

# Searches for beyond the Standard Model Higgs bosons in ATLAS & CMS

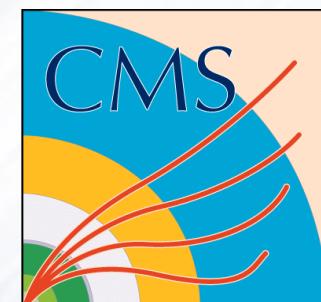
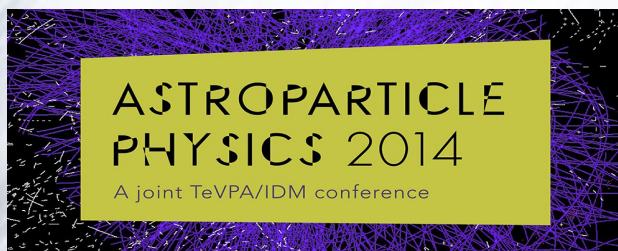
Johann Collot

LPSC

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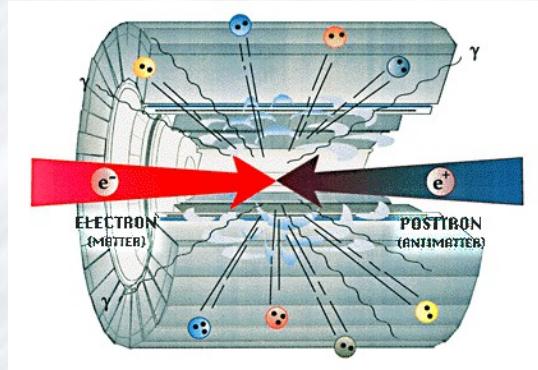
On behalf of the ATLAS & CMS collaborations





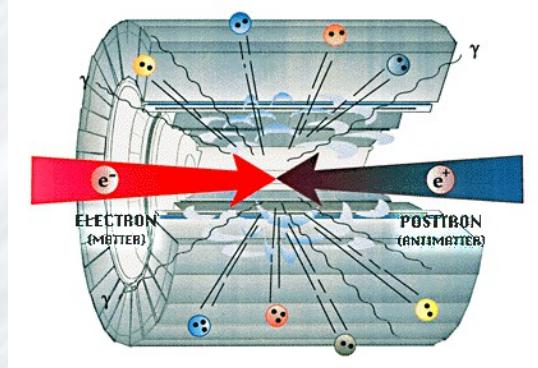
# After the discovery of 125 GeV $\text{EBH}$ boson

# After the discovery of 125 GeV EBH boson

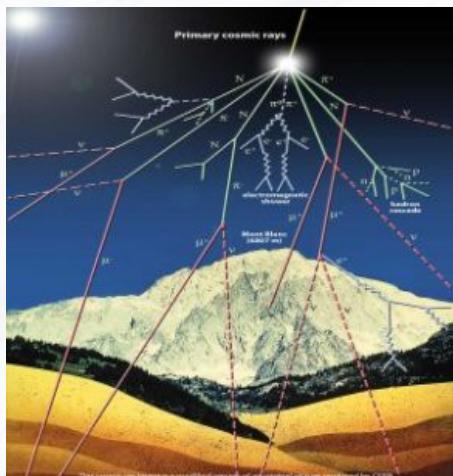


Collider physics : all Standard Model consistent  
A marvellous & much acclaimed achievement of  
modern science

# After the discovery of 125 GeV EBH boson



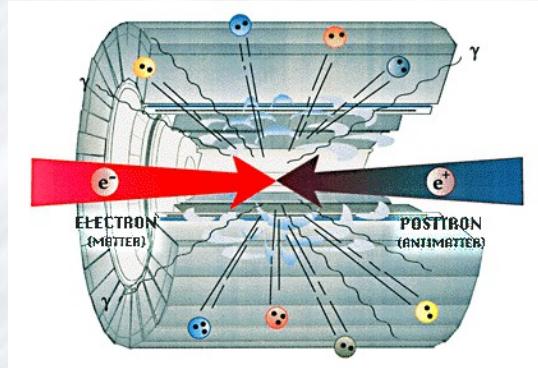
**Collider physics** : all Standard Model consistent  
A marvellous & much acclaimed achievement of modern science



**Cosmology, astrophysics & astroparticle** : many long-lasting puzzles

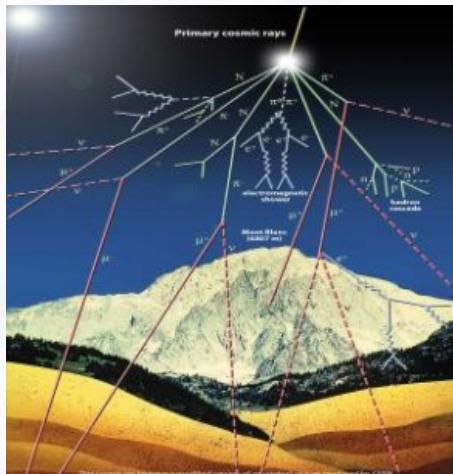
- matter-antimatter asymmetry
- dark matter
- dark energy
- cosmic inflation
- neutrino masses & oscillations

# After the discovery of 125 GeV EBH boson



Collider physics : all Standard Model consistent  
A marvellous & much acclaimed achievement of modern science

*Can the physics of fundamental scalar field(s) help to bridge these two domains ?*



Cosmology, astrophysics & astroparticle

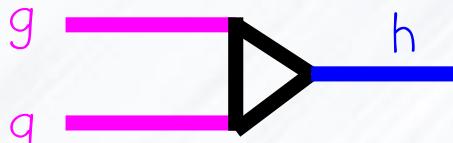
- matter-antimatter asymmetry
- dark matter
- dark energy
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# Many theorists say yes

- In general, any extension of SM comes with extra scalar fields and Higgses
- Extra scalar singlets
- More scalar doublets like in 2 Higgs Doublet Models (2HDM)
- Supersymmetry : MSSM, NMSSM...
- Extra scalar triplets
- Inflaton
- ...

# Search strategies

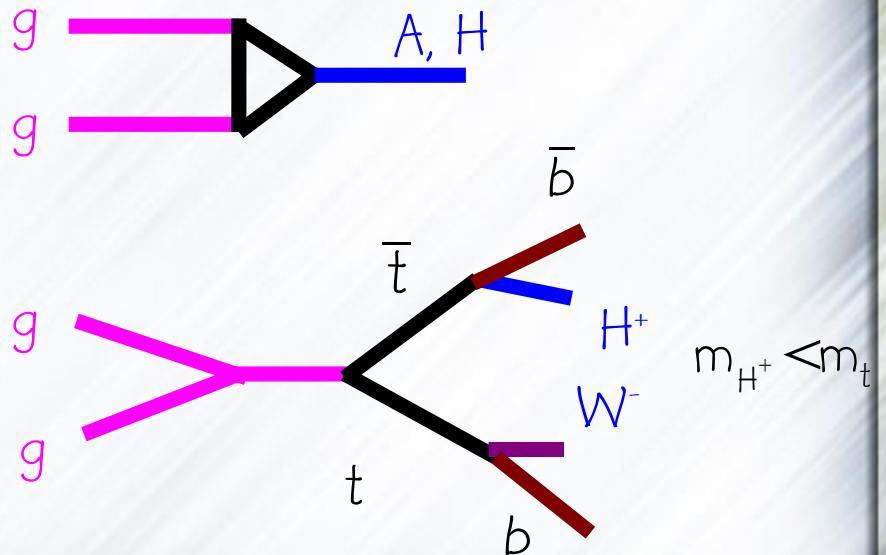
- Indirect search through careful studies of 125 GeV E<sub>BH</sub> boson properties & decays



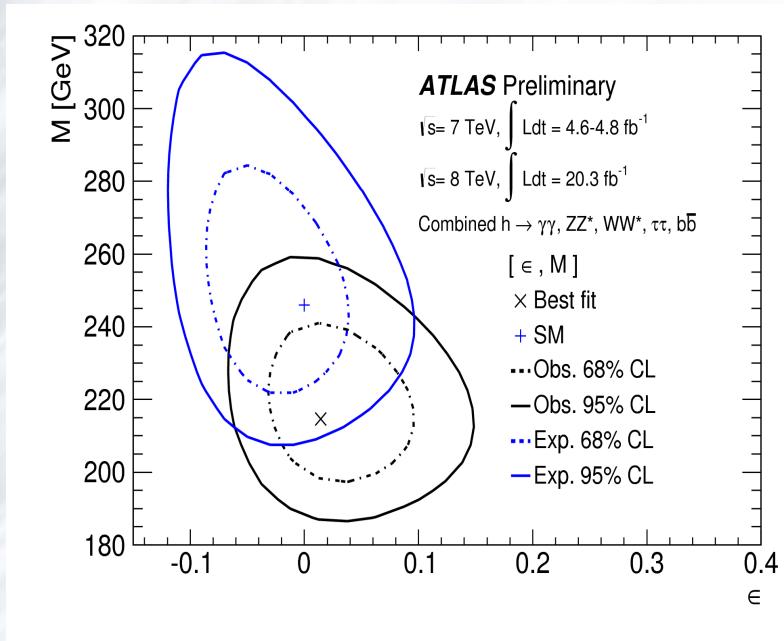
+ sub-dominant production channels : VBF, VH, ttH...

- Direct search from decays of extra neutral or charged Higgs(es)

(associated production may also significantly contribute)



# Probe mass scaling of observed h couplings



- General consistency test of observed h decay rates in all channels with respect to SM

Couplings to fermions ( $f$ ) and vector bosons ( $v$ ) scale like :

$$k_{f,i} = v \frac{m_{f,i}^\epsilon}{M^{1+\epsilon}} \quad \text{linear mass dependence}$$

$$k_{v,j} = v \frac{m_{v,j}^{2\epsilon}}{M^{1+2\epsilon}} \quad \text{quadratic mass dependence}$$

$$\text{SM} \Rightarrow \epsilon = 0, M = v = 246 \text{ GeV}$$

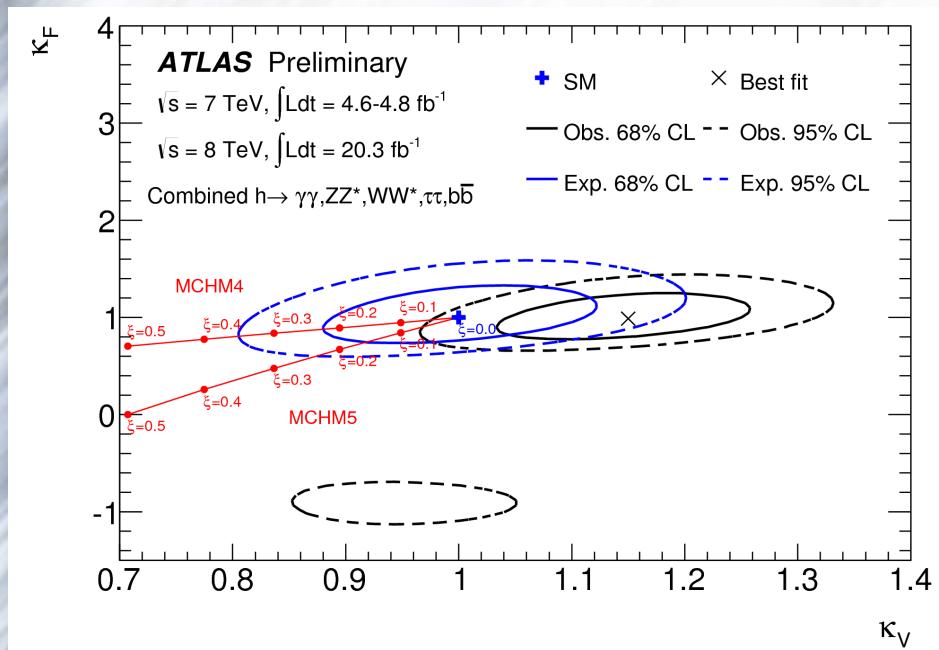
Best fit agrees with SM at  $1.5 \sigma$

