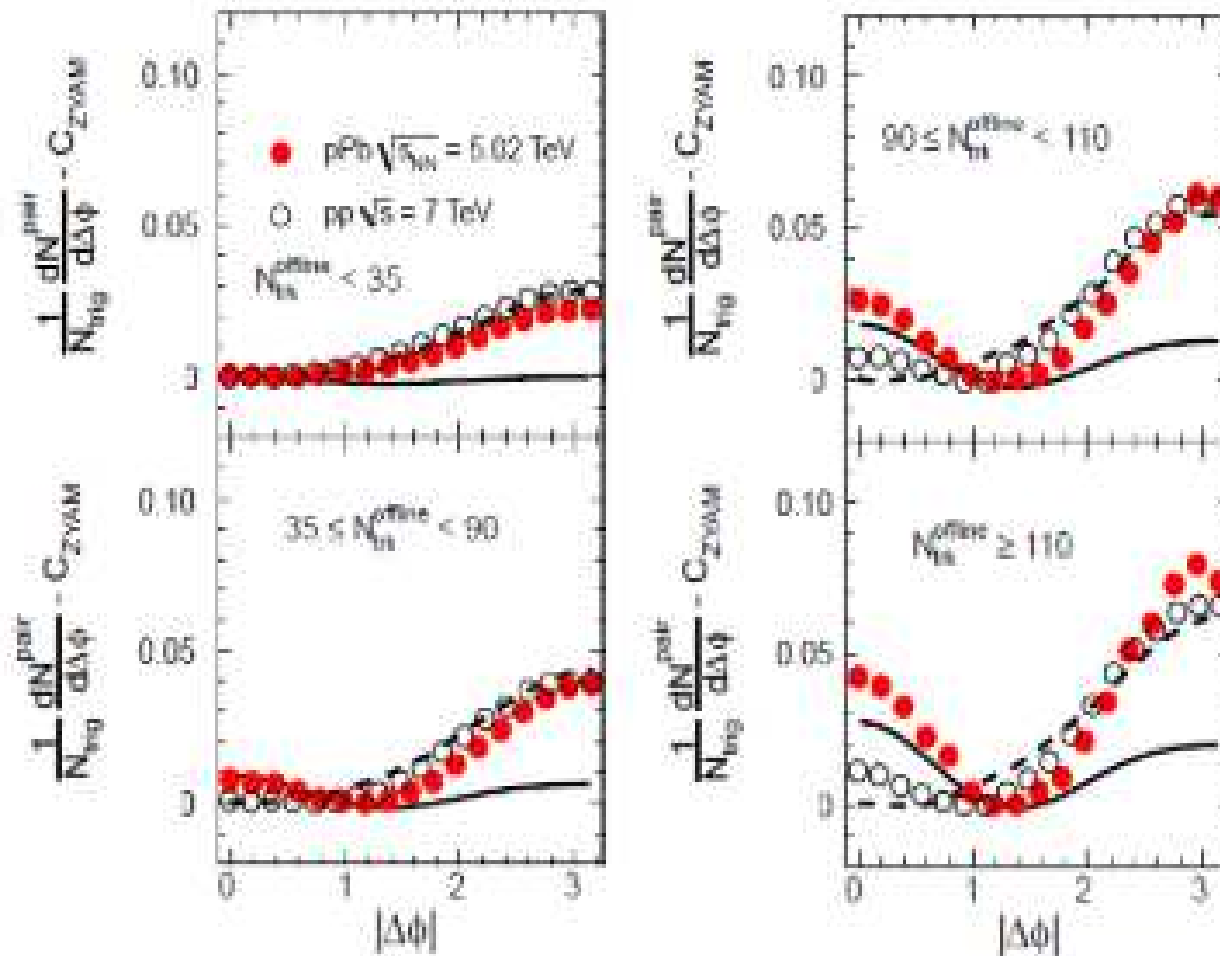
Anomalous centrality evolution of two-particle angular correlations from Au-Au collisions
at $\sqrt{s_{NN}} = 62$ and 200 GeV



Physics Letters B 718 (2013) 795–814

Observation of long-range, near-side angular correlations in pPb collisions at the LHC*

