

Dark Messages from the LHC

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Missing Energy Signals

- Missing energy signals are a big part of the new physics menu at colliders, largely because of the potential connection to dark matter.
- We still don't know what dark matter is, but we know it is at most weakly interacting:
 - We know it should look like "nothing" to a collider detector.
- The relic density suggests that it should have reasonably large couplings to at least some of the Standard Model, in order to explain its abundance in the Universe.



"Cold Dark Matter: An Exploded View" by Cornelia Parker



A Cartoon WIMP Theory

- A typical WIMP theory has a whole "layer" of new particles.
 - E.g. SUSY, UED, Little Higgs, ...
- The WIMP is the lightest of these new states, and must be neutral and ~stable to be viable dark matter.
- Most of the heavier "WIMP siblings" usually are colored and/or charged, and thus interact much more strongly with the Standard Model particles than the WIMP does.
- The details of the model determine how LHC signals translate into other properties of dark matter, such as annihilation and nuclear scattering.





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LHC WIMP Production



LHC can produce WIMP siblings, which decay into WIMPs and other SM particles.

LHC can directly produce WIMP pairs.



WIMP Sibling Production

Squarks and Gluinos



- Searches for missing energy plus various numbers of jets put bounds on squark and/or gluino ("colored sibling") production.
 - Gluinos decay to two jets + WIMP
 - Squarks into one jet + WIMP [Assuming degenerate "light" squarks]
- These are important constraints on SUSY. The specific message for dark matter depends very much on the model parameters.

3rd Generation Squarks



- Naturalness requires SUSY to have light(ish) stops. This should be balanced by the fact that in the MSSM, the Higgs mass is calculable, suggesting the stops aren't too light.
- Searches for stops are starting to reach 600-700 GeV, and carving out the natural regions of supersymmetry!

Simplified Model

- Moving away from more complete theories, we can also consider a model containing the dark matter as well as the most important particle(s) mediating its interaction with the SM.
- For example, if we are interesting in dark matter interacting with quarks, we can sketch a theory containing a colored scalar particle which mediates the interaction.
- This looks like part of the MSSM, but has more freedom to choose couplings, etc.
- There are basically three parameters to this model: the mass of the dark matter, the mass of the mediator, and the coupling strength.







Simplified Models

- This is a model that is used by the LHC collaborations as a way of presenting more generic searches for a colored particle which decays into a single jet and missing energy.
- If we exchange the LHC production cross section for the mediator coupling to quarks, we can translate the LHC bounds into dark matter properties.

Of course, we can also consider a wider variety of WIMP properties and mediators and get away from MSSM-like theories.





ũ_R Model

- For example, we can look at a model where a Dirac DM particle couples to right-handed up-type quarks.
- Motivated by MFV we set the couplings and mediator masses equal for all three generations.
 - Third generation could actually look much different. (But this would not change the results here much).
- In the parameter plane of the mass of the dark matter and mass of the mediators, we can determine a limit on the coupling strength.





QCD production saturates the LHC limits, resulting in no allowed value of g.

ũ_R Model

- A Dirac WIMP also has spinindependent scattering with nucleons.
 For most of the parameter space, there are bounds from the Xenon-100 experiment.
 - Not included here are the recent LUX results, which improve these bounds by a factor of ~2 at a DM mass of ~100 GeV.
- Elastic scattering does not rule out any point of the mass plane, but it does impose stricter constraints on the coupling in the regions CMS left as allowed.



DiFranzo, Nagao, Rajaraman, TMPT arXiv: I 308.2679 & JHEP



Traditional direct detection searches peter out for masses below about 10 GeV.

Majorana versus Dirac



PDF-friendly qq initial state.

Majorana versus Dirac



At colliders, t-channel exchange of a Majorana WIMP can produce two mediators, leading to a PDF-friendly qq initial state.