ATLAS-CONF-2014-010

Ellis & You, JHEP 1306 (2013) 103

# Could observed h be composite ?

- Minimal Composite Higgs Models – Scaling factors times SM couplings



ATLAS-CONF-2014-010

K. Agashe et al., Nucl.Phys. B719 (2005) 165

R. Contino, Phys. Rev. D75 (2007) 055014

M.S. Carena et al. Phys. Rev. D76 (2007) 035006

MCHM4 :  $k = k_v = k_f = \sqrt{1-\epsilon}$ , with  $\epsilon = v^2/f^2$

$f$  is the  $h$  compositeness scale

SM :  $f \rightarrow \infty$ ,  $\epsilon = 0$

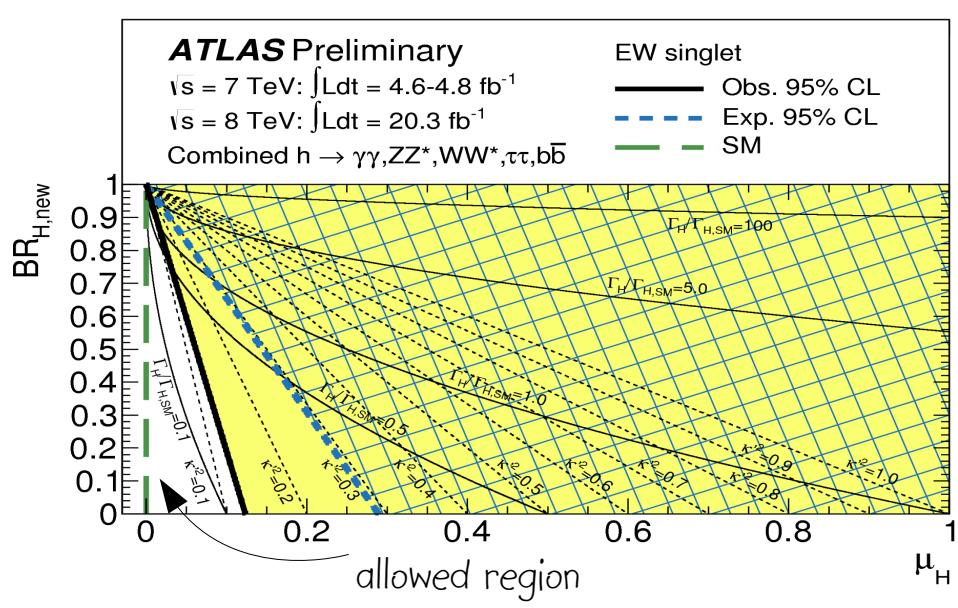
$f > 710 \text{ GeV}$  or  $0 \leq \epsilon < 0.12 @ 95 \text{ CL}$

MCHM5 :  $k_v = \sqrt{1-\epsilon}$ , and  $k_f = \frac{1-2\epsilon}{\sqrt{1-\epsilon}}$

$f > 640 \text{ GeV}$  or  $0 \leq \epsilon < 0.15 @ 95 \text{ CL}$

# Extra EW scalar singlet in data ?

- Assume the addition of a scalar singlet that mixes with doublet state after SSB to produce two neutral Higgses :  $h$  &  $H$
- Assume both couple to  $f$  and  $V$  in a similar way as SM but with a strength reduced by  $k$  &  $k'$  scale factors, for  $h$  &  $H$  respectively



ATLAS-CONF-2014-010

LHC Higgs XS WG, arxiv : 1307.1347 [hep-ph]

A. Hill & J van der Bij, Phys. Rev. D36

M. Veltman & F. Yndurain, Nucl. Phys. B325 (1989)

T. Binoth & J. van der Bij, Z. Phys. C75 (1997) 17

R. Schabinger & J.D. Wells, Phys. Rev. D 72 (2005) 093007

B. Patt & F. Wilczek, arxiv : hep-ph/0605188

G.M. Pruna & T. Robens, arxiv : 1303.1150 [hep-ph]

$$k^2 + k'^2 = 1$$

$$k^2 = \mu_h \quad \text{signal strength of } h$$

$$k'^2 = 1 - \mu_h$$

$$\mu_H = k'^2(1 - BR_{H,\text{new}}) \quad \text{signal strength of } H$$

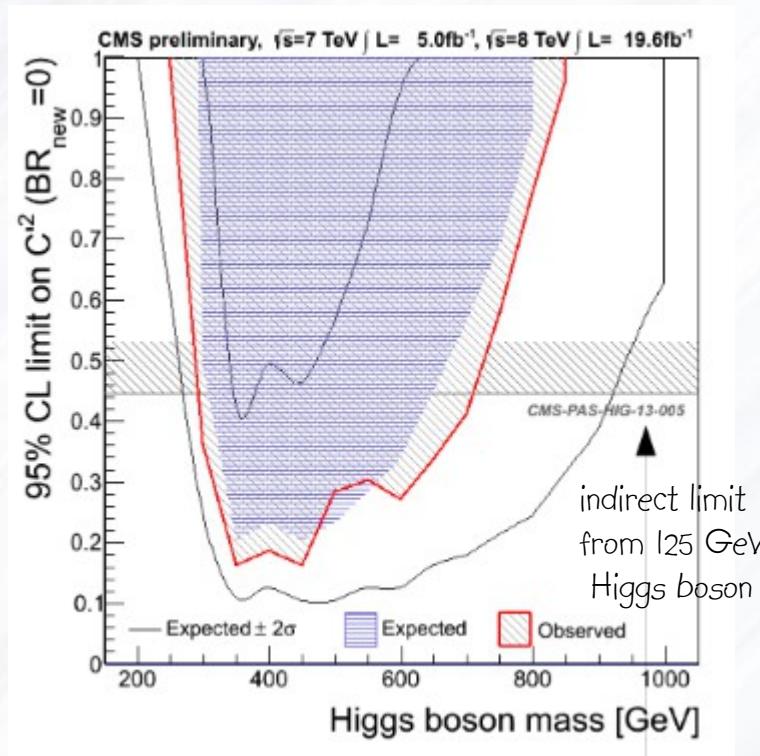
branching ratio of all new non-SM kinematically-allowed channels that open when  $m_H$  increases

$$\text{SM}: k' = 0 = \mu_H$$

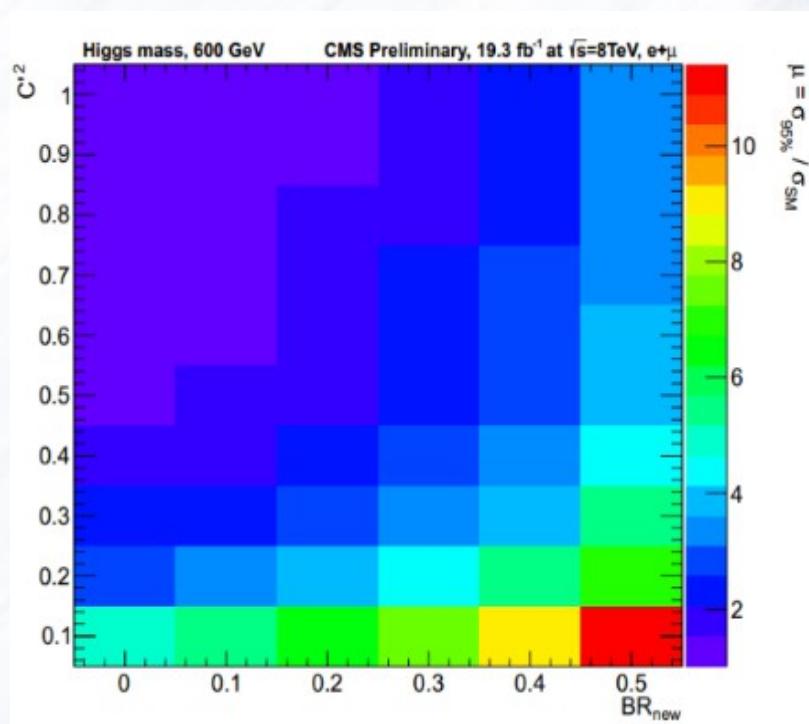
# Extra EW scalar singlet in data ?