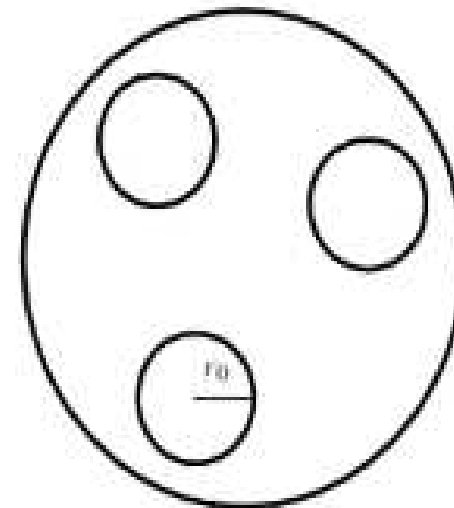
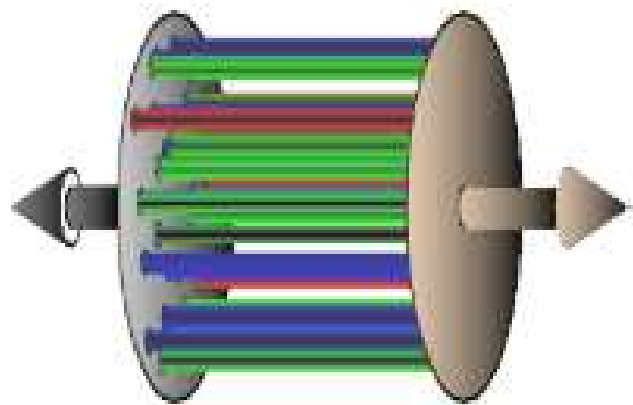
CMS Collaboration*



Outline

- 1 The string percolation model
- 2 Ridge structure
- 3 Results

Physical picture



$$r_0 = 0.2 - 0.25 \text{ fm}$$

- Projectile and target interact via color field created by the constituent partons of the nuclei.
- Color field is confined in a region with transverse size $r_0 \sim 0.2 \text{ fm}$.
- We can see them as small areas in transverse plane.
- These color “strings” break producing $q\bar{q}$ pairs (Schwinger mechanism) that subsequently lead to the observed hadrons.

String fusion

- A cluster of n strings behaves like a single string with a color field

$$\vec{Q}_n = \sum_1^n \vec{Q}_1$$

- The field is randomly oriented so

$$\langle \vec{Q}_n^2 \rangle = n \langle \vec{Q}_1^2 \rangle$$

- Using the Schwinger formula

$$\mu_n = \sqrt{\frac{nS_n}{S_1}} \mu_1, \quad \langle p_T^2 \rangle_n = \sqrt{\frac{nS_1}{S_n}} \langle p_T^2 \rangle_1$$

where μ_n and $\langle p_T^2 \rangle_n$ are, respectively, the multiplicity and the mean p_T^2 of the particles created by the fragmentation of a cluster of n strings occupying an area S_n .

Percolation



$$S_i = 3S_1$$



$$S_1$$



$$S_i < 4S_1$$



$$\rho = N_s \frac{S_1}{S_A}$$

- Limiting cases

a) $S_n = nS_1 \longrightarrow \mu_n = n\mu_1, \quad \langle \rho_T^2 \rangle_n = \langle \rho_T^2 \rangle_1$

b) $S_n = S_1 \longrightarrow \mu_n = \sqrt{n}\mu_1, \quad \langle \rho_T^2 \rangle_n = \sqrt{n} \langle \rho_T^2 \rangle_1$

- At a certain critical density 1.2-1.5 a macroscopic cluster appears which marks the percolation transition.

Homogeneous and high density case

- Mean fraction of the area covered by clusters

$$(1 - e^{-\rho}).$$

- So the basic equations concerning clusters are

$$\mu_n = N_s F(\rho) \mu_1, \quad \langle p_T^2 \rangle_n = \langle p_T^2 \rangle_1 / F(\rho).$$

where

$$F(\rho) = \sqrt{\frac{1 - e^{-\rho}}{\rho}}.$$

More realistic profile and also low density

- Mean fraction of the area covered by clusters

$$A(\rho) = \frac{1}{1 + ae^{-(\rho - \rho_0)/b}}.$$

String percolation and CGC

- Transverse size

Percolation

$$r_0 F(\rho)^{1/2}.$$

CGC

$$1/Q_s$$

- Effective number of clusters \rightarrow flux tubes

$$\langle N \rangle = \frac{(1 - e^{-\rho})R^2}{r_0^2 F(\rho)} = \sqrt{1 - e^{-\rho}} \sqrt{\rho} \left(\frac{R}{r_0} \right)^2. \quad \langle N \rangle = \frac{1}{\alpha_s} Q_s^2 R_A^2$$

