## Symmetry Energy: from Nuclei to Neutron Stars

#### Pawel Danielewicz<sup>1</sup>, Jenny Lee<sup>2</sup> and Jun Hong<sup>1</sup>

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3<sup>rd</sup> International Conference on New Frontiers in Physics

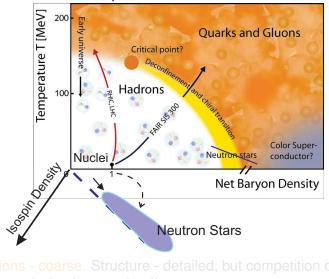
July 28 - August 6, 2014, Kolymbari, Crete, Greece



Symmetry Energy

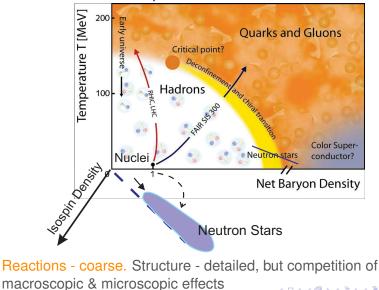
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#### Equation of State





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Symmetry Energy

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# Charge symmetry: invariance of nuclear interactions under $n \leftrightarrow p$ interchange

An isoscalar quantity *F* does not change under  $n \leftrightarrow p$ interchange. E.g. nuclear energy. Expansion in asymmetry  $\eta = (N - Z)/A$ , for smooth *F*, yields even terms only:  $F(\eta) = F_0 + F_2 \eta^2 + F_4 \eta^4 + ...$ 

An isovector quantity *G* changes sign. Example:  $\rho_{np}(r) = \rho_n(r) - \rho_p(r)$ . Expansion with odd terms only:  $G(\eta) = G_1 \eta + G_3 \eta^3 + \dots$ 

Note:  $G/\eta = G_1 + G_3 \eta^2 + \dots$ 

In nuclear practice, analyticity requires shell-effect averaging! Charge invariance: invariance of nuclear interactions under rotations in *n-p* space



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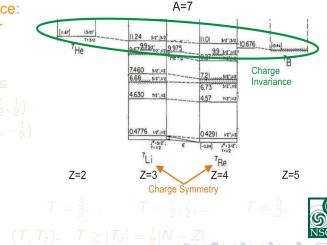
- Charge symmetry:  $n \leftrightarrow p$  invariance
- Charge invariance: symmetry under rotations in n-p space
- Isospin doublets
- $p:(\tau,\tau_z) = (\frac{1}{2},\frac{1}{2})$  $p:(\tau,\tau_z) = (1,\frac{1}{2})$

Net isospin

Net isospin

$$ec{T} = \sum^{A} ec{ au_i}$$

Isobars: Nuclei with the same A

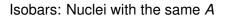


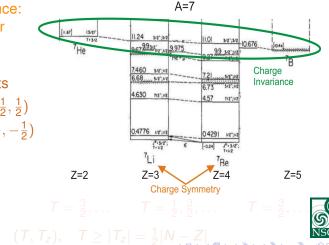


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 $\sum_{i=1}^{A} \vec{\tau}_{i}$ 

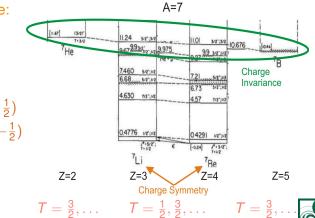




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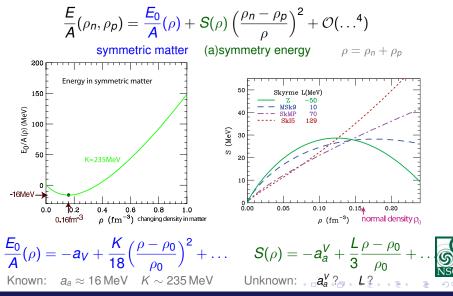
Nuclear states:  $(T, T_z), T \ge |T_z| = \frac{1}{2}|N - Z|$ 



Isobars: Nuclei with the same A

. S NSCL

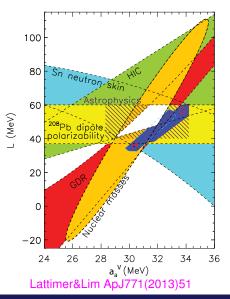
#### Energy in Uniform Matter



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#### Importance of Slope



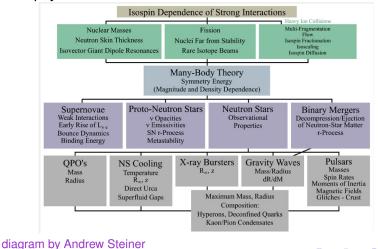
$$egin{split} rac{E}{A} &= rac{E_0}{A}(
ho) + S(
ho) \left(rac{
ho_n - 
ho_p}{
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ight)^2 \ S &\simeq -a_a^V + rac{L}{3}rac{
ho - 
ho_0}{
ho_0} \end{split}$$