Similar limits were obtained with CMS in direct & indirect searches

$H \rightarrow ZZ \rightarrow 2l\bar{l} 2\nu$



$H \rightarrow WW \rightarrow l\nu q\bar{q}'$



$$C'^2 = K'^2$$

CMS-PAS-HIG-13-014

CMS-PAS-HIG-13-008

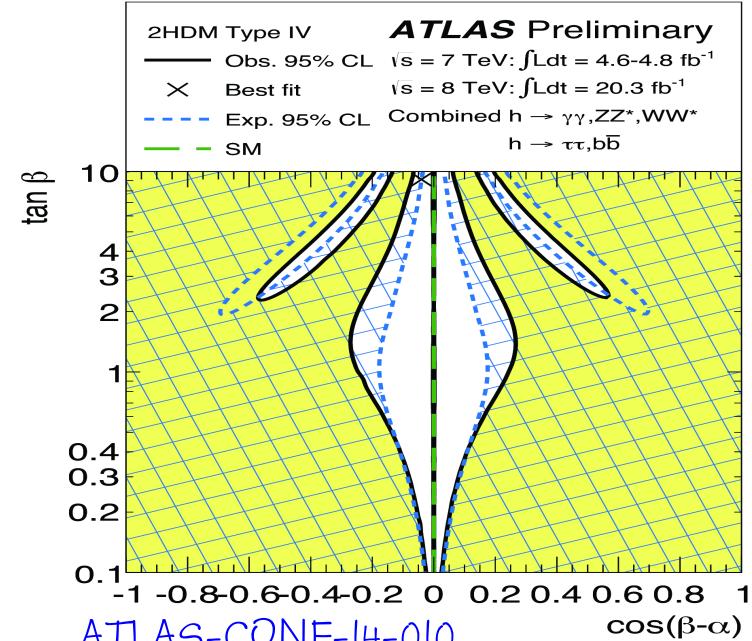
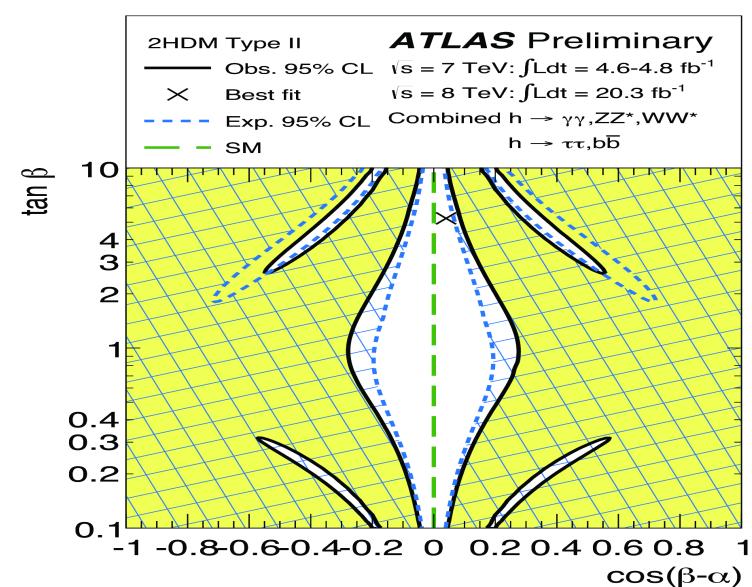
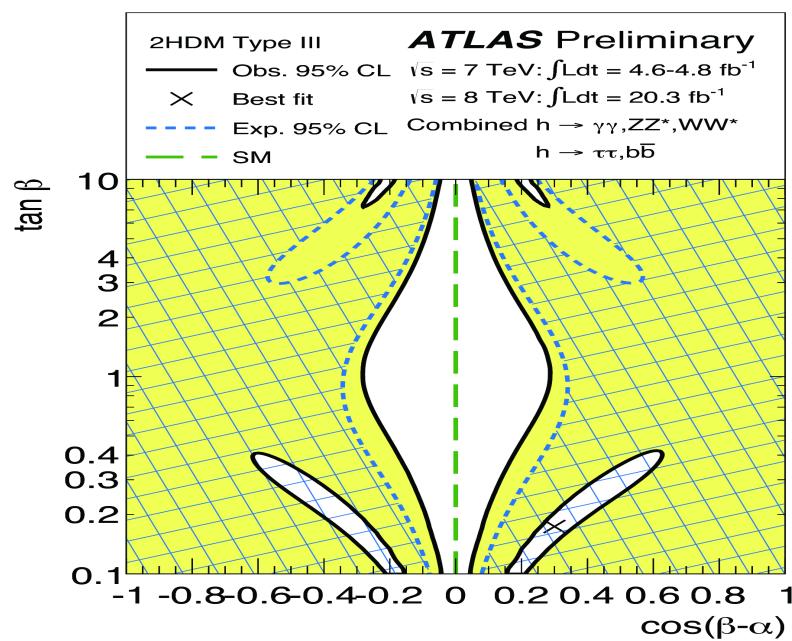
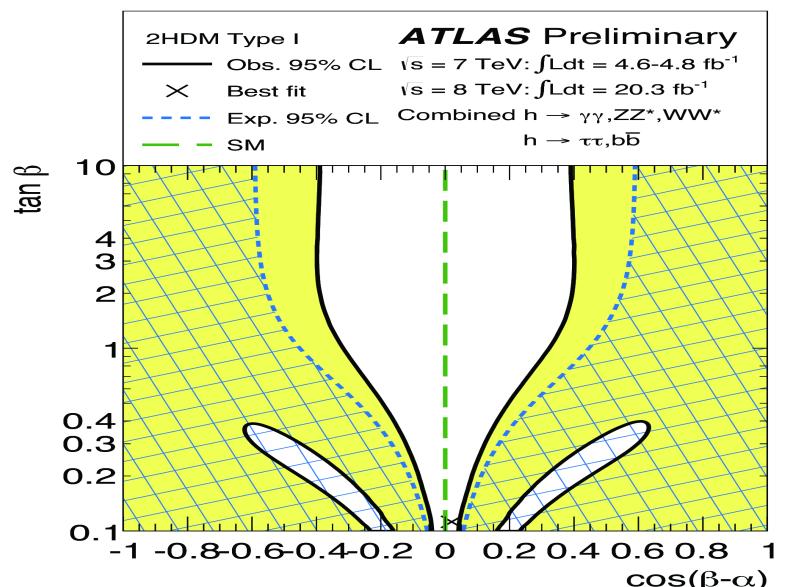
# 2HDM reminder

- 8 initial degrees of freedom
  - 3 eaten by  $W_L$ 's and  $Z_L$
  - 3 neutral Higgses and 2 charged ones ( $H^+$ ,  $H^-$ )
- CP conservation and Paschos-Glashow-Weinberg condition :
  - 2 CP-even neutrals :  $h$  &  $H$ , 1 CP-odd neutral  $A$ ,  $H^+$ ,  $H^-$
  - 4 Higgs masses,  $\tan \beta = v_u/v_d$ ,  $\alpha$  :  $h$  &  $H$  mixing angle
  - $v_u^2 + v_d^2 = v^2 = (246 \text{ GeV})^2$

	coupling scale factor of $h$ /SM	type I fermophobic	type II MSSM like	type III lepton-specific	type IV flipped
gauge bosons	$k_v$	$\sin(\beta-\alpha)$	$\sin(\beta-\alpha)$	$\sin(\beta-\alpha)$	$\sin(\beta-\alpha)$
up quarks	$k_u$	$\cos(\alpha)/\sin(\beta)$	$\cos(\alpha)/\sin(\beta)$	$\cos(\alpha)/\sin(\beta)$	$\cos(\alpha)/\sin(\beta)$
down quarks	$k_d$	$\cos(\alpha)/\sin(\beta)$	$-\sin(\alpha)/\cos(\beta)$	$\cos(\alpha)/\sin(\beta)$	$-\sin(\alpha)/\cos(\beta)$
charged leptons	$k_l$	$\cos(\alpha)/\sin(\beta)$	$-\sin(\alpha)/\cos(\beta)$	$-\sin(\alpha)/\cos(\beta)$	$\cos(\alpha)/\sin(\beta)$

G. Branco et al :arXiv:1106.0034 [hep-ph].

# Additional scalar doublet in data ?

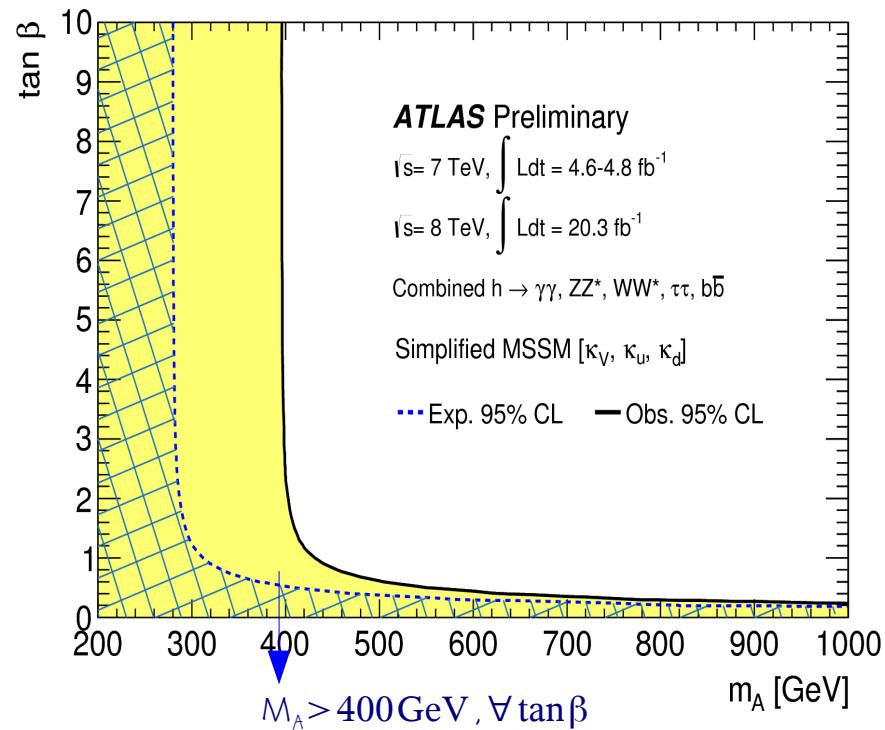


ATLAS-CONF-14-010

# Simplified MSSM

- In such a model, coupling scale factors are expressed as functions of  $m_A$  and  $\tan\beta$ .

ATLAS-CONF-14-010



$$\kappa_v = \frac{s_d(m_A, \tan\beta) + \tan\beta s_u(m_A, \tan\beta)}{\sqrt{1 + \tan^2\beta}}$$

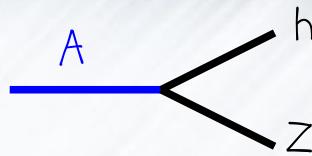
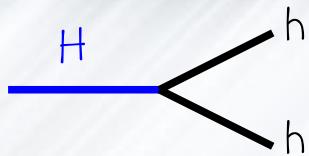
$$\kappa_u = s_u(m_A, \tan\beta) \frac{\sqrt{1 + \tan^2\beta}}{\tan\beta}$$

$$\kappa_d = \kappa_l = s_d(m_A, \tan\beta) \sqrt{1 + \tan^2\beta}$$

$$s_u = \left( 1 + \left( \frac{(m_Z^2 + m_A^2) \tan\beta}{m_Z^2 + m_A^2 \tan^2\beta - m_h^2 (1 + \tan^2\beta)} \right)^2 \right)^{-1/2}$$

$$s_d = \frac{(m_Z^2 + m_A^2) \tan\beta}{m_Z^2 + m_A^2 \tan^2\beta - m_h^2 (1 + \tan^2\beta)} s_u$$

# direct search for $H \rightarrow hh$ & $A \rightarrow Zh$

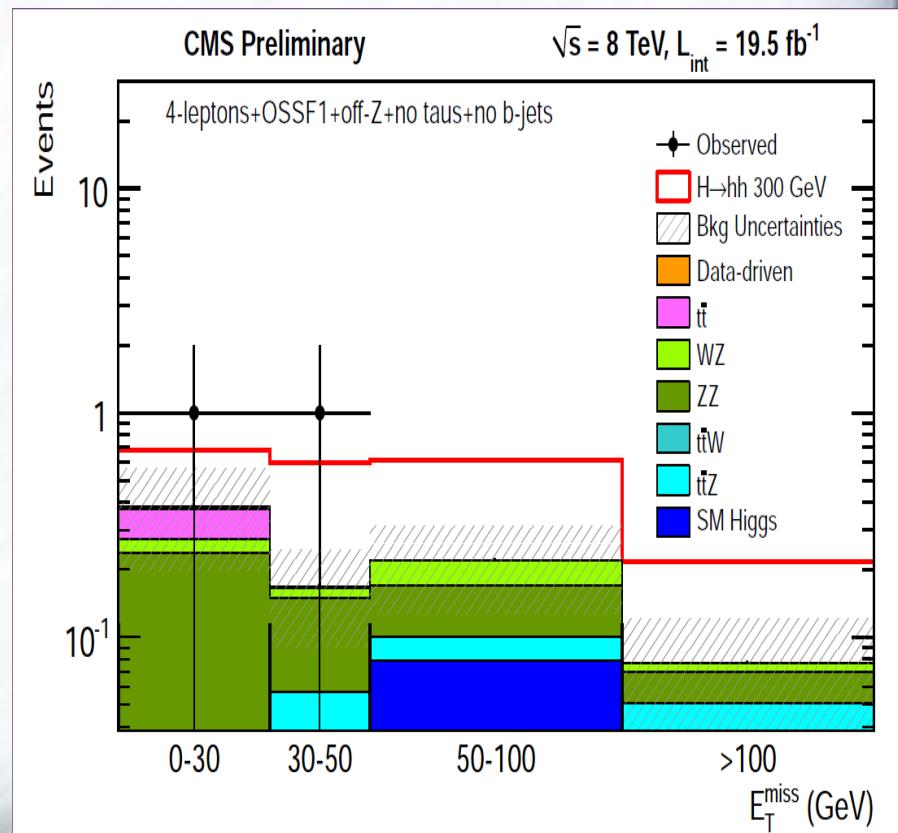


- for  $m_H > 2 m_h$  and  $2 m_h < m_A < 2 m_t$ , these decays are predominant
- $h$  decays like in SM, leads to multilepton and diphoton signatures further categorized as in table below