- Normalized 2-particle correlation function

$$\mathfrak{K} \equiv \frac{\langle n^2 \rangle - \langle n \rangle^2 - \langle n \rangle}{\langle n \rangle^2} = \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle^2} = \frac{1}{k}$$

Low density limit: n is basically Poisson-like $\implies k \rightarrow \infty$

High density limit: $\langle N^2 \rangle - \langle N \rangle^2 \approx \langle N \rangle \implies k \rightarrow \langle N \rangle \rightarrow \infty$

$$k = \frac{\sqrt{\rho} (R/r_0)^2}{1 - e^{-\rho}} = \frac{\langle N \rangle}{(1 - e^{-\rho})^{3/2}}.$$

The two particle correlations can be written in a factorized form,

$$\frac{\Delta\rho}{\sqrt{\rho_{ref}}} = \mathcal{R} \frac{dn}{dy} G(\phi), \quad \mathcal{R} \frac{dn}{dy} = \frac{\langle N \rangle}{k} = (1 - e^{-\rho})^{3/2}$$

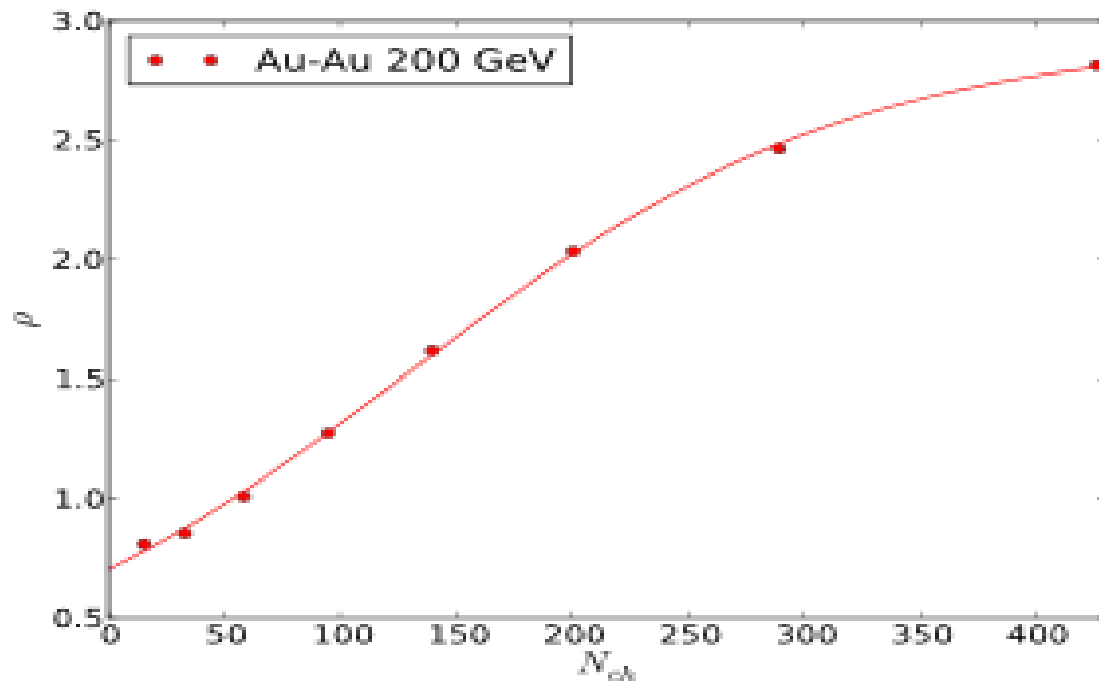
$$A_1 = cA(\rho)^{3/2} \quad A(\rho) = \frac{1}{1 + ae^{-(\rho-\rho_c)/b}}$$