ũ_R Model: Forecasts

- Now that we understand the current bounds, we can forecast what this implies for future searches.
- For example, we can plot the largest spin-dependent cross sections that are consistent with CMS and Xenon in this simplified model.
- Again, Dirac versus Majorana dark matter look very different from one another!

DiFranzo, Nagao, Rajaraman, TMPT arXiv: I 308.2679 & JHEP





ũ_R Model: Forecasts

- Similarly, we can forecast for the annihilation cross section.
- The Fermi LAT does not put very interesting constraints at the moment, but it is very close to doing so, and limits from dwarf satellite galaxies are likely to be relevant in the near future for Majorana DM.
- We can also ask where in parameter space this simple module would lead to a thermal relic with the correct relic density (σv ~ 10⁻²⁶ cm³/s).

DiFranzo, Nagao, Rajaraman, TMPT arXiv:1308.2679 & JHEP





Direct WIMP Production

Maverick WIMP Production

- Producing WIMPs directly requires there to be some initial radiation from the incoming quarks or gluons: a "monojet" event.
- We're not very sensitive to the details of how the WIMP couples to quarks and gluons: we can use effective field theories to parameterize all leading contributions.
- This kind of process works best for very light WIMPs, because they can be produced easily with a lot of kinetic energy, leading to large missing energy.





Birkedal, Matchev, Perelstein hep-ph/0403004 Feng, Su, Takayama hep-ph/0503117 Beltran, Hooper, Kolb, Krusberg, TMPT, 1009:037 Goodman, Ibe, Rajaraman, Shepherd, TMPT, Yu 1005.1286 & 1008.1783 Bai, Fox, Harnik 1005.3797



Example: Majorana WIMP

- As an example, we can write down operators of interest for a Majorana WIMP.
- There are 10 leading operators consistent with Lorentz and SU(3) x U(1)_{EM} gauge invariance coupling the WIMP to quarks and gluons.
- Each operator has a (separate) coefficient M* which parametrizes its strength.

Goodman, Ibe, Rajaraman, Shepherd, TMPT, Yu 1005.1286 & PLB

Name	Type	G_{χ}	Γ^{χ}	Γ^q
M1	qq	$m_q/2M_*^3$	1	1
M2	qq	$im_q/2M_*^3$	γ_5	1
M3	qq	$im_q/2M_*^3$	1	γ_5
M4	qq	$m_q/2M_*^3$	γ_5	γ_5
M5	qq	$1/2M_{*}^{2}$	$\gamma_5\gamma_\mu$	γ^{μ}
M6	qq	$1/2M_{*}^{2}$	$\gamma_5\gamma_\mu$	$\gamma_5\gamma^\mu$
M7	GG	$\alpha_s/8M_*^3$	1	-
M8	GG	$ilpha_s/8M_*^3$	γ_5	-
M9	$G\tilde{G}$	$\alpha_s/8M_*^3$	1	-
M10	$G\tilde{G}$	$i\alpha_s/8M_*^3$	γ_5	-

 $ec{q} \qquad ec{q} \qquad ec{g^2} \leftrightarrow ec{1} \ ec{q} \qquad ec{q} \qquad ec{g^2} \leftrightarrow ec{1} \ ec{M_{ ilde{q}}} \leftrightarrow ec{M_{ ilde{q}}}
ightarrow ec{q}$



Example: Majorana WIMP

Goodman, Ibe, Rajaraman, Shepherd, TMPT, Yu 1005.1286 & PLB

- The various types of interactions are accessible to different kinds of experiments.
 - Spin-independent elastic < scattering
 - Spin-dependent elastic scattering
 - Annihilation in the galactic halo
 - Collider Production –