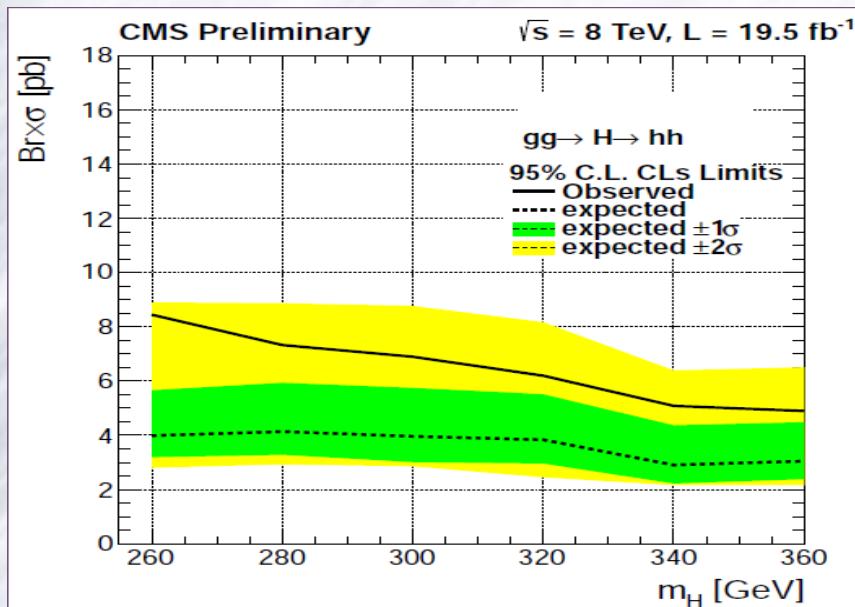
leptons	photons	opposite side same flavor pairs	hadronic tau	$b$ -tag
4	0	0,1 or 2	0 or 1	0 or 1
3	0	0 or 1	0 or 1	0 or 1
2	2	0 or 1	0	-
1	2	-	0	-
1	2	-	1	-
0	2	-	1 or 2	-

each channel further divided in bins of  $E_T^{\text{miss}}$

CMS-PAS-HIG-13-025



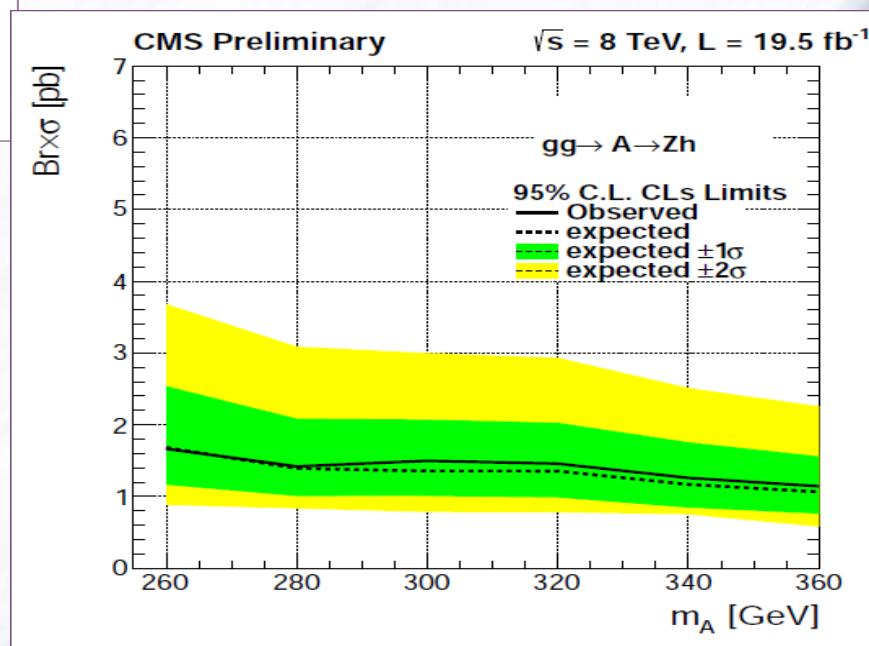
# direct search for H $\rightarrow$ hh & A $\rightarrow$ Zh



CMS-PAS-HIG-13-025

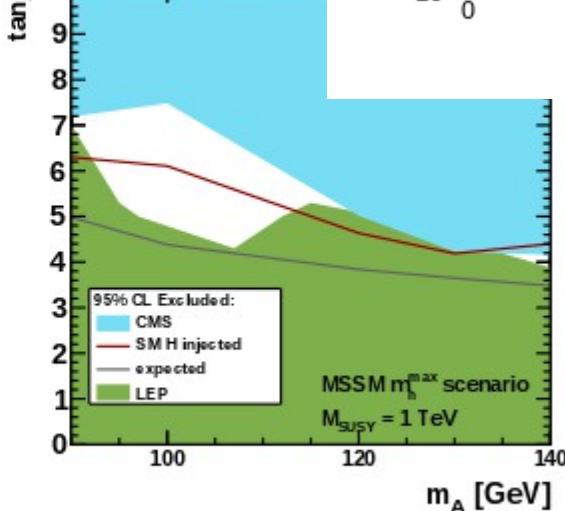
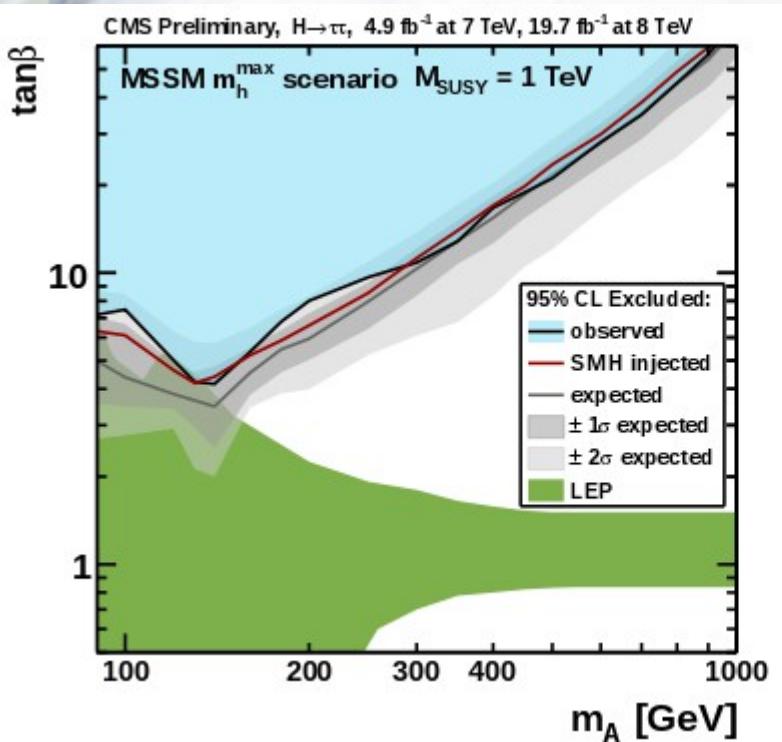
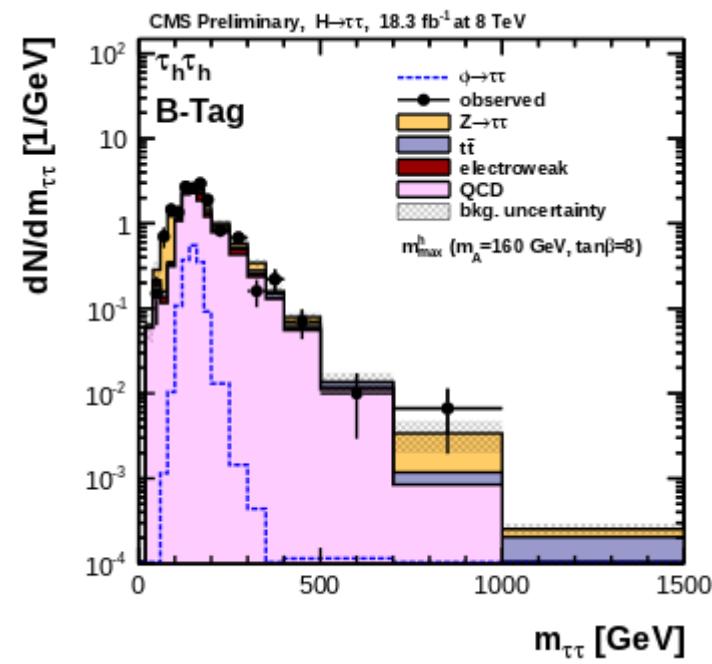
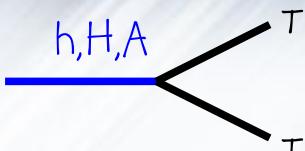
model-independent limits

$B \times \sigma$  H (A) production probed down to 4 (1.5) pb.



# Direct search for MSSM $h, H, A \rightarrow 2 \tau$ s

- $h$  is 125 GeV SM Higgs
- 0 or 1 b-tag for ggF or bbH production
- MSSM  $m_h^{\max}$  scenario (to push  $m_h$  to 125 GeV)

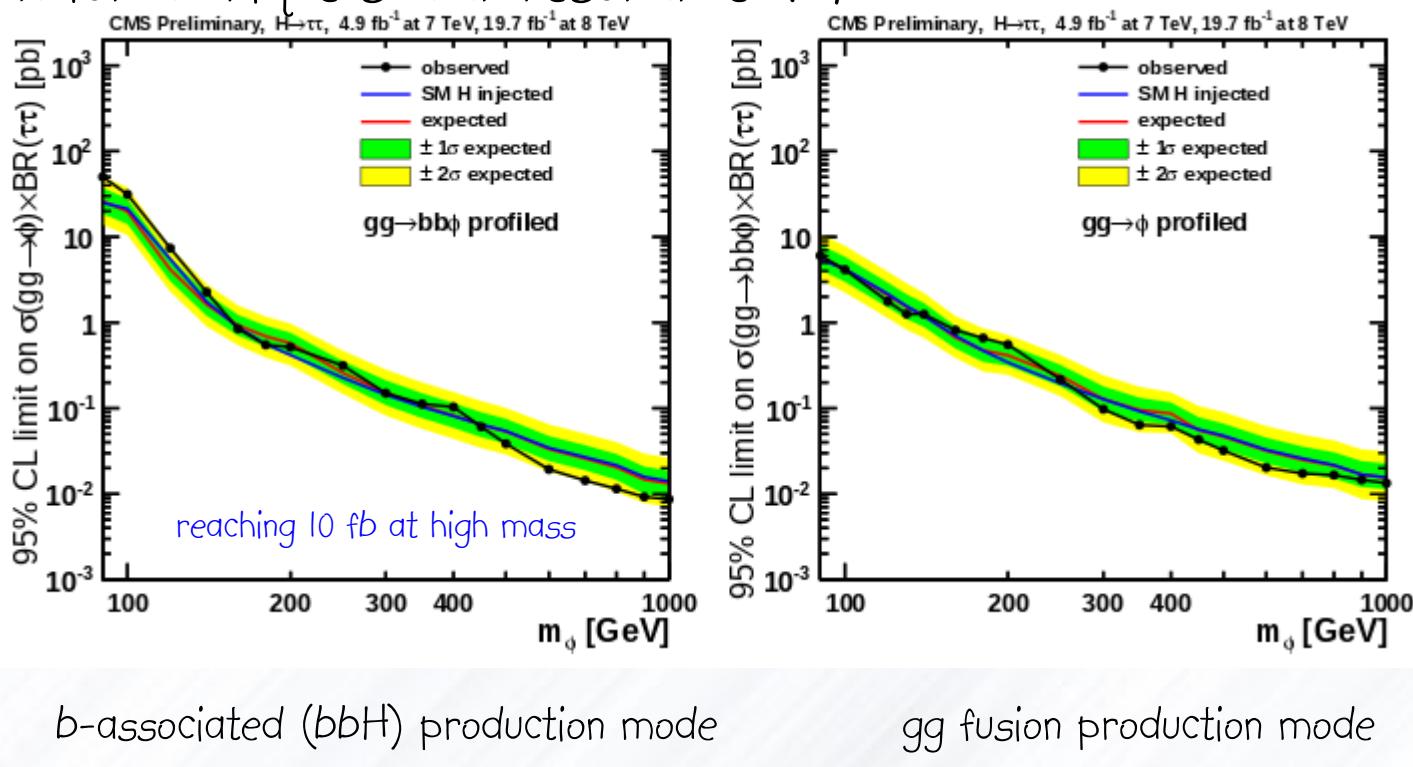


CMS-PAS-HIG-13-021  
MS. Carena et al.  
Eur. Phys. J.C. 45 (2006) 797  
Eur. Phys. J.C. 26 (2003) 601