In neutron matter:  $\rho_{\rho} \approx 0 \& \rho_{n} \approx \rho.$ Then,  $\frac{E}{A}(\rho) \approx \frac{E_{0}}{A}(\rho) + S(\rho)$ Pressure:  $P = \rho^{2} \frac{d}{d\rho} \frac{E}{A} \simeq \rho^{2} \frac{dS}{d\rho} \simeq \frac{L}{3\rho_{0}} \rho^{2}$ 



## Symmetry-Energy Connections

Symmetry energy ties research efforts in nuclear physics & astrophysics:



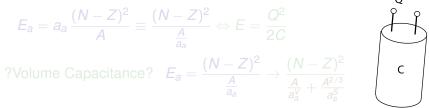


#### Symmetry Energy

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$$E = -a_V A + a_S A^{2/3} + a_C \frac{Z^2}{A^{1/3}} + a_a \frac{(N-Z)^2}{A} + E_{mic}$$

Symmetry energy: charge  $n \leftrightarrow p$  symmetry of interactions Analogy with capacitor:



ρ r

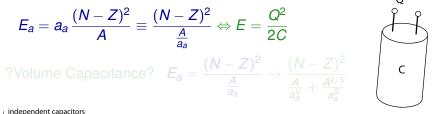
Thomas-Fermi (local density) approximation:

 $C' \equiv \frac{A}{a_a(A)} = \int \frac{\rho \, \mathrm{d} r}{S(\rho)} = \frac{A}{a_a^V}, \text{ for } S(\rho) \equiv a_a^V$ 

TF breaks in nuclear surface at  $ho < 
ho_0/4$  PD&Lee NPA\$18(2009)36

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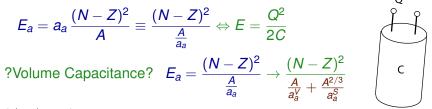
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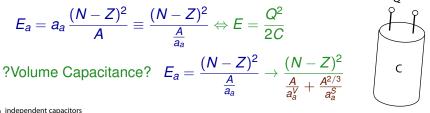


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Symmetry Energy

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### Mass Formula & Charge Invariance

Symmetry-energy details in a mass-formula are intertwined with details of other terms: Coulomb, Wigner & pairing + even those asymmetry-independent, due to (N - Z)/A - A correlations along stability line (PD)!

Best would be to study the symmetry energy in isolation from the rest of mass-formula! Absurd?!

Charge invariance to rescue: lowest nuclear states characterized by different isospin values  $(T,T_z)$ ,  $T_z = (Z - N)/2$ . Nuclear energy scalar in isospin space

sym energy 
$$E_a = a_a(A) \frac{(N-Z)^2}{A} = 4 a_a(A) \frac{T_Z^2}{A}$$

$$\rightarrow E_a = 4 a_a(A) \frac{T^2}{A} = 4 a_a(A) \frac{T(T+1)}{A}$$



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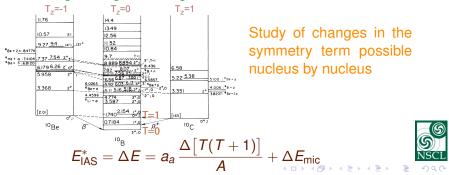
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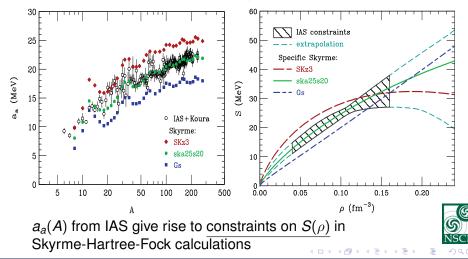
Symmetry Coefficient Nucleus-by-Nucleus Mass formula generalized to the lowest state of a given *T*:  $E(A, T, T_z) = E_0(A) + 4a_a(A) \frac{T(T+1)}{A} + E_{mic} + E_{Coul}$ In the ground state *T* takes on the lowest possible value  $T = |T_z| = |N - Z|/2$ . Through '+1' most of the Wigner term absorbed.

?Lowest state of a given *T*: isobaric analogue state (IAS) of some neighboring nucleus ground-state.