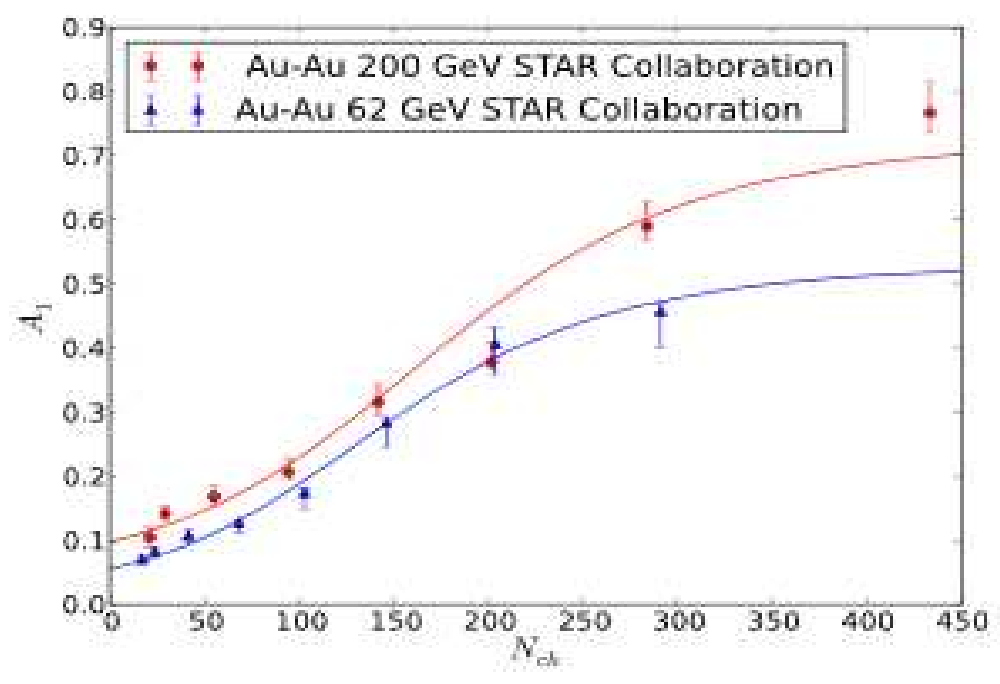
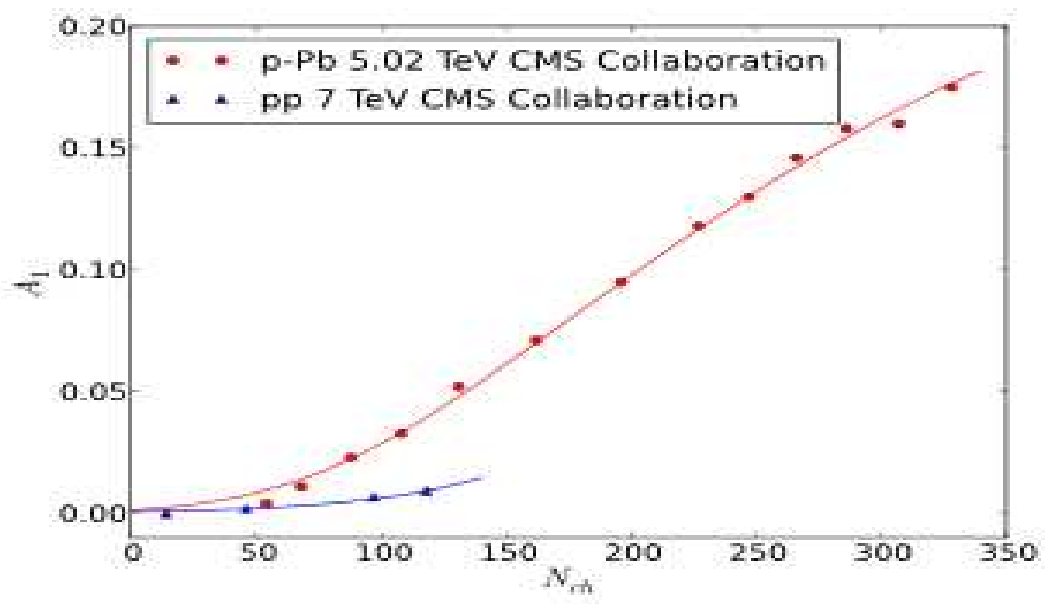
$$\rho_c = 1.5 \quad a = 1.5 \quad \begin{array}{ll} b = 0.75 & \text{for Au-Au} \\ b = 0.35 & \text{pPb and pp.} \end{array}$$