Complementary to ATLAS indirect search limits  
especially at high  $m_A$

# Direct search for $\phi \rightarrow 2 \tau$ s

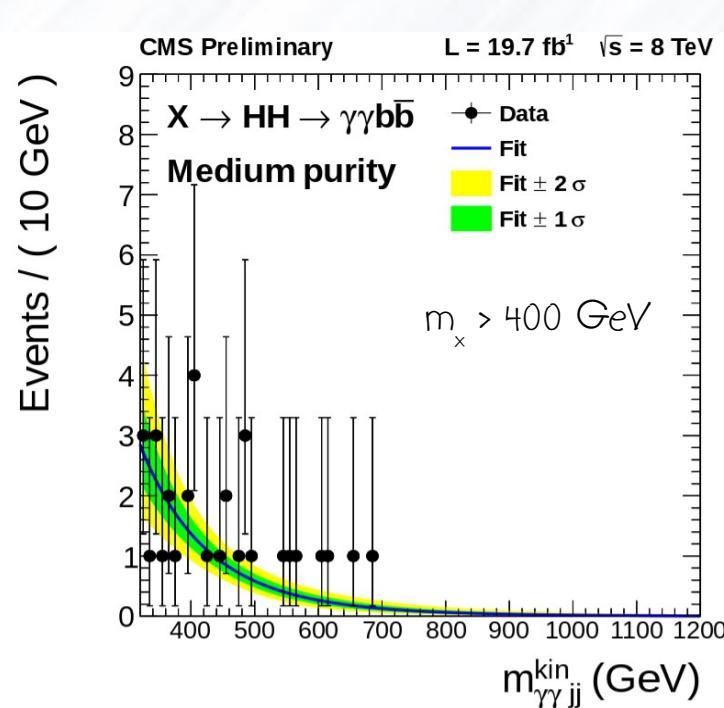
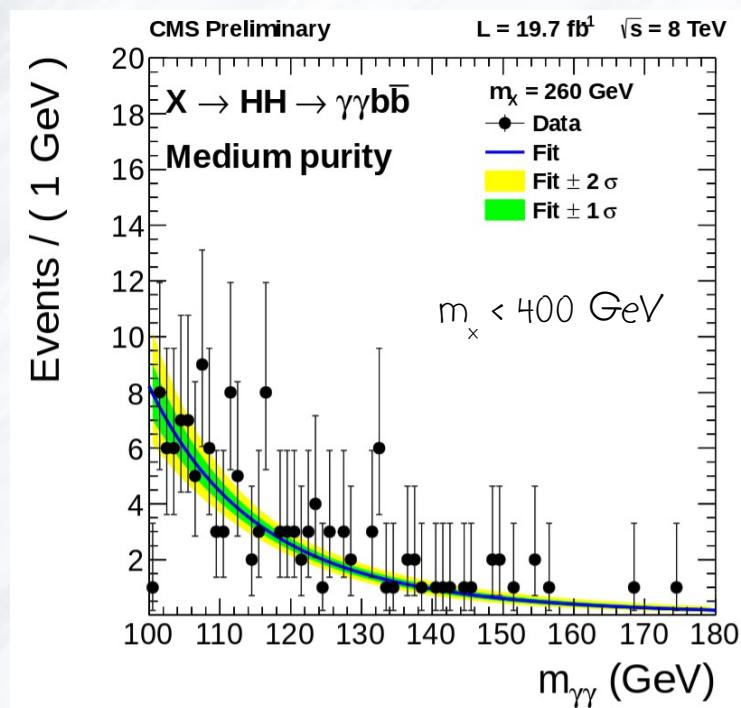
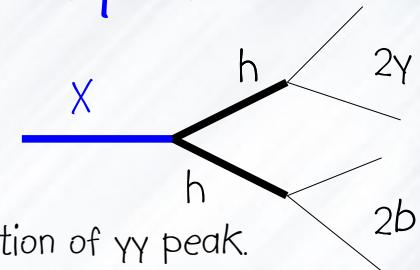
- same analyzed data interpreted in a model independent way
- search for a unique scalar resonance :  $\phi$



CMS-PAS-13-021

# direct search of $X \rightarrow hh \rightarrow 2\gamma 2b$

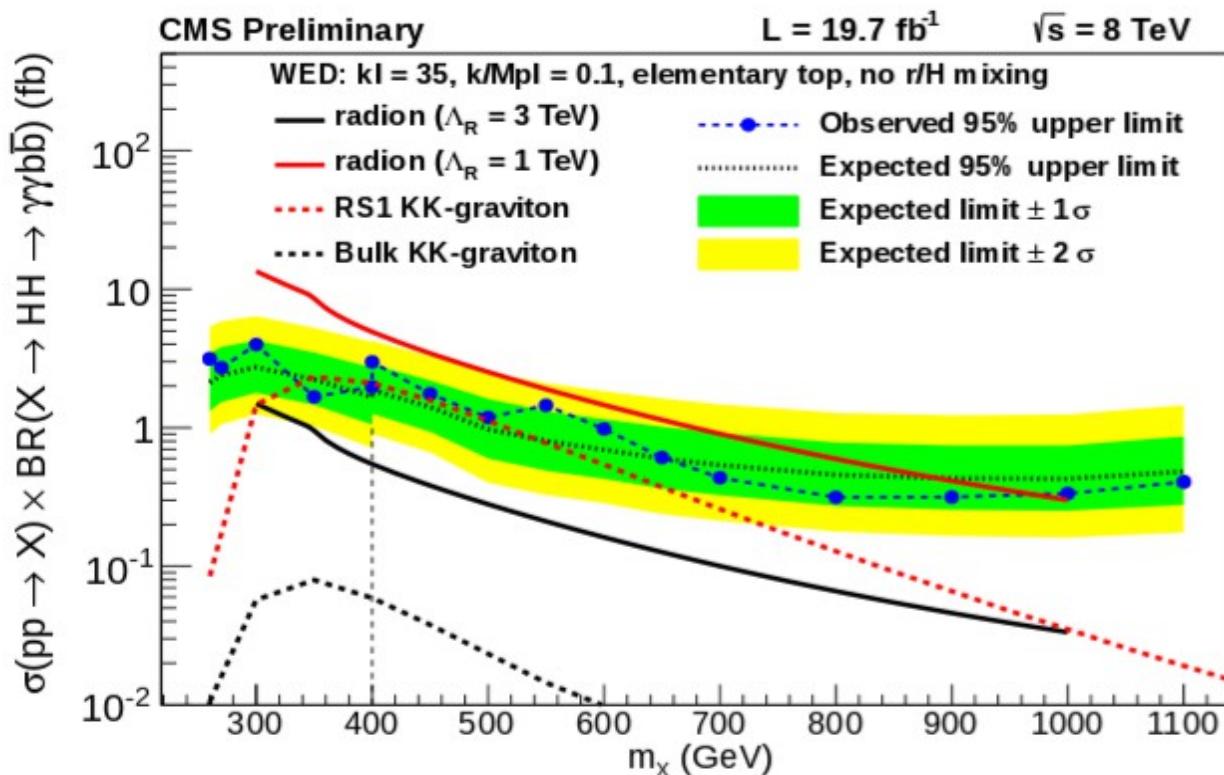
- search for a new  $hh$  resonance in the  $m_X$  region 260-1100 GeV
- combines high  $bb$  BR with better background rejection capability & resolution of  $\gamma\gamma$  peak.
- search in the  $m_{\gamma\gamma}$  spectrum for  $m_X < 400$  GeV and in  $m_{\gamma\gamma jj}^{\text{kin}}$  spectrum for  $m_X > 400$  GeV



CMS-PAS-13-032

# direct search of $X \rightarrow hh \rightarrow 2\gamma 2b$

- Model-independent upper limit
- Prediction lines of several Wrapped Extra Dimension models also shown