Name	Type	G_{χ}	Γ^{χ}	Γ^q
M1	qq	$m_q/2M_*^3$	1	1
M2	qq	$im_q/2M_*^3$	γ_5	1
/M3	qq	$im_q/2M_*^3$	1	γ_5
M4	qq	$m_q/2M_*^3$	γ_5	γ_5
M5	qq	$1/2M_{*}^{2}$	$\gamma_5\gamma_\mu$	γ^{μ}
M6	qq	$1/2M_{*}^{2}$	$\gamma_5\gamma_\mu$	$\gamma_5\gamma^\mu$
$\sim M7$	GG	$\alpha_s/8M_*^3$	1	-
M 8	GG	$i\alpha_s/8M_*^3$	γ_5	-
M9	$G\tilde{G}$	$\alpha_s/8M_*^3$	1	-
M10	$G\tilde{G}$	$i\alpha_s/8M_*^3$	γ_5	-

 $G_{\chi} \left[\bar{\chi} \Gamma^{\chi} \chi \right] G^{2}$ $\sum G_{\chi} \left[\bar{q} \Gamma^{q} q \right] \left[\bar{\chi} \Gamma^{\chi} \chi \right]$

Other operators may be rewritten in this form by using Fierz transformations.

Monojets

- In terms of the WIMP mass and its interaction with quarks and/or gluons, we can predict the rate of monojet production.
- There are SM backgrounds from producing a Z which decays into neutrinos plus a jet of hadrons as well as fakes.
- Other analyses such as e.g. razor can help improve sensitivity. Fox, Harnik, Primulando, Yu 1203.1662 & PRD
- The EFT provides a mapping between collider signals, direct, and indirect rates.



From Colliders to Direct Detection



Mono-Whatever

- We can go beyond mono-jets (and mono-photons).
- One can imagine similar searches involving other SM particles, such as mono-Ws (leptons), mono-Zs (dileptons), or even mono-Higgs.
- If we're just interested in the interactions of WIMPs with quarks and gluons, these processes are not going to add much.
- But they are also sensitive to interactions directly involving the bosons.
- And even for quarks, if we do see something, they can dissect the couplings to different quark flavors, etc.



(d coupling) = $\xi \times$ (u coupling)

Jet Substructure!

 Since the events of interest have boosted Ws, one can use substructure techniques to try to capture hadronically decaying Ws.

• This helps increase statistics, and ultimately gives a better limit than the lepton channel.

• A recent ATLAS study puts this idea into practice!



Annihilation

- We can also map interactions into predictions for WIMPs annihilating.
- For example, into continuum photons from a given tree level final state involving quarks/gluons.
- This allows us to consider bounds from indirect detection, and with assumptions, maps onto a thermal relic density.
- We see similar trends as were present before: Colliders do better for lighter WIMPs or p-wave annihilations whereas indirect detection is more sensitive to heavy WIMPs.



DM Complementarity, arXiv:1305.1605

Quarks & Leptons



DM Complementarity, arXiv:1305.1605

From Sketch to Life



Outlook

- LHC has a lot to tell us about dark matter!
- Already big statements are being made about missing energy, dark matter, and supersymmetric theories with R-parity conservation.
- The next years will get into very interesting territory, with sensitivity to scalar stops and gluinos which should cover the most well-motivated regions of SUSY parameter space.
 - (And to say nothing about the Higgs mass and the MSSM...)
- More direct maverick production of dark matter is less effective than traditional SUSY searches if we can produce coloured mediator particles directly. If they are too heavy, maverick production will be how we fall back to quantify limits on dark matter interactions, and make contact between accelerator data and (in)direct searches.
- Simplified models fill a niche between complete theories like the MSSM and effective field theories assuming the mediators are inaccessible.

Sketches of <u>....</u>



Bonus Material

How Effective a Theory?



"s-channel" mediators are not protected by the WIMP stabilization symmetry. They can couple to SM particles directly, and their masses can be larger or smaller than the WIMP mass itself. "t-channel" mediators are protected by the WIMP stabilization symmetry. They must couple at least one WIMP as well as some number of SM particles. Their masses are greater than the WIMP mass (or else the WIMP would just decay into them).

Where things can go wrong, and by how much, depends on the actual UV-completion.

We can understand some general features by imagining how one could resolve the contact interaction into a mediating particle.

From WIMPs to SIMPs...