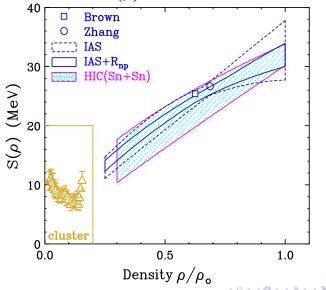


#### From $a_a(A)$ to $S(\rho)$ Strong $a_a(A)$ dependence [PD & Lee NPA922(14)1]:

lower  $A \Rightarrow$  more surface  $\Rightarrow$  lower  $\rho \Rightarrow$  lower S

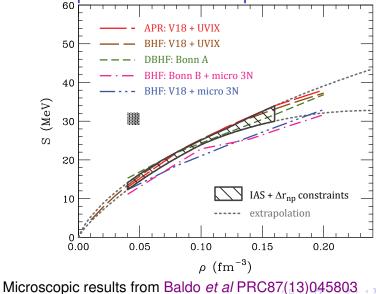


#### Subnormal $S(\rho)$ from Different Data



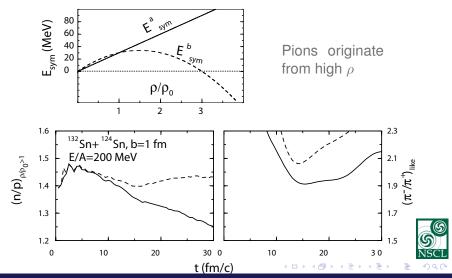


#### Comparison to Microscopic Calculations





#### Pions as Probe of High- $\rho$ Symmetry Energy B-A Li: $S(\rho > \rho_0) \Rightarrow n/p_{\rho > \rho_0} \Rightarrow \pi^-/\pi^+$

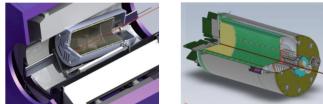


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## Dedicated Experimental Efforts

**SAMURAI-TPC Collaboration** (8 countries and 43 researchers): comparisons of near-threshold  $\pi^-$  and  $\pi^+$  and also *n-p* spectra and flows at RIKEN, Japan. NSCL/MSU, Texas A&M U Western Michigan U, U of Notre Dame GSI, Daresbury Lab, INFN/LNS U of Budapest, SUBATECH, GANIL China IAE, Brazil, RIKEN, Rikkyo U Tohoku U, Kyoto U

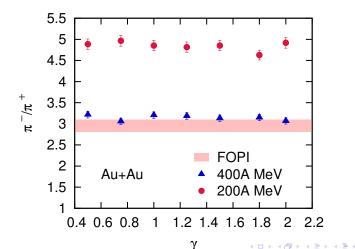
#### AT-TPC Collaboration (US & France)





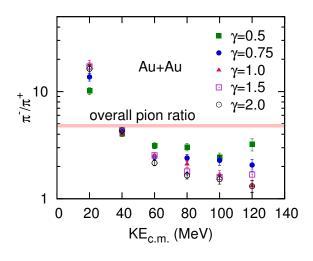
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FOPI:  $\pi^-/\pi^+$  at 400 MeV/nucl and above Hong & PD, PRC in press: measured ratios reproduced irrespectively of  $S_{int}(\rho) = S_0 (\rho/\rho_0)^{\gamma}$ :





### Original Idea Still Correct for High- $E \pi$ 's





ightarrow charge-exchange reactions blur the signal ightarrow

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### Conclusions

- Symmetry-energy term weakens as nuclear mass number decreases: from  $a_a \sim 23$  Mev to  $a_a \sim 9$  MeV for  $A \leq 8$ .
- Weakening of the symmetry term can be tied to the weakening of S(ρ) in uniform matter, with the fall of ρ.
- Convergence observed for conclusions on S(ρ) at moderately subnormal densities, from variety of data, including isobaric analog states, and from microscopic calcs testing mostly 2-body ints.
- In the region of ρ ≥ ρ<sub>0</sub>, S(ρ) is quite uncertain. One promising observable is the high-energy charged-pion yield-ratio around NN threshold.

PD&Lee NPA922(14)1; Hong&PD arXiv:1307.7654, in press NSF PHY-1068571 & PHY-1403906



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- Symmetry-energy term weakens as nuclear mass number decreases: from a<sub>a</sub> ~ 23 Mev to a<sub>a</sub> ~ 9 MeV for A ≤ 8.
- Weakening of the symmetry term can be tied to the weakening of S(ρ) in uniform matter, with the fall of ρ.
- Convergence observed for conclusions on S(ρ) at moderately subnormal densities, from variety of data, including isobaric analog states, and from microscopic calcs testing mostly 2-body ints.
- In the region of ρ ≥ ρ<sub>0</sub>, S(ρ) is quite uncertain. One promising observable is the high-energy charged-pion yield-ratio around NN threshold.

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