CMS-PAS-13-032

C. Csaki et al. : arXiv hep-th/0008151

L. Randall & R. Sundrum : arXiv hep-ph/9905221

T. Hapola and O. Antipin : <http://cp3-origins.dk/research/units/ed-tools>

P. Aquino : <http://feynrules.irmp.ucl.ac.be/wiki/RSmodel>

A. L. Fitzpatrick et al. : arXiv hep-ph/0701150

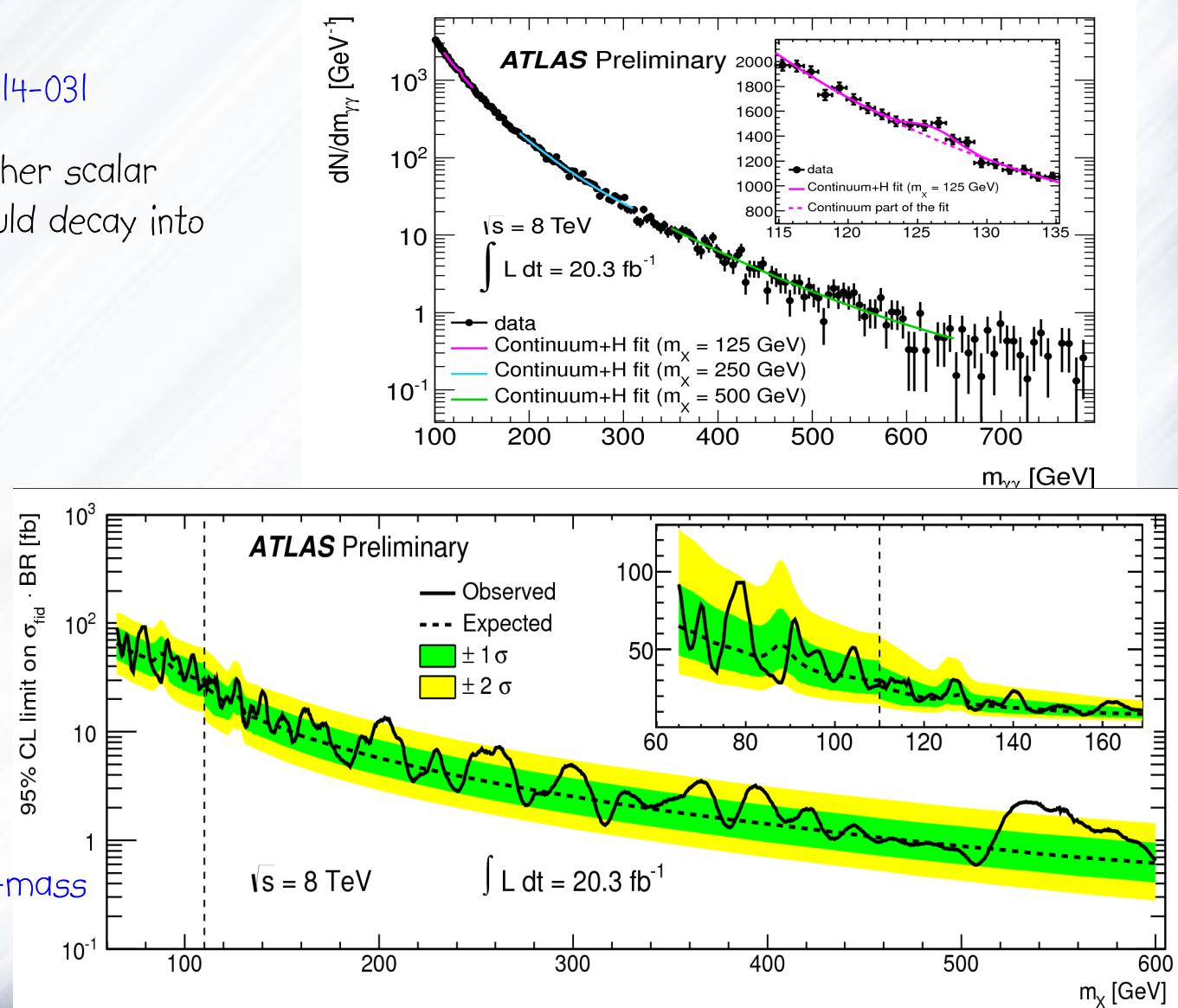
G. F. Giudice et al. : arXiv hep-ph/0002178

Ruling out one radion model up to  $m_X = 1 \text{ TeV}$

# Model-independent search for $H \rightarrow 2$ photons

ATLAS-CONF-2014-031

Searching for another scalar resonance that would decay into 2 photons

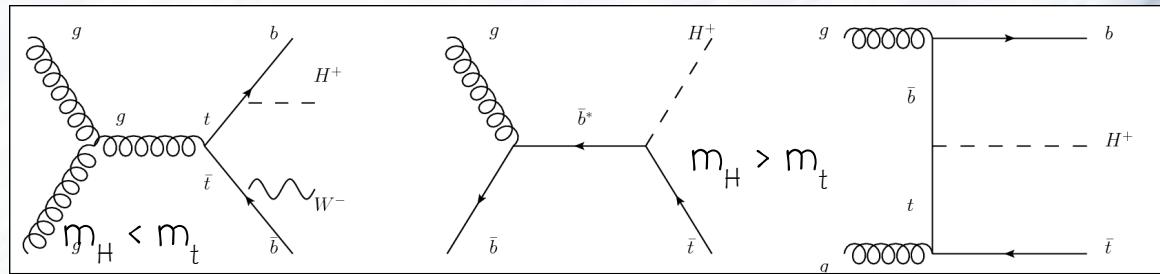


Model  
independent  
interpretation

reaching 1 fb at high-mass

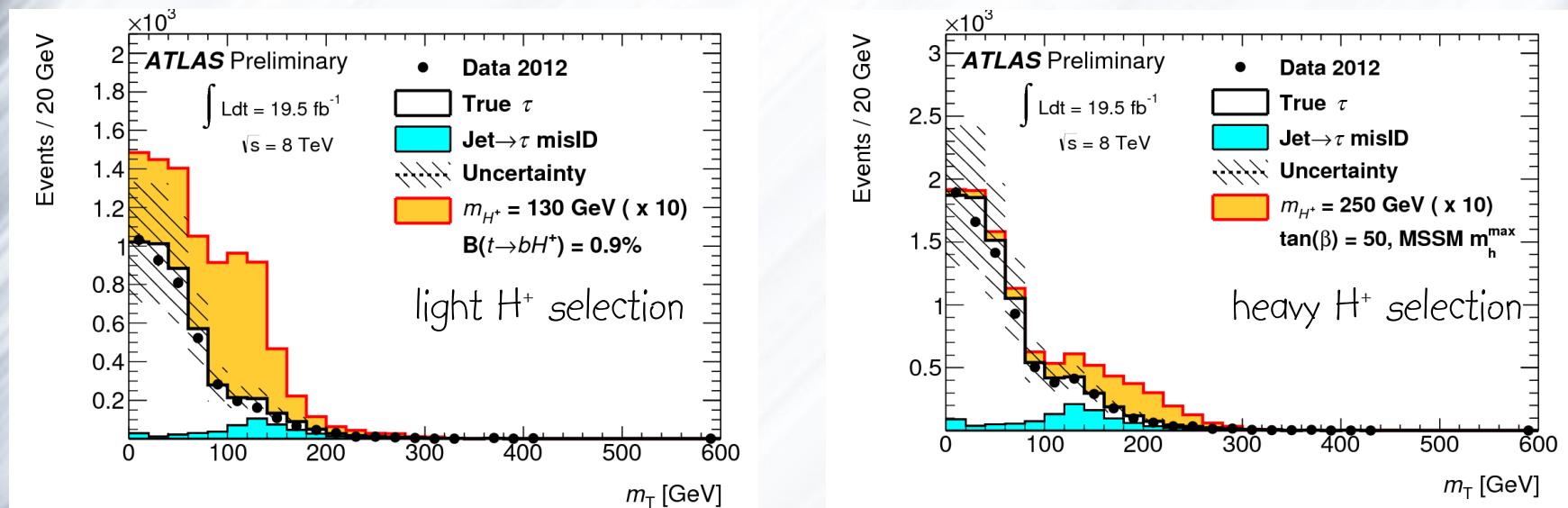
# Direct search for charged Higgs

- $H^+ \rightarrow T^+ \nu_T$
- tau (had) + jets channel



Measure :  $\text{BR}(t \rightarrow b H^+)$        $\sigma_{H^+}(\text{production cross section})$

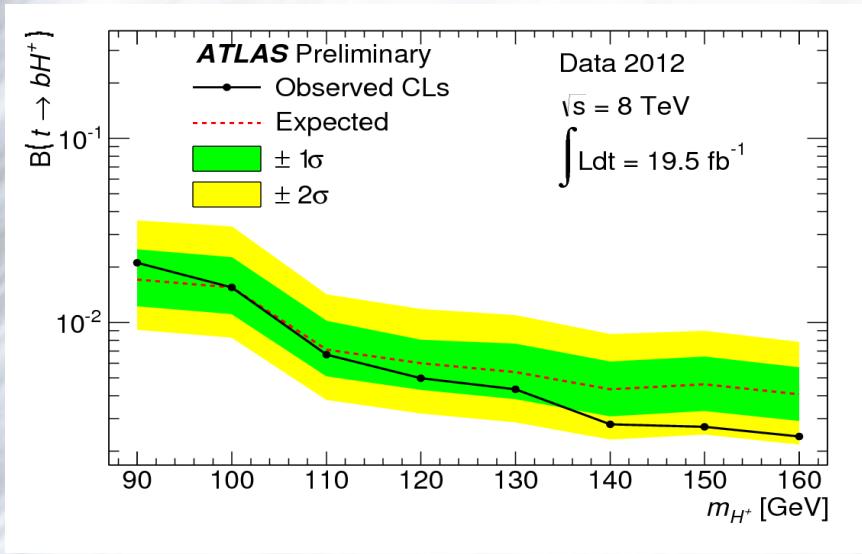
$$m_T = \sqrt{2 p_T^\tau E_T^{\text{miss}} (1 - \cos \Delta\Phi_{(\tau, \text{miss})})}$$