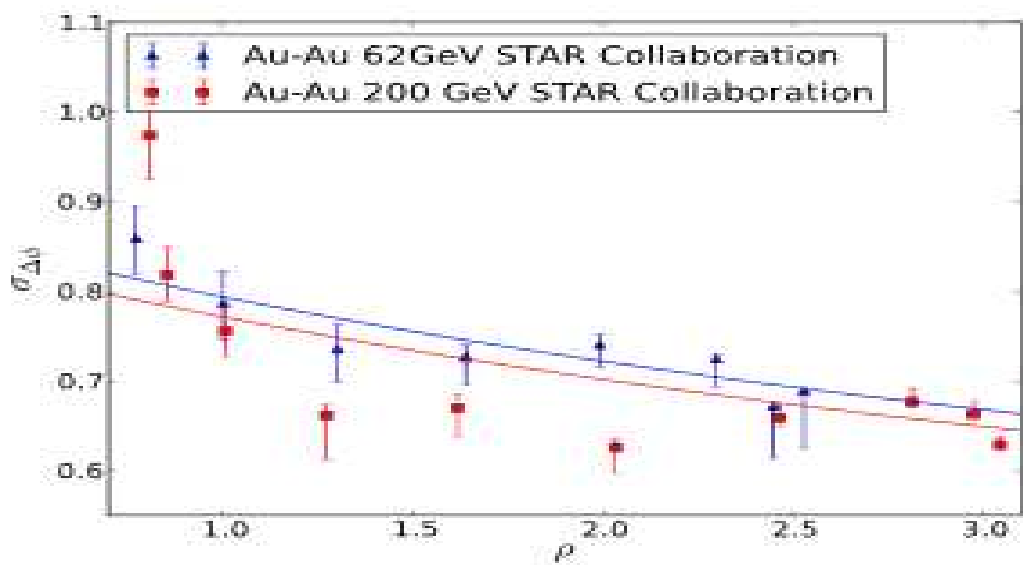
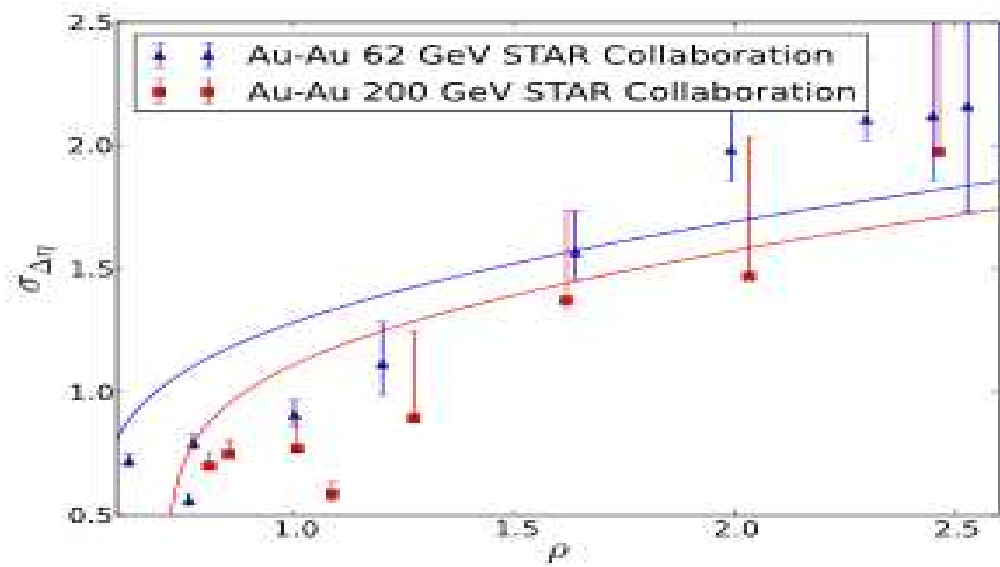
$$\sigma_{\Delta y} = c_1 \left(1 - \frac{N_A}{N_C} \right) \ln \frac{s}{s_0} + \ln \sqrt{\frac{\rho}{1 - e^{-\rho}}}$$

$$r_0 F(\rho)^{1/2}$$

$$\sigma_{\Delta\phi} = c_2 \left(\frac{1 - e^{-\rho}}{\rho} \right)^{1/4}$$







Conclusions

- The strength of the near side ridge is well described in string percolation for AA,pA and pp collisions
- This strength has to do with the fraction of the collision surface covered by strings(color sources) and therefore with the profile functions
- The widths of the pseudorapidity and azimuthal distributions are qualitatively described