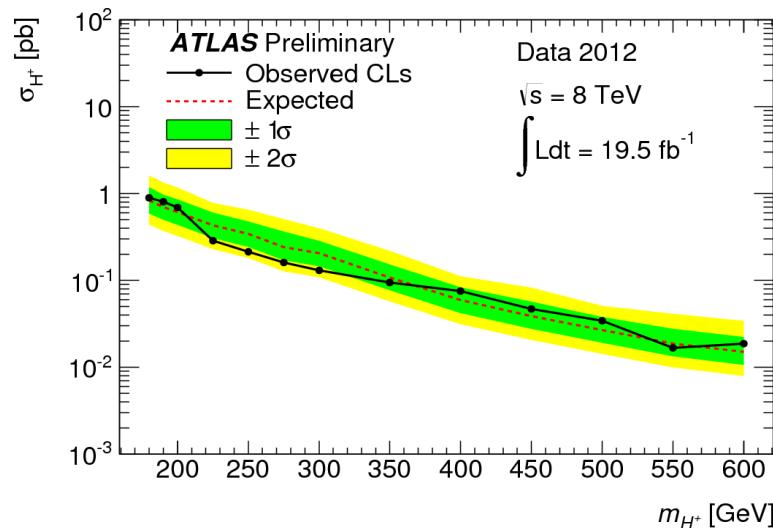
ATLAS-CONF-2013-090

# Direct search for charged Higgs

- interpreted in a model-independent way



$\text{BR}(t \rightarrow bH^+) < 1\%$  for a light charged Higgs



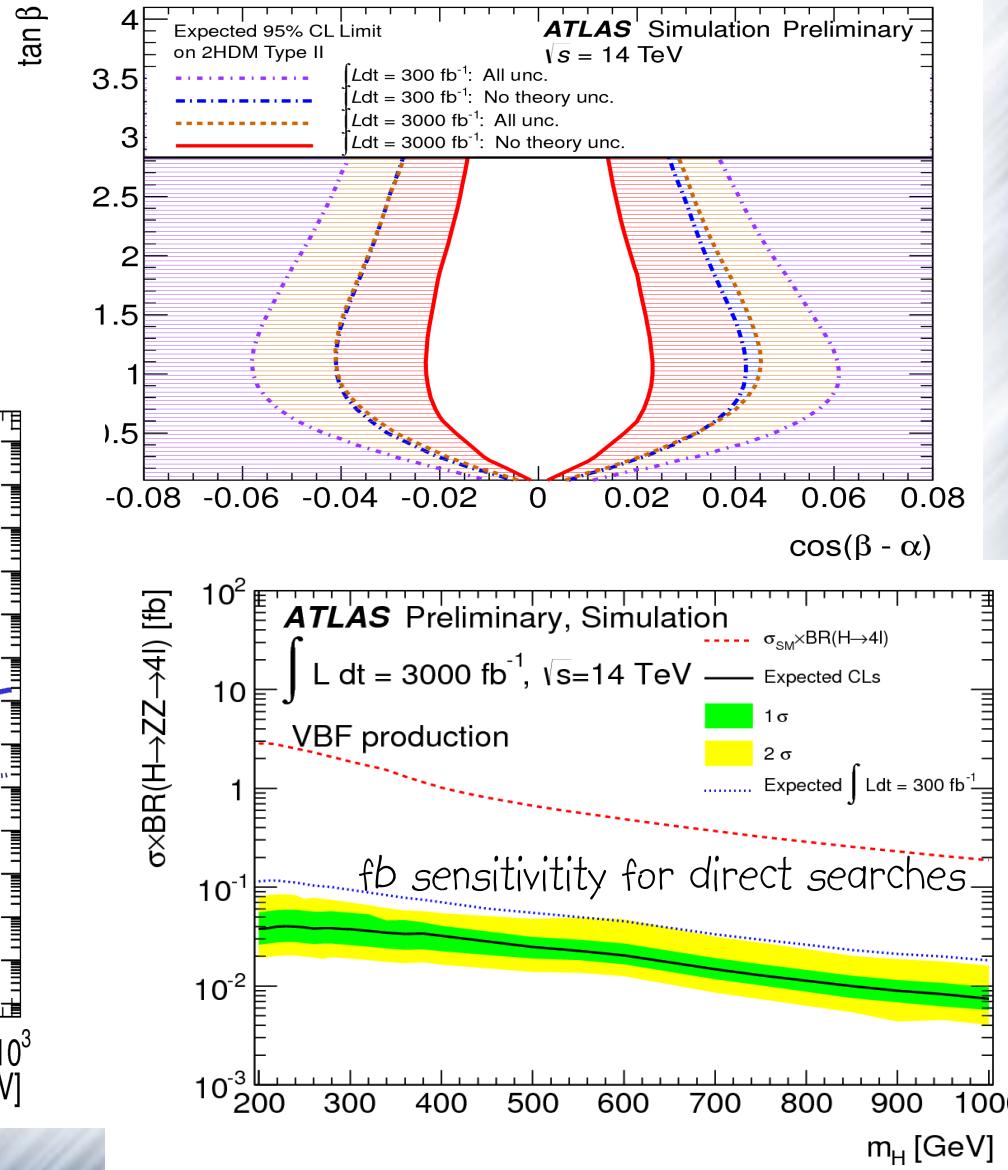
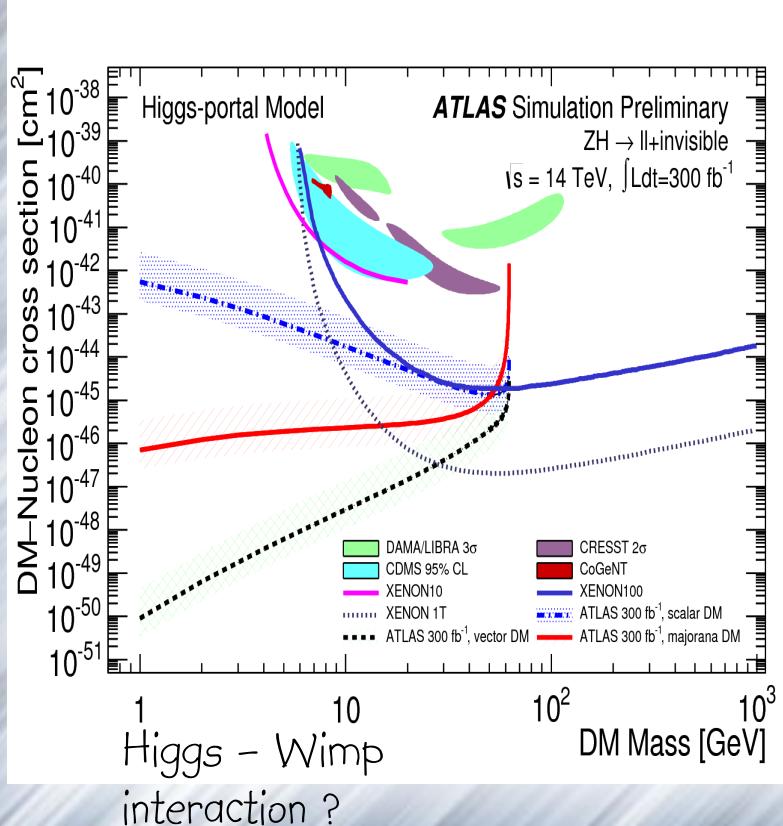
Production cross section of heavy charged Higgs  $< 1 \text{ pb}$

ATLAS-CONF-2013-090

# Run II prospects

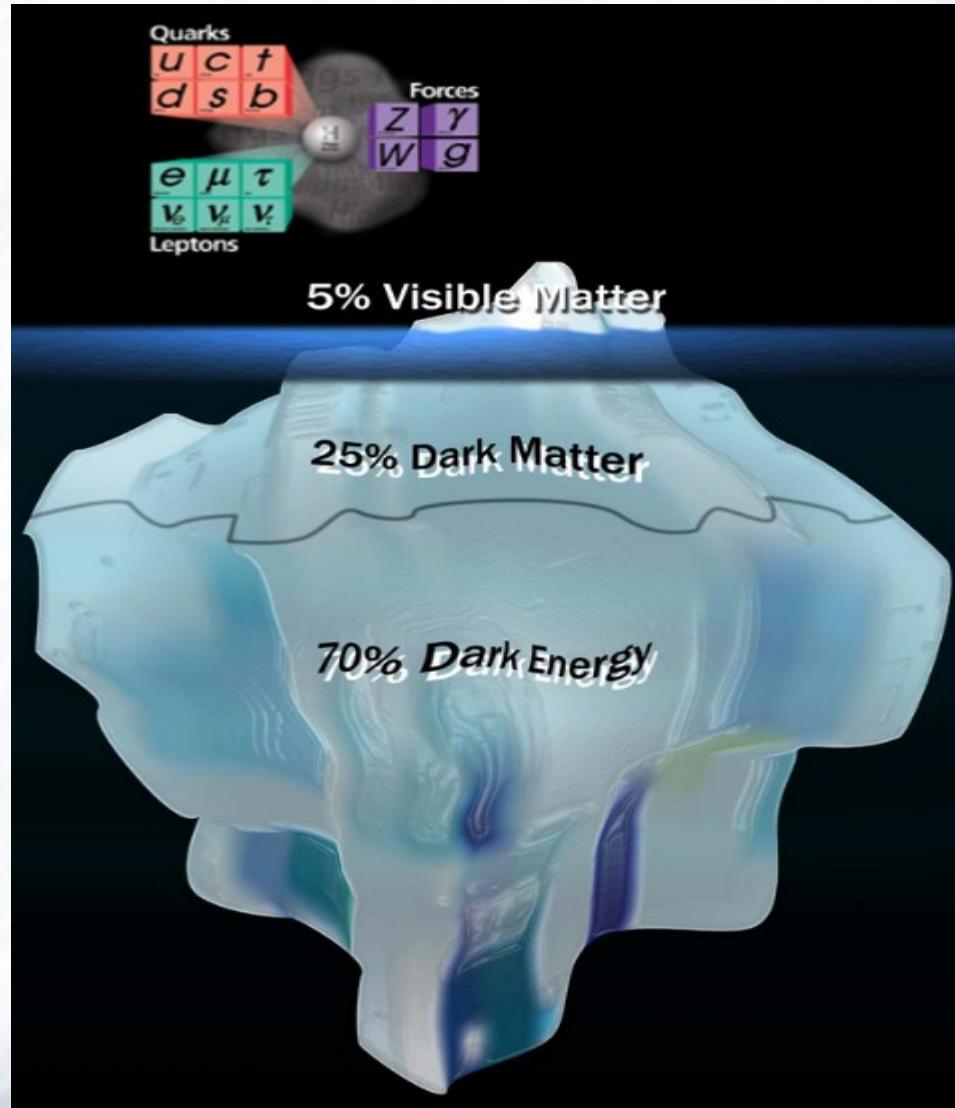
ATLAS-PHYS-PUB-2013-016  
 ATLAS-PHYS-PUB-2013-015  
 ATLAS-PHYS-PUB-2013-014

"Precision" coupling measurements



# Conclusion

- LHC Run I has promoted fundamental scalar fields to credible physical entities. It would be strange that nature has made the Higgs boson a unique child !
- Run II has the potential to help us bridge micro & macro physics.



# Backup slides

# Bibliography

- Limits on new phenomena via coupling measurements :  
ATLAS-CONF-2014-010 , CMS-PAS-HIG-13-005
- Direct searches :
  - CMS-PAS-HIG-14-001 , CMS-PAS-HIG-13-034 , CMS-PAS-HIG-13-025
  - CMS-PAS-HIG-13-032 , CMS-PAS-HIG-13-024 , CMS-PAS-HIG-13-010
  - ATLAS-CONF-2014-005 , ATLAS-CONF-2013-027 , ATLAS-CONF-2013-090
  - ATLAS-CONF-2012-079 , ATLAS-CONF-2012-013
  - PRL 108 (2012) 251801 , Phys. Lett. B721 (2013) 32, New J. Phys. 15(2013) 043009
  - Eur. Phys. J. C72 (2012) 2244 , JHEP02 (2013) 095
  - arXiv : 1312.1956 , 1210.5070 , 1207.2666

# direct search for H $\rightarrow$ hh & A $\rightarrow$ Zh

- $m_H > 2 m_h$  then H $\rightarrow$ hh ;  $2 m_h < m_A < 2 m_t$  then A $\rightarrow$ Zh predominantly
- h decays like in SM ; multilepton and diphoton signatures

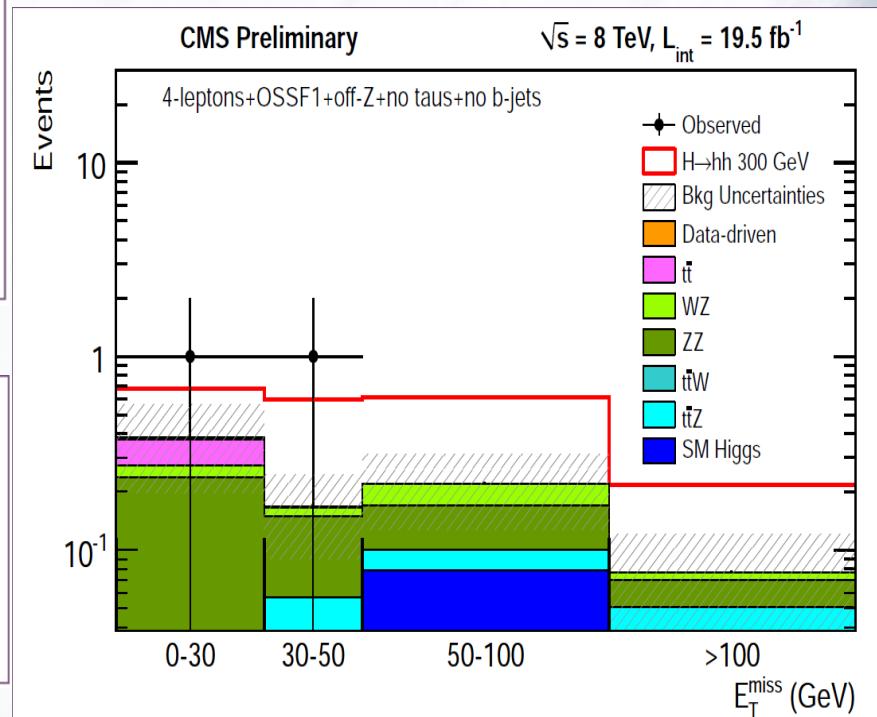
Leptons	Photons	OSSF pairs	Hadronic $\tau$	b-tag
4	0	0, 1 or 2	0 or 1	0 or 1
3	0	0 or 1	0 or 1	0 or 1
2	2	0 or 1	0	-
1	2	-	0	-
1	2	-	1	-
0	2	-	1 or 2	-

each channel further divided in bins of  $E_T^{\text{miss}}$

	$h \rightarrow WW^*$	$h \rightarrow ZZ^*$	$h \rightarrow \tau\tau$	$h \rightarrow bb$	$h \rightarrow \gamma\gamma$
$h \rightarrow WW^*$	✓	✓	✓	X	✓
$h \rightarrow ZZ^*$	-	✓	✓	✓	✓
$h \rightarrow \tau\tau$	-	-	✓	X	✓
$h \rightarrow bb$	-	-	-	X	X
$h \rightarrow \gamma\gamma$	-	-	-	-	X

✓ considered , X not considered

	$h \rightarrow WW^*$	$h \rightarrow ZZ^*$	$h \rightarrow \tau\tau$	$h \rightarrow \gamma\gamma$
$Z \rightarrow ll$	✓	✓	✓	✓
$Z \rightarrow qq$	X	✓	X	X
$Z \rightarrow \nu\nu$	X	✓	X	X



CMS-PAS-HIG-13-025

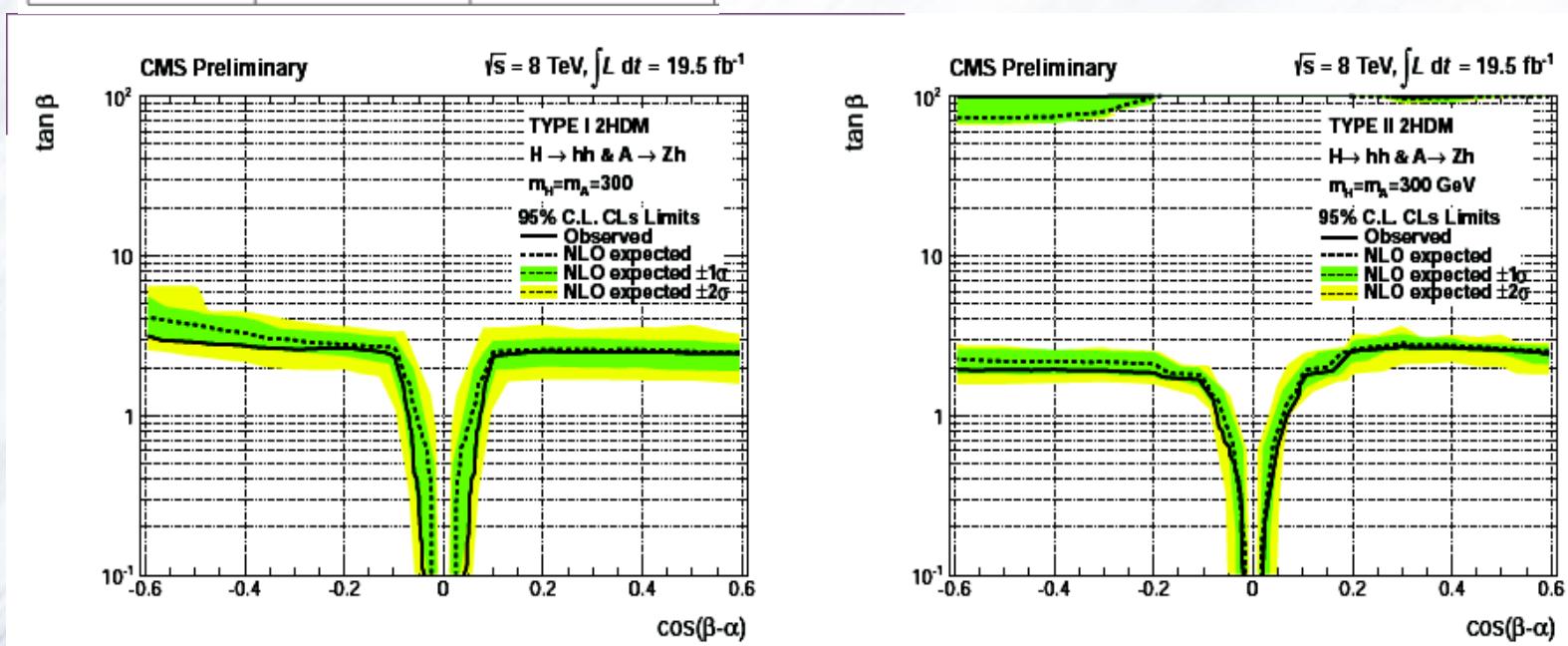
# direct search for $H \rightarrow hh$ & $A \rightarrow Zh$

$y_{\text{2HDM}}/y_{\text{SM}}$	2HDM I	2HDM II
$hVV$	$\sin(\beta - \alpha)$	$\sin(\beta - \alpha)$
$hQu$	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$
$hQd$	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$
$hLe$	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$
$HVV$	$\cos(\beta - \alpha)$	$\cos(\beta - \alpha)$
$HQu$	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$
$HQd$	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$
$HLe$	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$
$AVV$	0	0
$AQu$	$\cot \beta$	$\cot \beta$
$AQd$	$-\cot \beta$	$\tan \beta$
$ALe$	$-\cot \beta$	$\tan \beta$

CMS-PAS-HIG-13-025

N. Craig et al. : arXiv:1210.0559.

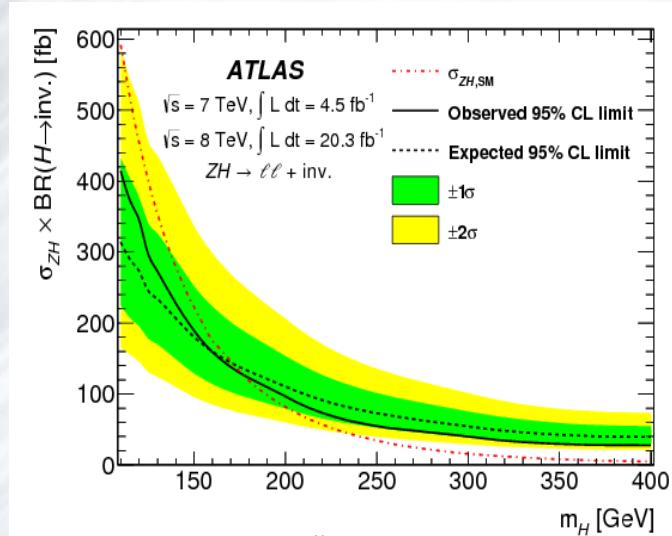
for  $m_H = m_A = 300$  GeV and type I & II 2HDMs



# Invisible decays of a BSM Higgs

Predicted branching ratio of SM Higgs invisible decay =  $1.3 \cdot 10^{-3}$  ( $h \rightarrow Z Z^* \rightarrow 4\nu$ )

Probing production of a Higgs boson produced in association with a  $Z$  that would predominantly decay into invisible particles :  $Z H \rightarrow 2 \text{ charged leptons} + E_T^{\text{miss}}$



when combined with all other observed decay channels : inv. BR of SM Higgs < 37% @ 95% CL

ATLAS-CONF-14-010 , arXiv : 1402.3244v2 , PRL 112(2014) 201802

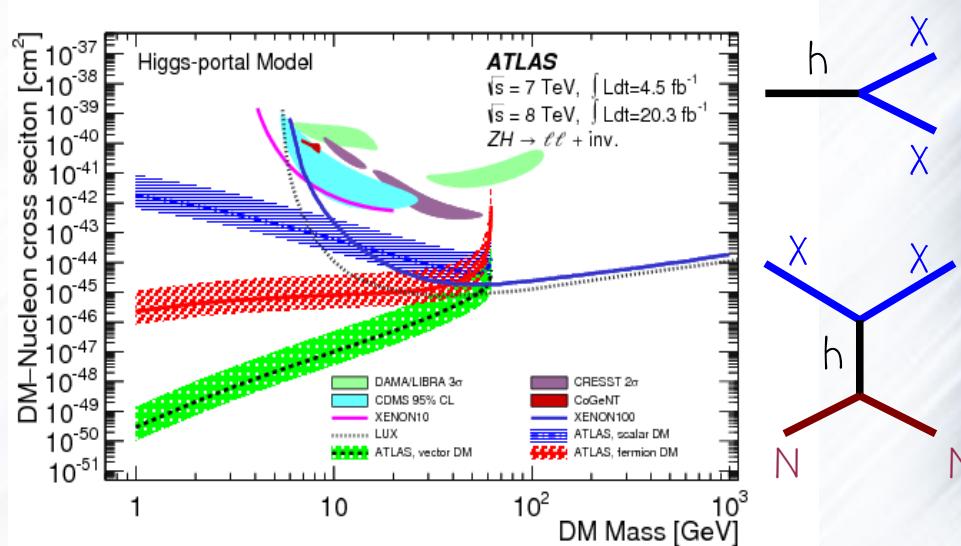
B. Patt & F. Wilczek : arXiv : hep-ph/0605188

S. Kanemura et al. : Phys.Rev. D82 (2010) 055026, arXiv : 1005.5661 [hep-ph]

P.J. Fox et al. : Phys. Rev. D85 (2012) 056011, arXiv : 1109.4398 [hep-ph]

A. Djouadi et al. : Phys. Lett. B709 (2012) 65, arXiv:1112.3288 [hep-ph]

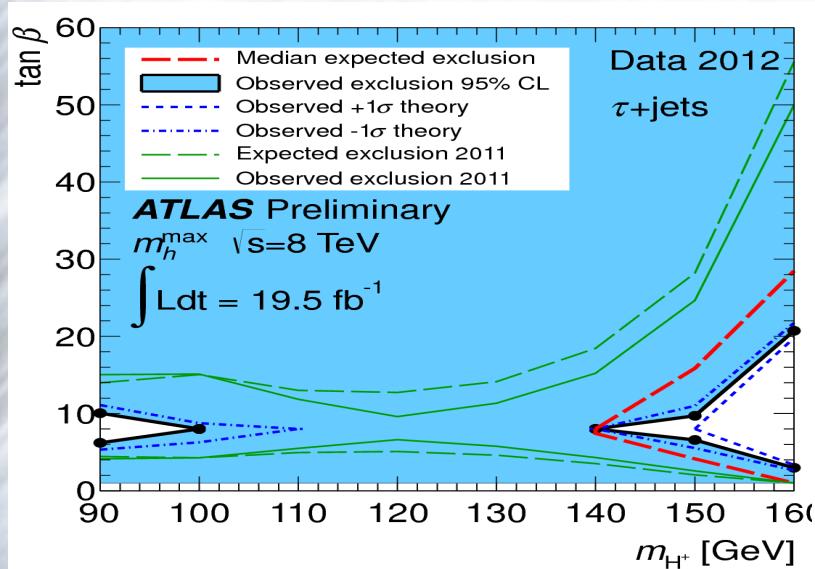
L. Lopez-Honorez et al. : Phys. Lett. B716 (2012) 179, arXiv:1203.2064 [hep-ph]



Interpreted in DM Higgs portal models in which wimps very weakly interact with SM particles except for the Higgs boson. In these models, wimps would predominantly interact with matter by Higgs boson exchange.

# Direct search for charged Higgs

- interpreted in MSSM  $m_h^{\max}$  scenario



Existence of a MSSM light charged Higgs almost ruled out, but still plenty of room for a heavy one (run II).

ATLAS-CONF-2013-090

M. Carena et al., arXiv : 1302.7033

